

¹ Observation of the Higgs boson in the WW^*
² channel and search for Higgs boson pair
³ production in the $b\bar{b}b\bar{b}$ channel with the
⁴ ATLAS detector

⁵ A DISSERTATION PRESENTED
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²⁰ **Observation of the Higgs boson in the WW^* channel and search
²¹ for Higgs boson pair production in the $b\bar{b}b\bar{b}$ channel with the
²² ATLAS detector**

²³ ABSTRACT

²⁴ This dissertation presents the observation and measurement of the Higgs boson in the $H \rightarrow WW^* \rightarrow$
²⁵ $\ell\nu\ell\nu$ channel at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 8$ TeV and a search for Higgs pair production in the $HH \rightarrow$
²⁶ $b\bar{b}b\bar{b}$ channel at $\sqrt{s} = 13$ TeV with the ATLAS detector in pp collisions at the Large Hadron Collider.

²⁷ First, the discovery of a particle consistent with the Higgs boson in 4.8 fb^{-1} at $\sqrt{s} = 7$ TeV and
²⁸ 5.8 fb^{-1} at $\sqrt{s} = 8$ TeV is discussed. Then, the measurement of the Higgs boson signal strength
²⁹ and cross section in both the gluon fusion and vector boson fusion (VBF) production modes using
³⁰ 20.3 fb^{-1} of $\sqrt{s} = 8$ TeV data combined with 4.8 fb^{-1} of 7 TeV data is shown. The combined signal
³¹ strength is measured to be $\mu = 1.09^{+0.23}_{-0.21}$. The total observed significance of the $H \rightarrow WW^*$ process
³² is observed to be 6.1σ (with 5.8σ expected). Advanced methods for background reduction and estima-
³³ tion, particularly in same-flavor lepton final states, are shown. The VBF signal strength is measured to
³⁴ be $\mu_{\text{VBF}} = 1.27^{+0.53}_{-0.45}$ with an observed significance of 3.2σ (with 2.7σ expected). In the VBF chan-
³⁵ nel, a selection requirement based method, the precursor to the final multivariate technique used for the
³⁶ result, is detailed.

³⁷ Finally, a search for Higgs pair production in the $b\bar{b}b\bar{b}$ final state with 3.2 fb^{-1} at $\sqrt{s} = 13$ TeV is
³⁸ presented. A particular focus is placed on a tailored signal region for resonant production of Higgs pairs
³⁹ at high masses. No significant excesses are observed, and upper limits on cross sections are placed for
⁴⁰ spin-2 Randall Sundrum gravitons (RSG) and narrow spin-0 resonances. The cross section of $\sigma(pp \rightarrow$
⁴¹ $G_{\text{KK}}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $k/\bar{M}_{\text{Pl}} = 1$ is constrained to be less than 70 fb for masses in the range
⁴² $600 < m_{G_{\text{KK}}^*} < 3000$ GeV. The cross section upper limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges
⁴³ from 30 to 300 fb in the mass range of $500 < m_H < 3000$ GeV.

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0

498

499

Introduction

500 The Higgs boson is often described as one of the cornerstones of particle physics. When the Standard
501 Model was first developed as a theory to describe the fundamental particles and forces of nature, physicists
502 were faced with a dilemma. The electroweak theory beautifully characterized both electromagnetism and
503 the weak force with a single underlying framework. However, the mass of the weak W and Z bosons
504 was puzzling given the fact that their electromagnetic counterpart, the photon, is massless. The Higgs
505 mechanism was developed as the leading theory for the origin of this electroweak symmetry breaking. It
506 predicted the existence of an additional spin-0 boson in the Standard Model, the Higgs boson. Generations
507 of collider experiments searched for this elusive particle. This dissertation presents research work on the
508 Higgs boson from its discovery to its use as a tool in the search for physics beyond the Standard Model
509 with the ATLAS detector at the Large Hadron Collider (LHC).

510 One of the first priorities for the LHC when it began colliding proton beams in 2010 was the search
511 for the Higgs boson. This search was initially tackled in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel, followed by

512 the $H \rightarrow \gamma\gamma$ and $H \rightarrow ZZ^* \rightarrow 4\ell$ channels. Each channel has its own merits, but the WW^* mode is
513 particularly suited to searching over a wide range of masses. The $H \rightarrow WW^*$ branching ratio is large and it
514 is the primary decay channel above the $2m_W$ mass threshold. Despite the fact that the full Higgs invariant
515 mass cannot be reconstructed in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel, its signal to background ratio makes
516 it ideal for measurement of Higgs properties such as the production cross section and couplings.

517 In 2012, the ATLAS and CMS experiments announced the discovery of a new particle consistent with
518 the Higgs boson [1, 2]. In ATLAS, this discovery was made with 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$
519 and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. This dissertation first presents the search for gluon fusion production
520 of the Higgs in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel, which played an important role in this discovery.
521 Selection requirements which were optimized to maximize the discovery significance in this channel, as
522 well as background estimation procedures, are discussed.

523 After its discovery, interest in the Higgs shifted to focus on the measurement of its properties. As a result,
524 extensions of the initial discovery analysis in larger datasets had two main goals. Improvement of signal to
525 background ratio was important to allow for precision measurements. Also, searches for rarer production
526 modes of the Higgs were a priority. The first such extension presented in this dissertation is a tailored
527 selection for $\ell\nu\ell\nu$ final states with same flavor leptons. Novel variables for the reduction of the $Z+\text{jets}$
528 background that could remain robust under increasing LHC instantaneous luminosities are shown. The
529 second post-discovery result shown is the first evidence of Vector Boson Fusion (VBF) production of the
530 Higgs boson.

531 VBF production of the Higgs boson is particularly interesting in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ final state.
532 In this combination of production and decay modes, the Higgs boson couples exclusively to W bosons,
533 allowing for precise measurement of the Higgs- W coupling constant. However, it is challenging to observe
534 VBF Higgs production because its cross section at the LHC is an order of magnitude lower than gluon
535 fusion production. The large $H \rightarrow WW^*$ branching ratio thus presents another advantage over other
536 final states. Additionally, VBF production of the Higgs boson creates two forward jets in addition to the
537 Higgs, and these jets can be used to isolate VBF Higgs events from other production modes. The VBF
538 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis first created a selection requirement based signal region using variables

539 constructed specifically for the VBF Higgs production topology. This “cut-based” analysis is presented
540 in detail in this dissertation. These VBF topology variables, once validated in the cut-based analysis, were
541 then input into a multivariate boosted decision tree discriminant to achieve the first evidence of VBF Higgs
542 production with the full 20.3 fb^{-1} of $\sqrt{s} = 8 \text{ TeV}$ data in ATLAS.

543 After a two year shutdown, the LHC restarted in 2015 with a center of mass energy of $\sqrt{s} = 13 \text{ TeV}$.
544 This increase improved the LHC’s ability to probe for physics beyond the Standard Model, and the Higgs
545 sector remained one of the largest regions of unprobed phase space where such new physics could be dis-
546 covered. Production of high mass resonances benefited most from the center of mass energy increase. In
547 particular, the cross section for a generic gluon-initiated 2 TeV resonance increased tenfold with the in-
548 crease from 8 to 13 TeV. Therefore, a natural next step in studies of the Higgs was a search for a new
549 heavy resonance which decays into a pair of Higgs bosons. The final result shown in this dissertation is
550 a search for resonant di-Higgs production in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ final state with 3.2 fb^{-1} recorded
551 by ATLAS at $\sqrt{s} = 13 \text{ TeV}$. This search has the unique advantage that it can both probe new physics
552 and gain further understanding of the Higgs potential through constraints on SM pair production of the
553 Higgs. It also extends the previous ATLAS results at $\sqrt{s} = 8 \text{ TeV}$ and probes higher mass resonances
554 that were not previously accessible. Additionally, it is an informative precursor to di-Higgs analyses at the
555 future High Luminosity LHC (HL-LHC), where a projected dataset of 3000 fb^{-1} at $\sqrt{s} = 14 \text{ TeV}$ will
556 begin to become sensitive to the SM Higgs self coupling.

557 As mentioned above, this dissertation begins by discussing the discovery of the Higgs and the role of
558 the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel. It then presents the first evidence for the VBF production mode using
559 the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel with the full ATLAS Run 1 dataset. It also shows the final combined
560 Run 1 measurements of gluon fusion Higgs production from this channel. Finally, it presents a search for
561 Higgs pair production in the $HH \rightarrow b\bar{b}b\bar{b}$ channel. It is organized into four parts.

562 Part 1 presents the theoretical and experimental background required for the subsequent parts. Chap-
563 ter 1 gives an overview of Higgs physics, particularly single and double Higgs production in the Standard
564 Model and beyond. Chapter 2 presents details regarding the Large Hadron Collider and the ATLAS experi-
565 ment. The evolution of machine conditions, descriptions of the ATLAS sub-detectors, and an overview of

object reconstruction in ATLAS are all shown. A brief interlude on the ATLAS Muon New Small Wheel upgrade is also given, as this upgrade has been a focus of my graduate work and will have an important impact on ATLAS' ability to study the Higgs at the High Luminosity LHC.

Part 2 discusses the observation and measurement of the Higgs in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel in the ATLAS Run 1 dataset at $\sqrt{s} = 7$ and 8 TeV. Because I worked in this channel from before the discovery through to the final analysis of the Run 1 dataset, Part 2 is organized in such a way to allow easy presentation of multiple analyses on different subsets of the full Run 1 dataset. Chapter 3 presents a general overview of the $H \rightarrow WW^*$ analysis strategy and defines many of the variables and common elements used in the rest of Part 2. Chapter 4 presents the discovery and subsequent measurements of the Higgs boson, focusing on the role of the WW^* channel in this discovery. Chapter 5 presents the first evidence for the VBF production mode of the Higgs, a result from the WW^* channel in the full Run 1 ATLAS dataset. In this chapter, the focus is mainly on the cut-based VBF analysis. The cut-based analysis was an important first step to the final VBF result which used a boosted decision tree. Where appropriate, connections between the cut-based and BDT analyses are shown and their compatibility is discussed. Finally, the VBF analysis was an important input into the combined Run 1 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ result, which used both the gluon fusion and VBF channels in a combined fit to infer properties of the Higgs, including its couplings to the gauge bosons and its production cross section. This is the topic of Chapter 6.

Part 3 presents a search for Higgs pair production in the $HH \rightarrow b\bar{b}b\bar{b}$ channel. Chapter 7 presents an overview of this search in the boosted regime, where the Higgs pairs are the result of the decay of a heavy resonance. Chapter 8 shows the combined results between the boosted regime and the resolved regime, which is sensitive to lower mass resonances and non-resonant Higgs pair production. Finally, Part 4 presents a conclusion and brief outlook of future Higgs physics with ATLAS.

Part I

589

590

Theoretical and Experimental Background

In modern physics, there is no such thing as “nothing.”

Richard Morris

1

591

592

The Physics of the Higgs Boson

593 This chapter presents an overview of the Standard Model of Particle Physics and in particular the physics
594 of the Higgs boson. First, a brief overview of the Standard Model is presented. Then, a description of the
595 Higgs mechanism of electroweak symmetry breaking is given. Next, the physics of single Higgs boson
596 production and decay is described. The Standard Model also allows for production of two Higgs bosons
597 and this is detailed as well. Finally, di-Higgs production in two beyond the Standard Model (BSM) theories
598 - Randall-Sundrum gravitons (RSG) and Two Higgs Doublet Models (2HDM) - is shown.

599 *I.I THE STANDARD MODEL OF PARTICLE PHYSICS*

600 The Standard Model (SM) of Particle Physics is a quantum field theory describing the fundamental
601 particles of nature and the forces that govern their interactions. Several comprehensive pedagogical treat-
602 ments of the SM already exist in the literature [3–8] and this section will not rehash those. Rather, this
603 section presents a brief overview of the SM particles and forces in order to define them for subsequent

604 discussions.

605 The Standard Model consists of two primary categories of fundamental particles: fermions (spin 1/2
606 particles) and bosons (integer spin particles). The SM also describes three forces: electromagnetism, the
607 weak nuclear force, and the strong nuclear force. Gravity is not included in the theory and is largely irrele-
608 evant at the scales currently probed by collider experiments. Within the fermions, there are both quarks
609 (which interact via all three forces) and leptons. The charged leptons interact via electromagnetic and weak
610 interactions, while neutrinos (neutral leptons) interact only via the weak force. Within the bosons, there
611 are the W^\pm and Z bosons (the mediators of the weak force), the gluon (g , the mediator of the strong
612 force), and the photon (γ , the mediator of the electromagnetic force). Finally, there is the Higgs boson,
613 a fundamental spin zero particle resulting from the Higgs mechanism of electroweak symmetry breaking.

614 Figure 1.1 summarizes the fermions and bosons of the SM.

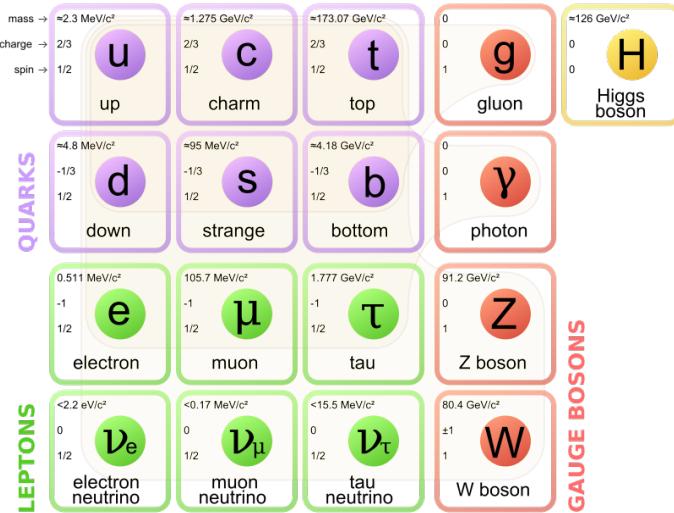


Figure 1.1: The particles of the Standard Model and their properties [6].

615 The Standard Model coalesced into a unified theoretical framework in the 1960s through the work
616 of Glashow, Weinberg, Salam, and others on the theory of electroweak interactions [9–12]. This theory
617 characterized both the electromagnetic and weak interactions as unified under a single gauge symmetry
618 group, namely $SU(2) \times U(1)$. At low enough energy scales (on the order of the W and Z masses), the
619 electroweak symmetry is broken, as evidenced by the fact that the weak bosons have mass while the photon
620 does not. The discovery of the Higgs boson in 2012 confirmed the Higgs mechanism as the most likely

621 candidate for this electroweak symmetry breaking [1, 2]. The complete SM consists of this electroweak
 622 theory combined with the theory of quantum chromodynamics (which models the strong sector as a non-
 623 Abelian $SU(3)$ gauge group)¹.

624 **I.2 ELECTROWEAK SYMMETRY BREAKING AND THE HIGGS**

625 In the Standard Model Lagrangian, it is difficult to include mass terms for the W and Z bosons without
 626 breaking the fundamental gauge symmetry of the Lagrangian. A traditional mass term does not preserve
 627 the $SU(2) \times U(1)$ symmetry. Additionally, scattering of massive W and Z bosons violate unitarity and
 628 these diagrams diverge at high energy scales. In the 1960s, Higgs, Brout, Englert, Guralnik, Kibble, and
 629 Hagen developed a mechanism for spontaneous symmetry breaking via the addition of a complex scalar
 630 doublet to the SM. Three of the four real degrees of freedom of this complex field would go to the lon-
 631 gitudinal modes of the W^\pm and Z , thus allowing them to have mass [14–17]. The remaining degree of
 632 freedom would manifest as an additional scalar, known now as the Higgs boson.

633 The mechanism works by introducing a Lagrangian for the newly introduced field that still respects the
 634 symmetry of the Standard Model inherently, but with a minimum at a non-zero vacuum expectation value
 635 for the field. In this minimum of the potential, the electroweak symmetry is broken. Specifically, consider
 636 a complex scalar doublet Φ with four degrees of freedom, as shown in equation I.1.

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1^+ + i\phi_2^+ \\ \phi_1^0 + i\phi_2^0 \end{pmatrix} \quad (\text{I.1})$$

637 The simplest potential of a self-interacting Higgs that still respects the SM symmetry is given in equa-
 638 tion I.2.

$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \quad (\text{I.2})$$

639 If the μ^2 term of this potential is positive, then the potential has a minimum at $\Phi = 0$ and the electroweak

¹For a pedagogical treatment of the physics of quantum chromodynamics, see reference [13].

⁶⁴⁰ symmetry is preserved. However, if instead $\mu^2 < 0$, then the minimum is at a finite value of Φ , namely

$$\Phi_{\min} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \quad (1.3)$$

⁶⁴¹ where $v = \sqrt{\mu^2/\lambda}$. Because this is the location of the minimum, it corresponds to the vacuum expecta-
⁶⁴² tion value for the field ($\langle \Phi \rangle = \Phi_{\min}$). The excitations of the Higgs can then be parameterized as

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix} \quad (1.4)$$

⁶⁴³ The full scalar Lagrangian, including the kinetic term, is then given as

$$\mathcal{L}_s = (D^\mu \Phi)^\dagger (D_\mu \Phi) - V(\Phi) \quad (1.5)$$

⁶⁴⁴ where the covariant derivative is defined as

$$D_\mu = \partial_\mu + \frac{ig}{2} \tau^a W_\mu^a + ig' Y B_\mu \quad (1.6)$$

⁶⁴⁵ and W^1, W^2, W^3 and B are the $SU(2)$ and $U(1)$ gauge fields of the electroweak theory, respectively. g
⁶⁴⁶ and g' are the corresponding coupling constants. The Pauli matrices are represented with τ . With the
⁶⁴⁷ scalar Lagrangian in place, the physical gauge fields can then be written as

$$W_\mu^\pm = \frac{1}{\sqrt{2}} (W_\mu^1 \mp i W_\mu^2) \quad (1.7)$$

⁶⁴⁸

$$Z_\mu = \frac{-g' B_\mu + g W_\mu^3}{\sqrt{g^2 + g'^2}} \quad (1.8)$$

⁶⁴⁹

$$A_\mu = \frac{g B_\mu + g' W_\mu^3}{\sqrt{g^2 + g'^2}} \quad (1.9)$$

650 Equation 1.7 corresponds to the charged W^+ and W^- bosons, equation 1.8 corresponds to the neutral Z
 651 boson, and equation 1.9 corresponds to the neutral photon. The masses of the particles also arise from the
 652 Lagrangian. The photon has zero mass, while the masses of the W and Z bosons are given in equation 1.10.

653

$$\begin{aligned} M_W^2 &= \frac{1}{4}g^2v^2 \\ M_Z^2 &= \frac{1}{4}(g^2 + g'^2)v^2 \end{aligned} \quad (1.10)$$

654 The fermion masses also arise through a coupling with the Higgs via the Yukawa interaction (for a detailed
 655 description, see [8]). In this case the coupling between the Higgs and the fermions goes as

$$g_{hf\bar{f}} = \frac{m_f}{v} \quad (1.11)$$

656 The full Lagrangian of Higgs interactions can be written as

$$\mathcal{L}_{\text{Higgs}} = -g_{hf\bar{f}}\bar{f}fh + \frac{g_{hhh}}{6}h^3 + \frac{g_{hhhh}}{24}h^4 + \delta_V V_\mu V^\mu \left(g_{hVV}H + \frac{g_{hhVV}}{2}h^2 \right) \quad (1.12)$$

657 with

$$\begin{aligned} g_{hVV} &= \frac{2m_V^2}{v} & g_{hhVV} &= \frac{2m_V^2}{v^2} \\ g_{hhh} &= \frac{3m_h^2}{v} & g_{hhHH} &= \frac{3m_h^2}{v^2} \end{aligned} \quad (1.13)$$

658 The last term of the Lagrangian appears twice, once for W bosons and once for Z bosons. V refers to
 659 the W^\pm and Z , and $\delta_W = 1$ while $\delta_Z = 1/2$. Phenomenologically, there are a few features of this
 660 Lagrangian that are useful to note. First, note that the Higgs mass is a free parameter of the theory that
 661 must be determined experimentally. Second, note that the coupling of the Higgs to the vector bosons and
 662 fermions scales as a function of the masses of these particles, a fact that is important when considering
 663 both the production and decays of the Higgs. Finally, note the presence of the cubic and quartic Higgs self
 664 interaction terms, which can lead to final states with multiple Higgs bosons produced.

665 1.3 HIGGS BOSON PRODUCTION AND DECAY

666 This section discusses the properties of Higgs production and decay mechanisms. The details presented
667 here will focus on the properties of a 125 GeV Higgs boson, as this is the mass closest to that of the newly
668 discovered Higgs.

669 1.3.1 HIGGS PRODUCTION

670 The Higgs is produced by four main production modes at the Large Hadron Collider - gluon-gluon
671 fusion (ggF), vector boson fusion (VBF), associated production with a W or Z boson, or associated pro-
672 duction with top quarks ($t\bar{t}H$). Figure 1.2 shows the Feynman diagrams for these four modes.

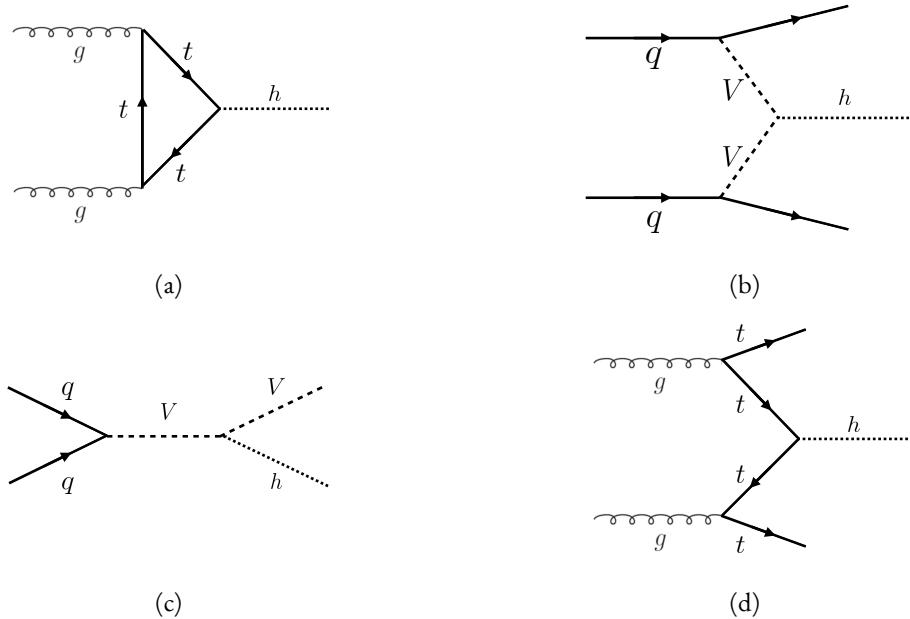


Figure 1.2: The four most common Higgs boson production modes at the LHC: (a) gluon-gluon fusion, (b) vector boson fusion, (c) $W/Z + H$ production, (d) $t\bar{t}H$ production

673 In gluon-gluon fusion, gluons from the incoming protons fuse via a top-quark loop to produce a Higgs.
674 The top quark is the dominant contribution in the loop due to its heavy mass and the fact that the Higgs-
675 fermion coupling constant scales with fermion mass. In vector boson fusion, the incoming quarks each
676 radiate a W or Z boson which fuse to produce the Higgs. This production mode results in a final state
677 with a Higgs boson and two additional jets which tend to be forward because they carry the longitudinal

678 momentum of the incoming partons. The Higgs can also be produced in association with a W or Z boson.
 679 The W/Z is produced normally and then radiates a Higgs². Finally, the Higgs can be produced in associa-
 680 tion with two top quarks. Each incoming gluon splits into a $t\bar{t}$ pair, and one of the top pairs combines to
 681 create a Higgs. Figure 1.3 shows the production cross section for a 125 GeV Higgs boson in each of these
 modes at a pp collider as a function of center of mass energy. In figure 1.3, note that gluon fusion has the

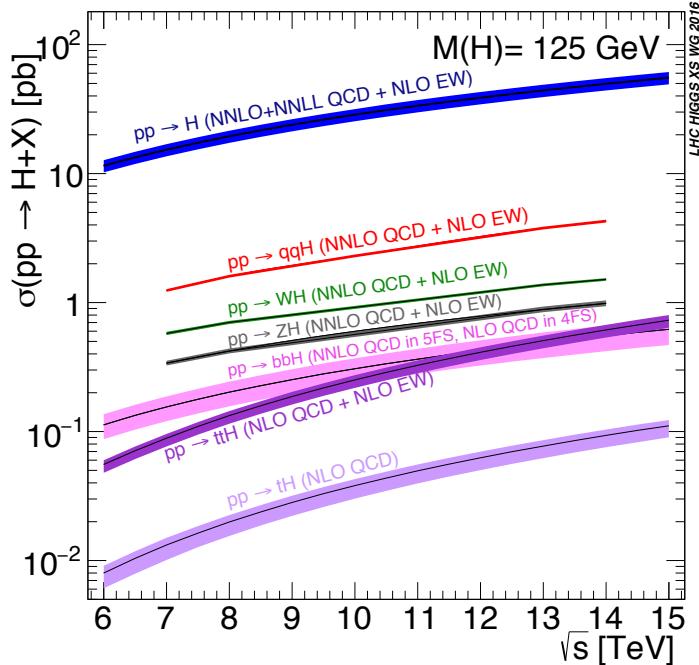


Figure 1.3: Higgs production cross sections as a function of center of mass energy (\sqrt{s}) at a pp collider [18].

682
 683 largest cross section, while VBF is the second largest at approximately a factor of 10 smaller. The figure also
 684 includes the less commonly studied $b\bar{b}H$ and tH modes. While the $b\bar{b}H$ mode has a larger cross section
 685 than $t\bar{t}H$, it also has larger backgrounds and is thus less sensitive. The tH mode is not as sensitive as $t\bar{t}H$
 686 due to its lower cross section. At $\sqrt{s} = 8$ TeV, ggF production of a 125 GeV Higgs has a cross section of
 687 $19.47^{+1.54}_{-1.67}$ pb, while VBF has a cross section of $1.601^{+0.036}_{-0.035}$ pb [18]. The cross sections of all of the main
 688 Higgs production modes at this center of mass energy, as well as their uncertainties from varying the QCD
 689 renormalization and factorization scales and PDFs, are summarized in table 1.1 for a 125 GeV Higgs. The
 690 relative uncertainty of the gluon fusion mode is larger than the relative uncertainty in the vector boson

²This mode is also sometimes known as “Higgs-strahlung”.

⁶⁹¹ fusion mode due to the fact that gluon fusion production happens through a loop.

Production mode	σ (pb)	QCD scale uncert. (%)	PDF + α_s uncert. (%)
Gluon fusion	19.47	+7.3 / - 8.0	3.1
Vector boson fusion	1.601	+0.3 / - 0.2	2.2
WH	0.7026	+0.6 / - 0.9	2.0
ZH	0.4208	+2.9 / - 2.4	1.7
bbH	0.2021	+20.7 / - 22.3	
$t\bar{t}H$	0.1330	+4.1 / - 9.2	4.3
tH (t -channel)	0.01869	+7.3 / - 16.5	4.6
tH (s -channel)	1.214×10^{-3}	+2.8 / - 2.4	2.8

Table 1.1: Production cross sections for a 125 GeV Higgs boson at $\sqrt{s} = 8$ TeV with scale and PDF uncertainties [18].

⁶⁹² **1.3.2 HIGGS BRANCHING RATIOS**

⁶⁹³ The fact that the Higgs couples more strongly to more massive particles is crucial for understanding its
⁶⁹⁴ branching ratios. The width for Higgs decays to fermions is given by equation 1.14 [5].

$$\Gamma(H \rightarrow f\bar{f}) = \frac{N_c \sqrt{2} G_F m_f^2 m_H}{8\pi} \quad (1.14)$$

⁶⁹⁵ In this case, N_c is the number of colors, G_F is the Fermi constant, m_f is the mass of the fermion, and
⁶⁹⁶ m_H is the mass of the Higgs. Note that the width scales with the square of the fermion mass. (This also
⁶⁹⁷ assumes that the Higgs mass is large enough to decay with both the fermions on shell.)

⁶⁹⁸ The decay width to WW , in the case where both W bosons are produced on shell ($m_H \geq 2m_W$), is
⁶⁹⁹ given in equation 1.15 [5].

$$\Gamma(H \rightarrow W^+ W^-) = \frac{\sqrt{2} G_F M_W^2 m_H}{16\pi} \frac{\sqrt{1-x_W}}{x_W} (3x_W^2 - 4x_W + 4) \quad (1.15)$$

⁷⁰⁰ where m_W is the mass of the W and $x_W = 4M_W^2/m_H^2$. To get the branching ratio to ZZ (in the regime
⁷⁰¹ where $m_H \geq 2m_Z$), the equation is divided by 2 to account for identical particles in the final state, and

⁷⁰² x_W is replaced with $x_Z = 4M_Z^2/m_H^2$. This is shown in equation 1.16 [5].

$$\Gamma(H \rightarrow ZZ) = \frac{\sqrt{2}G_F M_Z^2 m_H}{32\pi} \frac{\sqrt{1-x_Z}}{x_Z} (3x_Z^2 - 4x_Z + 4) \quad (1.16)$$

⁷⁰³ The more general formula for Higgs branching into WW or ZZ , taking into account the case where one
⁷⁰⁴ or both vector bosons is off-shell, is shown in equation 1.17 [19].

$$\Gamma(H \rightarrow V^*V^*) = \frac{1}{\pi^2} \int_0^{M_H^2} \frac{dq_1^2 M_V \Gamma_V}{(q_1^2 - M_V^2)^2 + M_V^2 \Gamma_V^2} \int_0^{(M_H - q_1)^2} \frac{dq_2^2 M_V \Gamma_V}{(q_2^2 - M_V^2)^2 + M_V^2 \Gamma_V^2} \Gamma_0 \quad (1.17)$$

⁷⁰⁵ Here, q_1^2 and q_2^2 are the invariant masses of the virtual gauge bosons, M_V is the W or Z mass, and Γ_V is
⁷⁰⁶ the W or Z width. Γ_0 is the squared matrix element, which is given in equation 1.18 [19].

$$\Gamma_0 = \frac{G_F M_H^3}{8\sqrt{2}\pi} \delta_V \sqrt{\lambda(q_1^2, q_2^2, M_H^2)} \left[\lambda(q_1^2, q_2^2, M_H^2) + \frac{12q_1^2 q_2^2}{M_H^4} \right] \quad (1.18)$$

⁷⁰⁷ The function λ is defined as $\lambda(x, y, z) = (1 - x/z - y/z)^2 - 4xy/z^2$. The integral in the general
⁷⁰⁸ off-shell boson case is much more difficult to interpret than the simpler on-shell branching ratios, but it
⁷⁰⁹ can be evaluated numerically. These branching ratio formulas can also be visualized as a function of Higgs
⁷¹⁰ mass, as shown in figure 1.4. There are a few interesting features to note in this figure. First, note that at
⁷¹¹ high Higgs masses, once on-shell production of both W and Z bosons is possible, these two decays are
⁷¹² dominant due to the large masses of the W/Z . Also note that the branching ratio to W s is twice that of
⁷¹³ Z s at these large masses due to the fact that there are two charged W bosons (W^\pm) and only one Z boson³.
⁷¹⁴ At 125 GeV, the Higgs is accessible through many different decay modes. The largest branching ratio is
⁷¹⁵ the decay $H \rightarrow b\bar{b}$ at 58.24% [18]. This branching is larger than the WW/ZZ decays because one of
⁷¹⁶ the two bosons must be produced off-shell for $m_h = 125$ GeV. The second largest branching ratio is
⁷¹⁷ to WW^* at 21.37 % (before taking into account the branching ratios of the W). Table 1.2 summarizes
⁷¹⁸ the theoretical branching ratios for a Higgs with a mass of 125 GeV. Note that there is a Higgs branching
⁷¹⁹ ratio to $\gamma\gamma$ even though photons are massless. This decay happens through a loop, which suppresses the

³In the Higgs Lagrangian, this extra symmetry factor is quantified by the δ_V noted in equation 1.12.

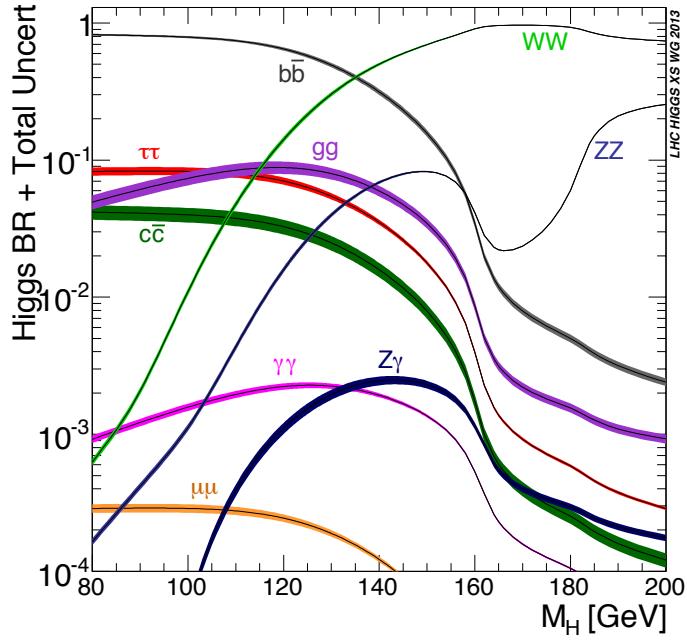


Figure 1.4: Higgs boson branching ratios as a function of m_H [18].

⁷²⁰ branching ratio⁴.

Decay	Branching ratio (%)	Relative uncertainty (%)
bb	58.24	+0.25 / - 0.25
WW^*	21.37	+0.99 / - 0.99
gg	8.187	+3.40 / - 3.41
$\tau\tau$	6.272	+1.17 / - 1.16
cc	2.891	+1.20 / - 1.20
ZZ^*	2.619	+0.99 / - 0.99
$\gamma\gamma$	0.2270	+1.73 / - 1.72
$Z\gamma$	0.1533	+5.71 / - 5.71
$\mu\mu$	0.02176	+1.23 / - 1.23

Table 1.2: Theoretical branching ratios for a 125 GeV Higgs boson, quoted as a percentage of the total width of the Higgs. Uncertainties shown are relative to the branching ratio value [18].

⁷²¹ Note that the branching ratios alone do not tell the full story of which Higgs channels are the most
⁷²² sensitive. For example, the $H \rightarrow b\bar{b}$ channel in gluon fusion production is incredibly difficult to observe
⁷²³ due to the large QCD dijet background at the LHC. However, in associated production of the Higgs,

⁴The largest contributions to the loop are the top quark and W boson.

724 where a W or Z gives additional final state particles that can be used to reduce background, a search for
 725 $H \rightarrow b\bar{b}$ can be sensitive. The combinations of production and decay modes that are most commonly
 726 studied at the LHC are summarized in table 1.3 [5].

Decay	Inclusive (incl. ggF)	VBF	WH/ZH	$t\bar{t}H$
$H \rightarrow \gamma\gamma$	✓	✓	✓	✓
$H \rightarrow bb$			✓	✓
$H \rightarrow \tau^+\tau^-$		✓		
$H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$	✓	✓	✓	
$H \rightarrow ZZ \rightarrow 4\ell$	✓			
$H \rightarrow Z\gamma \rightarrow \ell\ell\gamma$	very low			

Table 1.3: Possible channels for Higgs searches. Checkmarks denote the most sensitive production modes for each decay channel [5].

727 1.4 HIGGS PAIR PRODUCTION IN THE STANDARD MODEL

728 The Standard Model also allows for processes that produce two Higgs bosons in the final state, known
 729 as Higgs pair production or di-Higgs production. The two main production mechanisms are shown in
 figure 1.5. The two diagrams in figure 1.5 interfere destructively with one another, resulting in a low overall

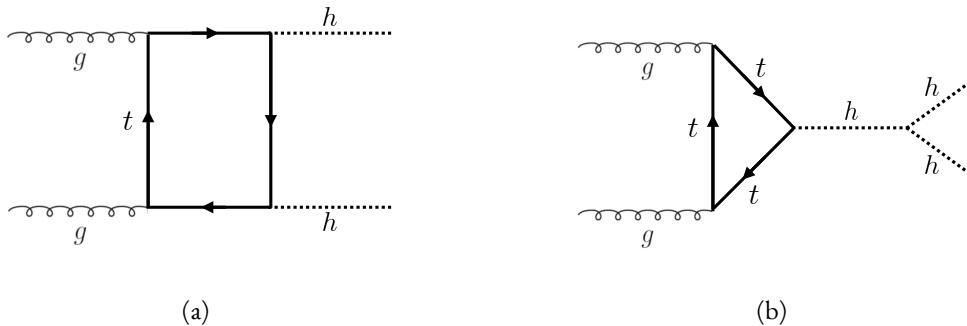


Figure 1.5: The two leading diagrams for Standard Model di-Higgs production at the LHC: (a) box diagram, (b) Higgs self coupling.

730 cross section for di-Higgs production at the LHC. Nevertheless, Higgs pair production is quite interesting
 731 to study because it gives direct access to the λ parameter of the Higgs potential, also known as the Higgs
 732 self coupling. The diagram in figure 1.5(b) is sensitive to this coupling through the triple Higgs vertex.

734 One can substitute the gluon fusion production of diagram 1.5(b) with any of the other production
 735 modes previously discussed. These other modes do not suffer from interference with the box diagram in
 736 figure 1.5(a) due to the presence of additional particles in the final state. They still have a lower cross section
 737 than the gluon fusion mode, however. The cross sections for di-Higgs production in the different modes,
 738 as well as their uncertainties, are shown in table 1.4 [20]. These are shown for $\sqrt{s} = 14$ TeV as this is the
 739 expected center of mass energy for the High Luminosity LHC and this energy is more sensitive to di-Higgs
 production. Note that the scale of cross section quoted is now in fb rather than pb.

Production mode	σ (fb)	Total uncert. (%)
Gluon fusion	33.89	+37.2 / - 27.8
Vector boson fusion	2.01	+7.6 / - 5.1
$W H H$	0.57	+3.7 / - 3.3
$Z H H$	0.42	+7.0 / - 5.5
$t \bar{t} H$	1.02	-

Table 1.4: Production cross sections for pair production of a 125 GeV Higgs boson at $\sqrt{s} = 14$ TeV with total uncertainty [20]. The uncertainties include QCD scale and PDF variations as well as uncertainties on α_S .

740

741 1.5 HIGGS PAIR PRODUCTION IN THEORIES BEYOND THE STANDARD MODEL

742 The Higgs pair production cross section in the Standard Model is rather small, and datasets on the
 743 scale of the full 3000 fb^{-1} expected from the High Luminosity LHC will be required to obtain sensitive
 744 measurements of the Higgs self-coupling [20]. However, the discovery of the Higgs also gives particle
 745 physicists a new tool that can be exploited in the search for new physics beyond the Standard Model. In
 746 particular, Higgs pair production is a promising channel in the search for new physics. The cross section for
 747 di-Higgs production can be altered through both resonant and non-resonant production of Higgs pairs. In
 748 non-resonant production, di-Higgs production vertices can arise from the presence of a new strong sector
 749 and additional colored particles [21–23]. Figure 1.6 shows examples of the types of vertices that can arise. In
 750 the resonant case, new heavy particle can decay to Higgs pairs. Such new particles can include heavy Higgs
 751 bosons arising in two Higgs doublet models (2HDM) or Higgs portal models as well as heavy gravitons in
 752 Randall-Sundrum theories [21, 24–30]. Figure 1.7 shows a generic diagram for a heavy resonance decaying

753 to two Higgs bosons. In the 2HDM, X corresponds to the heavy CP-even scalar H . In the Randall-Sundrum model, X corresponds to a heavy spin-2 graviton G_{KK}^* . The next sections provide more detail

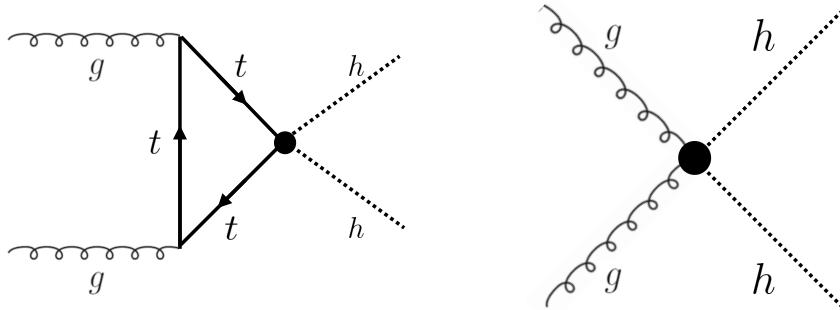


Figure 1.6: Diagrams with new vertices for non-resonant Higgs pair production arising in composite Higgs models.

754 on the phenomenology of resonant Higgs production in Randall-Sundrum and 2HDM models, as these
755 models will later be tested in a dedicated search for resonant production of boosted Higgs pairs.

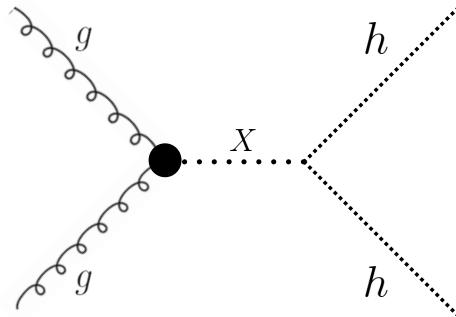


Figure 1.7: Generic Feynman diagram for resonant Higgs pair production in BSM theories.

756

757 1.5.1 RANDALL-SUNDRUM GRAVITONS

758 The Randall-Sundrum model is a proposed solution to the hierarchy problem that posits a five-dimensional
759 warped spacetime that contains two branes: one where the force of gravity is very strong and a second brane
760 at the TeV scale corresponding to the known Standard Model sector [24]. In the theory, the branes are
761 weakly coupled and the graviton probability function drops exponentially going from the gravity brane
762 to the SM brane, rendering gravity weak on the SM brane. The experimental consequence of this theory

763 is a tower of widely spaced (in mass) Kaluza-Klein graviton resonances. In theories where the fermions
 764 are localized to the SM brane, production of gravitons from fermion pairs is suppressed and the primary
 765 mode of production is gluon fusion [25]. These gravitons have a substantial branching fraction to Higgs
 766 pairs, ranging from 6.43% for gravitons with a mass of 500 GeV to 7.66% at 3 TeV. Figure 1.8 shows the
 767 branching ratios of the spin-2 Randall Sundrum graviton (RSG) as a function of its mass. The predomi-
 768 nant decays are to $t\bar{t}$ above the mass threshold for that channel.

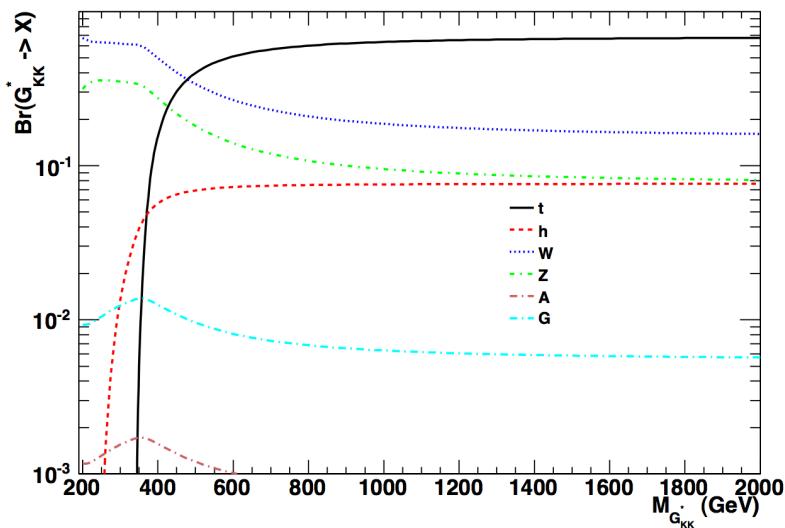


Figure 1.8: Branching ratios for a spin-2 Randall-Sundrum graviton as a function of mass computed in MadGraph with the CP3-Origins implementation [25, 31, 32].

769 Randall-Sundrum models have two free parameters - the mass of the graviton and a curvature parameter
 770 k . Typically, rather than k , the theory is parameterized using $c \equiv k/\bar{M}_{\text{pl}}$, where \bar{M}_{pl} is the reduced
 771 Planck mass. The cross section for production of the RSG decreases as a function of mass and is strongly
 772 dependent on the gluon PDF. The increase in center of mass energy from 8 to 13 TeV in LHC Run 2
 773 greatly increases the cross section at higher mass. Figure 1.9 shows the cross section as a function of graviton
 774 mass at $\sqrt{s} = 13$ TeV for RSG models with $c = 1.0$ and $c = 2.0$.

775 Another interesting feature of the theory is that the width of the graviton increases with both c and
 776 $m_{G_{KK}^*}$. Figure 1.10 shows the graviton width for both $c = 1.0$ and $c = 2.0$ as a function of mass. In
 777 $c = 1.0$, the width starts at 8.365 GeV for a mass of 300 GeV and increases to 187.2 GeV at a mass of
 778 3 TeV. Similarly, with $c = 2.0$, the width starts at 33.46 GeV for $m_G = 300$ GeV and increases to

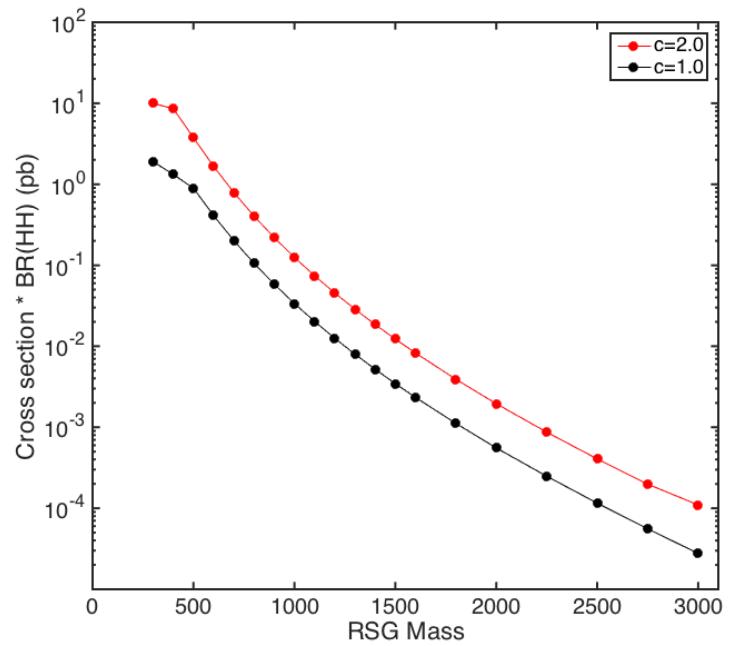


Figure 1.9: $\sigma \times \text{BR}(HH)$ for Randall-Sundrum gravitons as a function of mass computed in MadGraph with the CP3-Origins implementation [25, 31, 32].

⁷⁷⁹ 748.8 GeV at a mass of 3 TeV.

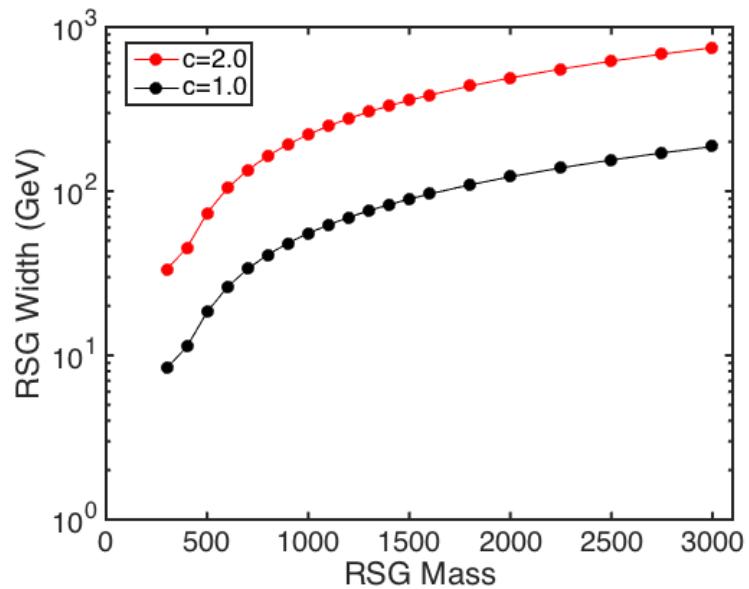


Figure 1.10: Randall-Sundrum graviton width as a function of mass computed in MadGraph with the CP3-Origins implementation [25, 31, 32]

780 1.5.2 TWO HIGGS DOUBLET MODELS

781 In Two Higgs Doublet Models (2HDM), a second complex scalar doublet is added to the Standard
782 Model [27–29]. In this case, all four degrees of freedom in the second doublet correspond to new particles,
783 meaning that there are five total scalars from the two Higgs doublets - h (light CP-even Higgs), H (heavy
784 CP-even Higgs), A (heavy CP-odd Higgs), and H^\pm (charged Higgs). The model is parameterized by two
785 main parameters. The first, $\tan \beta \equiv \frac{v_2}{v_1}$, is the ratio of the vacuum expectation values of the two Higgs
786 doublets (where v_1 corresponds to the v in the SM Higgs model described above). The second parameter
787 is α , a mixing angle between the heavy and light Higgs fields. Models are also often parameterized with
788 $\cos(\beta - \alpha)$ rather than α directly. The limit where $\cos(\beta - \alpha) = 0$ is called the alignment limit, and in
789 this limit the light Higgs h has the same couplings as a Standard Model Higgs.

790 2HDM models are usually separated into two main types - Type I and Type II. In Type I models, the
791 charged fermions only couple to the second Higgs doublet, leading to a fermiophobic light Higgs. In
792 Type II models, up-type quarks couple to the first doublet while down-type quarks couple to the second
793 doublet. One specific realization of a Type II 2HDM is the Minimal Supersymmetric Standard Model
794 (MSSM).

795 Resonant di-Higgs production in 2HDM models can proceed through decays of the heavy CP-even
796 Higgs $H \rightarrow hh$. The branching ratio for $H \rightarrow hh$ depends on the model type as well as the values of
797 $\tan \beta$ and $\cos \beta - \alpha$. Figure 1.II shows the branching ratios as a function of the mass of the heavy scalar
798 H for both Type I and Type II models. Depending on the type of model hh can be a substantial fraction
799 of the decays of H .

800 1.6 CONCLUSION

801 Studying the Higgs sector is essential for understanding the details of how mass arises in the Standard
802 Model and how the electroweak symmetry is broken. The discovery of the Higgs boson also opens the
803 door for its use as a tool to search for new physics, and Higgs pair production is an ideal candidate for
804 this study. Even if no BSM physics is found in Higgs pair production, searches for Higgs pairs will put
805 constraints on the Higgs self coupling and thus improve knowledge of the Standard Model and the details

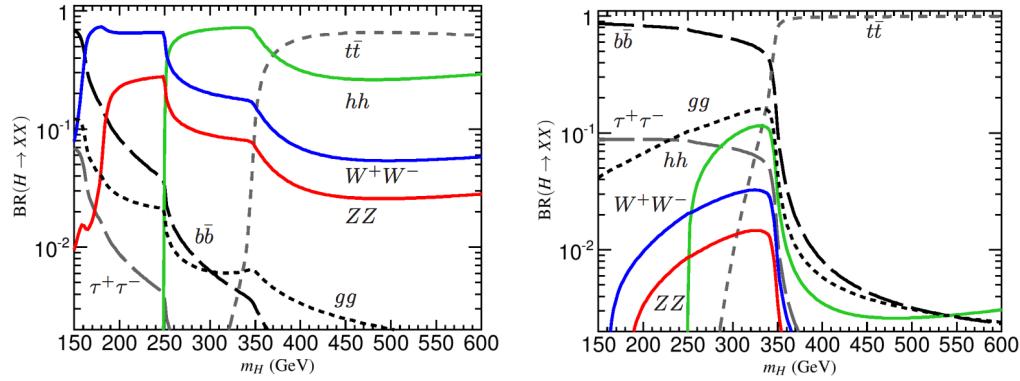


Figure 1.II: Branching ratios for heavy Higgs H in Type I (left) and Type II (right) 2HDM models with $\tan \beta = 1.5$ and $\cos(\beta - \alpha) = 0.1$ (0.01) for Type I (Type II) [29].

of the Higgs potential.

*The enthusiasm and motivation to explore particle physics
at the high-energy frontier knows no borders between the
nations and regions of the planet.*

Peter Jenni

2

807

808

809

The ATLAS detector and the Large Hadron Collider

810 This chapter presents an overview of the experimental systems used to conduct the measurements pre-
811 sented in this thesis. First, a brief overview of the accelerator, the Large Hadron Collider, will be given. In
812 this section, the accelerator conditions relevant to data-taking are presented as well. Next, an overview of
813 the ATLAS experiment is given. The basics of each sub-detector's role are summarized, as well as the details
814 of the datasets accumulated. Then, a brief interlude on the ATLAS Muon New Small Wheel upgrade is
815 presented. While this new detector does not have a direct impact on any of the datasets recorded so far,
816 it will have an impact on future analyses and the work done on it is briefly summarized here. Finally, an
817 overview of object reconstruction in ATLAS is given. While the details of all of the algorithms will not be
818 presented in detail, aspects of the reconstruction performance are shown as these are relevant to the results
819 presented later in this thesis.

820 2.1 THE LARGE HADRON COLLIDER

821 The Large Hadron Collider (LHC) is a proton-proton collider at the CERN laboratory in Geneva,
822 Switzerland [33]. It is designed for a maximum collision center of mass energy of $\sqrt{s} = 14$ TeV and has a
823 circumference of 26.7 kilometers. Four main experiments are located at the interaction points (IP) of the
824 accelerator: ATLAS (A Toroidal LHC Apparatus), CMS (the Compact Muon Solenoid), ALICE (A Large
825 Ion Collider Experiment), and LHCb [34–37]. The results presented in this thesis were all completed with
826 the ATLAS detector. Figure 2.1 shows a schematic of the LHC ring and its experiments.

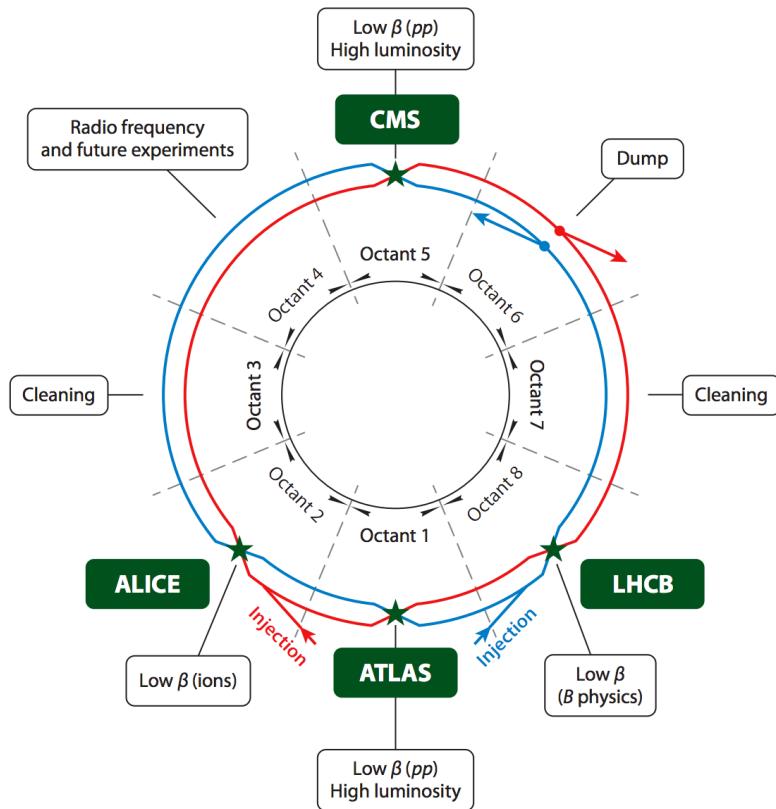


Figure 2.1: A schematic view of the LHC ring [38].

827 One of the most interesting features of the LHC is its magnet design. Because the tunnel does not have
828 room for separate superconducting magnets for each of the beam pipes, the LHC employs a twin-bore
829 magnet design. Each magnet must hold an 8.3 Tesla magnetic field in order to bend the proton beams at
830 $\sqrt{s} = 14$ TeV. The superconducting magnets are cooled to a temperature of 1.9 Kelvin with superfluid

831 helium.

832 2.1.1 INSTANTANEOUS LUMINOSITY

833 The rate of physics events expected from the accelerator is dependent on the instantaneous luminosity
834 of the machine and the cross section of the physics process, $R_{\text{events}} = L\sigma$. Here, R_{events} is the num-
835 ber of events per second, L is the instantaneous luminosity of the machine, and σ is the cross section for
836 the physics process being measured. The instantaneous luminosity of the LHC is determined by numer-
837 ous factors related to beam conditions. Equation 2.1 gives the equation for instantaneous luminosity of a
838 Gaussian beam profile [38].

$$L = \frac{N_b^2 n_b f_{\text{rev}} \gamma_r}{4\pi \epsilon_n \beta^*} F \quad (2.1)$$

839 The LHC collides protons in bunches, and in the above equation N_b is the number of protons per bunch
840 while n_b is the number of bunches per beam. Nominally, the LHC can hold up to 2808 proton bunches.
841 f_{rev} is the revolution frequency. ϵ_n is the normalized transverse beam emittance, a measurement of the
842 average spread of the particles in position-momentum space which has the dimension of length. β^* is the
843 value of the β function for the beam at the interaction point. It relates the emittance to the Gaussian
844 width of the beam with $\sigma_{\text{beam}} = \sqrt{\epsilon \cdot \beta}$. F is a reduction factor that corrects for the fact that the beams
845 are colliding at an angle at the IP.

846 Another way of writing the instantaneous luminosity is shown in equation 2.2. In this case, the instan-
847 taneous luminosity is written as the ratio of the rate of inelastic collisions to the inelastic cross section [39].

848

$$L = \frac{R_{\text{inel}}}{\sigma_{\text{inel}}} = \frac{\mu n_b f_{\text{rev}}}{\sigma_{\text{inel}}} \quad (2.2)$$

849 In this case, μ is the average number of interactions per bunch crossing in the accelerator. μ is a useful
850 parameter for characterizing the amount of activity recorded in an experiment. As the instantaneous lu-
851 minosity and thus μ increase, there are more interactions per bunch crossing and more activity is present
852 in the detector. The level of activity is often characterized with $\langle \mu \rangle$, the measured per bunch crossing μ
853 value averaged over all bunch crossings. The interactions inside each bunch crossing that are not the main

854 physics process of interest are often referred to as “pileup” interactions, and $\langle \mu \rangle$ is a measurement of the
855 level of pileup in the detector.

856 **2.1.2 EVOLUTION OF MACHINE CONDITIONS**

857 This thesis uses datasets taken at three different center of mass energies: $\sqrt{s} = 7\text{ TeV}$ data taken in the
858 year 2011, $\sqrt{s} = 8\text{ TeV}$ data taken in the year 2012, and $\sqrt{s} = 13\text{ TeV}$ data taken in the year 2015. In
859 addition to increasing center of mass energy, the instantaneous luminosity and parameters that determine
860 it were evolving. Table 2.1 summarizes that machine conditions in each of these datasets.

	2011	2012	2015	Design
$\sqrt{s} [\text{TeV}]$	7	8	13	14
Number of bunches	1380	1380	1825	2808
Max. protons per bunch	1.45×10^{11}	1.7×10^{11}	1.2×10^{11}	1.15×10^{11}
Bunch spacing [ns]	50	50	25	25
Max. instantaneous luminosity [$\text{cm}^{-2}\text{s}^{-1}$]	3.7×10^{33}	7.7×10^{33}	5×10^{33}	10^{34}
$\beta^* [\text{m}]$	1.0	0.6	0.8	0.55
$\langle \mu \rangle$	11.6	20.7	13.7	-

Table 2.1: Evolution of LHC machine conditions [40, 41].

861 **2.2 THE ATLAS DETECTOR**

862 The ATLAS detector is a multi-purpose particle detector experiment at the LHC’s Point 1 [34]. It has
863 nearly 4π coverage in solid angle around the interaction point. It consists of an inner detector for mea-
864 suring charged particles, electromagnetic and hadronic calorimeters, and a muon spectrometer. Figure 2.2
865 gives an overview of the detector.

866 **2.2.1 COORDINATE SYSTEM**

867 Before defining the properties of the individual detectors, it is important to establish the coordinate
868 system used. Figure 2.3 shows a schematic of the coordinate system. The azimuthal plane (perpendicular
869 to the beam line) is defined as the x - y plane. The angle in this plane is referred to as ϕ . The angle relative

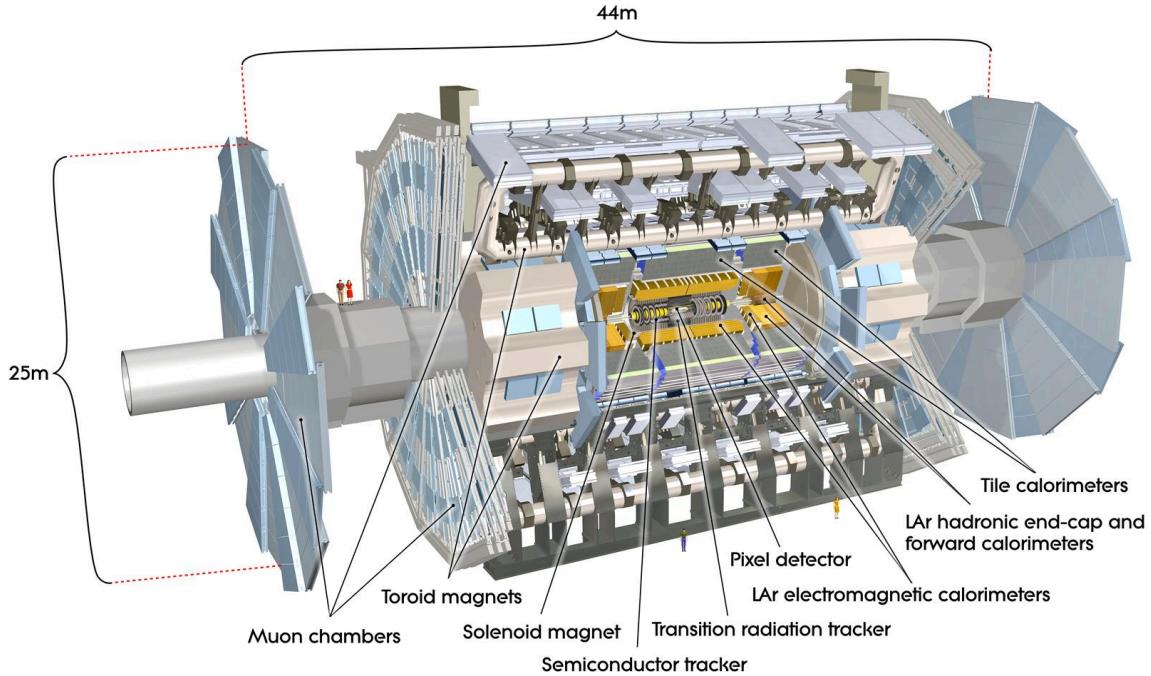


Figure 2.2: A full diagram of the ATLAS detector [34]

870 to the beam axis is referred to as θ . Rather than using θ directly as a coordinate, the experiment often uses
 871 the pseudorapidity η . η is defined in equation 2.3.

$$\eta = -\ln \left(\tan \left(\frac{\theta}{2} \right) \right) \quad (2.3)$$

872 Pseudorapidity is the massless approximation of rapidity, the angle used to parameterize boosts in spe-
 873 cial relativity. This coordinate is useful in particle physics for two reasons. First, it means that differences
 874 in η are Lorentz invariant. Second, particle production is roughly constant in pseudorapidity. Particles
 875 with η close to zero are referred to as “central”, while those at high $|\eta|$ are called “forward”. In general,
 876 two main detector configurations can be seen in figure 2.2. There are “barrel” elements, which surround
 877 the beam line cylindrically and are in the central region of the detector. In the forward region, there are
 878 “endcap” regions which are arranged as disks perpendicular to the beam line.

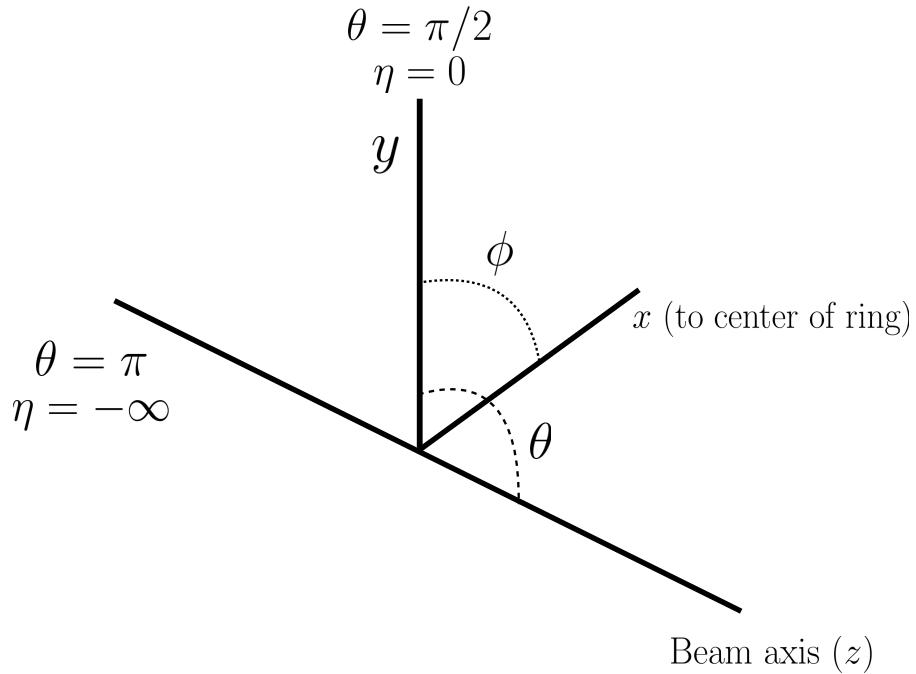


Figure 2.3: The ATLAS coordinate system.

879 2.2.2 INNER DETECTOR

880 The ATLAS Inner Detector (ID) system is built for precision tracking of charged particles. It covers
 881 the range $|\eta| < 2.5$. In this range, approximately 1000 particles are generated every bunch crossing in the
 882 detector. This requires having fine granularity to achieve the resolutions required for good momentum
 883 measurement and vertex reconstruction.

884 The ID consists of three sub-components: the pixel detector, semiconductor tracker (SCT), and trans-
 885 sition radiation tracker (TRT). It is surrounded by a solenoid providing a 2 T axial magnetic field which
 886 bends particles in the transverse plane to allow for momentum measurement. Figure 2.4 shows the layout
 887 of each of these components.

888 PIXEL DETECTOR

889 The pixel detector is the first detector particles traverse after being generated in proton collisions and
 890 is the most granular detector. Its operation is crucial for precision tracking and vertex reconstruction as



Figure 2.4: Layout of the ATLAS Inner Detector system [42].

well as higher level object reconstruction like tagging of jets from b -quarks. The basic sensing element in this subdetector is a silicon pixel detector. The operating principle for the silicon pixels is that of a p - n junction. When a charged particle passes through, it creates electron-hole pairs that are then separated by the electric field. The sensors are $250\ \mu\text{m}$ thick and use oxygenated n -type wafers with readout pixels on the n^+ side of the detector [34]. Overall, the pixel detector has 1744 sensors and 80.4 million readout channels.

In the barrel region, the pixel detector has three concentric layers of sensors surrounding the beamline. In the endcap region, it consists of disks perpendicular to the beam axis. The detector is segmented in the R - ϕ plane and in z . Usually, three pixel layers are crossed by a charged particle track. The intrinsic accuracies of the sensors are $10\ \mu\text{m}$ in R - ϕ and $115\ \mu\text{m}$ in z (or R for the endcap).

901 INSERTABLE B-LAYER

In Run 2, a new innermost pixel layer, known as the insertable B-layer (IBL), was added to the Inner Detector [43]. This layer was added to cope with the higher luminosities planned in LHC Run 2 and at the

904 high luminosity HL-LHC. Additionally it improves tracking position resolution which in turn improves
905 the vertexing and b -tagging capabilities in ATLAS. The detector sits directly on a new beam pipe, only
906 33.25 mm away from the collision points in the azimuthal plane.

907 **SEMICONDUCTOR TRACKER (SCT)**

908 The semiconductor tracker (SCT) consists of silicon microstrips and comprises the next four layers
909 of the ID. This sub-detector has 6.4 cm long sensors that are daisy-chained into strips with a strip pitch
910 of $80 \mu\text{m}$ [34]. Some of the strips have a small stereo angle to allow for measurement of both angular
911 coordinates. In total there are 6.3 million readout channels. The intrinsic accuracies are $17 \mu\text{m}$ in $R\phi$
912 and $580 \mu\text{m}$ in z (or R in the endcap).

913 **TRANSITION RADIATION TRACKER (TRT)**

914 The transition radiation tracker (TRT) serves two purposes. First, it consists of 4mm diameter straw
915 tubes filled with a 70/27/3% gas mixture of xenon, carbon dioxide, and oxygen to provide tracking of
916 charged particles. Particles typically have 36 TRT straw tube hits per track. The material in between
917 the straws is designed to induce transition radiation which can be useful for particle identification. As
918 particles pass between media with different dielectric constants, they emit transition radiation that can
919 cause additional showers in the TRT. In particular it is useful for discrimination between electrons and
920 pions or other charged hadrons, as the amount of transition radiation is proportional to the Lorentz factor
921 of the particle.

922 **2.2.3 CALORIMETERS**

923 The calorimeter system consists of two main sub-components: a fine granularity electromagnetic calorime-
924 ter tailored for the measurement of photons and electrons and multiple coarser hadronic calorimeters ded-
925 icated to the measurement of hadronic showers [34]. The calorimeter system has broader coverage than
926 the inner detector, covering the region out to $|\eta| < 4.9$. It is also designed to deliver good containment of
927 showers so as to limit leakage into the muon system. Figure 2.5 shows the layout of the calorimeter system.

928 Both the electromagnetic and hadronic calorimeters are sampling calorimeters. They alternate active
 929 material for energy measurement with passive material for energy absorption. The materials used for each
 930 purpose vary based on the type of calorimeter and its location in the detector.

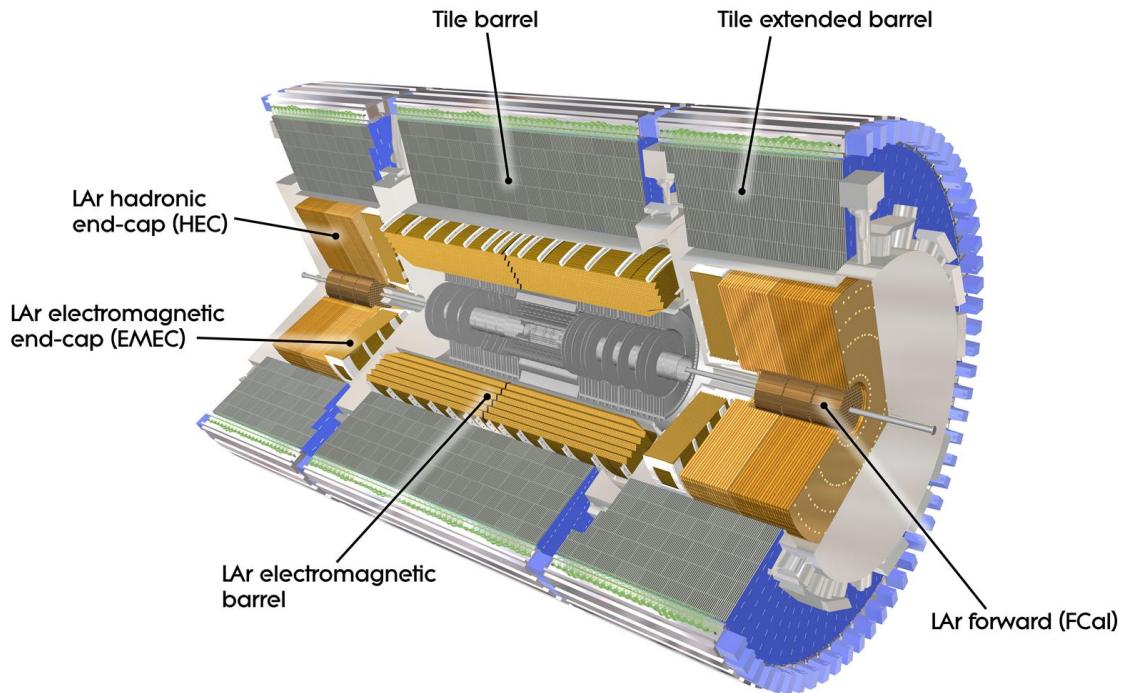


Figure 2.5: Layout of the ATLAS calorimeter system [34].

931 ELECTROMAGNETIC CALORIMETER

932 The electromagnetic calorimeter (EM calorimeter) use liquid Argon (LAr) as its active material and lead
 933 as its passive material. It is arranged in an accordion geometry to increase the absorption area while still
 934 allowing it to have no azimuthal cracks (complete symmetry in ϕ). The EM calorimeter is divided into a
 935 barrel portion that extends to $|\eta| < 1.475$ and an endcap portion going from $1.375 < |\eta| < 3.2$. The
 936 region where these two units overlap is called the “transition region”.

937 In order to provide good containment the calorimeter depth must be optimized. Typically, for elec-
 938 tromagnetic calorimeters the depth is measured in radiation lengths. In general, the intensity of a par-
 939 ticle beam attenuates exponentially in distance with a constant equal to the radiation length. That is,

940 $I(x) = I_0 e^{-x/X_0}$, where I is the intensity, x is the distance traveled, and X_0 is the radiation length.
941 The ATLAS EM calorimeter is designed to have > 22 radiation lengths in the barrel and > 24 in the
942 endcap [34].

943 **HADRONIC CALORIMETERS**

944 There are three types of hadronic calorimeters present in ATLAS: the tile calorimeter (TileCal), hadronic
945 endcap (HEC), and forward calorimeter (FCal). Each one is optimized for stopping of hadronic showers
946 and the materials chosen are specific to their placement in the detector.

947 The TileCal is a scintillating tile calorimeter placed directly outside the EM calorimeter. It uses steel as
948 the absorber and plastic scintillator tiles as the active material. It has coverage in the barrel at $|\eta| < 1.0$
949 and in the “extended barrel” region of $0.8 < |\eta| < 1.7$.

950 The HEC had two wheels perpendicular to the beam line per endcap and is located directly behind the
951 EM calorimeter endcap modules. The HEC covers the region from $1.5 < |\eta| < 3.2$, overlapping slightly
952 with both the tile calorimeter and the forward calorimeter. Like the EM calorimeter, it uses liquid Argon
953 as the active material, but it uses copper as the absorber.

954 The FCal covers the most forward regions of the calorimeter system, extending to the region of $3.1 <$
955 $|\eta| < 4.9$. It again uses liquid argon as its active material. For absorber, it consists of an innermost module
956 made of copper followed by a module made of tungsten.

957 The hadronic equivalent of radiation length is called the interaction length and is denoted as λ . In the
958 barrel, the hadronic calorimeter depth is approximately 9.7λ , while in the endcap is 10λ . The outer
959 supports contribute an additional 1.3λ . This is been shown to be sufficient to limit punch-through of
960 showers to the muon system [34].

961 **2.2.4 MUON SPECTROMETER**

962 The muon spectrometer is dedicated to measuring the momentum and position of muons. It consists
963 of tracking and trigger chambers which are unique in the barrel and endcap regions. The magnetic field
964 for bending of muons is provided by a system of three large air-core toroid magnets (from which ATLAS

965 derives its name.) These magnets provide 1.5 to 5.5 Tm of bending power at $0 < |\eta| < 1.4$ and approx-
 966 imately 1 to 7.5 Tm in the endcap region of $1.6 < |\eta| < 2.7$. The entire muon system covers the range
 967 $0 < |\eta| < 2.7$. Monitored drift tubes (MDTs) are used for tracking in the barrel and the two outer layers
 968 of the endcap, while cathode strip chambers (CSCs) are used to provide tracking in the innermost endcap
 969 wheel. In the barrel, resistive plate chambers (RPCs) are used as trigger chambers while thin gap chambers
 970 (TGCs) are used in the endcap. Figure 2.6 shows the layout of the ATLAS muon system. The entire muon
 971 system is designed with the specification of providing a 10% momentum resolution for a 1 TeV muon.

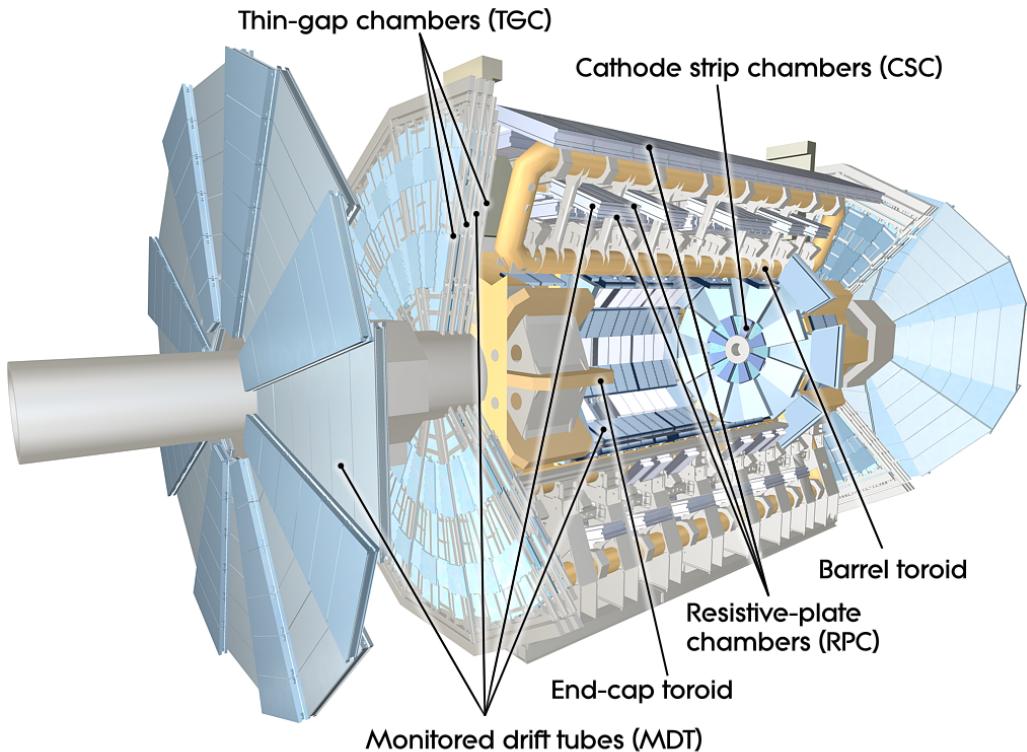


Figure 2.6: Layout of the ATLAS muon system [34].

972 MONITORED DRIFT TUBES (MDTs)

973 The monitored drift tubes (MDTs) are aluminum 3 cm diameter tubes filled with a 93/7 % mixture of
 974 Argon and CO₂, with trace amounts of water. As a charged particle traverses the tube, it ionizes the gas
 975 and the ions drift to a wire at the center of the tube. The radial distance of traversal of the particle in the

976 tube is determined by the drift time of the electrons, allowing for fine position resolution. The tubes have
977 an average resolution of $80 \mu\text{m}$ per tube and a maximum drift time of approximately 700ns. The tubes
978 are oriented so that they give precision measurement in η and run along ϕ . They cover $|\eta| < 2.7$, except
979 in the innermost layer of the endcap where they only go to $|\eta| < 2.0$ [34].

980 CATHODE STRIP CHAMBERS (CSCs)

981 The cathode strip chambers cover a narrow window of the innermost endcap region at $2.0 < |\eta| <$
982 2.7. In this region the background rates in the cavern are particularly high and the CSCs are designed to
983 handle these higher rates. The CSCs are multiwire proportional chambers with wires pointing in the radial
984 direction (away from the beam pipe). The wire serves as an anode and there are two types of segmented
985 cathode strip, one perpendicular to the wires which gives the precision measurement and one parallel which
986 provides the transverse coordinate. It has an 80/20% gas mixture of Argon and CO₂ [34].

987 RESISTIVE PLATE CHAMBERS (RPCs)

988 The resistive plate chambers (RPCs) are gaseous electrode-plate detectors covering the region $|\eta| <$
989 1.05. They consist of two resistive plates separated by a distance of 2 mm. The gas mixture used is a
990 94.7/5/0.3% mixture of C₂H₂F₄, Iso-C₄H₁₀, and SF₆. It has readout strips with a pitch of 23-35 mm
991 for both η and ϕ measurement and thus provides measurement of the azimuthal coordinate in the barrel.
992 The thin gas gap allows for a quick response time which makes it ideal for use in the trigger. There are three
993 layers of RPCs which are referred to as the three trigger stations. They allow for both a low p_T and high
994 p_T trigger. The coincidence of hits in the innermost chambers allows for triggering of muons between 6
995 and 9 GeV, while the outermost layer allows the trigger to select high momentum tracks in the range of 9
996 to 35 GeV [34].

997 THIN GAP CHAMBERS (TGCs)

998 The thin gap chambers (TGCs) are multiwire proportional chambers where the wire to cathode dis-
999 tance (1.4mm) is smaller than the wire-to-wire distance (1.8 mm). They contain a gas mixture of CO₂
1000 and *n*-pentane and use a high electric field to gain good time resolution. They serve two functions in the

end-cap system. First, they serve as the trigger chambers. Second, they also provide azimuthal coordinate measurement. They sit on the inner and middle layers of the endcap. The outermost layer's azimuthal coordinate is determined by extrapolation [34].

2.2.5 MAGNET SYSTEM

As mentioned previously, there are two independent magnet systems in ATLAS. The first is a 2 T solenoid field in the inner detector which provides bending in the azimuthal plane. The second is an approximately 0.5 T toroidal field in the muon system which provides bending in η . Figure 2.7 shows the predicted field integral as a function of $|\eta|$ [34].

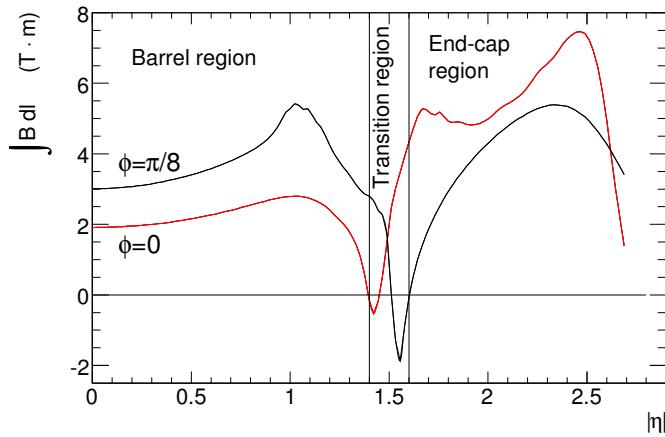


Figure 2.7: Predicted field integral as a function of $|\eta|$ for the ATLAS magnet system [34].

2.2.6 TRIGGER SYSTEM

The ATLAS trigger system searches for signatures of muons, electrons, photons, hadronically decaying τ leptons, and jets in order to save these events for further analysis. The trigger system in ATLAS is designed to reduce the maximum LHC event rate of 40 MHz to a more reasonable rate that can be recorded. The trigger first consists of a fast, hardware based system called the Level-1 (L1) trigger. The L1 trigger consists of independent dedicated detector sub-components that can seed regions of interest (RoIs) for further analysis downstream. For muons, the RPCs and TGCs are used, while in the calorimeter coarsely grained sections of calorimeter cells called towers are used. Once regions of interest are seeded, a software

1017 based system called the High Level Trigger (HLT) is used to reconstruct objects and integrate information
 1018 from different parts of the detector. In Run 1 of ATLAS, the HLT consisted of two separate stages: the
 1019 level 2 (L2) trigger and the event filter (EF).

1020 The maximum trigger rate that the L1 trigger can handle is 75 kHz. In the HLT, the rate of events
 1021 written to disk is approximately 200 Hz. Figure 2.8 shows the trigger rates for different L1 triggers in 2012
 1022 and 2015 for ATLAS [44].

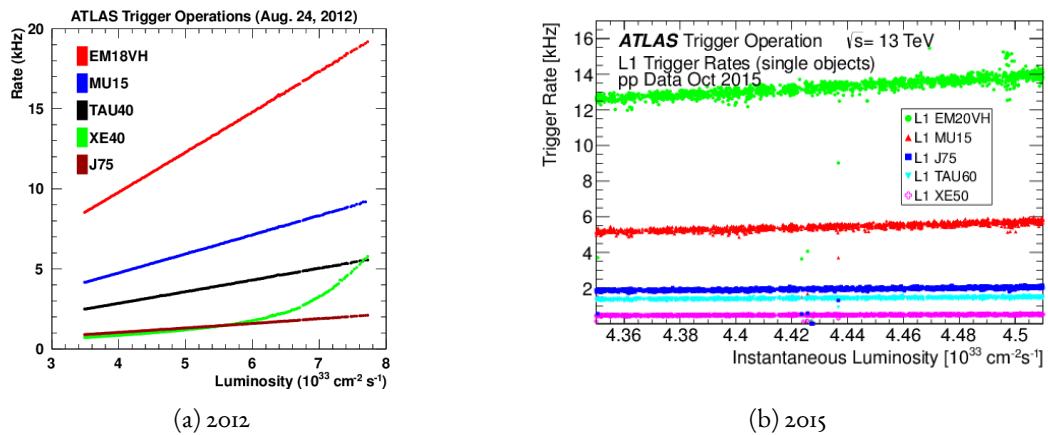


Figure 2.8: ATLAS trigger rates for Level-1 triggers as a function of instantaneous luminosity in 2012 and 2015 operation. These are single object triggers for electromagnetic clusters (EM), muons (MU), jets (J), missing energy (XE), and τ leptons (TAU). The threshold of the trigger is given in GeV [44].

1023 2.2.7 ATLAS DATASETS

1024 ATLAS has collected data at center of mass energies of 7, 8, and 13 TeV. Figure 2.9 shows the integrated
 1025 luminosity as a function of time for each of the three collected datasets. At $\sqrt{s} = 7$ TeV, ATLAS recorded
 1026 5.08 fb^{-1} . Increased instantaneous luminosity in 2012 led to a larger dataset of 21.3 fb^{-1} recorded at
 1027 $\sqrt{s} = 8$ TeV. After Long Shutdown 1 (LS1) of the LHC and a restart in 2015, ATLAS recorded 3.9 fb^{-1}
 1028 of data at $\sqrt{s} = 13$ TeV [45, 46].

1029 2.2.8 DETECTOR PERFORMANCE

1030 Table 2.2 summarizes the design requirements for each of the different sub-detectors. This table shows
 1031 the energy and momentum resolution of tracking, calorimetry, and muon measurements.

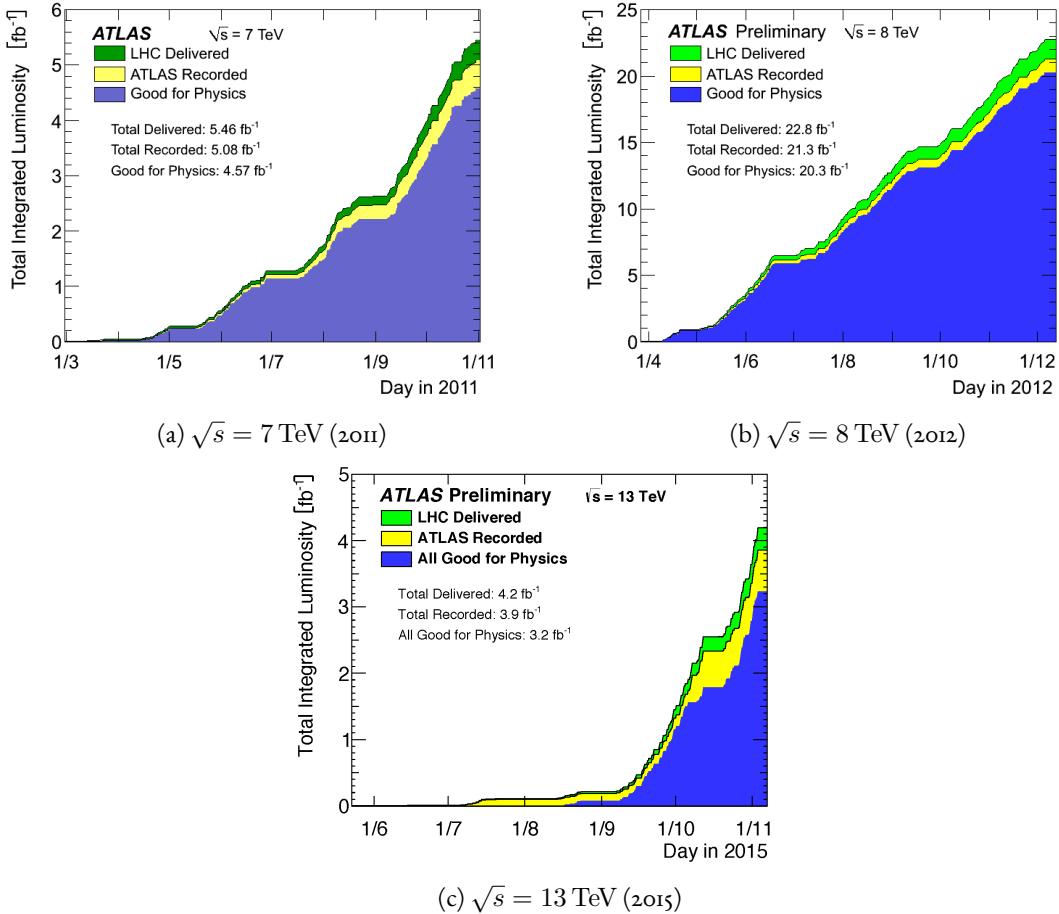


Figure 2.9: Instantaneous luminosity as a function of time for data recorded by ATLAS at different center of mass energies [45, 46].

	Required resolution
Tracking	$\sigma_{p_T}/p_T = 0.05\% p_T \oplus 1\%$
EM calorimetry	$\sigma_E/E = 10\%/\sqrt{E} \oplus 0.7\%$
Hadronic calorimetry	
Barrel and end-cap	$\sigma_E/E = 50\%/\sqrt{E} \oplus 3\%$
Forward	$\sigma_E/E = 100\%/\sqrt{E} \oplus 10\%$
Muon spectrometer	σ_{p_T}/p_T at $p_T = 1 \text{ TeV}$

Table 2.2: Performance requirements for the ATLAS detector [34].

2.3 THE ATLAS MUON NEW SMALL WHEEL UPGRADE

As the LHC continues operation, it is scheduled to be upgraded in several phases to allow it to reach higher instantaneous luminosities and thus collect larger datasets. These conditions will open new doors

1035 for study of rare physics processes but will also present interesting challenges that must be faced. ATLAS
 1036 will require new detector technologies to cope with the increased background rates in the cavern in these
 1037 high luminosity conditions. One such upgrade, scheduled to be installed during Long Shutdown 2 (LS2)
 1038 of the LHC in 2018, is the ATLAS Muon New Small Wheel (NSW) upgrade [47]. The NSW will replace
 1039 the innermost end-cap wheel of the muon system with new technologies. This is the part of the muon
 1040 detector closest to the beam line and thus experiences the highest rates of particle flux in the muon system.

1041 2.3.1 MOTIVATION

1042 The motivation of the NSW is two-fold. The first objective is to alleviate the decreased tracking ef-
 1043 ficiency that comes in a high rate environment. As figure 2.10, at the LHC design luminosity both the
 1044 efficiency of recording hits and reconstructing track segments in the MDTs decreases.

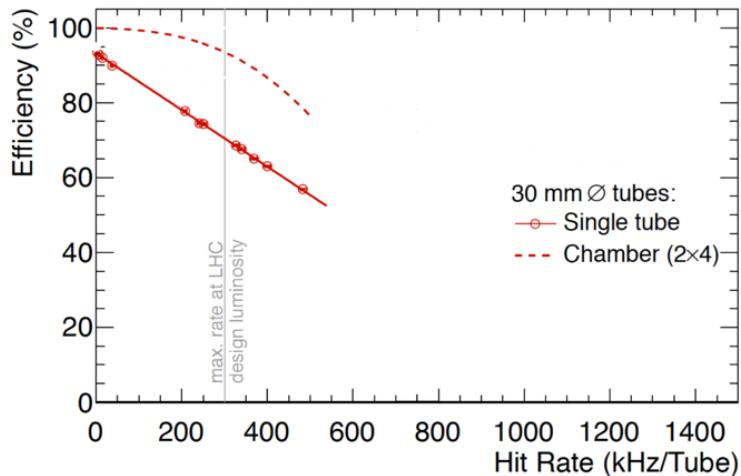


Figure 2.10: MDT tube hit (solid) and segment (dashed) efficiency as a function of hit rate per tube [47].

1045 The NSW will also work to alleviate the rate of fake triggers arising in the endcap. Figure 2.11 shows
 1046 the extrapolated trigger rates as a function of the p_T threshold with and without the NSW upgrade. As
 1047 the figure shows, the NSW upgrade will reduce the trigger rate by an order of magnitude compared to the
 1048 current endcap trigger system. This reduction allows the p_T thresholds on muons to remain low, increasing
 1049 the phase space of possible physics studies and in particular maintaining good acceptance for Higgs physics.

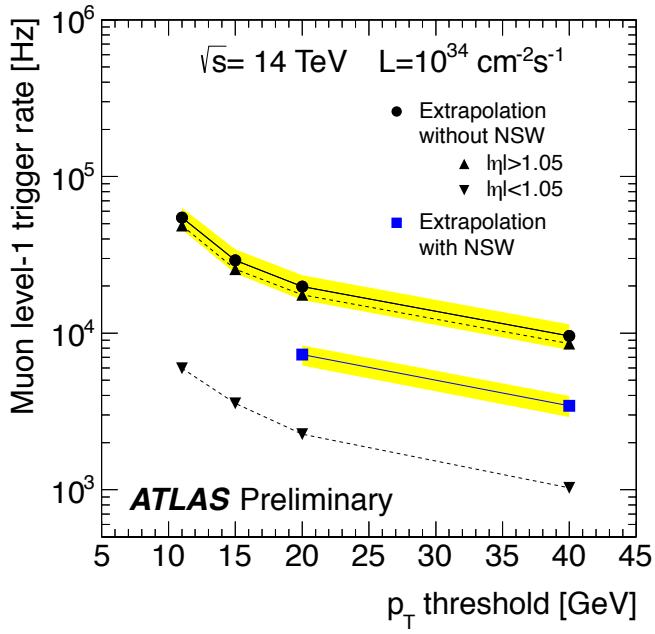


Figure 2.11: Trigger rate as a function of p_T threshold with and without the NSW upgrade [47].

1050 2.3.2 NSW DETECTOR TECHNOLOGIES

1051 The NSW will use two new detector technologies - micromesh gaseous structure detectors (micromegas)
 1052 and small-strip thin gap chambers (sTGCs) [47, 48]. The micromegas is more suited to tracking because
 1053 of its good spatial resolution, while the sTGCs have better time resolution and are more suited for the
 1054 trigger. However, both systems are capable of providing tracking and trigger information. To maintain
 1055 full redundancy in cases of detector failure, both technologies will be used for tracking and trigger in the
 1056 NSW.

1057 MICROMEGAS

1058 Micromegas detectors operate using a thin metallic mesh that sits approximately $100\ \mu\text{m}$ away from the
 1059 readout electrodes to create the amplification region. Above this mesh, there is a drift region on the order
 1060 of a few mm in length capped by a drift electrode. As a charged particle traverses the detector, it ionizes gas
 1061 and the electrons drift down towards readout strips. The timing of the drift can be used to reconstruct the
 1062 angle of traversal of the particle. This is illustrated in figure 2.12. The micromegas used in ATLAS will be

1063 resistive micromegas, where the readout electrodes are topped with resistive strips [49]. This alleviates the
1064 risk of sparking in the large area detectors that ATLAS will use.

1065 In ATLAS, the micromegas drift gap will be 5 mm and the amplification gap will be $128 \mu\text{m}$. They are
1066 filled with the same gas mixture as the MDTs. They will be stacked in an octuplet in an XXUV-UVXX
1067 geometry, where X refers to straight strips and U and V refer to stereo strips at an angle of $\pm 1.5^\circ$. This
1068 arrangement allows for measurement of the azimuthal coordinate and gives a large lever arm between the
1069 straight strips for triggering purposes. Figure 2.12 shows the geometry of a single micromegas detector as
1070 well as its operating principle [47].

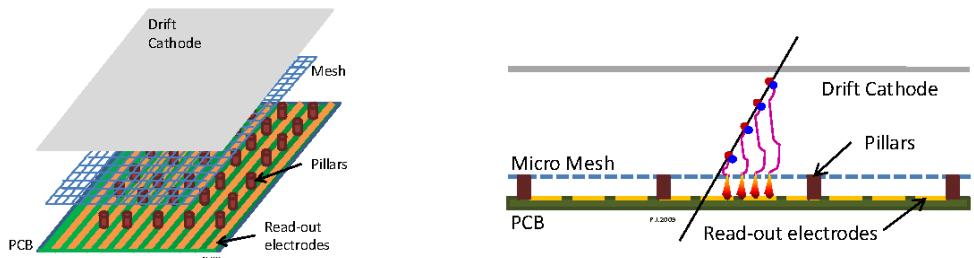


Figure 2.12: Illustrations of the geometry (left) and operating principle (right) of the micromegas detector [47].

1071 sTGCs

1072 The sTGCs are similar to the TGCs currently in the ATLAS endcap muon system [34]. They consist
1073 of gold-plated tungsten wires (with a 1.8 mm pitch) between two cathode planes 1.4 mm away from the
1074 wire plane. One cathode plane consists of strips with a 3.2 mm pitch (much smaller pitch than the TGCs),
1075 while the other consists of coarser pads that are used for defining regions of interest in the sTGC trigger
1076 algorithm. Figure 2.13 shows the basic detector geometry.

1077 2.3.3 PHYSICS IMPACT

1078 Maintaining low p_T thresholds for muons while still staying within the trigger rate budget at Level 1
1079 (20 kHz) for the muon system is crucial for physics analyses to be successful in high luminosity condi-
1080 tions. One realm where the lepton trigger threshold is especially important is in Higgs physics. In the
1081 $H \rightarrow WW^*$ analysis, one of the W bosons is off shell and tends to decay to soft leptons. In associated

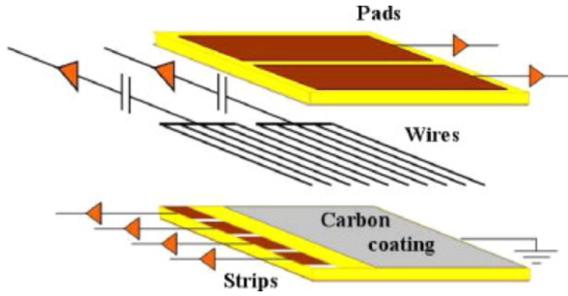


Figure 2.13: Geometry of the sTGC detector [47].

1082 production of a Higgs with a W , the lepton is also important because it provides the main handle which
 1083 allows the event to be triggered. Without the NSW, analyses would be required to either raise the muon
 1084 p_T threshold or only use muons triggered from the barrel muon system. Table 2.3 shows that both of these
 1085 alternatives significantly reduce the Higgs signal efficiency. With the NSW, the signal efficiency is largely
 1086 maintained and the triggers can remain unprescaled at lower p_T thresholds.

Threshold	$H \rightarrow bb$ (%)	$H \rightarrow WW^*$ (%)
$p_T > 20$ GeV	93	94
$p_T > 40$ GeV	61	75
$p_T > 20$ GeV (barrel only)	43	72
$p_T > 20$ GeV (with NSW)	90	92

Table 2.3: Signal efficiencies for WH production with $H \rightarrow b\bar{b}$ and $H \rightarrow WW^* \rightarrow \mu\nu qq$ under different trigger configurations [47].

1087 2.4 OBJECT RECONSTRUCTION IN ATLAS

1088 ATLAS analyses first start by requiring the presence of certain reconstructed physics objects in the event.
 1089 This section will present a brief overview of the algorithms used to reconstruct electrons, muons, jets (in-
 1090 cluding b -jets), and missing energy¹. The performance of physics object reconstruction and identification
 1091 will also be discussed as these are relevant to the analyses presented later. Figure 2.14 gives an overview of
 1092 the different sub-detectors that each type of particle will interact with in ATLAS.

¹Reconstruction algorithms for other objects, such as photons and hadronically decaying τ leptons, are not detailed here as these objects are not used in the results presented in this dissertation.

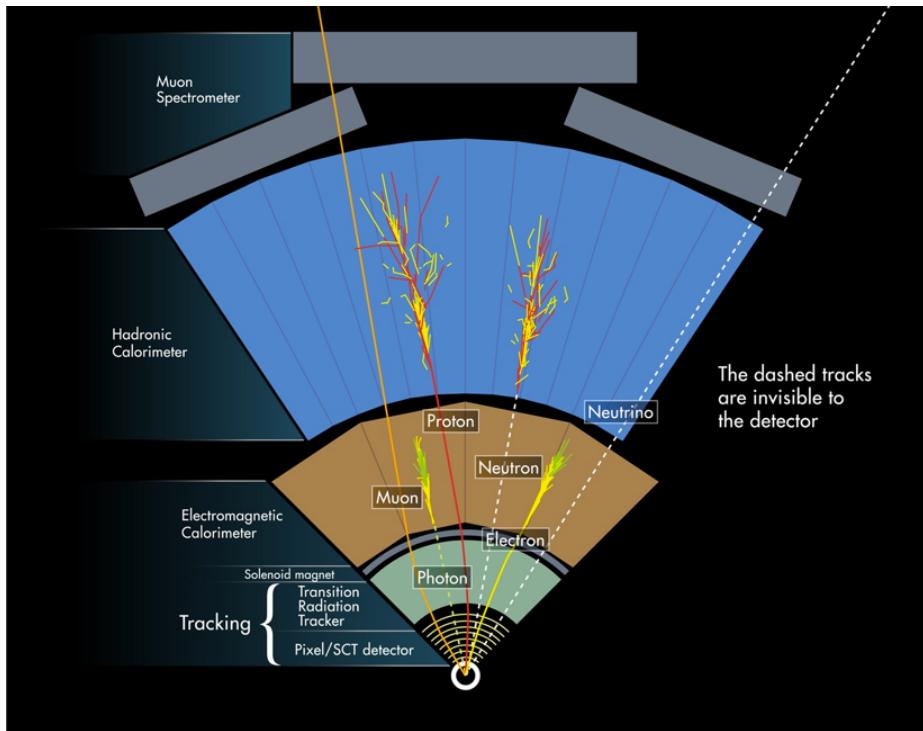


Figure 2.14: Illustration of particle interactions in ATLAS [50]

1093 2.4.I ELECTRONS

1094 Electrons in ATLAS will leave tracks in the inner detector and energy deposits in the electromagnetic
 1095 calorimeter. The algorithm for recognizing the signature of electrons proceeds in two steps: reconstruction
 1096 and identification.

1097 In reconstruction, an electron candidate is formed by matching EM calorimeter deposits with ID tracks.
 1098 The algorithm first chooses seed clusters in the EM calorimeter by using a sliding window algorithm that
 1099 searches for towers with transverse energy larger than 2.5 GeV. In addition to seed clusters, track candi-
 1100 dates must be identified in the ID. The algorithm selects seed tracks with $p_T > 1$ GeV that do not fit well
 1101 with a pion hypothesis. Once candidate tracks are selected, they are re-fit with a Gaussian Sum Filter (GSF)
 1102 algorithm to estimate electron parameters [51]. Finally, an electron candidate is formed if at least one track
 1103 matches to a seed cluster in the calorimeter. The full details of the reconstruction algorithm can be found
 1104 in reference [52].

1105 Once an electron candidate is present, identification criteria must be applied in order to reject fake elec-

trons from background. Many different variables are used for this identification. They include information about the shower shape in the EM calorimeter and the amount of leakage into the hadronic calorimeter, as well as information from the ID and in particular the TRT. There are both selection requirement based and likelihood-based criteria that range from “loose” to “very tight”. For details, see reference [52].

Figure 2.15 shows the algorithm’s reconstruction efficiency for true electrons with different identification criteria as well as the electron energy resolution in simulation [52, 53]. The reconstruction efficiency is measured using both Z and J/ψ tag and probe techniques.

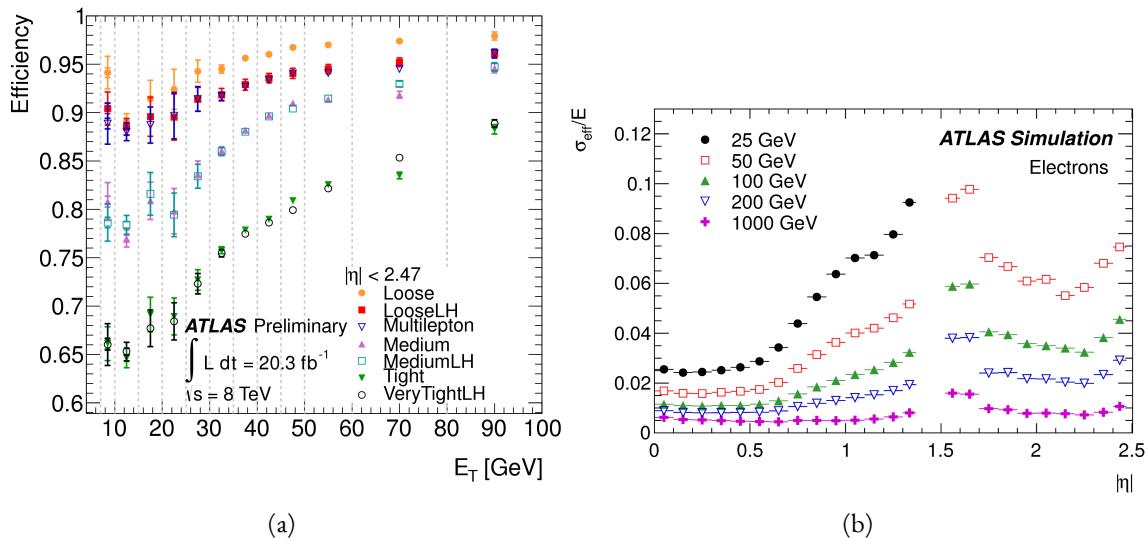


Figure 2.15: Electron performance: (a) reconstruction efficiency as a function of electron E_T [52] (b) energy resolution in simulation as a function of $|\eta|$ for different energy electrons [53].

2.4.2 MUONS

The ATLAS detector is designed to stop most particles before they reach the muon spectrometer. Muons, however, are minimum ionizing particles, meaning that they will not lose a significant amount of energy through interactions with the detector and will thus pass through. Therefore, the muon reconstruction works to match tracks in the muon spectrometer with tracks in the inner detector.

The first step of reconstruction is to reconstruct local straight line tracks, called segments, in each muon chamber. Segments are then fit to larger tracks that traverse the entire muon spectrometer. Such muon tracks are referred to as “standalone” tracks (SA) as they only use information from the muon spectrometer.

1121 The standalone tracks are then matched to tracks in the inner detector to form “combined” (CB) muons,
 1122 where the combined ID and MS fit are used to determine the momentum and direction of the muon. To
 1123 improve acceptance, segment-tagged and calorimeter-tagged muons are also reconstructed. In these cases,
 1124 ID tracks are matched to segments in the MS and calorimeter deposits consistent with a minimum ionizing
 1125 particle, respectively. The details of the reconstruction can be found in reference [54].

1126 As with electrons, once muon candidates are reconstructed they have identification criteria applied to
 1127 reduce background. These criteria include the χ^2 match between the ID and MS tracks, the number of
 1128 hits in the ID, overall ID and MS track fit quality, and additional variables [54]. The criteria range from
 1129 “loose” to “tight” as with electrons. Figure 2.16 shows the muon reconstruction efficiency (measured with
 1130 Z and J/ψ tag and probe) and invariant mass resolution [54].

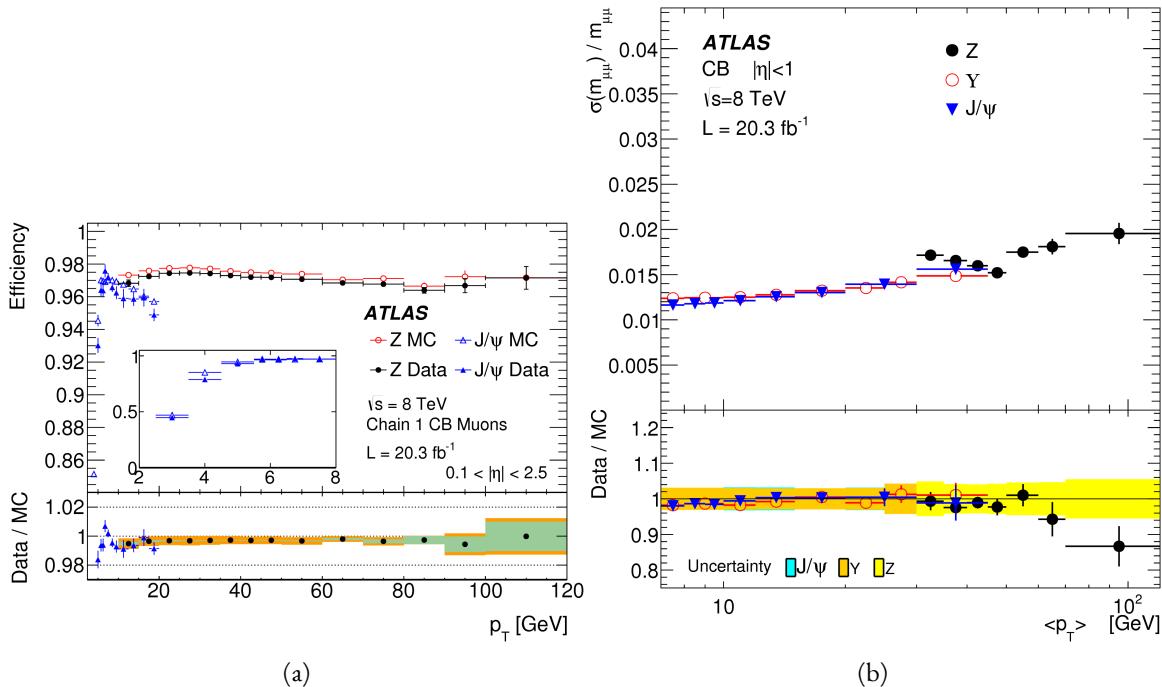


Figure 2.16: Muon performance: (a) reconstruction efficiency as a function of muon p_T (b) dimuon mass resolution as a function of average p_T [54].

1131 2.4.3 JETS

1132 When a quark or gluon is produced in collisions, it is not measured directly in ATLAS. Rather, due to
 1133 QCD effects, it produces a collimated spray of hadrons in the direction of the original parton, which is

known as a jet. Jets are reconstructed in ATLAS using energy deposits in the hadronic calorimeter. The first step is build “topological clusters” out of energy deposits in calorimeter cells [55, 56]. This is done using strategy where seed cells are chosen by picking cells whose energy measurements are four times the amount of noise expected for that cell. Adjacent cells with at least 2σ energy measurements are added to the cluster, then a final layer of clusters with energy above 0σ are added. Once calorimeter clusters are formed, they are clustered further into jet candidates using the anti- k_T jet clustering algorithm [57]. This algorithm defines a parameter R that appears in the denominator of the clustering distance metric and defines the radial size of the jet in η - ϕ space.

The energy response of the calorimeter must be properly characterized in order to reconstruct the true jet energy. Calorimeter clusters can be calibrated either with the EM calibration, where each cluster is assumed to have come from the energy deposit of an electron or photon, or the LCW calibration, where local cluster weights are computed to allow for local calibration of clusters as hadronic or electromagnetic. The details of the jet energy calibration are not discussed here and are presented in reference [58]. Figure 2.17 shows the jet energy response after calibration in Monte Carlo as a function of the true p_T of the jet [58].

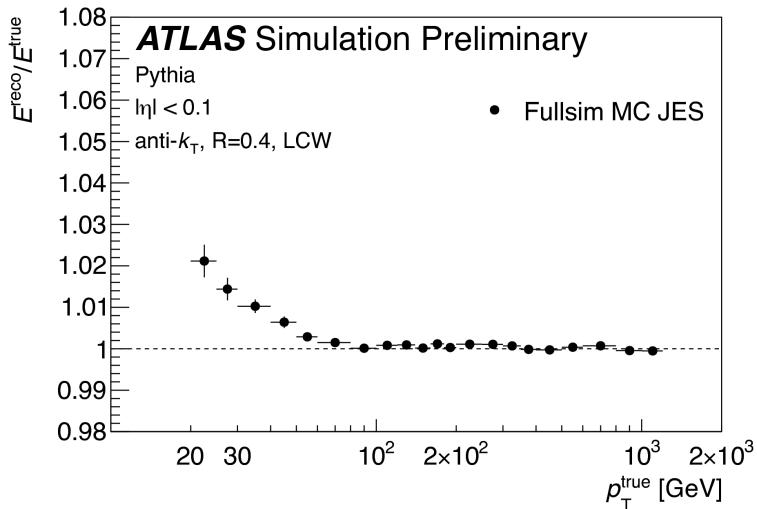


Figure 2.17: Jet energy response after calibration as a function of true p_T in simulation [58].

Analyses often need to know how consistent a particular jet is with the primary vertex of the event in order to avoid contamination from pileup interactions. One measure of this consistency is known as the

1150 jet vertex fraction (JVF). The JVF is the ratio of tracks associated with a primary vertex to the total number
1151 of tracks inside a jet. Jets from the primary interaction in the event should have a large fraction of tracks
1152 consistent with the primary vertex and therefore have a large JVF value.

1153 **2.4.4 b -TAGGING**

1154 One important aspect of jet physics is the task of identifying the flavor of parton that produced the
1155 measured jet. While in general this is very difficult, jets from b -quarks offer an interesting case where such
1156 identification is possible. B mesons have a lifetime on the order of 10^{-12} seconds, which makes a $c\tau$ on
1157 the order of millimeters [6]. This type of displaced decay vertex can be identified in detectors like ATLAS
1158 and allows b -jets to be distinguished from other flavors of jets².

1159 ATLAS uses a multivariate machine learning algorithm to identify jets from b -quarks. The inputs to this
1160 algorithm are determined from lower level reconstruction algorithms. There are three distinct algorithms
1161 that reconstruct variables which are used as input to the multivariate technique.

1162 The first family of algorithms is referred to as IPxD (where the x can either be 2 or 3). These algorithms
1163 use the transverse and longitudinal impact parameters d_0 and z_0 of the tracks inside a jet to determine
1164 their consistency with the primary vertex. They two or three dimensional (hence the x) templates for light
1165 flavor, charm, and bottom jets and then evaluate the likelihood of the jet coming from each of these types.
1166 The likelihood ratios are used as inputs to the multivariate algorithm.

1167 The next two algorithms used as input are referred to as the secondary vertex (SV) and JetFitter (JF)
1168 algorithms. The SV algorithm uses tracks inside the jet to fit for vertices that are displaced from the pri-
1169 mary vertex. The JF algorithm attempts to reconstruct the full flight path of the b by looking for multiple
1170 displaced vertices along the same line (as B decays often result in subsequent charm meson decays).

1171 In Run 1, the multivariate b -tagging algorithm used a neural network and was referred to as MV1.
1172 The details of this algorithm and its inputs are given in reference [59]. In Run 2, the number of inputs
1173 was simplified and a boosted decision tree with 24 input variables was used, referred to as MV2 [60]. The
1174 MV2 algorithm is a boosted decision tree incorporating twenty-four input variables constructed from three

²Jets from charm quarks can also be detected in this way but they do not live quite as long so the displacement of the vertex is harder to distinguish

1175 lower level input algorithms described above. Figure 2.18 summarizes the inputs to MV2. Figure 2.19 shows
1176 the performance of each of these algorithms.

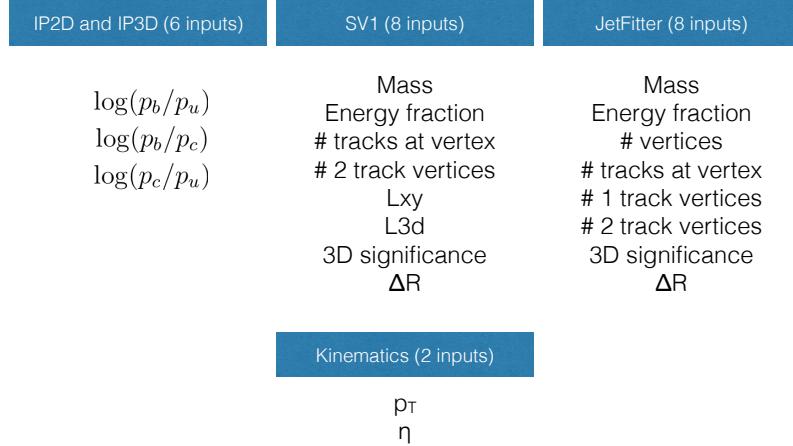


Figure 2.18: Summary of the inputs to the MV2 b -tagging algorithm.

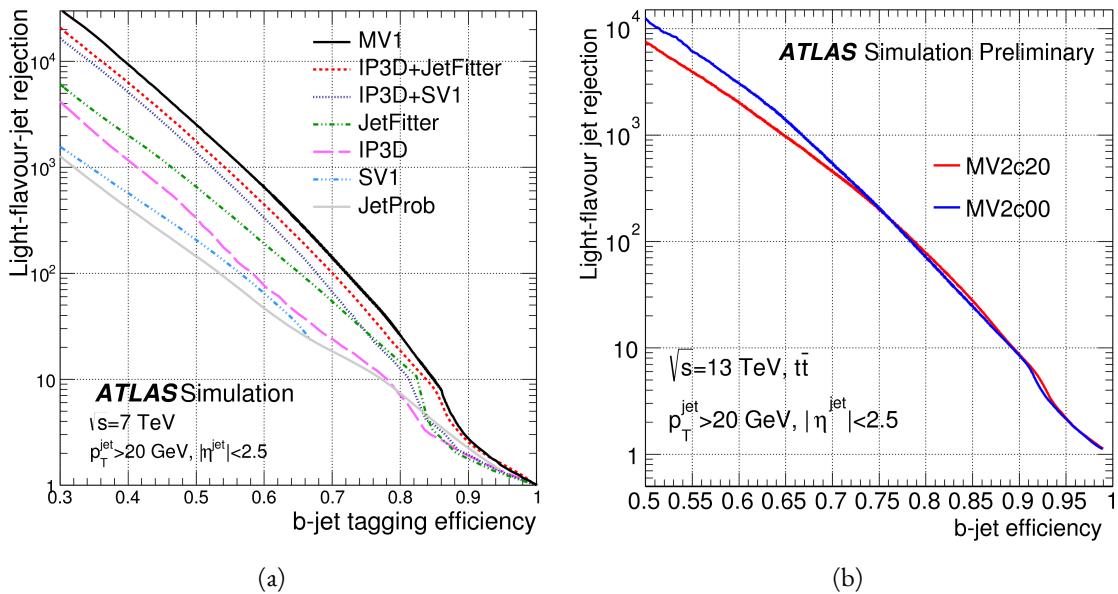


Figure 2.19: Light jet rejection (1/efficiency) vs. b -jet efficiency for MV1 and its input algorithms (a) [59] and MV2 (b) [60] in simulated $t\bar{t}$ events. The numbers in the algorithm names in (b) refer to the fraction of charm events used in the MV2 training.

1177 2.4.5 MISSING TRANSVERSE ENERGY

1178 As noted in figure 2.14, neutrinos produced in ATLAS will pass through the detector without interact-
1179 ing. The only way of detecting the presence of weakly interacting particles like neutrinos (or BSM particles
1180 that are long-lived) is to use missing transverse momentum. The basic principle of missing transverse en-
1181 ergy is to use the momentum balance of the incoming protons to infer the presence of missing particles.
1182 The net longitudinal momentum of the incoming partons that collide is not known (since each carries
1183 an unknown fraction of the proton's momentum). However, the protons (and thus incoming partons)
1184 have no net momentum in the plane transverse to the beam line (the x - y) plane. Therefore, if there are no
1185 undetected particles in the final state, the transverse momenta of all of the final state particles should bal-
1186 ance. The magnitude of the imbalance in the transverse plane is known as missing transverse momentum
1187 (E_T^{miss}).

1188 The basic calculation of missing transverse momentum from calorimeter cells is given in equation 2.4 [61].

1189

$$\begin{aligned} E_x^{\text{miss}} &= - \sum_{i=1}^{N_{\text{cell}}} E_i \sin \theta_i \cos \phi_i \\ E_y^{\text{miss}} &= - \sum_{i=1}^{N_{\text{cell}}} E_i \sin \theta_i \sin \phi_i \end{aligned} \quad (2.4)$$

1190 The E_T^{miss} calculation is separated into different terms based on the objects that the calorimeter clusters
1191 are associated with. This way, each cell's contribution is calibrated appropriately according to the object.
1192 This separation of terms is shown in equation 2.5 [61].

$$\begin{aligned} E_{x(y)}^{\text{miss,calo}} &= E_{x(y)}^{\text{miss},e} + E_{x(y)}^{\text{miss},\gamma} + E_{x(y)}^{\text{miss},\tau} + E_{x(y)}^{\text{miss,jets}} \\ &\quad + E_{x(y)}^{\text{miss,softjets}} + E_{x(y)}^{\text{miss},\mu} + E_{x(y)}^{\text{miss,CellOut}} \end{aligned} \quad (2.5)$$

1193 The CellOut term of the above equation corresponds to calorimeter cells with energy deposits that are
1194 not associated with other objects. The soft jets term comes from cells associated to jets with p_T between
1195 7 and 20 GeV, while the jets term comes from jets with $p_T > 20$ GeV. Because muons do not deposit
1196 significant energy in the calorimeter, the muon momentum is used for the muon term [61]. The final

₁₁₉₇ E_T^{miss} is calculated using equation 2.6.

$$E_T^{\text{miss}} = \sqrt{(E_x^{\text{miss}})^2 + (E_y^{\text{miss}})^2} \quad (2.6)$$

₁₁₉₈ Figure 2.20 shows the resolution of the components of the E_T^{miss} with different pileup suppression techniques [62].

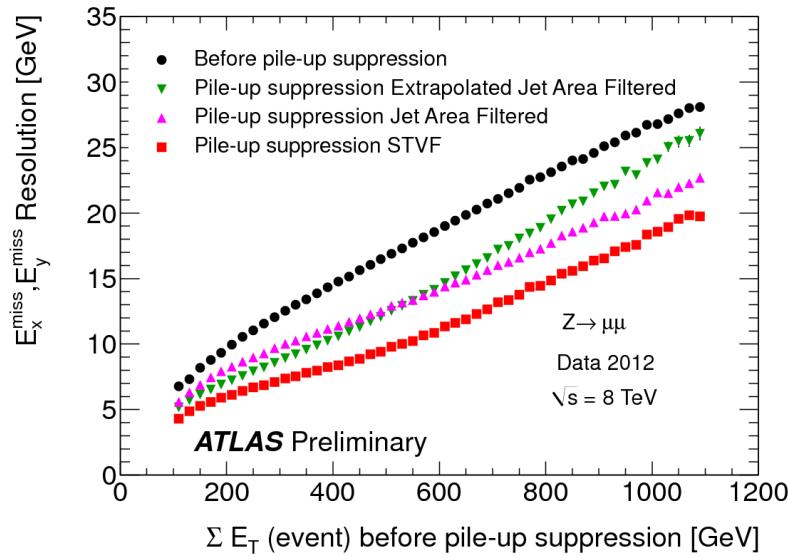


Figure 2.20: Resolution of E_T^{miss} components as a function of $\sum E_T$ before pileup suppression with different pileup techniques [62].

₁₁₉₉

1200

Part II

1201

Observation and measurement of Higgs

1202

boson decays to WW^* in LHC Run I at

1203

$\sqrt{s} = 7$ and 8 TeV

*Basic research is what I am doing when I don't know what
I am doing.*

Wernher von Braun

3

1204

1205 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ Analysis Strategy

1206 3.1 INTRODUCTION

1207 This chapter presents an overview of the strategy for searching for a Higgs boson in the $H \rightarrow WW^* \rightarrow$
1208 $\ell\nu\ell\nu$ decay topology. Its purpose is to define in broad terms how the search and measurement are under-
1209 taken, before going into details on the specific sub-categories within the larger analysis. First, the properties
1210 of the Higgs signal are discussed and the associated backgrounds are presented. Next, the observables used
1211 to enhance the signal to background ratio are defined. Finally, the parameters of interest in the search
1212 and measurement will be shown, along with a brief overview of the statistical treatment of the final Higgs
1213 candidates.

1214 Following this chapter, three different results from the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel are shown.
1215 Chapter 4 presents a search for Higgs boson production in gluon fusion mode and the role of the $H \rightarrow WW^*$
1216 channel in its discovery. Chapter 5 shows the search and first evidence in ATLAS for the Vector Boson Fu-
1217 sion (VBF) production mode of the Higgs. Finally, chapter 6 shows the combined Run 1 $H \rightarrow WW^*$

1218 results for the measurement of the Higgs cross section and relative coupling strengths to other SM parti-
1219 cles.

1220 3.2 THE $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ SIGNAL IN ATLAS

1221 The signal studied in this and subsequent chapters is the Higgs boson in the WW^* final state, where
1222 each W boson subsequently decays into a charged lepton and a neutrino. In the simplest decay path, the
1223 final state consists of two neutrinos and two charged leptons, each of which can be either an electron or a
1224 muon. If a W decays to a τ lepton, only leptonic decays of the τ are considered. The τ lepton produce
1225 additional neutrinos in the final state but still yield two charged leptons (where each lepton is an electron or
1226 muon). Neutrinos are not detected in ATLAS, so the final state ultimately consists of two reconstructed
1227 leptons and missing transverse momentum (denoted as E_T^{miss}). Final states where both of the charged
1228 leptons are electrons or muons are referred to as the “same flavor” ($ee/\mu\mu$) final states, while those with
1229 one electron and one muon are referred to as “different flavor” ($e\mu$ or μe).

1230 There can be additional jets produced in association with the Higgs boson. As described in detail in
1231 Chapter 1, if the Higgs is produced via vector boson fusion production, there will be two additional forward
1232 jets in the event. In gluon fusion, one or more jets can be produced through initial state radiation from
1233 the incoming gluons. Because of the varying background composition as a function of jet multiplicity,
1234 each bin in this variable has its own dedicated requirements applied in the search and measurement. The
1235 $n_j = 0$ and $n_j = 1$ bins are dedicated to gluon fusion production, while the $n_j \geq 2$ bin has separate
1236 dedicated searches for ggF and VBF production.

1237 Figure 3.1 shows the relative branching fractions for the $H \rightarrow WW^*$ process, calculated from the Par-
1238 ticle Data Group values for the W and τ branching ratios [6]. The largest branching ratio corresponds
1239 to both W bosons decaying to quark pairs at 45.44%. The second largest ratio is for one W decaying lep-
1240 tonically and the other decaying to quarks, a branching ratio of 34.18%. In all cases, ℓ denotes either an
1241 electron or muon, and the leptonic branching ratios of the τ are included. For example, the $\ell\nu qq$ final
1242 state includes one W decaying to $e\nu$, $\mu\nu$, or $\tau\nu$. In the case of the $W \rightarrow \tau\nu$ decay, the τ lepton then
1243 decays to an electron or muon via $\tau \rightarrow \nu_\tau \ell \nu_\ell$. Final states with a τ_h refer to hadronic decays of the τ . The

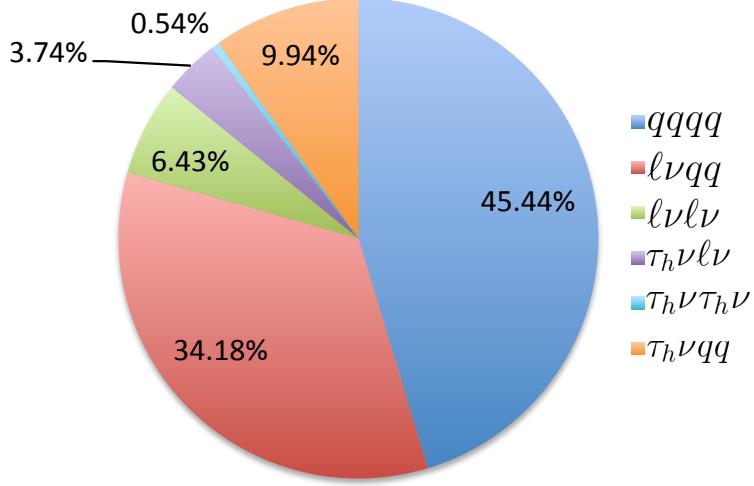


Figure 3.1: Branching ratios for a WW system. q refers to quarks. ℓ can be either an electron or muon, and the leptonic branching ratios of the τ are included. For example, the $\ell\nu qq$ final state includes one W decaying to $e\nu$, $\mu\nu$, or $\tau\nu$. τ_h refer to hadronic decays of the τ .

branching ratio of the $\ell\nu\ell\nu$ final state is 6.43%.

While the $\ell\nu\ell\nu$ final state is not a large fraction of the branching ratio, there are significant advantages to using this channel in an analysis. First, both the $qqqq$ and $\ell\nu qq$ channels suffer from a large QCD multijet background, which is often difficult to model. Second, events in the the $\ell\nu\ell\nu$ channel in data can be triggered more efficiently due to the presence of two leptons.

3.3 BACKGROUND PROCESSES

Many processes from the Standard Model can also produce a final state with two leptons and missing transverse momentum. This section describes the dominant backgrounds to Higgs production and further explains how they can be reduced. Table 3.1 summarizes the different background processes.

3.3.1 STANDARD MODEL WW PRODUCTION

Non-resonant Standard Model diboson production, as shown in figure 3.2, is an irreducible background to Higgs boson production in the WW final state. It produces the same exact final state objects, namely

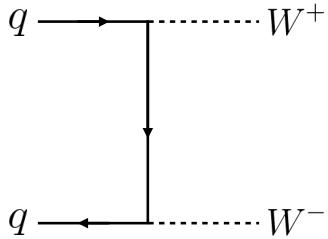


Figure 3.2: Feynman diagram for Standard Model WW production

leptonically decaying W bosons. There are no additional objects in the final state that allow for background reduction. Therefore the analysis solely relies on the correlations between the leptons to reduce this background.

3.3.2 TOP QUARK PRODUCTION

Top quark production can mimic the Higgs in the WW^* final state as well. Top quarks can be produced either in pairs ($t\bar{t}$ production) or singly (s -channel, t -channel, or associated production Wt). The dominant top background are $t\bar{t}$ and Wt production.

Because top quarks decay via $t \rightarrow Wb$, top pair production can produce a final state with two W bosons that then decay leptonically. In Wt production, there are two real W bosons produced, as with $t\bar{t}$. In both cases, there is at least one b -jet in the final state. By vetoing on the presence of b -jets, these top quark backgrounds can be reduced. Figure 3.3 shows the Feynman diagrams for $t\bar{t}$ and Wt production.

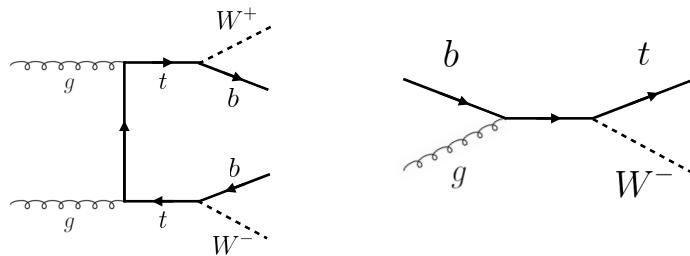


Figure 3.3: Feynman diagrams for top pair production (left) and Wt production (right)

1267 3.3.3 W +JETS BACKGROUND

1268 Single W boson production in association with jets is a unique background to Higgs production. The
1269 other backgrounds considered thus far have all included two prompt leptons, each decaying from a W
1270 boson, in the final state. In W +jets production, however, only one reconstructed lepton originates from
1271 a W . The second reconstructed lepton is either an algorithmic “fake” or the result of non-prompt decays.
1272 In the first case, the lepton is a jet misidentified as a lepton by either the electron or muon reconstruction
1273 algorithms. In the second case, the lepton may be a real lepton but coming from semi-leptonic decays of
1274 particles inside the shower of the jet. This background can be reduced by requiring that the reconstructed
1275 lepton have little activity in the surrounding region of the calorimeter (also known as an “isolation”). Fig-
1276 ure 3.4 shows the Feynman diagram for W +jets production.

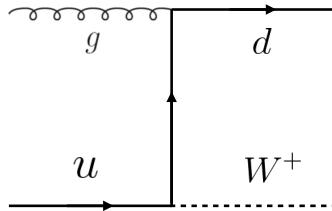


Figure 3.4: An example Feynman diagram of W +jets production

1277 3.3.4 Z/γ^* +JETS BACKGROUND

1278 Production of a Z boson or virtual photon (also known as Drell-Yan and denoted with Z/γ^*) in as-
1279 sociation with jets is also a background to Higgs production. The Z boson decays to two leptons of the
1280 same flavor. When the Z/γ^* decays directly to electrons or muons, the background enters the same flavor
1281 final state. When the Z decays to two τ leptons the background can enter the different flavor final state as
1282 well. Figure 3.5 shows the production of a Z in association with one jet. Because there are no neutrinos in
1283 this final state, variables like E_T^{miss} can be used to reduce the background¹.

¹The E_T^{miss} cut is much more effective for the reduction of Z/γ^* production in the same flavor final state. If the background enters the different flavor final state through τ decays, there will be neutrinos present. Other requirements on the lepton invariant mass are made to reduce the $Z/\gamma^* \rightarrow \tau\tau$ background.



Figure 3.5: An example Feynman diagram of Z +jets production

1284 3.3.5 SUBDOMINANT BACKGROUNDS

1285 There are additional processes which contribute to the background composition. These backgrounds
 1286 are subdominant and contribute less to the total background estimate than those discussed previously.
 1287 The first process is referred to as VV or “Other diboson” processes and includes multiple Standard Model
 1288 diboson processes, including WZ , ZZ , $W\gamma$, $W\gamma^*$, and $Z\gamma$ production. Additionally, there is a back-
 1289 ground contribution from QCD multijet production. While the cross section for this process is large, its
 1290 contribution to the WW^* final state is small because two jets must be misidentified as leptons.

Category	Process	Description
SM WW	$WW \rightarrow \ell\nu\ell\nu$	Real leptons and neutrinos
Top quark production	$t\bar{t} \rightarrow WbWb \rightarrow \ell\nu b\bar{b}\nu\bar{b}$	Real leptons, untagged b s
	$tW \rightarrow WbW \rightarrow \ell\nu\ell\nu b$	Real leptons, untagged b
	$t\bar{b}, t\bar{q}\bar{b}$	Untagged b , jet misidentified as lepton
Drell-Yan	$Z/\gamma^* \rightarrow ee, \mu\mu$	“Fake” E_T^{miss}
	$Z/\gamma^* \rightarrow \tau\tau \rightarrow \ell\nu\ell\nu\nu$	Real leptons and neutrinos
Other dibosons	$ZZ \rightarrow \ell\ell\nu\nu$	Real leptons and neutrinos
	$W\gamma^*, WZ \rightarrow \ell\nu\ell\ell, ZZ \rightarrow \ell\ell\ell\ell$	Unreconstructed leptons
	$W\gamma, Z\gamma$	γ reconstructed as e , unreconstructed lepton
W +jets	$Wj \rightarrow \ell\nu j$	Jet reconstructed as lepton
QCD multijet	jj	Jets reconstructed as leptons

Table 3.1: A summary of backgrounds to the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ signal

1291 3.4 SHARED SIGNAL REGION SELECTION REQUIREMENTS

1292 As presented in section 3.2, there are many different combinations of physics objects that can define a
 1293 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ final state. The multiplicity of jets and the flavor combinations of the leptons

1294 both lead to many potential signal regions. Additionally, signal regions can be optimized separately to be
 1295 sensitive to the distinct production modes of the Higgs. Gluon fusion, vector boson fusion, and associated
 1296 production of a Higgs all lead to unique final state topologies. Figure 3.6 delineates the different signal
 1297 regions used in the gluon fusion and vector boson fusion $H \rightarrow WW^*$ analyses. While there are different
 1298 optimizations possible in each signal region, there are also some commonly shared selections that will be
 1299 described here.

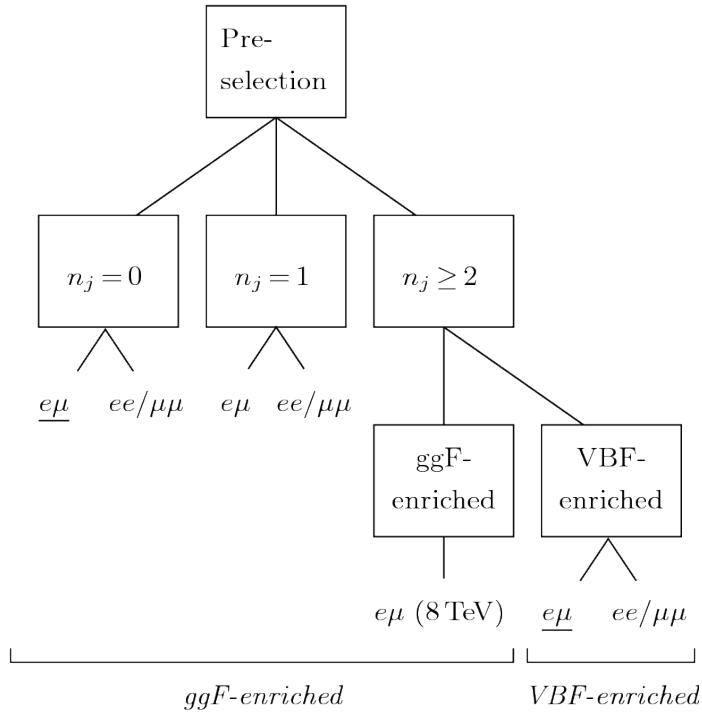


Figure 3.6: An illustration of the unique analysis signal regions [63]. The most sensitive regions for both gluon fusion and vector boson fusion production are underlined.

1300 3.4.I EVENT PRE-SELECTION

1301 Before being sorted into the distinct signal regions, basic requirements are applied to the reconstructed
 1302 objects in the event to select Higgs-like event candidates. First, two oppositely charged leptons are required.
 1303 Once the leptons are selected, the last requirement for event pre-selection is the presence of neutrinos. As
 1304 neutrinos cannot be detected directly in ATLAS, E_T^{miss} can be used as a proxy for the combined neutrino
 1305 momentum in the transverse plane.

1306 In general, it is expected that the signal should have a harder E_T^{miss} spectrum than backgrounds, espe-
 1307 cially if these backgrounds do not contain neutrinos in the final state. When using E_T^{miss} , it is possible
 1308 mis-measurements of objects in the detector can lead to imbalances in the transverse plane. When such a
 1309 mis-measurement occurs, the E_T^{miss} vector in the transverse plane will often point in the same direction as
 1310 the mis-measured object. Therefore, a new variable, $E_{T,\text{rel}}^{\text{miss}}$, is used in the pre-selection. $E_{T,\text{rel}}^{\text{miss}}$ is defined
 1311 in equation 3.1.

$$E_{T,\text{rel}}^{\text{miss}} = \begin{cases} E_T^{\text{miss}} \sin \Delta\phi_{\text{near}} & \text{if } \Delta\phi_{\text{near}} < \pi/2 \\ E_T^{\text{miss}} & \text{otherwise,} \end{cases} \quad (3.1)$$

1312 If the closest object to the E_T^{miss} vector is within $\pi/2$ radians in the transverse plane, the E_T^{miss} is projected
 1313 away from this object. Otherwise, the normal E_T^{miss} vector is used. Figure 3.7 shows a graphical illustration
 of this concept.

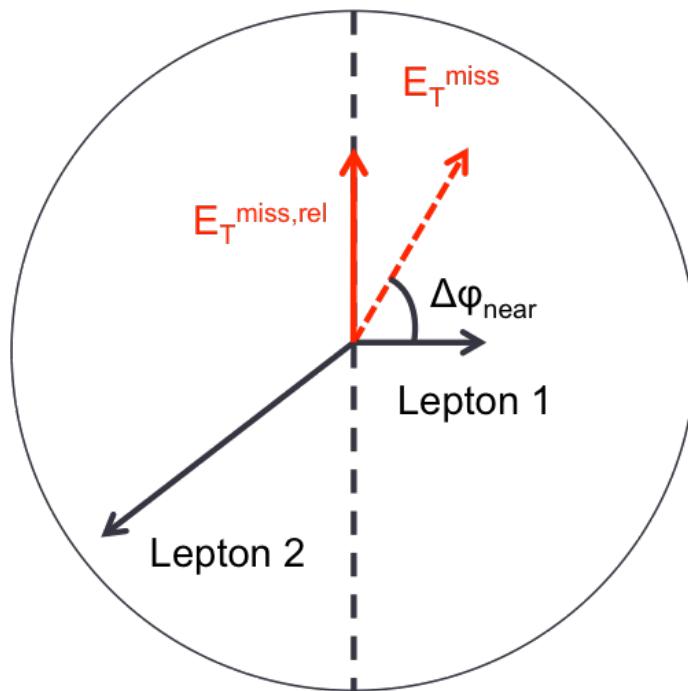


Figure 3.7: A graphical illustration of the $E_{T,\text{rel}}^{\text{miss}}$ calculation

1314
 1315 Once the lepton and E_T^{miss} pre-selections are made, the analysis is divided into different regions accord-
 1316 ing to jet multiplicity.

1317 3.4.2 JET MULTIPLICITY

1318 Jet multiplicity, denoted as n_j , is used to sub-divide the analysis into distinct signal regions. By creating
 1319 separate signal regions, each bin in jet multiplicity becomes sensitive to different modes of Higgs produc-
 1320 tion and different backgrounds.

1321 For example, the $n_j \geq 2$ region is more sensitive to VBF production because of the two high momen-
 1322 tum jets produced at matrix element level. For gluon fusion production to enter this bin, two initial state
 1323 radiation jets must be emitted.

1324 Figure 3.8 shows the jet multiplicity in both the different flavor and same flavor regions after the pre-
 1325 selection. It also shows the background composition in the bins of n_b . A few trends from this distribution
 1326 are worth noting. The first is that the Drell-Yan background dominates in the same flavor channels for
 1327 $n_j \leq 1$. Second, the top background becomes a clear contributor to the total background for $n_j \geq$
 1328 1. Lastly, the SM WW production dominates in the $n_j = 0$ bin, as it is an irreducible background to
 1329 $H \rightarrow WW^*$ production. Because of these distinct features, each jet multiplicity bin is treated separately.

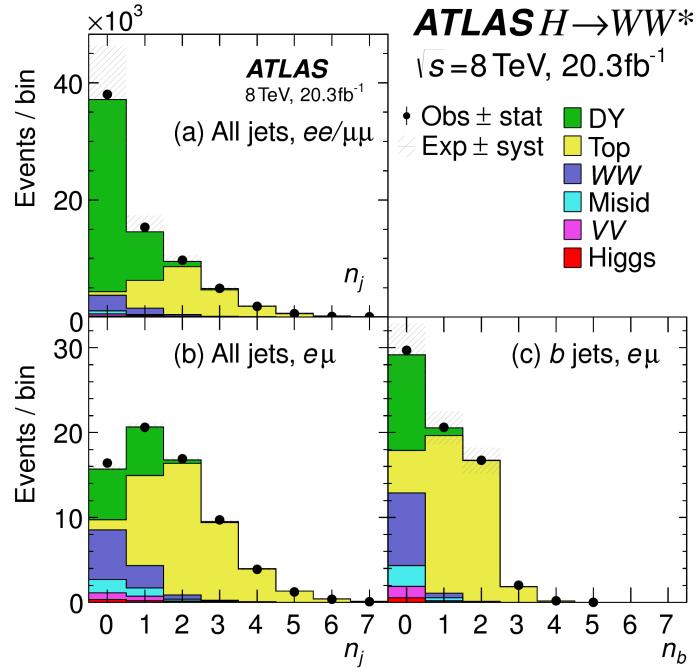


Figure 3.8: Predicted backgrounds (compared with data) as a function of n_j (a and b) and n_b (c) after pre-selection requirements

1330 3.5 BACKGROUND REDUCTION IN SAME-FLAVOR FINAL STATES

1331 As described in section 3.4.2, the background composition of the same flavor final states is different
1332 from that of the different flavor states. In particular, Drell Yan processes play a much larger role because
1333 the Z/γ^* decays to same flavor leptons. Because real neutrinos are absent in the Z/γ^* decays to ee and $\mu\mu$,
1334 a requirement on E_T^{miss} should largely reduce the background. However, as this section will demonstrate,
1335 with increasing pileup conditions the resolution of the calorimeter-based E_T^{miss} degrades greatly. There-
1336 fore, two new variables for Z/γ^* background reduction are constructed and described in this section.

1337 3.5.1 PILEUP AND E_T^{miss} RESOLUTION

1338 Secondary interactions of protons in the colliding bunches of the LHC (known as pileup interactions,
1339 described in detail in Chapter 2) deposit energy into the ATLAS calorimeter in addition to the energy that
1340 comes from the hard scatter process of interest. The calculation of E_T^{miss} is fundamentally Poissonian.
1341 Summing up all of the energy deposits in individual calorimeter cells or clusters is similar to a counting
1342 experiment. The error on a mean of N in a Poisson distribution is \sqrt{N} , so the energy resolution scales
1343 as \sqrt{E} . As more energy is deposited in the calorimeter, the E_T^{miss} resolution degrades, meaning that the
1344 E_T^{miss} resolution is particularly sensitive to LHC instantaneous luminosity conditions.

1345 Figure 3.9 shows an event display of a $Z/\gamma^* + \text{jets}$ event candidate with the twenty-five reconstructed
1346 primary vertices. This display illustrates that while the interaction of interest only has tracks coming from
1347 the hardest primary vertex, all of the secondary interactions deposit energy in the calorimeter as well.

1348 Figure 3.10 shows the RMS of the E_T^{miss} distribution in $Z \rightarrow \mu\mu$ events (where there are no real neu-
1349 trinos) as a function of the number of the average number of interactions. Under 2011 LHC conditions,
1350 this RMS was approximately 9 GeV, while under 2012 running conditions the resolution worsened to 12
1351 GeV. The increase in pileup dilutes the E_T^{miss} variable's ability to reduce the Z/γ^* background.

1352 3.5.2 TRACK-BASED DEFINITIONS OF MISSING TRANSVERSE MOMENTUM

1353 Because the increasing number of secondary proton-proton interactions degrades calorimeter-based
1354 E_T^{miss} resolution, a new variable using only contributions from the primary interaction vertex is necessary

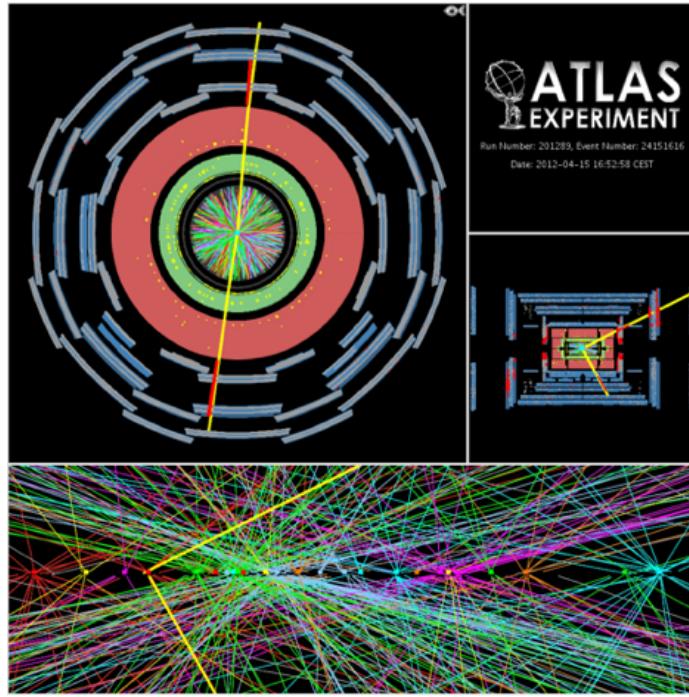


Figure 3.9: An event display of a Z/γ^* + jets event illustrating the effect of pileup interactions

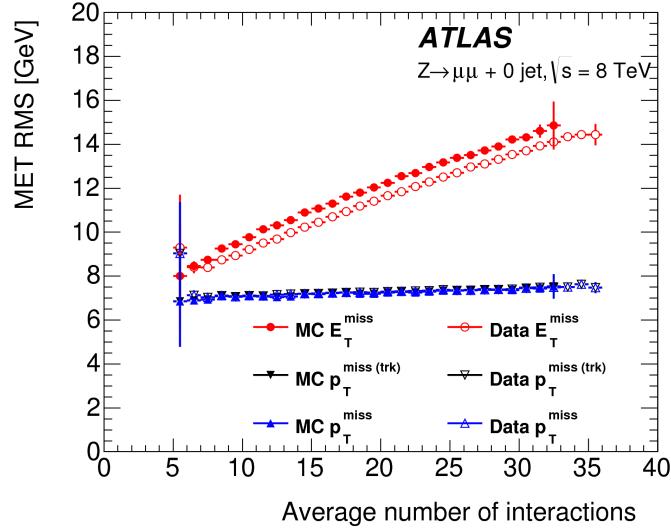


Figure 3.10: The RMS of different missing transverse momentum definitions as a function of the average number of interactions per bunch crossing

to further reduce the Z/γ^* background. While it is not possible to associate calorimeter energy deposits with a particular vertex, individual charged particle tracks in the Inner Detector are associated to unique vertices. Thus, two track-based definitions of missing transverse momentum , using only tracks coming

1358 from the primary vertex in the event, are used in the analysis. The simplest variable, $p_T^{\text{miss}(\text{trk})}$, is the vec-
 1359 torial sum of the p_T of all of the tracks from the primary vertex and the selected leptons (excluding the
 1360 tracks associated with the selected leptons to avoid double counting). Equation 3.2 defines $p_T^{\text{miss}(\text{trk})}$.

$$p_T^{\text{miss}(\text{trk})} = - \left(\sum_{\text{selected leptons}} p_T + \sum_{\text{other tracks}} p_T \right), \quad (3.2)$$

1361 To further improve the resolution on the missing transverse momentum, the variable p_T^{miss} is used as de-
 1362 fined in equation 3.3. For selected leptons and jets, the nominal p_T measurements are used. Tracks are used
 1363 to estimate the soft component of the missing transverse momentum instead of calorimeter measurements.

1364

$$p_T^{\text{miss}} = - \left(\sum_{\text{selected leptons}} p_T + \sum_{\text{selected jets}} p_T + \sum_{\text{other tracks}} p_T \right), \quad (3.3)$$

1365 Figure 3.10 illustrates that these two new variables accomplish their intended purpose. The resolution as a
 1366 function of mean number of interactions for both $p_T^{\text{miss}(\text{trk})}$ and p_T^{miss} is much flatter than the dependence
 1367 for E_T^{miss} . Figure 3.11a shows the difference between the true and reconstructed values of missing transverse
 1368 momentum using both the track-based p_T^{miss} and calorimeter based E_T^{miss} . The RMS of the distribution
 1369 improves by 3.5 GeV when using p_T^{miss} .

1370 3.5.3 DISTINGUISHING Z/γ^* +JETS AND $H \rightarrow WW^*$ TOPOLOGIES

1371 In addition to measuring missing transverse momentum, another variable can be constructed to exploit
 1372 kinematic and topological differences between the Z/γ^* background and $H \rightarrow WW^*$ signal. Because
 1373 there are no real neutrinos in the final state (in the case of $Z/\gamma^* \rightarrow ee, \mu\mu$ decays), the dilepton system will
 1374 be balanced with the jets produced in the hard scatter. A new variable, f_{recoil} , is constructed to estimate
 1375 the balance between the dilepton system and recoiling jets and is defined in equation 3.4. The transverse
 1376 plane is divided into four sections, or quadrants, with one quadrant centered on the dilepton vector. The
 1377 numerator of f_{recoil} is the magnitude of the vectorial sum of the p_T of jets in the quadrant opposite the
 1378 dilepton system, weighted by each jet's Jet Vertex Fraction (JVF, described in chapter 2). The denominator

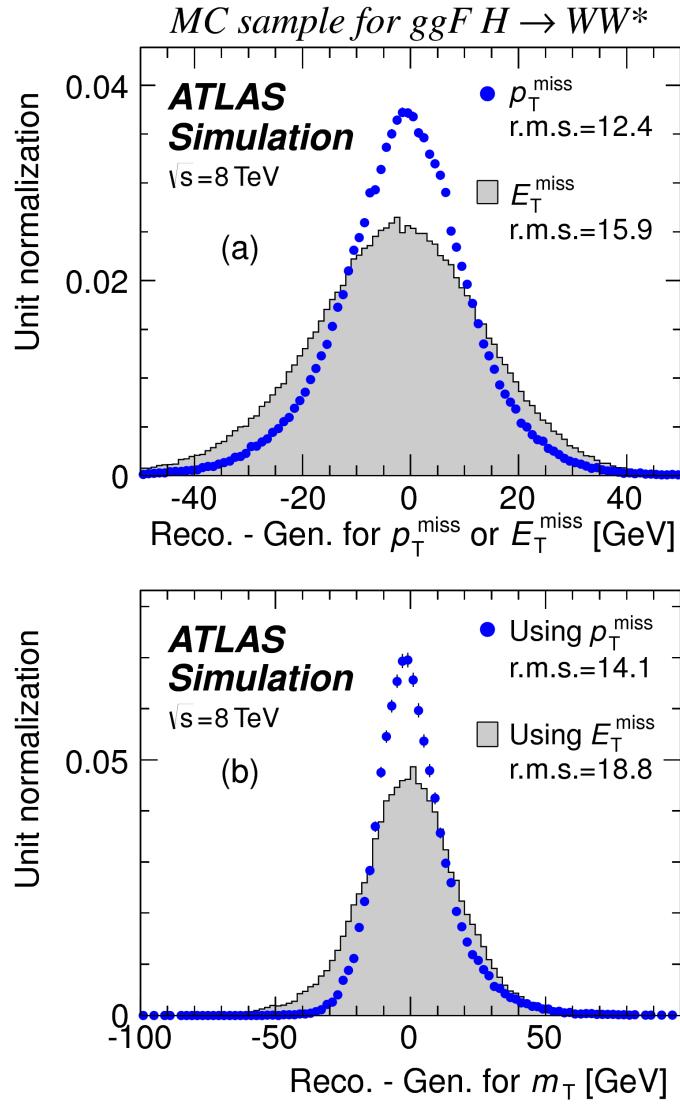


Figure 3.11: The difference between the true and reconstructed values of the missing transverse momentum (a) and m_T (b) in a gluon fusion signal sample

¹³⁷⁹ is the magnitude of the dilepton p_T .

$$f_{\text{recoil}} = \left| \sum_{\text{jets } j \text{ in } \wedge} \text{JVF}_j \cdot \mathbf{p}_T^j \right| / p_T^{\ell\ell}. \quad (3.4)$$

¹³⁸⁰ Figure 3.12 shows a shape comparison of the f_{recoil} distribution in a simulated $Z/\gamma^* + \text{jets}$ sample, a
¹³⁸¹ $H \rightarrow WW^*$ signal sample, and other backgrounds that contain real neutrinos. The $Z/\gamma^* + \text{jets}$ events

1382 tend to be more balanced between the dilepton system and recoiling jets, while the processes containing
 1383 real neutrinos are less balanced in the transverse plane. Thus, a requirement on f_{recoil} will reduce the Z/γ^*
 1384 + jets background while maintaining a good signal efficiency.

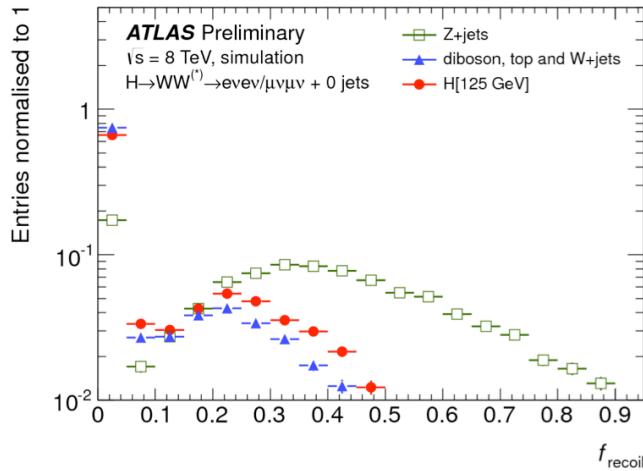


Figure 3.12: Comparison of f_{recoil} distributions for $Z/\gamma^* + \text{jets}$, $H \rightarrow WW^*$, and other backgrounds with real neutrinos.

1385 3.5.4 OPTIMIZING BACKGROUND REDUCTION SELECTION REQUIREMENTS

1386 The requirements on $p_T^{\text{miss}(\text{trk})}$ and f_{recoil} used to reduce the $Z + \text{jets}$ background must be optimized
 1387 to maximize expected signal significance in the same flavor channels. Figure 3.13 shows an optimization of
 1388 the combination of the two requirements in the gluon fusion zero jet bin. Each bin shows the expected
 1389 signal significance if the $p_{T,\text{rel}}^{\text{miss}(\text{trk})}$ is required to be greater than the left edge of the bin and the f_{recoil} is
 1390 required to be less than the top edge of the bin. The figure shows that the best signal significance comes
 1391 from requiring low values of f_{recoil} (< 0.05) and $p_{T,\text{rel}}^{\text{miss}(\text{trk})}$ values greater than 45 GeV.

1392 3.6 PARAMETERS OF INTEREST AND STATISTICAL TREATMENT

1393 As with any search or measurement, there are particular parameters of the Higgs that the $H \rightarrow WW^*$
 1394 analysis is interested in measuring. In this case, the parameters of interest are the mass of the Higgs boson
 1395 and its production cross section. Because the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ process does not have a closed final

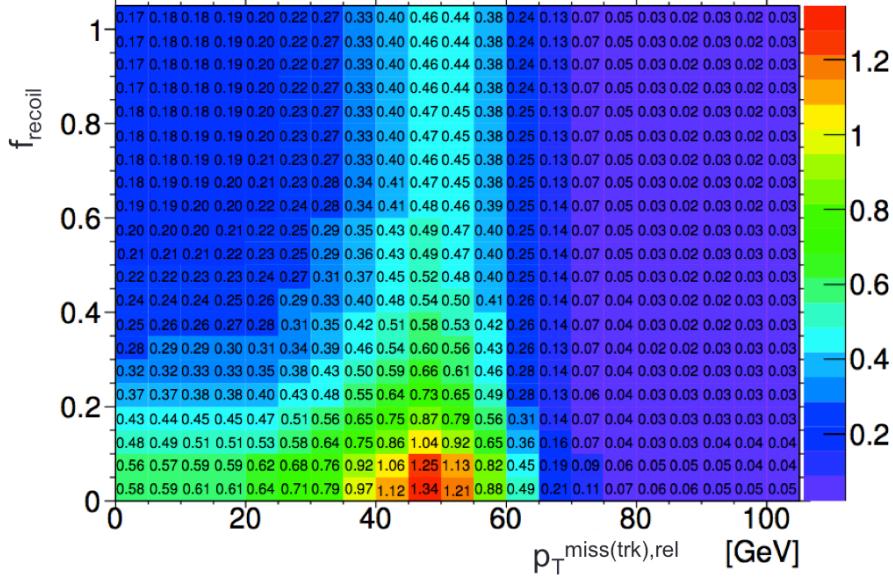


Figure 3.13: Signal significance as a function of required value for f_{recoil} and $p_{T,\text{rel}}^{\text{miss}(\text{trk})}$ in the ggF $H \rightarrow WW^*$ with $n_j = 0$

1396 state, it is not possible to measure the full invariant mass of the particle that may have produced the final
 1397 state. However, a proxy for the invariant mass is defined using transverse plane information and detailed
 1398 in section 3.6.1. The second parameter of interest is the ratio of the measured cross section to that expected
 1399 from the Standard Model Higgs, which is denoted a μ . This is defined in equation 3.5.

$$\mu = \frac{\sigma}{\sigma_{\text{SM}}} \quad (3.5)$$

1400 All of the likelihoods used in the statistical analysis of the final signal region events are parameterized as a
 1401 function of μ . μ is a natural variable for hypothesis testing, as $\mu = 0$ corresponds to a background only
 1402 hypothesis and $\mu = 1$ corresponds exactly to a Standard Model Higgs.

1403 3.6.1 TRANSVERSE MASS

1404 The $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis cannot reconstruct the full invariant mass of the Higgs because
 1405 of the neutrinos in the final state. The transverse mass serves as a proxy for the full invariant mass by

1406 exploiting information from the transverse plane. The transverse mass is defined in equation 3.6.

$$m_T = \sqrt{(E_T^{\ell\ell} + p_T^{\text{miss}})^2 - |\vec{p}_T^{\ell\ell} + \vec{p}_T^{\text{miss}}|^2}, \quad (3.6)$$

1407 Here the $E_T^{\ell\ell}$ and $p_T^{\ell\ell}$ are the transverse energy and momentum of the dilepton system, while p_T^{miss} is a
1408 proxy for the transverse momentum of the di-neutrino system. The track-based p_T^{miss} is used in the m_T
1409 rather than the calorimeter based E_T^{miss} because it has a better resolution on the true transverse mass.
1410 Figure 3.11b shows the improvement in the RMS of the difference between the true and reconstructed
1411 transverse mass in a ggF signal sample. The RMS improves by 4.7 GeV using p_T^{miss} in the m_T calculation.

1412 **3.6.2 STATISTICAL TREATMENT²**

1413 **LIKELIHOOD FUNCTION**

1414 The statistical analysis of final event candidates is framed as a hypothesis test, where the null hypoth-
1415 esis is background-only (no Standard Model Higgs). The first step in the analysis is to form a likelihood
1416 function for the data. In its simplest form, this likelihood is the probability of observing the number of
1417 events seen in the final signal region given knowledge of the signal strength. Because observation of events
1418 is fundamentally a Poisson counting experiment, this simple likelihood can be expressed as a Poisson prob-
1419 ability of observing N events given a total number of predicted signal and background events. This basic
1420 likelihood is shown in equation 3.7.

$$\mathcal{L}(\mu) = P(N|\mu S + B) \quad (3.7)$$

1421 Here, P is the Poisson probability density function, N is the total number of observed events, μ is the
1422 signal strength, S is the predicted number of signal events, and B is the predicted number of background
1423 events.

1424 In particle physics, certain background estimates are commonly normalized in so-called “control” re-
1425 gions and those predictions are scaled by the same normalization factor in the signal region. This leads to a

²Many thanks to Aaron Armbruster, whose thesis [64] inspired parts of this section.

1426 slightly more complicated likelihood, which is a function of both the signal strength and the background
1427 normalization. This is shown in equation 3.8.

$$\mathcal{L}(\mu, \theta) = P(N|\mu S + \theta B) P(N_{\text{CR}}|\theta B_{\text{CR}}) \quad (3.8)$$

1428 Here, θ serves as a “nuisance parameter”, or a parameter that is not of primary interest but still enters the
1429 likelihood. The second Poisson term enforces that the background normalization be consistent with the
1430 number of observed events in data in the control region, N_{CR} .

1431 So far, these two formulations of likelihoods have assumed a single signal region and do not take into
1432 account any shape information of potential discriminating variables. The $H \rightarrow WW^*$ analysis is divided
1433 into many different categories, the counting experiment described above can be performed in each individ-
1434 ual category. As mentioned in section 3.6.1, the transverse mass is used as the primary discriminating vari-
1435 able in many of the $H \rightarrow WW^*$ sub-analyses. The same counting experiment can be performed in each
1436 bin of the m_T distribution to incorporate some shape information. Thus, the total likelihood becomes a
1437 product over signal regions and bins of the m_T distribution. Finally, there are usually many backgrounds
1438 that are normalized in control regions. The new formulation of the likelihood takes this into account by
1439 including a product over control regions in the second Poisson term. All of these modifications are shown
1440 in equation 3.9.

$$\mathcal{L}(\mu, \theta) = \prod_{\substack{\text{SRs } i \\ \text{bins } b}} P\left(N_{ib} \middle| \mu S_{ib} + \sum_{\text{bkg } k} \theta_k B_{kib}\right) \prod_{\text{CRs } l} P\left(N_l \middle| \sum_{\text{bkg } k} \theta_k B_{kl}\right) \quad (3.9)$$

1441 The final step to obtain the full likelihood used in the analysis is to add nuisance parameters for the
1442 systematic uncertainties. In cases where the uncertainty does not affect the shape of m_T bin-by-bin, each
1443 systematic uncertainty ϵ is allowed to affect the expected event yields through a linear response function
1444 of the nuisance parameter, namely $\nu(\theta) = (1 + \epsilon)\theta$. If instead the uncertainty does affect the shape, the
1445 effect is instead parameterized by $\nu_b(\theta) = 1 + \epsilon_b\theta$. The value of the nuisance parameters for the systematic
1446 uncertainty are constrained with a Gaussian term that is added to the likelihood as well. This is of the form
1447 $g(\delta|\theta) = e^{-(\delta-\theta)^2/2}/\sqrt{2\pi}$, where δ is the central value and θ is a nuisance parameter. Finally, a last term is

¹⁴⁴⁸ added to account for the statistical uncertainty in the Monte Carlo samples used, which adds an additional
¹⁴⁴⁹ poisson term. The full likelihood used in the final statistical analysis is defined in equation 3.10.

$$\begin{aligned} \mathcal{L}(\mu, \boldsymbol{\theta}) = & \prod_{\substack{\text{SRs i} \\ \text{bins b}}} P \left(N_{ib} \middle| \mu S_{ib} \cdot \prod_{\substack{\text{sig.} \\ r}} \nu_{br}(\theta_r) + \sum_{\text{bkg k}} \theta_k B_{kib} \cdot \prod_{\substack{\text{bkg.} \\ s \\ \text{syst.}}} \nu_{bs}(\theta_s) \right) \\ & \cdot \prod_{\text{CRs l}} P \left(N_l \middle| \sum_{\text{bkg k}} \theta_k B_{kl} \right) \\ & \cdot \prod_{\substack{\text{syst} \\ t}} g(\delta_t | \theta_t) \cdot \prod_{\text{bkg k}} P(\xi_k | \zeta_k \theta_k) \end{aligned} \quad (3.10)$$

¹⁴⁵⁰ The fourth term of the equation quantifies the uncertainty due to finite Monte Carlo sample size. Here,
¹⁴⁵¹ ξ represents the central value of the background prediction, θ is the associated nuisance parameter, $\zeta =$
¹⁴⁵² $(B/\delta B)^2$, where δB is the statistical uncertainty of B .

¹⁴⁵³ The best fit value of the signal strength μ is determined by finding the values of μ and $\boldsymbol{\theta}$ that maximize
¹⁴⁵⁴ the likelihood, while setting $\delta = 0$ and $\xi = \zeta$. Once the likelihood is defined, a test statistic must be built
¹⁴⁵⁵ for use in hypothesis testing.

¹⁴⁵⁶ TEST STATISTIC

¹⁴⁵⁷ To distinguish whether the data match a background only or background and signal hypothesis, a test
¹⁴⁵⁸ statistic must be used. The $H \rightarrow WW^*$ analysis uses the profile likelihood technique [65]. The first step
¹⁴⁵⁹ in formulating this test statistic is to define the profile likelihood ratio, shown in equation 3.11.

$$\lambda(\mu) = \frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(\hat{\mu}, \hat{\theta})} \quad (3.11)$$

¹⁴⁶⁰ Here $\hat{\theta}_\mu$ is the value of θ that maximizes the likelihood for the choice of μ being tested. Additionally, $\hat{\theta}$
¹⁴⁶¹ and $\hat{\mu}$ represent the values of θ and μ that gives the overall maximum value of the likelihood.

₁₄₆₂ Once this is defined, a test statistic q_μ is constructed. This is shown in equation 3.12.

$$q_\mu = -2 \ln \lambda(\mu) \quad (3.12)$$

₁₄₆₃ A higher value of q_μ indicates that the data are more incompatible with the hypothesized value of μ , and
₁₄₆₄ q_0 then corresponds to the value of the test statistic for the background only hypothesis. A p_0 value is
₁₄₆₅ then defined to quantify the compatibility between the data and the null hypothesis. The p_0 value is the
₁₄₆₆ probability of obtaining a value of q_0 larger than the observed value, and this is shown in equation 3.13.

$$p_0 = \int_{q_0^{\text{obs}}}^{\infty} f(q_\mu | \mu = 0) dq_\mu \quad (3.13)$$

₁₄₆₇ Here $f(q_\mu)$ is the probability distribution function of the test statistic. Finally, the p_0 value can be con-
₁₄₆₈ verted into a signal significance, using the formula in equation 3.14, or the one-sided tail of the Gaussian
₁₄₆₉ distribution.

$$Z_0 = \sqrt{2} \operatorname{erf}^{-1}(1 - 2p_0) \quad (3.14)$$

₁₄₇₀ The threshold for discovery used in particle physics is $Z_0 \geq 5$, more commonly known as a value of 5σ .

The real voyage of discovery consists not in seeking new landscapes, but in having new eyes.

Marcel Proust

4

1471

1472 The discovery of the Higgs boson and the role 1473 of the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel

1474 4.1 INTRODUCTION

1475 This chapter presents the results of the search for the Higgs boson in 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$
1476 and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. The results of three searches at $\sqrt{s} = 8 \text{ TeV}$ in the $H \rightarrow WW^* \rightarrow$
1477 $\ell\nu\ell\nu$, $H \rightarrow \gamma\gamma$, and $H \rightarrow ZZ \rightarrow 4\ell$ channels are shown. These results at 8 TeV are combined
1478 with the results of searches at $\sqrt{s} = 7 \text{ TeV}$ in the same channels along with $H \rightarrow \tau\tau$ production and
1479 associated production searches for $H \rightarrow b\bar{b}$. The results of this combination are a 5.9σ detection of a
1480 new particle consistent with a Higgs boson produced via gluon fusion. Rather than going into detail for
1481 all of the different Higgs decay searches, this chapter will discuss the three most sensitive channels and in
1482 particular focus on $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$. While the focus is on WW^* , some of the ZZ^* and $\gamma\gamma$ results

1483 are shown for completeness. The results not discussed here can be found in the ATLAS Higgs discovery
1484 publication [1].

1485 **4.2 DATA AND SIMULATION SAMPLES**

1486 The data sample used for the following results was taken in 2011 and 2012 at center of mass energies
1487 of 7 and 8 TeV, respectively, with 4.8 fb^{-1} collected at 7 TeV and 5.8 fb^{-1} collected at 8 TeV. Higgs
1488 production in the gluon fusion and vector boson fusion modes is modeled with **POWHEG** for the hard
1489 scattering event and **PYTHIA** for the showering and hadronization. Associated production of a Higgs with a
1490 vector boson or top quarks is modeled via **PYTHIA**. Table 4.1 shows the Monte Carlo generators used for
1491 modeling the signal and background processes relevant for the three analyses to be discussed.

Process	Generator
ggF, VBF H	POWHEG + PYTHIA
$WH, ZH, t\bar{t}H$	PYTHIA
$W + \text{jets}, Z/\gamma^* + \text{jets}$	ALPGEN + HERWIG
$t\bar{t}, tW, tb$	MC@NLO + HERWIG
tqb	ACERMC + PYTHIA
$q\bar{q} \rightarrow WW$	MC@NLO + HERWIG
$gg \rightarrow WW$	GG2WW+ HERWIG
$q\bar{q} \rightarrow ZZ$	POWHEG + PYTHIA
$gg \rightarrow ZZ$	GG2ZZ+ HERWIG
WZ	MADGRAPH+ PYTHIA , HERWIG
$W\gamma + \text{jets}$	ALPGEN + HERWIG
$W\gamma^*$	MADGRAPH+ PYTHIA
$q\bar{q}/gg \rightarrow \gamma\gamma$	SHERPA

Table 4.1: Monte carlo generators used to model signal and background for the Higgs search [1].

1492 **4.3 $H \rightarrow WW \rightarrow e\nu\mu\nu$ SEARCH**

1493 The $H \rightarrow WW \rightarrow e\nu\mu\nu$ search is unique compared to the ZZ and $\gamma\gamma$ channels. The Higgs mass
1494 cannot be fully reconstructed due to the presence of neutrinos in the final state, so the transverse mass m_T
1495 is used as the final discriminating variable. This channel also has a wider variety of backgrounds compared
1496 to other channels, as discussed in chapter 3. The same flavor final states are excluded from the 8 TeV dataset

¹⁴⁹⁷ due to high pileup conditions¹. These final states are later included in results with the full Run 1 dataset,
¹⁴⁹⁸ as discussed in chapters 5 and 6.

¹⁴⁹⁹ **4.3.1 EVENT SELECTION**

¹⁵⁰⁰ The analysis requires two opposite charge isolated leptons, with the leading (sub-leading) lepton required
¹⁵⁰¹ to have $p_T > 25(15)$ GeV. The events are separated into different signal regions depending on which
¹⁵⁰² flavor of lepton is leading ($e\mu$ for leading electron, μe for leading muon). Strict lepton quality cuts are
¹⁵⁰³ applied to the sample to reduce backgrounds from mis-reconstructed leptons.

¹⁵⁰⁴ Jets are reconstructed with the anti- k_T algorithm with a radius parameter $R = 0.4$. The jets are re-
¹⁵⁰⁵ quired to have $p_T > 25$ GeV and $|\eta| < 4.5$, with jets in the tracking volume required to have a jet vertex
¹⁵⁰⁶ fraction of 0.5 and jets in the forward region required to have $p_T > 30$ GeV. The analysis is separated
¹⁵⁰⁷ into three different signal regions based on jet multiplicity: $n_j = 0, 1, \geq 2$.

¹⁵⁰⁸ To indicate the presence of neutrinos in the event, a requirement of $E_{T,\text{rel}}^{\text{miss}} > 25$ GeV is made². This
¹⁵⁰⁹ requirement significantly reduces the QCD multijet and Z/γ^* + jets backgrounds. Figure 4.1 shows the
¹⁵¹⁰ distribution of n_j in data and simulation after applying these “pre-selection” requirements.

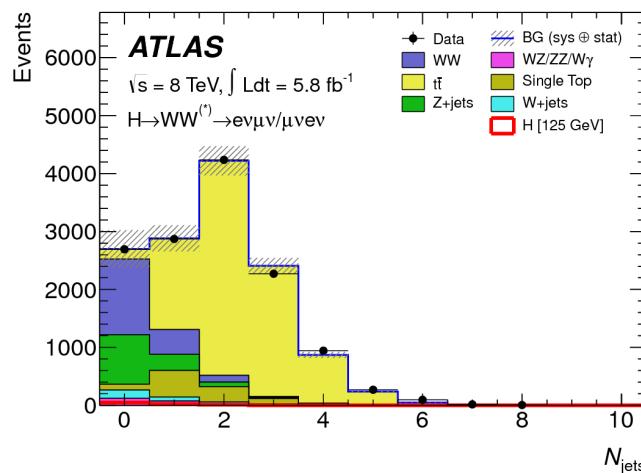


Figure 4.1: Jet multiplicity distribution in data and MC after applying lepton, jet, and $E_{T,\text{rel}}^{\text{miss}}$ selections. The WW and top backgrounds have been normalized using control samples, and the hashed band indicates the total uncertainty on the prediction [1].

¹The less sensitive 7 TeV search result includes both different flavor and same flavor final states.

²For the definition of $E_{T,\text{rel}}^{\text{miss}}$, see section 3.4.1.

1511 Additional selections are applied to require the dilepton topology to correspond to that of a Standard
1512 Model Higgs boson. The requirements are presented here - more detailed discussion on the motivation
1513 for each requirement is saved for Chapter 5. In all of the jet multiplicity channels, the dilepton system
1514 is required to have a small gap in azimuthal angle, $\Delta\phi_{\ell\ell} < 1.8$. Similarly, the dilepton invariant mass,
1515 $m_{\ell\ell}$, is required to be less than 50 GeV in the lower jet multiplicity channels and less than 80 GeV in the
1516 $n_j \geq 2$ channel. In the $n_j = 0$ channel, the magnitude of the dilepton p_T , $p_T^{\ell\ell}$, is required to be greater
1517 than 30 GeV.

1518 In the higher jet multiplicity channels ($n_j \geq 1$), the top background is a larger fraction of the total
1519 background and must be reduced more carefully. The total transverse momentum p_T^{sum} is thus required
1520 to be less than 30 GeV. Additionally, the di- τ invariant mass $m_{\tau\tau}$ (dilepton mass computed under the
1521 assumption that the neutrinos from the τ decay are emitted collinear to the charged leptons [66]) is used
1522 to reject $Z \rightarrow \tau\tau$ events by requiring $|m_{\tau\tau} - m_Z| > 25$ GeV. These variables are also discussed in more
1523 detail in Chapter 5.

1524 In the $n_j \geq 2$ channel, requirements are made to isolate the VBF contribution to Higgs production.
1525 The kinematics of the two leading jets are used to make these requirements. In particular, the event must
1526 have $\Delta y_{jj} > 3.8$ and $m_{jj} > 500$ GeV, along with a veto on having any additional jets with rapidity
1527 between the two leading jets.

1528 **4.3.2 BACKGROUND ESTIMATION**

1529 The details of the background estimation techniques used in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis are
1530 discussed in section 5.5. The dominant backgrounds are SM WW production and top (both pair and
1531 single) production, and these backgrounds have their normalizations estimated from dedicated control
1532 regions while their shapes are taken from simulation.

1533 The control sample for the Standard Model WW background is defined by making the same require-
1534 ments as the signal region with the $m_{\ell\ell}$ requirement inverted (now requiring $m_{\ell\ell} > 80$ GeV) and remov-
1535 ing the $\Delta\phi_{\ell\ell}$ requirement. This creates a control sample that is 70% (40%) pure in the 0(1)-jet region. The
1536 correction to the pure MC-based background estimate is quantified by defining a normalization factor β

1537 which is the ratio of the data yield to the MC yield ($N_{\text{data}}/N_{\text{MC}}$) in this control sample. Table 4.2 shows
 1538 the WW normalization factors in the $n_j = 0$ and $n_j = 1$ bins (the $n_j \geq 2$ estimate is taken directly
 1539 from MC).

n_j	β_{WW}	β_t
= 0	1.06 ± 0.06	1.11 ± 0.06
= 1	0.99 ± 0.15	1.11 ± 0.05
≥ 2	-	1.01 ± 0.26

Table 4.2: Normalization factors (ratio of data and MC yields in a control sample) for the Standard Model WW and top backgrounds in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis [1]. Only statistical uncertainties are shown.

1540 The top background estimate is also computed separately in each jet multiplicity bin. In the $n_j = 0$
 1541 channel, the background is first normalized using data after pre-selection requirements with no selection
 1542 on n_j . Then, a dedicated b -tagged control sample is used to evaluate the ratio of one-jet to two-jet events in
 1543 data. The details of this technique are shown in reference [67]. In the $n_j = 1$ and the $n_j \geq 2$ regions, the
 1544 top background is normalized in a control sample where the signal region selections are applied, but the
 1545 b -jet veto is reversed and the Higgs topology requirements on $m_{\ell\ell}$ and $\Delta\phi_{\ell\ell}$ are removed. The resulting
 1546 normalization factors for these techniques are shown in table 4.2.

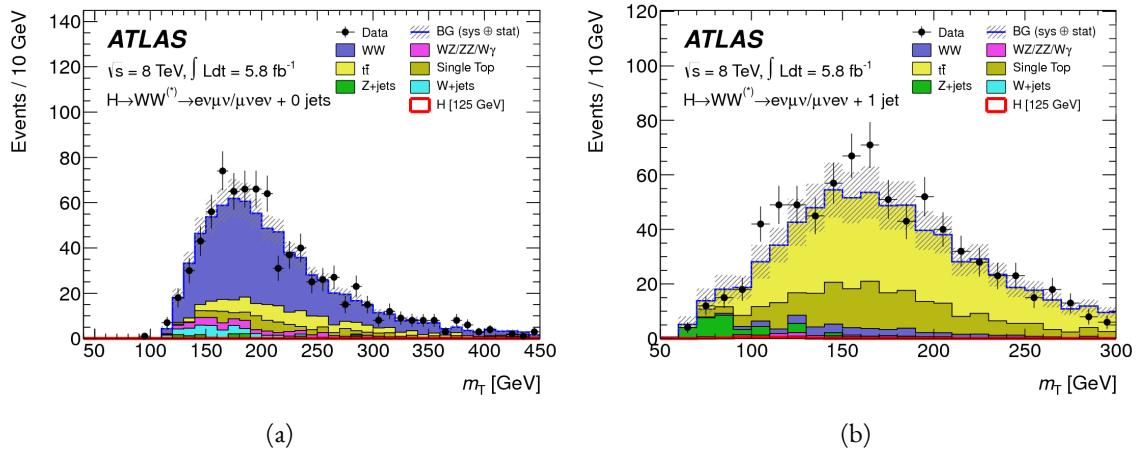


Figure 4.2: Comparison of m_T between data and simulation in the $n_j = 0$ WW (a) and $n_j = 1$ top (b) control samples [1].

1547 The control samples which are used for background normalization can also be used to validate the mod-
 1548 eling of the m_T distribution for each background. Figure 4.2 shows the comparison between data and MC

1549 for the m_T distribution after correcting the normalization of the backgrounds in the WW and top control
1550 regions. Good agreement between data and simulation is seen in both cases.

1551 The $W + \text{jets}$ background estimate is taken entirely from data using a control sample with one well recon-
1552 structed lepton and one anti-identified lepton. All other backgrounds are taken purely from simulation.

1553 4.3.3 SYSTEMATIC UNCERTAINTIES

1554 The systematic uncertainties that have the largest impact on the analysis are the theoretical uncertainties
1555 associated with the signal cross section. These are shared with the ZZ^* and $\gamma\gamma$ channels. The uncertainties
1556 resulting from variations of the QCD scale are $+7\% / -8\%$ on the final signal yield. Those coming from
1557 variations of the parton distribution function (PDF) used in the simulation add a $\pm 8\%$ uncertainty on
1558 the yield. The uncertainties on the branching ratios of the Higgs are $\pm 5\%$.

1559 The main experimental uncertainties come from variations of the jet energy scale (JES), jet energy reso-
1560 lution, pile-up, E_T^{miss} , b -tagging efficiency, $W + \text{jets}$ background estimate, and integrated luminosity. For
1561 more details, see reference [1].

1562 4.3.4 RESULTS

1563 Table 4.3 shows the signal and background yields in the final signal region after normalizing the back-
1564 grounds according to the methods described above.

	$n_j = 0$	$n_j = 1$	$n_j \geq 2$
Signal	20 ± 4	5 ± 2	0.34 ± 0.07
WW	101 ± 13	12 ± 5	0.10 ± 0.14
Other dibosons	12 ± 3	1.9 ± 1.1	0.10 ± 0.10
$t\bar{t}$	8 ± 2	6 ± 2	0.15 ± 0.10
Single top	3.4 ± 1.5	3.7 ± 1.6	-
$Z/\gamma^* + \text{jets}$	1.9 ± 1.3	0.10 ± 0.10	-
$W + \text{jets}$	15 ± 7	2 ± 1	-
Total background	142 ± 16	26 ± 6	0.35 ± 0.18
Observed in data	185	38	0

Table 4.3: Data and expected yields for signal and background in the final $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ signal region.
Uncertainties shown are both statistical and systematic [1].

1565 Figure 4.3 shows the m_T distribution in the $n_j \leq 1$ channels for 8 TeV data. (No events are observed
 1566 in data in the $n_j \geq 2$ channels in this dataset). The excess shown here relatively flat as a function of
 1567 hypothesized Higgs mass. The combined 7 and 8 TeV data gives an excess with local significance of 2.8σ
 1568 with an expected significance of 2.3σ , corresponding to a μ measurement of 1.3 ± 0.5 .

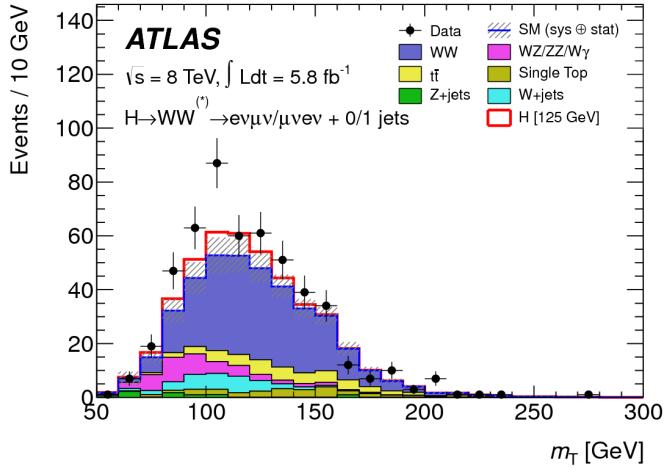


Figure 4.3: m_T distribution in the $H \rightarrow WW \rightarrow e\nu\mu\nu$ $n_j \leq 1$ channels for 8 TeV data [1].

1569 4.4 $H \rightarrow \gamma\gamma$ SEARCH

1570 The $H \rightarrow \gamma\gamma$ search is a search for a peaked excess above the falling SM diphoton mass spectrum, with
 1571 $m_{\gamma\gamma}$ as the ultimate discriminating variable. Events are selected by requiring two isolated photons, with
 1572 the leading (sub-leading) photon required to have $E_T > 40(30)$ GeV. In the 8 TeV data, the photons are
 1573 required to pass identification criteria consistent with a photonic shower in the electromagnetic calorimeter
 1574 and little leakage in the hadronic calorimeter.

1575 The main challenges for this analysis are accurate mass reconstruction and background estimation. In
 1576 order to accurately reconstruct the invariant mass of the di-photon system, both the energy and direction
 1577 of the photons must be measured well. Therefore, the identification of the primary vertex of the hard
 1578 interaction is particularly important, and is done using a multivariate likelihood which combines informa-
 1579 tion about the photon direction and vertex position. The background is modeled with a falling spectrum
 1580 in $m_{\gamma\gamma}$ that is parameterized by different functions depending on the category of the event.

1581 4.4.1 RESULTS

1582 The resulting diphoton mass spectrum is shown in figure 4.4. The best fit mass value in the $\gamma\gamma$ channel
 1583 alone in the combined 7 and 8 TeV data is 126.5 GeV. The local significance at this point is 4.5σ , with
 1584 an expected significance of 2.5σ . Therefore, the measured signal strength μ is 1.8 ± 0.5 in this channel.

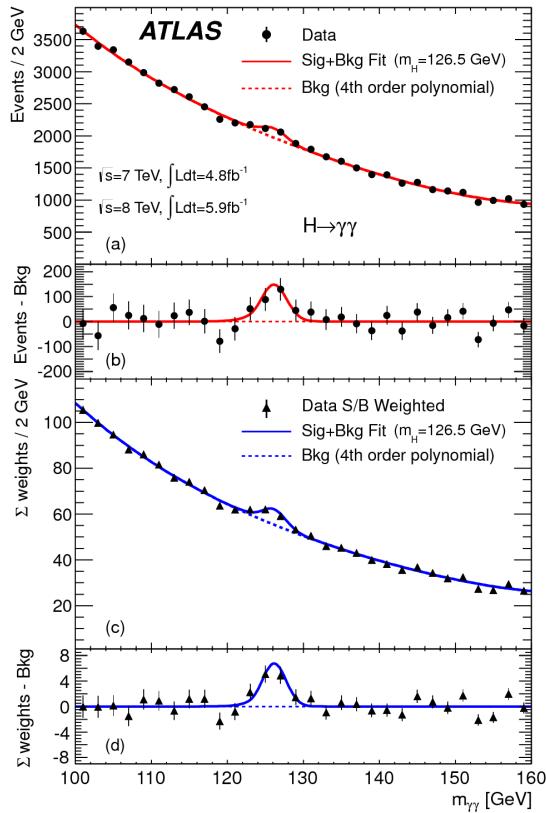


Figure 4.4: Diphoton mass spectrum in 7 and 8 TeV data. Panel a) shows the unweighted data distribution superimposed on the background fit, while panel c) shows the data where each event category is weighted by its signal to background ratio. Panels b) and d) show the respective distributions with background subtracted [1].

1585 4.5 $H \rightarrow ZZ \rightarrow 4\ell$ SEARCH

1586 The $H \rightarrow ZZ \rightarrow 4\ell$ analysis searches for a Standard Model Higgs boson decaying to two Z bosons,
 1587 each of which decays to a pair of same flavor, opposite charge isolated leptons. The ultimate discriminating
 1588 variable is $m_{4\ell}$, or the invariant mass of the four selected leptons. The ℓ denotes an e or μ as with the
 1589 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis.

1590 Four distinct signal regions are constructed depending on the flavors of the final state, additionally sep-
 1591 arated by the flavor of the leading lepton pair. These are referred to as $4e$, $2e2\mu$, $2\mu2e$, 4μ .

1592 The main backgrounds in the $H \rightarrow ZZ \rightarrow 4\ell$ search are continuum ZZ^* production, $Z + \text{jets}$ pro-
 1593 duction, and $t\bar{t}$. The $m_{4\ell}$ distribution for background is estimated from simulation. The normalization
 1594 of the SM ZZ^* background is also taken from MC simulation, while the $Z + \text{jets}$ and $t\bar{t}$ normalizations are
 1595 taken from data-driven methods.

1596 4.5.1 RESULTS

1597 Figure 4.5 shows the $m_{4\ell}$ spectrum measured in the 7 and 8 TeV datasets. The total number of events
 1598 observed in the window between 120 and 130 GeV is 13, with 6 events in the 4μ channel, 2 events in
 1599 the $4e$ channel, and 5 events in the $2e2\mu/2\mu2e$. The best fit μ value in the combined 7 and 8 TeV data
 1600 occurs at 125 GeV and is measured to be 1.2 ± 0.6 . The observed significance at this mass is 3.6σ , with
 1601 an expected significance of 2.7σ .

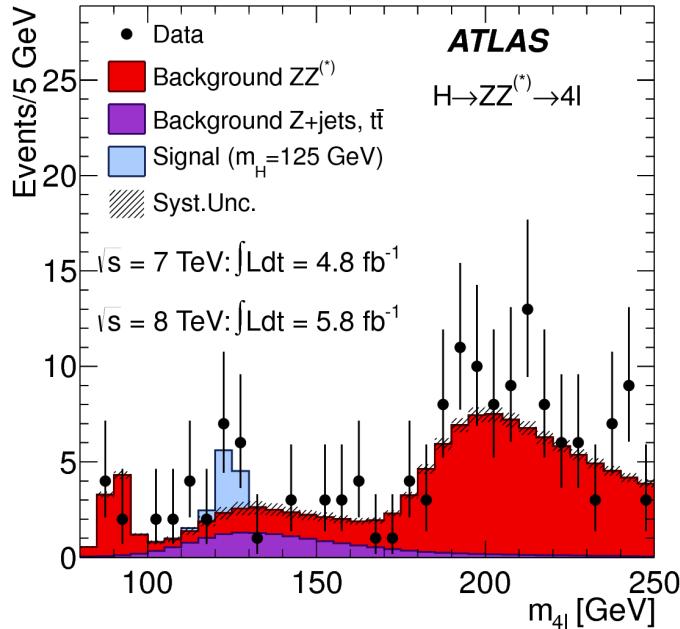


Figure 4.5: Four lepton invariant mass spectrum ($m_{4\ell}$) in 7 and 8 TeV data compared to background estimate. A 125 GeV SM Higgs signal is shown in blue [1].

1602 4.6 COMBINED RESULTS

1603 The statistical interpretation of the combined results is undertaken as described in section 3.6.2, with a
 1604 hypothesis test based on a likelihood ratio parameterized by the Higgs signal strength μ . The null hypoth-
 1605 esis corresponds to $\mu = 0$, while the SM Higgs corresponds to $\mu = 1$.

1606 Table 4.4 summarizes the properties of the individual channels as well as the significances of the excesses
 1607 seen. The most significant observed local excess comes from the $\gamma\gamma$ channel. Figure 4.6 shows a com-
 1608 parison of the observed local p_0 values as a function of hypothesized mass for the three different search
 1609 channels. Both the ZZ^* and $\gamma\gamma$ channels have very peaked excesses, while the WW^* excess can be seen as
 1610 very broad because the m_T distribution does not provide detailed information about the true Higgs mass.
 1611 Figure 4.7 shows the combined exclusion limit, p_0 , and signal strength. The highest local excess comes at
 1612 a value of 126.5 GeV and corresponds to a 6.0σ observed excess.

Channel	Fit var.	Observed Z_l	Expected Z_l	$\hat{\mu}$
$H \rightarrow ZZ^* \rightarrow 4\ell$	$m_{4\ell}$	3.6	2.7	1.2 ± 0.6
$H \rightarrow \gamma\gamma$	$m_{\gamma\gamma}$	4.5	2.5	1.8 ± 0.5
$H \rightarrow WW^* \rightarrow e\nu\mu\nu$	m_T	2.8	2.3	1.3 ± 0.5
Combined	-	6.0	4.9	1.4 ± 0.3

Table 4.4: Summary of the expected and observed significance and measured signal strengths in the combined 7 and 8 TeV datasets for the Higgs discovery analysis [1].

1613 Figure 4.8 shows a comparison of the measured signal strengths between the different Higgs search
 1614 channels. All measured μ are consistent with unity within their uncertainty, and the combined μ mea-
 1615 surement is 1.4 ± 0.3 .

1616 The likelihood can also be computed in a two-dimensional plane of m_H and μ , and this is shown in
 1617 figure 4.9. The figure shows that while the $\gamma\gamma$ and ZZ^* channels have very good mass resolution, the
 1618 excess in WW^* covers a broad mass range. The banana shape of the WW^* result is due to the fact that
 1619 the excess in this channel can either be explained by increasing the signal strength or by changing the mass
 1620 (and thus the cross section). The two parameters are correlated due to the lack of mass sensitivity in this
 1621 channel.

1622 Because multiple Higgs mass points are searched for, the local significance must be corrected for a look-

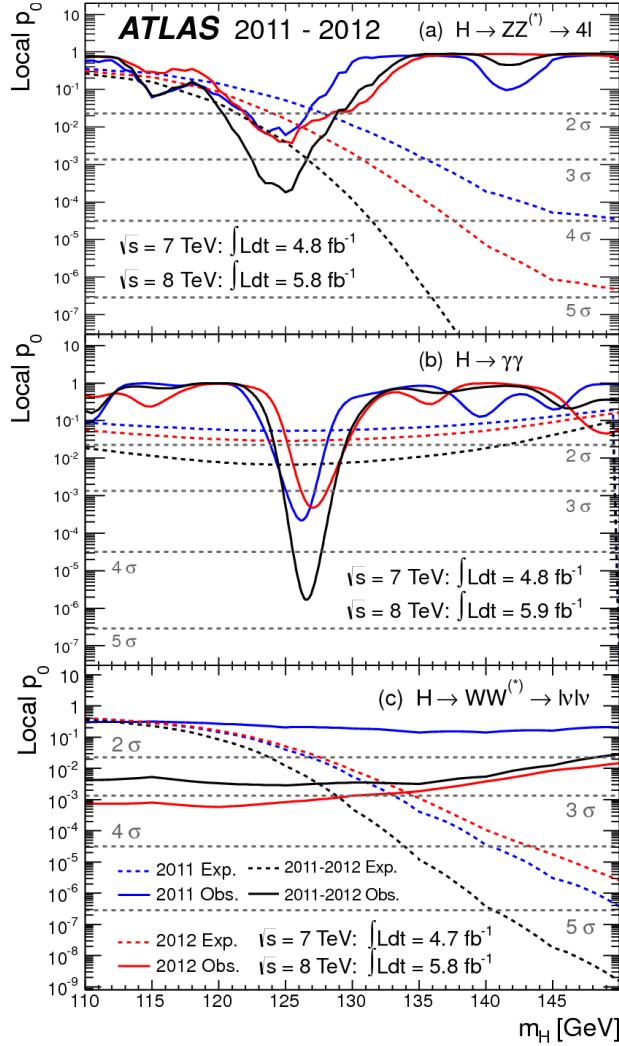


Figure 4.6: Local p_0 distribution as a function of hypothesized Higgs mass for the $H \rightarrow ZZ^* \rightarrow 4\ell$ (a), $H \rightarrow \gamma\gamma$ (b), and $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ (c) channels. Dashed curves show expected results, while solid curves show observed. Red curves are from 7 TeV data, blue curves from 8 TeV, and black curved combined [1].

¹⁶²³ elsewhere effect to compute a true global significance. The global significance for finding a Higgs anywhere
¹⁶²⁴ in the mass range of 110 GeV to 600 GeV is 5.1σ . This increases slightly to 5.3σ if only mass range from
¹⁶²⁵ 110 to 150 GeV.

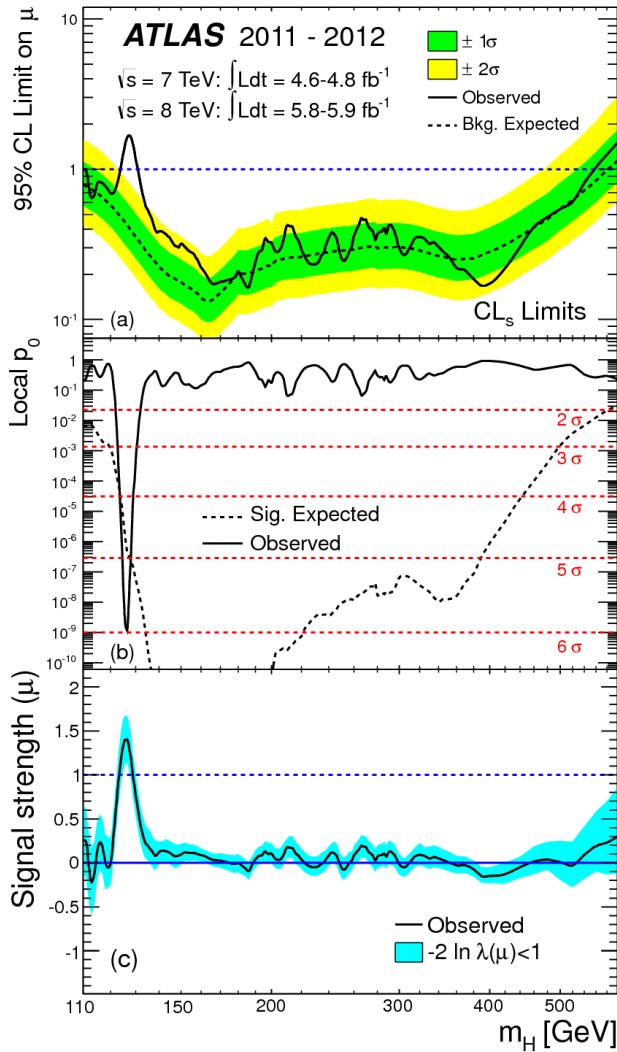


Figure 4.7: Combined 95% CL limits (a), local p_0 values (b), and signal strength measurement (c) as a function of Higgs mass [1].

1626 4.7 CONCLUSION

1627 A search for the production of a Standard Model Higgs boson was conducted in 4.8 fb^{-1} collected at
 1628 $\sqrt{s} = 7 \text{ TeV}$ and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. A new particle consistent with the Higgs boson was observed,
 1629 with a mass of 126.5 GeV and a global (local) significance of $5.1(6.0)\sigma$. This is the first discovery level
 1630 observation of a particle consistent with the Higgs.

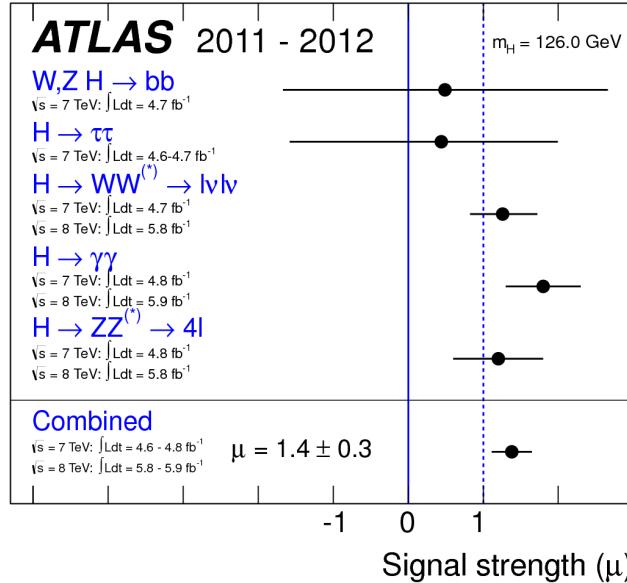


Figure 4.8: Comparison of measured signal strength μ for a 126 GeV Higgs in the 7 and 8 TeV datasets [1].

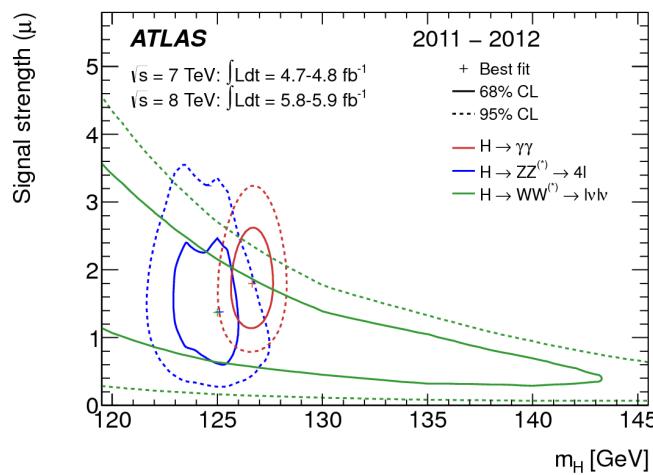


Figure 4.9: Two dimensional likelihood as a function of signal strength μ and Higgs mass m_H [1].

*The imagination of nature is far, far greater than the
imagination of man.*

Richard Feynman

5

1631

1632 Evidence for Vector Boson Fusion production

1633

$$\text{of } H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$$

1634 5.1 INTRODUCTION

1635 After the discovery of a particle consistent with the Higgs boson, the $H \rightarrow WW^*$ analysis had two main
1636 goals. The first goal was to increase the sensitivity of the analysis to fully confirm that the $H \rightarrow WW^*$
1637 process did indeed exist. The second goal was to characterize the particle as much as possible, including
1638 searching for the lower cross-section production modes. This chapter presents a dedicated search for Vec-
1639 tor Boson Fusion (VBF) production of a Higgs boson decaying via the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ mode. First,
1640 the data and Monte Carlo samples are detailed, along with trigger and physics object selections. Then, the
1641 details of the analysis are shown, including signal region definition, background estimation techniques,
1642 and systematic uncertainties. Finally, the results of the analysis are presented. As will be shown, this anal-

ysis is the first and most sensitive evidence of the VBF production mode of the Higgs on ATLAS.

In the VBF channel, there are both a selection requirement based signal region analysis (known as the “cut-based”) and a multivariate analysis which uses a boosted decision tree (known as the BDT analysis). The focus of this chapter will be on the cut-based signal region, as this is an important component of the VBF analysis and in particular acts as strong validation for the final BDT result. Connections between the cut-based and BDT analyses will be discussed where appropriate.

5.2 DATA AND SIMULATION SAMPLES

The results presented here are with 20.3 fb^{-1} taken at $\sqrt{s} = 8 \text{ TeV}$ and 4.5 fb^{-1} taken at $\sqrt{s} = 7 \text{ TeV}$. The details of the LHC and detector conditions during this period are given in Chapter 2. The trigger selection defining the dataset is discussed in section 5.2.1. The simulation samples used for signal and background modeling are given in section 5.2.2.

5.2.1 TRIGGERS

The analysis uses a combination of single lepton and dilepton triggers to allow lowering of the p_T thresholds and increased signal acceptance. The p_T threshold on the leptons is a particularly important consideration for this signal. Because the W^* produced in the decay is off-shell, it tends to produce lower momentum leptons. Thus, being able to lower the p_T threshold while still maintaining a low background rate is critical. Figure 5.1 shows an example of the subleading lepton p_T for a VBF $H \rightarrow WW^*$ signal compared to the corresponding $t\bar{t}$ background. Note that the lepton p_T spectrum is considerably softer in the signal sample.

As discussed in Chapter 2, there are multiple levels in the ATLAS trigger system, and there are different p_T thresholds imposed for the leptons at each level. Additionally, some triggers have a loose selection on the isolation of the lepton (looser than that applied offline in the analysis object selection). Table 5.1 shows the p_T thresholds used for single lepton triggers, while table 5.2 shows the p_T thresholds coming from di-lepton triggers. The single lepton trigger efficiency for muons that pass the analysis object selection is 70% for muons in the barrel region ($|\eta| < 1.05$) and 90% in the endcap region. The electron trigger



Figure 5.1: A comparison of the subleading lepton p_T spectrum between VBF $H \rightarrow WW^*$ production and $t\bar{t}$ background.

efficiency increases with electron p_T but the average is approximately 90%. These efficiencies are measured by combined performance and trigger signature groups [68, 69].

	Level-I threshold	High-level threshold
Electron	18	$24i$
	30	60
Muon	15	$24i$
		36

Table 5.1: Single lepton triggers used for electrons and muons in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis. A logical “or” of the triggers listed for each lepton type is taken. Units are in GeV, and the i denotes an isolation requirement in the trigger.

	Level-I threshold	High-level threshold
ee	10 and 10	12 and 12
$\mu\mu$	15	18 and 8
$e\mu$	10 and 6	12 and 8

Table 5.2: Di-lepton triggers used for different flavor combinations in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis. The two thresholds listed refer to leading and sub-leading leptons, respectively. The di-muon trigger only requires a single lepton at level-I.

The combination of all triggers shown gives good efficiency for signal events. This efficiency is summarized in table 5.3. The relative improvement in efficiency by adding the dilepton triggers is also shown

1672 in the same table. The largest gain comes in the $\mu\mu$ channel. Overall the trigger selection shows a good
1673 efficiency for $H \rightarrow WW^*$ signal events.

Channel	Trigger efficiency	Gain from 2ℓ trigger
ee	97%	9.1%
$\mu\mu$	89%	18.5%
$e\mu$	95%	8.3%
μe	81%	8.2%

Table 5.3: Trigger efficiency for signal events and relative gain of adding a dilepton trigger on top of the single lepton trigger selection. The first lepton is the leading, while the second is the sub-leading. Efficiencies shown here are for the ggF signal in the $n_j = 0$ category but are comparable for the VBF signal.

1674 **5.2.2 MONTE CARLO SAMPLES**

1675 Modeling of signal and background processes in the signal region, in particular for the m_T distribution,
1676 is an important consideration for the final interpretation of the analysis. Therefore, careful consideration
1677 must be paid to which Monte Carlo (MC) generators are used for specific processes. With the exception
1678 of the $W + \text{jet}$ and multijet backgrounds, the m_T shape used as the final discriminant is taken from simu-
1679 lation¹.

1680 Table 5.4 shows the MC generators used for the signal and background processes, as well as the cross
1681 sections of each process. In order to include corrections up to next-to-leading order (NLO) in the QCD
1682 coupling constant α_s , the POWHEG [70] generator is often used. In some cases, only leading order gener-
1683 ators like ACERMC [71] and GG2VV [72] are available for the process in question. If the process requires
1684 good modeling for very high parton multiplicities, the SHERPA [73] and ALPGEN [74] generators are used
1685 to provide merged calculations for five or fewer additional partons. These matrix element level calculations
1686 must then be additionally matched to models of the underlying event, hadronization, and parton shower.
1687 There are four generators used for this purpose: SHERPA , PYTHIA 6 [75], PYTHIA 8 [76], or HERWIG
1688 [77] + JIMMY [78]. The simulation additionally requires an input parton distribution function (PDF).
1689 The CT10 [79] PDFs are used for SHERPA and POWHEG simulated samples, while CTEQ6L1 [80] is used

¹Many backgrounds are normalized from data, as described in section 5.5.

Process	MC generator	$\sigma \cdot \mathcal{B}$ (pb)
Signal		
ggF $H \rightarrow WW^*$	POWHEG +PYTHIA 8	0.435
VBF $H \rightarrow WW^*$	POWHEG +PYTHIA 8	0.0356
VH $H \rightarrow WW^*$	PYTHIA 8	0.0253
WW		
$q\bar{q} \rightarrow WW$ and $qg \rightarrow WW$	POWHEG +PYTHIA 6	5.68
$gg \rightarrow WW$	GG2VV +HERWIG	0.196
$(q\bar{q} \rightarrow W) + (q\bar{q} \rightarrow W)$	PYTHIA 8	0.480
$q\bar{q} \rightarrow WW$	SHERPA	5.68
VBS $WW + 2$ jets	SHERPA	0.0397
Top quarks		
$t\bar{t}$	POWHEG +PYTHIA 6	26.6
Wt	POWHEG +PYTHIA 6	2.35
$t\bar{q}\bar{b}$	ACERMC +PYTHIA 6	28.4
$t\bar{b}$	POWHEG +PYTHIA 6	1.82
Other dibosons (VV)		
$W\gamma$ ($p_T^\gamma > 8$ GeV)	ALPGEN +HERWIG	369
$W\gamma^*$ ($m_{\ell\ell} \leq 7$ GeV)	SHERPA	12.2
WZ ($m_{\ell\ell} > 7$ GeV)	POWHEG +PYTHIA 8	12.7
VBS $WZ + 2$ jets	SHERPA	0.0126
($m_{\ell\ell} > 7$ GeV)		
$Z\gamma$ ($p_T^\gamma > 8$ GeV)	SHERPA	163
$Z\gamma^*$ (min. $m_{\ell\ell} \leq 4$ GeV)	SHERPA	7.31
ZZ ($m_{\ell\ell} > 4$ GeV)	POWHEG +PYTHIA 8	0.733
$ZZ \rightarrow \ell\ell\nu\nu$ ($m_{\ell\ell} > 4$ GeV)	POWHEG +PYTHIA 8	0.504
Drell-Yan		
Z ($m_{\ell\ell} > 10$ GeV)	ALPGEN +HERWIG	16500
VBF $Z + 2$ jets	SHERPA	5.36
($m_{\ell\ell} > 7$ GeV)		

Table 5.4: Monte Carlo samples used to model the signal and background processes [63].

for ALPGEN +HERWIG and ACERMC simulations. The Drell-Yan samples are reweighted to the MRST [81] PDFs, as these are found to give the best agreement between data and simulation.

Once the basic hard scattering process is simulated, it must be passed through a detector simulation

and additional pile-up events must be overlaid. The pile-up events are modeled with PYTHIA 8, and the ATLAS detector is simulated with GEANT4 [82]. Because of the unique phase space of the $H \rightarrow WW^*$ analysis, events are sometimes filtered at generator level to allow for more efficient generation of relevant events. The efficiency of the trigger in MC simulation does not always match the measured efficiency in data, so trigger scale factors are applied to correct the MC efficiency to the data. These are derived by the combined performance groups [68, 69].

5.3 OBJECT SELECTION

In order to define the signal region, the analysis must first select the reconstructed physics objects to be considered. The details of the object reconstruction algorithms are discussed in Chapter 2, while this section gives specific selection requirements used in the $H \rightarrow WW^*$ analysis. The first step in this process is to select a primary vertex candidates. The event's primary vertex is the vertex with the largest sum of p_T^2 for associated tracks and is required to have at least three tracks with $p_T > 450$ MeV. Many of the object selection cuts are then made relative to this chosen primary vertex.

5.3.1 MUONS

The analysis uses combined muon candidates, where a track in the Inner Detector has been matched to a standalone track in the Muon Spectrometer. The track parameters are combined statistically in the muon reconstruction algorithm [54]. The muons are required to be within $|\eta| < 2.5$ and have a $p_T > 10$ GeV. To reduce backgrounds coming from mis-reconstructed leptons, there are requirements on the impact parameter of the muon relative to the primary vertex. The transverse impact parameter d_0 is required to be small relative to its estimated uncertainty, the exact cut value being $d_0/\sigma_{d_0} < 3$. The longitudinal impact parameter z_0 must satisfy $|z_0 \sin \theta| < 1$ mm.

As discussed previously, the muons must also be isolated. There are two types of lepton isolations that are calculated: track-based and calorimeter-based. For muons, the track-based isolation is defined using the scalar sum $\sum p_T$ for tracks with $p_T > 1$ GeV (excluding the muon's track) within a cone of $\Delta R = 0.3$ (0.4) for muon with $p_T > 15$ GeV ($10 < p_T < 15$ GeV). The final isolation requirement is made my

1718 requiring that this scalar sum be no more than a certain fraction of the muon's p_T . This requirement varies
1719 with muon p_T and the exact requirements are defined in table 5.5.

1720 The calorimeter-based muon isolation is defined using the $\sum E_T$ calculated from calorimeter cells with
1721 the same cone size as the track-based isolation but excluding cells within $\Delta R < 0.05$ around the muon.
1722 This isolation is also defined as a requirement on the ratio of the sum to the muon p_T and varies with
1723 muon p_T . The requirement values as a function of p_T are also given in table 5.5.

1724 The isolation requirements loosen as a function of p_T to allow for larger signal acceptance. At low p_T ,
1725 the isolation is tightened to reduce the $W + \text{jets}$ background which arises from a misidentified lepton.

p_T range (GeV)	Calorimeter isolation	Track isolation
10 – 15	0.06	0.06
15 – 20	0.12	0.08
20 – 25	0.18	0.12
> 25	0.30	0.12

Table 5.5: p_T dependent isolation requirements for muons. Muons are required to have their calorimeter based or track based cone sums be less than this fraction of their p_T .

1726 5.3.2 ELECTRONS

1727 Electrons are identified by matching reconstructed clusters in the electromagnetic calorimeter with tracks
1728 in the inner detector. The electrons are identified using a likelihood based method [[51](#), [52](#)] which takes into
1729 account the shower shapes in the calorimeter, the matching of tracks to clusters, and the amount of transi-
1730 tion radiation in the TRT. The electrons are required to have $|\eta| < 2.47$, and candidates in the transition
1731 region between the barrel and endcap ($1.37 < |\eta| < 1.52$) are excluded. As the muons, the electrons
1732 are required to have transverse impact parameter significance < 3 , while in the longitudinal direction
1733 they must have $|z_0 \sin \theta| < 0.4$ mm. Some electron requirements also vary with electron E_T , and these
1734 requirements are summarized in table 5.6.

1735 The isolation for electrons is defined similarly to the muons but with unique requirements on the ob-
1736 jects included. The track-based isolation is constructed using tracks with $p_T > 400$ MeV with cone sizes
1737 as defined for the muons. The calorimeter-based isolation also uses the same cone size as the muon, but

1738 here the cells within a 0.125×0.175 area in $\eta \times \phi$ around the electron cluster's barycenter are excluded.
 1739 The other difference with respect to muons is that the denominator of the isolation ratio is the electron's
 1740 E_T rather than p_T . The isolation cuts very with electron E_T and are defined in table 5.6. The electron is
 1741 also required to not be consistent with a vertex coming from a photon conversion.

p_T range (GeV)	Quality cut	Calorimeter isolation	Track isolation
10 – 15	Very tight LH	0.20	0.06
15 – 20	Very tight LH	0.24	0.08
20 – 25	Very tight LH	0.28	0.10
> 25	Medium	0.28	0.10

Table 5.6: p_T dependent requirements for electrons. Electrons are required to have their calorimeter based or track based cone sums be less than this fraction of their E_T .

1742 5.3.3 JETS

1743 Jets are clustered with the anti- k_T reconstruction algorithm using a radius parameter of $R = 0.4$. They
 1744 are required to have a jet vertex fraction (jVF) of at least 50%, meaning that half of the tracks associated with
 1745 the jet originated from the primary vertex. Jets with no tracks associated (i.e. those outside the acceptance
 1746 of the ID) do not have this requirement applied. Jets are required to have $p_T > 25$ GeV if they are within
 1747 the tracking acceptance ($|\eta| < 2.4$). Jets with $2.4 < |\eta| < 4.5$ are required to have $p_T > 30$ GeV. This
 1748 tighter requirement reduces jets from pileup in the region where jVF requirements cannot be applied. The
 1749 two highest p_T jets in the event are referred to as the “VBF” jets and used to compute variables used in the
 1750 analysis selection.

1751 Identification of b -jets is done using the MV1 algorithm and is limited to the acceptance of the ID
 1752 ($|\eta| < 2.5$) [59]. The operating point of MV1 that is used is 85% efficient for identifying true b -jets.
 1753 This operating point has a 10.3% of mis-tagging a light quark jet as a b -jet. In order to improve the rejec-
 1754 tion of b -jets, a lower threshold than the nominal p_T threshold described above is used. For the purposes
 1755 of counting the number of b -jets, jets with p_T down to 20 GeV are used.

1756 5.3.4 OVERLAP REMOVAL

1757 There are some cases where reconstructed objects will overlap and one will have to be chosen (for exam-
1758 ple, an electron and a jet in the calorimeter). First, the case of lepton overlap is dealt with. If an electron
1759 candidate extends into the muon spectrometer, it is removed. If a muon and electron are within $\Delta R < 0.1$
1760 of each other, the electron is removed and the muon is kept. If two electron candidates overlap within the
1761 same radius, then the higher E_T electron is kept. Next, the overlap between leptons and jets is considered.
1762 If an electron and jet are within $\Delta R < 0.3$ of one another, the electron is kept and the jet is removed.
1763 However, if a muon and jet overlap within $\Delta R < 0.3$, the jet is kept (as it is likely that the muon is the
1764 result of a semileptonic decay inside the jet). Once the overlap removal is complete, the final set of objects
1765 used in the analysis is defined.

1766 5.4 ANALYSIS SELECTION

1767 The VBF analysis uses two distinct selections. The first is a more standard selection, referred to as “cut-
1768 based”, that applies requirements on the VBF variables and uses m_T as the final discriminating variable.
1769 The second is a looser selection that uses an algorithm known as a Boosted Decision Tree (BDT). A BDT is
1770 a multivariate technique that uses an ensemble of decision trees to split the phase space of input variables
1771 into signal-like and background-like regions in order to provide separation power [83–85]. The output
1772 score of a BDT trained to distinguish the VBF Higgs signal from background processes is used as the final
1773 discriminating variable in order to take advantage of the detailed correlations between the VBF variables.
1774 While the BDT-based analysis is ultimately more sensitive, the cut-based serves as an important component
1775 of the analysis. First, the cut-based allows for confirming the modeling and validity of the variables used
1776 as input to the BDT. Second, because this is the first use of such an MVA technique in the $H \rightarrow WW^*$
1777 analysis, the cut-based selection allows confirmation of the final BDT result with a more traditional anal-
1778 ysis. The cut-based techniques are the focus of this chapter, but connections to the BDT result will be
1779 illustrated when appropriate.

1780 One important note is that because this analysis is dedicated to the measurement of the VBF pro-
1781 duction mode of the Higgs, events coming from gluon fusion production with the Higgs decaying via

1782 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ are treated as background events. This will be seen throughout the background
1783 predictions shown below.

1784 **5.4.1 COMMON PRE-SELECTION**

1785 Both the cut-based and BDT analyses have a common pre-selection that is applied before the signal
1786 region requirements. The requirements on leptons are common to all n_j bins. The analysis requires two
1787 oppositely charged leptons, with the leading lepton required to have $p_T > 22$ GeV while the subleading
1788 lepton must have $p_T > 10$ GeV. Next, to remove low mass Z/γ^* events, a requirement on the dilepton
1789 mass $m_{\ell\ell} > 10(12)$ GeV is applied in the different (same) flavor channel. In the same flavor channels,
1790 there is an additional veto placed on the region around the Z peak, requiring that $|m_{\ell\ell} - m_Z| > 15$ GeV.

1791 There are also requirements on the amount of missing transverse momentum in the event. These are
1792 only applied in the same flavor channels, as in the different flavor channels $t\bar{t}$ is the dominant background
1793 in $n_j \geq 2$. The BDT analysis requires $p_T^{\text{miss}} > 40$ GeV and $E_T^{\text{miss}} > 45$ GeV. The cut-based analysis
1794 must select more tightly on these variables to have maximal sensitivity and thus requires $p_T^{\text{miss}} > 50$ GeV
1795 and $E_T^{\text{miss}} > 55$ GeV.

1796 Finally, because this analysis is focused on VBF Higgs production, a requirement on the jet multiplicity
1797 is placed, with $n_j \geq 2$. Additionally, the analysis requires that there are no jets identified as b-quarks in
1798 the event, or $n_b = 0$.

1799 **5.4.2 CUT-BASED SELECTION**

1800 The cut-based selection places sequential requirements on variables reconstructed from the VBF jets in
1801 order to increase the signal to background ratio.

1802 **GENERAL BACKGROUND REDUCTION**

1803 Top pair production is the primary background in the $n_j \geq 2$ bin. Even though $n_b = 0$ is required, an
1804 additional variable is constructed to further suppress the top background. There is often additional QCD
1805 radiation that accompanies the $t\bar{t}$ system when it is produced. Therefore, a variable which tests for the

1806 presence of this additional radiation, p_T^{sum} , is constructed. It is defined in equation 5.1.

$$p_T^{\text{sum}} = p_T^{\ell\ell} + p_T^{\text{miss}} + \sum p_T^j \quad (5.1)$$

1807 After pre-selection, the cut-based analysis requires the event to have $p_T^{\text{sum}} < 15 \text{ GeV}$ to further suppress
1808 $t\bar{t}$ production.

1809 In the different flavor channels, a requirement is made to reduce the contamination from $Z \rightarrow \tau\tau$
1810 decays. The di- τ invariant mass, $m_{\tau\tau}$, is constructed by assuming that the neutrinos from the τ decays
1811 were collinear with the leptons [66]. The analysis requires that this mass satisfy $m_{\tau\tau} < m_Z - 25 \text{ GeV}$ so
1812 that it is not consistent with the mass of the Z boson.

1813 VBF TOPOLOGICAL CUTS

1814 The characteristic feature of VBF production of the Higgs is the presence of two additional forward
1815 jets coming from the incoming partons which radiate the vector bosons that make the Higgs. These jets
1816 are forward because the outgoing partons still carry the longitudinal momentum of the incoming partons.

1817 Figure 5.2 shows the distribution of the η for the leading jet in a VBF event compared to a background top
1818 pair production event. As can be seen, the VBF jets tend to be more forward in η , while the $t\bar{t}$ jets are more
1819 central. Because the cross section for VBF production is an order of magnitude smaller than gluon fusion
1820 production, these forward jets must be used in order to reduce background and achieve a good signal to
1821 background ratio. The dedicated VBF search selection requirements are constructed to maximally exploit
1822 the features of the unique VBF topology.

1823 Requirements on the VBF jets are collectively referred to as the “VBF topological cuts”. First, a require-
1824 ment on the dijet invariant mass of the VBF jets, m_{jj} , is placed, requiring $m_{jj} > 600 \text{ GeV}$. Next, the
1825 event is required to have a large gap in rapidity between the two VBF jets, or $\Delta y_{jj} > 3.6$. Both of these
1826 are tight requirements on the presence of two forward, high p_T jets moving in opposite directions in the
1827 longitudinal plane.

1828 Beyond requiring the presence of the two forward VBF jets, the analysis also vetoes on the presence of
1829 any additional jets that fall between the two VBF jets. This requirement is referred to as the central jet

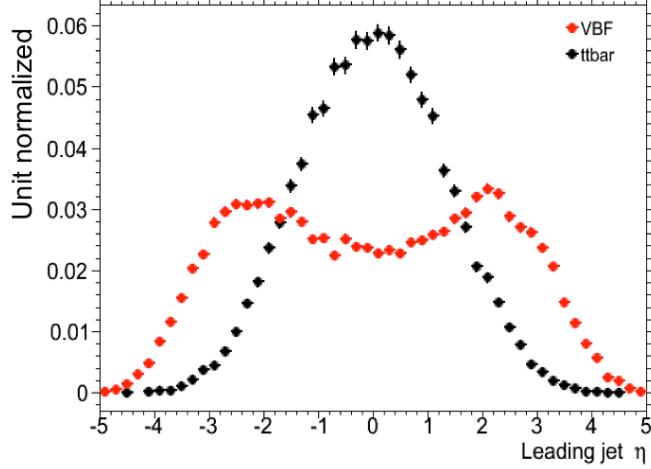


Figure 5.2: Leading jet η in VBF $H \rightarrow WW^*$ (red) and $t\bar{t}$ (black)

1830 veto, or CJV. Events are vetoed if they have a third jet with $p_T > 20$ GeV whose rapidity is between the
 1831 region defined by the two VBF jets. This requirement can be expressed in terms of a variable called the jet
 1832 centrality, defined in equation 5.2.

$$C_{j3} = \left| \eta_{j3} - \frac{\eta_{j1} + \eta_{j2}}{2} \right| / \frac{|\eta_{j1} - \eta_{j2}|}{2}, \quad (5.2)$$

1833 Here, η_{j1} and η_{j2} are the pseudorapidities of the leading and subleading jets, respectively, while η_{j3} is
 1834 the pseudorapidity of the extra jet in the event (if one exists). Intuitively, C_{j3} is zero when η_{j3} is directly
 1835 centered between the two jets and unity when η_{j3} is aligned with either of the VBF jets. Thus, the CJV
 1836 can be expressed as a requirement that $C_{j3} > 1$.

1837 The decay products of the Higgs tend to be central as well. Thus, the analysis also requires that both
 1838 leptons in the analysis fall within the rapidity gap defined by the jets. This cut is referred to as the outside
 1839 lepton veto, or OLV. Stated another way, leptons are required to have a centrality (defined analogously to
 1840 that of the third jet in equation 5.2) within the jet rapidity gap, or $C_\ell < 1$ for both leptons.

1841 Figure 5.3a-c shows the m_{jj} , Δy_{jj} , and $C_{\ell 1}$ variables at the stage where all previous requirements in the
 1842 sequence have been made. The agreement between data and Monte Carlo is good, and the bottom panels

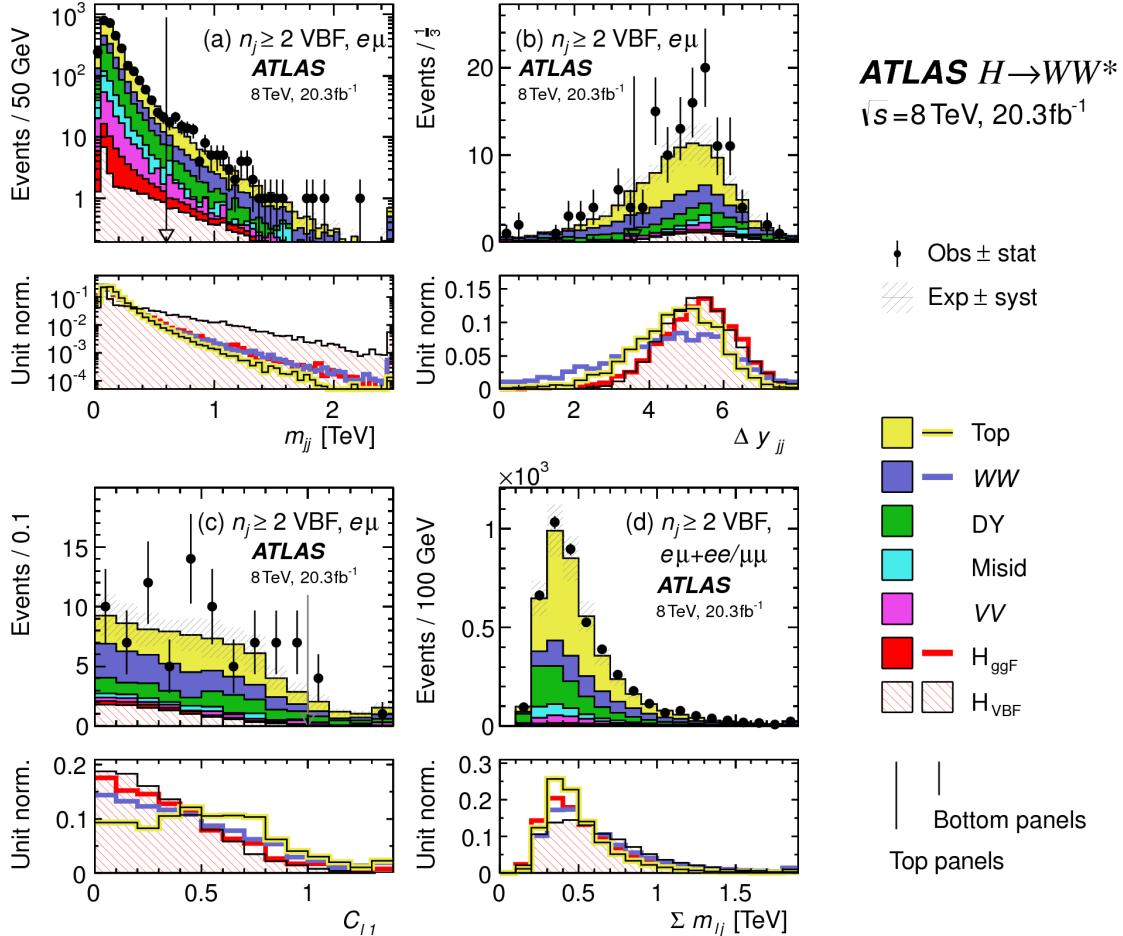


Figure 5.3: Distributions of (a) m_{jj} , (b) Δy_{jj} , (c) $C_{\ell 1}$, and (d) $\Sigma m_{\ell j}$, for the cut-based VBF analysis. The top panels compare simulation and data, while the bottom panels show normalized distributions for all background processes and signal for shape comparisons [63].

1843 show their power in discriminating the VBF signal from the background processes.

1844 The final signal region is also split into two bins of m_{jj} , with the first bin corresponding to $600\text{ GeV} <$
 1845 $m_{jj} < 1\text{ TeV}$ and the second bin corresponding to $m_{jj} > 1\text{ TeV}$. The first bin has more events but also
 1846 a larger contribution from background, while the second bin has a lower expected number of events but a
 1847 1:1 signal to background ratio.

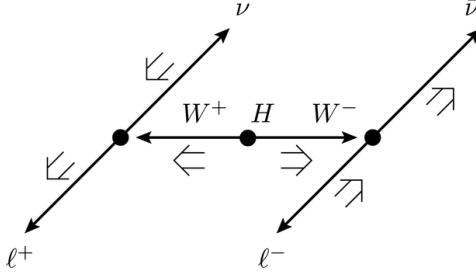


Figure 5.4: A cartoon of the WW final state. Momenta are represented with thin arrows, spins with thick arrows [63].

1848 HIGGS TOPOLOGICAL CUTS

1849 The final state leptons will exhibit unique correlations due to the fact that they arise from the decay
 1850 of a spin zero resonance. In particular, the spins of the final state leptons and neutrinos must all cancel,
 1851 as shown in figure 5.4. Because the neutrino has a left handed chirality and the anti-neutrino has a right
 1852 handed chirality (in the massless neutrino approximation), the spin and momentum of the particles will
 1853 be anti-aligned and aligned, respectively. In the transverse plane, the momenta of all four final state objects
 1854 must cancel as well. With the constraint of having both the momenta and the spin alignments cancel, the
 1855 final state kinematics strongly prefer having a small angle between the leptons in the transverse plane (low
 1856 $\Delta\phi_{\ell\ell}$). This angular correlation will also lead to low values of the di-lepton invariant mass $m_{\ell\ell}$. These
 1857 unique signal final state kinematic correlations are exploited to define the ultimate signal region.

1858 Two requirements on dilepton kinematics are made that are common with lower multiplicity jet bins
 1859 as well. The angle between leptons in the transverse plane, $\Delta\phi_{\ell\ell}$, is required to be less than 1.8 radians.
 1860 Additionally, the dilepton invariant mass, $m_{\ell\ell}$, is required to be less than 50 GeV.

1861 The cut-based analysis uses m_T as the final discriminating variable as in the ggF focused analysis. The
 1862 optimal number of bins in m_T was found to be three bins, with the bin boundaries at 80 and 130 GeV.

1863 Table 5.7 shows a summary of the data and estimated signal and background yields from simulation
 1864 as each requirement described above is made. The table shows how the overall signal to background ra-
 1865 tio grows through the various selection requirements. Table 5.8 shows the background composition after
 1866 each selection requirement, illustrating which backgrounds are reduced most by certain requirements. Fig-

¹⁸⁶⁷ ure 5.5 shows an ATLAS event display of a candidate event in the final signal region.

Selection	Summary					
	$\bar{N}_{\text{obs}}/\bar{N}_{\text{bkg}}$	\bar{N}_{obs}	\bar{N}_{bkg}	\bar{N}_{signal}		
				N_{ggF}	N_{VBF}	N_{VH}
$e\mu$ sample	1.00 ± 0.00	61434	61180	85	32	26
$n_b = 0$	1.02 ± 0.01	7818	7700	63	26	16
$p_T^{\text{sum}} < 15$	1.03 ± 0.01	5787	5630	46	23	13
$m_{\tau\tau} < m_Z - 25$	1.05 ± 0.02	3129	2970	40	20	9.9
$m_{jj} > 600$	1.31 ± 0.12	131	100	2.3	8.2	—
$\Delta y_{jj} > 3.6$	1.33 ± 0.13	107	80	2.1	7.9	—
$C_{j3} > 1$	1.36 ± 0.18	58	43	1.3	6.6	—
$C_{\ell 1} < 1, C_{\ell 2} < 1$	1.42 ± 0.20	51	36	1.2	6.4	—
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T$	2.53 ± 0.71	14	5.5	0.8	4.7	—
$ee/\mu\mu$ sample	0.99 ± 0.01	26949	27190	31	14	10.1
$n_b, p_T^{\text{sum}}, m_{\tau\tau}$	1.03 ± 0.03	1344	1310	13	8.0	4.0
$m_{jj}, \Delta y_{jj}, C_{j3}, C_\ell$	1.39 ± 0.28	26	19	0.4	2.9	0.0
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T$	1.63 ± 0.69	6	3.7	0.3	2.2	0.0

Table 5.7: Summary of event selection for the $n_j \geq 2$ VBF analysis in the 8 TeV cut-based analysis [63].

	Composition of N_{bkg}									
	N_{WW}		N_{top}		N_{misid}		N_{VV}		$N_{\text{Drell-Yan}}$	
	N_{WW}^{QCD}	N_{WW}^{EW}	$N_{t\bar{t}}$	N_t	N_{Wj}	N_{jj}	N_{VV}	$N_{ee/\mu\mu}$	$N_{\tau\tau}^{\text{QCD}}$	$N_{\tau\tau}^{\text{EW}}$
$e\mu$ sample	1350	68	51810	2970	847	308	380	51	3260	46
$n_b = 0$	993	43	3000	367	313	193	273	35	2400	29
$p_T^{\text{sum}} < 15$	781	38	1910	270	216	107	201	27	2010	23
$m_{\tau\tau} < m_Z - 25$	484	22	1270	177	141	66	132	7.6	627	5.8
$m_{jj} > 600$	18	8.9	40	5.3	1.8	2.4	5.1	0.1	15	1.0
$\Delta y_{jj} > 3.6$	11.7	6.9	35	5.0	1.6	2.3	3.3	—	11.6	0.8
$C_{j3} > 1$	6.9	5.6	14	3.0	1.3	1.3	2.0	—	6.8	0.6
$C_{\ell 1} < 1, C_{\ell 2} < 1$	5.9	5.2	10.8	2.5	1.3	1.3	1.6	—	5.7	0.6
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T$	1.0	0.5	1.1	0.3	0.3	0.3	0.6	—	0.5	0.2
$ee/\mu\mu$ sample	594	37	23440	1320	230	8.6	137	690	679	16
$n_b, p_T^{\text{sum}}, m_{\tau\tau}$	229	12.0	633	86	26	0.9	45	187	76	1.5
$m_{jj}, \Delta y_{jj}, C_{j3}, C_\ell$	3.1	3.1	5.5	1.0	0.2	0.0	0.7	3.8	0.7	0.1
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_T$	0.4	0.2	0.6	0.2	0.2	0.0	0.1	1.5	0.3	0.1

Table 5.8: Background composition after each requirement in the $n_j \geq 2$ VBF analysis in the 8 TeV cut-based analysis [63].

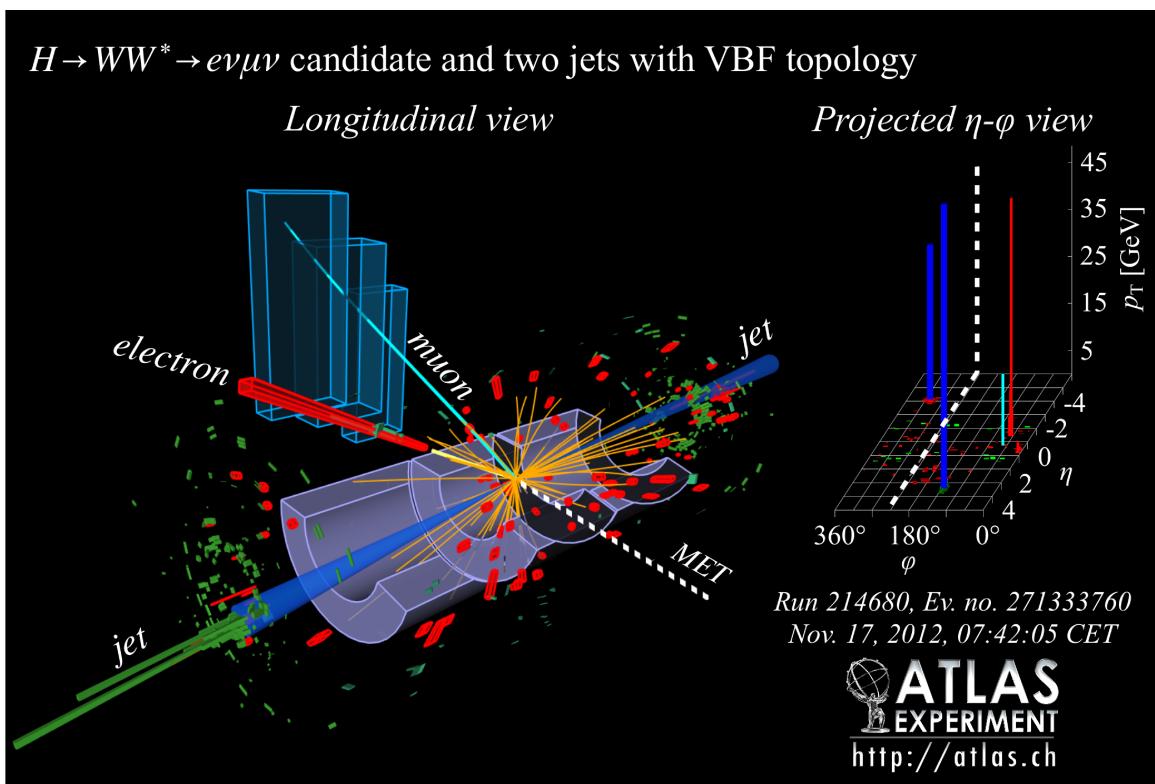


Figure 5.5: Event display of a VBF candidate event [63].

1868 5.4.3 BDT-BASED SELECTION

1869 The boosted decision tree based analysis takes a different philosophy compared to the cut-based. Rather
 1870 than making sequential requirements on many variables, the BDT analysis uses many of these variables as
 1871 inputs to the BDT. The output BDT score (O_{BDT}) is used as the final discriminant rather than m_T^2 .
 1872 The BDT is trained with the VBF $H \rightarrow WW^*$ simulation as the signal samples and all other processes as
 1873 background, including ggF $H \rightarrow WW^*$ production. While the BDT based analysis is treated as a separate
 1874 result, it has significant overlap with the cut-based selection.

1875 PRE-TRAINING SELECTION AND BDT INPUTS

1876 Before training, the common pre-selection cuts described in section 5.4.1 are applied. Additionally, the
 1877 central jet veto and outside lepton veto described in section 5.4.2 are applied. The BDT has eight input

²For the final discriminant analysis, the O_{BDT} distribution is divided into four bins, with boundaries at $[-1, -0.48, -0.3, 0.78, 1]$.

variables, six of which are also variables that are used in the cut-based analysis. The six shared variables are p_T^{sum} , m_{jj} , Δy_{jj} , $m_{\ell\ell}$, $\Delta\phi_{\ell\ell}$, and m_T . The seventh variable input in the BDT is a combination of the variables used to define the OLV in the cut-based analysis. The BDT uses as input the sum of lepton centralities, or $\sum C_\ell = C_{\ell 1} + C_{\ell 2}$. The final BDT input variable, $\Sigma m_{\ell j}$, is constructed to account for the correlations between the jets and leptons in the event. It is the sum of the invariant masses of all four possible lepton-jet combinations.

Figure 5.3d shows the agreement between data and simulation for the $\Sigma m_{\ell j}$ variable, as well as showing its discriminating power. Figure 5.6 shows the distributions of the Higgs topological variables that are shared between the cut-based and BDT analyses. Figure 5.7 shows the distributions of the VBF topological variables shared between the cut-based and BDT analyses. In both cases, the VBF yield has been scaled by a factor of 50 to better show the shape difference compared to the backgrounds.

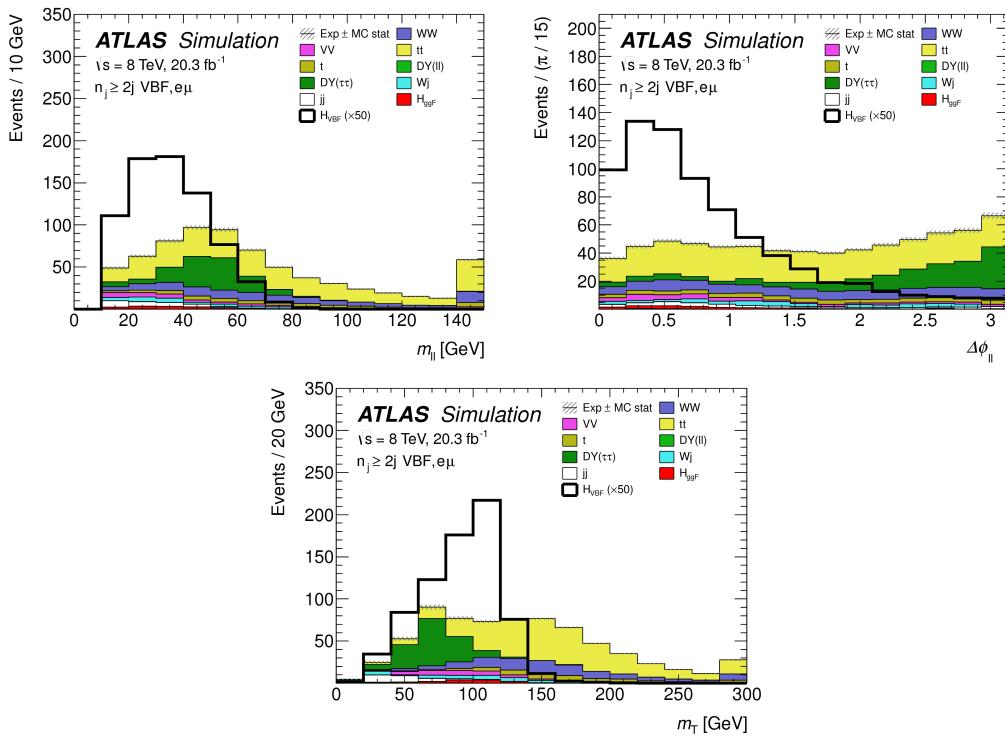


Figure 5.6: Distributions of $m_{\ell\ell}$ (top left), $\Delta\phi_{\ell\ell}$ (top right), and m_T (bottom), the Higgs topology variables used in the selection requirements of the cut-based signal region and as inputs to the BDT result. These are plotted after all of the BDT pre-training selection cuts [63].

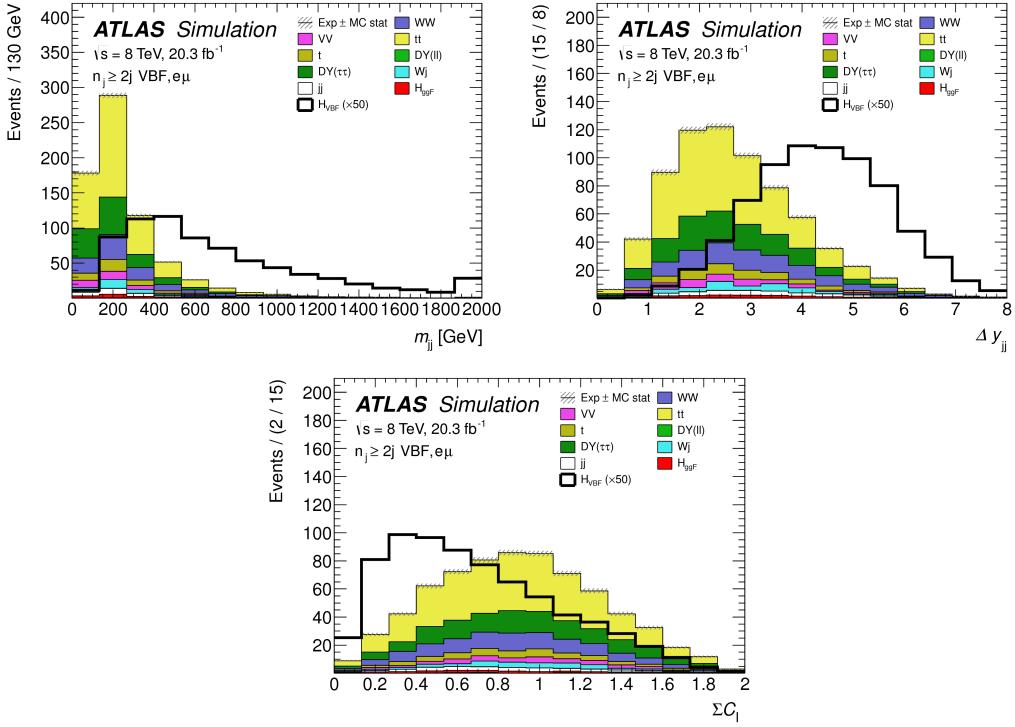


Figure 5.7: Distributions of m_{jj} (top left), Δy_{jj} (top right), $\sum C_\ell$ (bottom), the VBF topology variables used in the selection requirements of the cut-based signal region and as inputs to the BDT result. These are plotted after all of the BDT pre-training selection cuts [63].

1889 5.5 BACKGROUND ESTIMATION

1890 This section describes the procedures used to estimate backgrounds for the VBF analysis in both the
 1891 cut-based and BDT analyses.

1892 5.5.1 GENERAL STRATEGY

1893 Most of the backgrounds in the VBF Higgs analysis have shapes estimated from Monte Carlo simula-
 1894 tion but normalizations derived from control regions in data. In essence, a normalization factor (denoted
 1895 with β or abbreviated as NF) is derived by scaling the MC yield in the control region to the corresponding
 1896 yield in data. Once this factor is derived, it can be used to scale the MC estimate of the background in the

1897 signal region. This is illustrated in equation 5.3.

$$B_{\text{SR}}^{\text{est}} = B_{\text{SR}} \times \frac{N_{\text{CR}}}{B_{\text{CR}}} \equiv B_{\text{SR}} \times \beta \quad (5.3)$$

1898 Here, B is the MC yield prediction in the denoted region, while N is the observed number of events in
1899 data in the denoted region.

1900 There is an alternative way of writing the same equation in terms of an extrapolation factor α rather
1901 than a normalization factor β . The overall calculation is exactly the same. However, when phrased in
1902 this way, it shows how the uncertainty on the background estimation can be reduced. This is shown in
1903 equation 5.4.

$$B_{\text{SR}}^{\text{est}} = N_{\text{CR}} \times \frac{B_{\text{SR}}}{B_{\text{CR}}} \equiv N_{\text{CR}} \times \alpha \quad (5.4)$$

1904 Phrased this way, the equation shows that with enough events in the control region, a large theoretical
1905 uncertainty on the overall background yield in the signal region can be replaced by a small statistical un-
1906 certainty coming from the number of data events in the CR and a smaller theoretical uncertainty on the
1907 extrapolation from the control region to the signal region.

1908 5.5.2 TOP BACKGROUND

1909 The normalization factor β_t for the top background in the VBF analysis is derived in a region required to
1910 have one b-tagged jet, or $n_b = 1$. In the cut-based analysis, normalization factors are computed after every
1911 selection requirement by making the same requirements in the CR. These NF are then applied to the $t\bar{t}$ and
1912 single top event yields in the SR. In the BDT analysis, a single normalization factor is computed for each
1913 bin of O_{BDT} after applying the BDT pre-training cuts described previously. The computed normaliza-
1914 tion factors are derived with all flavor combinations combined in order to decrease statistical uncertainty.
1915 Additionally, in the BDT analysis, BDT bins 2 and 3 are merged for the same reason.

1916 Table 5.9 shows the evolution of the β_t through the cut-based selection. Table 5.10 shows the value of
1917 the β_t in each bin of O_{BDT} . In all cases, the computed factors are relatively consistent with unity, with
1918 the largest discrepancy coming in bin 1 of O_{BDT} . The normalization factors in the bins of O_{BDT} are also

¹⁹¹⁹ consistent with those derived in the cut-based signal region, increasing confidence in the BDT estimation.

Figure 5.8 shows the m_{jj} and O_{BDT} distributions in the top control region. Overall the modeling looks

Cut	β_t
$p_T^{\text{sum}} < 15 \text{ GeV}$	1.03 ± 0.01
$m_{\tau\tau} < m_Z - 25$	1.05 ± 0.01
$m_{jj} > 600 \text{ GeV}$	0.96 ± 0.06
$\Delta y_{jj} > 3.6$	1.02 ± 0.08
CJV	1.13 ± 0.16
OLV	1.01 ± 0.19
$m_{jj} < 1 \text{ TeV}$	0.94 ± 0.19
$m_{jj} > 1 \text{ TeV}$	1.48 ± 0.66

Table 5.9: Top normalization factors computed at each stage of the cut-based selection. Uncertainties are statistical only.

O_{BDT}	β_t
Bin 0	1.09 ± 0.02
Bin 1	1.58 ± 0.15
Bin 2	0.95 ± 0.31
Bin 3	0.95 ± 0.31

Table 5.10: Top normalization factors computed for each bin of O_{BDT} . Uncertainties are statistical only.

¹⁹²⁰

¹⁹²¹ consistent with the data. While these normalization factors can be computed and applied to the expected
¹⁹²² background yields listed in tables like table 5.8, the final normalization of the top background is profiled
¹⁹²³ (meaning there is a dedicated Poisson constraint) and allowed to float in the final statistical fit.

¹⁹²⁴ 5.5.3 $Z/\gamma^* \rightarrow \tau\tau$ BACKGROUND

¹⁹²⁵ In the different flavor channels, the $Z/\gamma^* \rightarrow \tau\tau$ background is an important one. Di-tau production
¹⁹²⁶ can produce an $e\mu$ final state if each τ lepton decays to a different flavor lepton.

¹⁹²⁷ In the BDT analysis, a single normalization factor for the background is derived. A control region
¹⁹²⁸ is defined using the pre-training selection cuts, except requiring that $|m_{\tau\tau} - m_Z| < 25 \text{ GeV}$ so that
¹⁹²⁹ the region is enriched in $Z/\gamma^* \rightarrow \tau\tau$ background. Additional requirements of $m_{\ell\ell} < 80(75) \text{ GeV}$
¹⁹³⁰ in the different (same) flavor channel, as well as $O_{\text{BDT}} > -0.48$ are applied to increase the purity of the

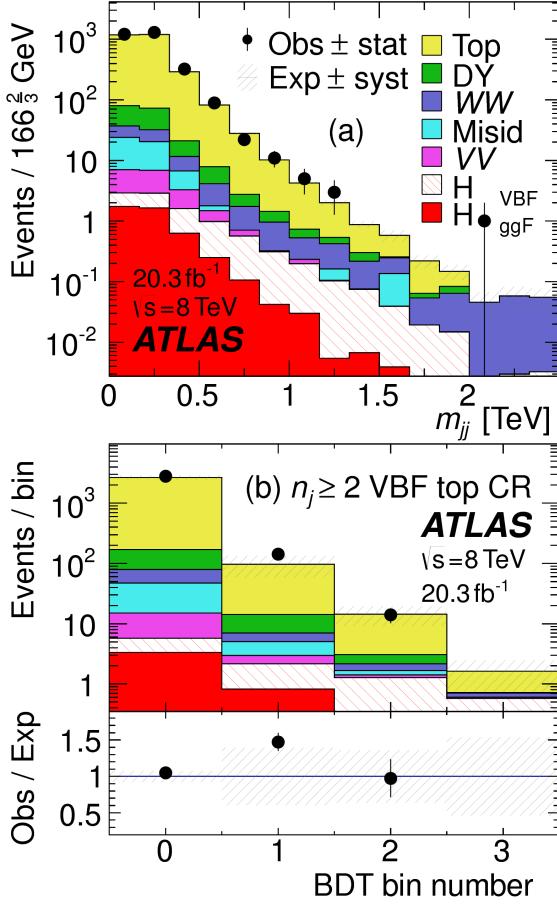


Figure 5.8: Distributions of m_{jj} (a) and O_{BDT} (b) in the VBF $n_b = 1$ top CR [63].

region. The final $\beta_{Z/\gamma^* \rightarrow \tau\tau}$ is calculated to be 0.9 ± 0.3 (statistical uncertainty only). Because of the small contribution of this background in the BDT analysis and the large statistical uncertainty, no additional systematics are calculated. The final SR estimate is scaled by this β and not allowed to float in the fit.

The cut-based corrections are a bit more involved because they need to be applied selection by selection, as well as in the final signal region for the fit. The control region is defined including all SR requirements up to the $Z/\gamma^* \rightarrow \tau\tau$ veto, which is instead turned into a Z mass peak requirement as for the BDT region. The $m_{\ell\ell}$ cut from the BDT region is included as well. The cut-based approach aims to correct the normalization of the $Z/\gamma^* \rightarrow \tau\tau$ background in two ways. First, an overall normalization factor is computed from the control region. However, the VBF topological cuts are not included in this region, and applying them as is done in the top CR is not feasible due to limited statistics. So, instead, correction

1941 factors (CF) to the cut efficiencies of the VBF cuts are derived in a same flavor $Z \rightarrow \ell\ell$ control region,
 1942 which has significantly more statistics. The CF is simply the ratio of the cut efficiencies in data and MC
 1943 derived in this region. In the end, the overall background estimate is given by equation 5.5.

$$N_{Z/\gamma^* \rightarrow \tau\tau}^{\text{est}} = B_{Z/\gamma^* \rightarrow \tau\tau}^{\text{SR}} \times \beta_{\tau\tau} \times \frac{\epsilon_{\text{VBF cuts}}^{\text{data}}}{\epsilon_{\text{VBF cuts}}^{\text{MC}}} \quad (5.5)$$

1944 The hypothesis is that while the normalization correction must be derived in a dedicated region, the effi-
 1945 ciency of the VBF topology requirements should not be sensitive to the type of Z/γ^* process and thus the
 1946 higher number of events can be exploited to derive the CF. Figure 5.9 shows a shape comparison for the
 1947 m_{jj} variable in $Z \rightarrow \tau\tau$ events in the signal region and $Z \rightarrow \ell\ell$ events in the control region. The figure
 1948 shows that the shapes are indeed comparable and thus any CF derived in the same flavor control region
 1949 can reliably be applied in the signal region.

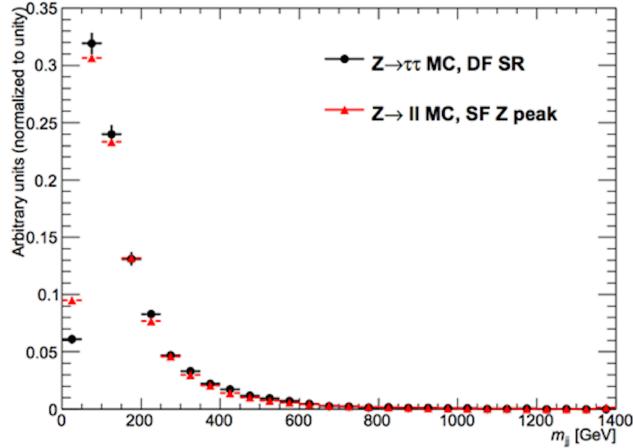


Figure 5.9: Comparison of m_{jj} shape in a same flavor $Z \rightarrow \ell\ell$ control region and the VBF cut-based signal region.

1950 Table 5.11 shows the overall normalization factor $\beta_{\tau\tau}$ and the efficiency correction factors for the various
 1951 VBF topological cuts. In general, the statistical uncertainties on the cut efficiency corrections are quite
 1952 good, and the MC tends to underestimate the efficiency of the VBF cuts for the $Z/\gamma^* \rightarrow \tau\tau$ background.
 1953 The overall normalization factor is also consistent with that calculated for the BDT analysis.

$\beta_{\tau\tau}$	0.97 ± 0.04
Cut	Correction factors
$m_{jj} > 600 \text{ GeV}$	1.09 ± 0.01
$\Delta y_{jj} > 3.6$	1.14 ± 0.02
CJV	1.20 ± 0.02
OLV	1.17 ± 0.03
$m_{jj} < 1 \text{ TeV}$	1.17 ± 0.06
$m_{jj} > 1 \text{ TeV}$	1.18 ± 0.13

Table 5.II: $Z/\gamma^* \rightarrow \tau\tau$ correction factors for the VBF cut-based analysis. Uncertainties are statistical only.

1954 5.5.4 $Z/\gamma^* \rightarrow \ell\ell$ BACKGROUND

1955 In the same flavor channels, the $Z/\gamma^* \rightarrow \ell\ell$ background is dominant and thus must be estimated cor-
 1956 rectly. In both the BDT and cut-based analyses, the background is estimated using the so-called “ABCD”
 1957 method. The ABCD method creates four different regions by defining requirements on two variables.
 1958 One of the regions (A) is the signal region, while the other regions are defined by inverting one of both of
 1959 the requirements. in this case, the two variables used are $m_{\ell\ell}$ and E_T^{miss} , because inverting either of the
 1960 SR cuts on these variables will give regions rich in the $Z/\gamma^* \rightarrow \ell\ell$ background. Figure 5.10 illustrates the
 1961 definitions of each region.

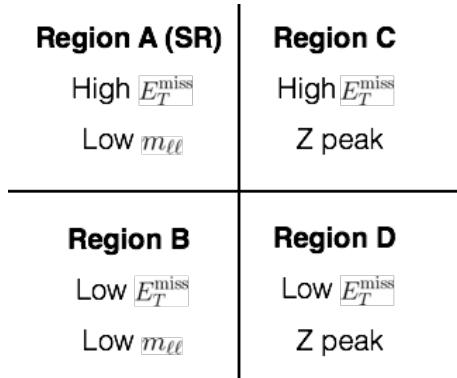


Figure 5.10: General illustration of the ABCD region definitions for $Z/\gamma^* \rightarrow \ell\ell$ background estimation.

1962 In both of the cut-based and BDT analyses, the Z peak region is defined with $|m_{\ell\ell} - m_Z| < 15 \text{ GeV}$.
 1963 In the cut-based analysis, low $m_{\ell\ell}$ corresponds to $m_{\ell\ell} < 50 \text{ GeV}$ (this defines the cut-based SR) while
 1964 in the BDT it is $m_{\ell\ell} < 75 \text{ GeV}$. In the cut-based, high and low E_T^{miss} are defined as opposite ends of

1965 the 55 GeV cut applied for the signal region definition. The BDT low E_T^{miss} region is between 25 and
 1966 45 GeV, while the high E_T^{miss} region is $E_T^{\text{miss}} > 45$ GeV.

1967 Once the regions are defined, the background in the signal region is estimated by extrapolating the
 1968 estimate in region B to region A. This extrapolation is done by multiplying the number of events in region
 1969 B by the ratio of the number of events in regions C and D. Effectively, the Z peak region is used to estimate
 1970 the efficiency of the E_T^{miss} requirement in data, and then this efficiency is applied in the low $m_{\ell\ell}$ region.
 1971 An additional correction is also applied for the non-closure of the method in MC. This is summarized in
 1972 equations 5.6 and 5.7.

$$N_{Z/\gamma^* \rightarrow \ell\ell}^{\text{SR}} = N_{Z/\gamma^* \rightarrow \ell\ell}^B \times \frac{N_{Z/\gamma^* \rightarrow \ell\ell}^C}{N_{Z/\gamma^* \rightarrow \ell\ell}^D} \times f_{\text{corr}} \quad (5.6)$$

1973

$$f_{\text{corr}} = \frac{B_{\text{MC}}^A / B_{\text{MC}}^B}{B_{\text{MC}}^C / B_{\text{MC}}^D} \quad (5.7)$$

1974 Here, the N refer to data yields in each region with the non Z/γ^* backgrounds subtracted, while B refer
 1975 to the Z/γ^* yields in MC in each region.

1976 A normalization factor $\beta_{\ell\ell}$ is computed for each analysis as the ratio of the predicted data yield to the
 1977 MC yield in the SR. The shape of the BDT distribution is taken from data region B, while the shape of
 1978 the m_T distribution in the cut-based analysis is taken from Z/γ^* MC in the SR. The values of $\beta_{\ell\ell}$ in the
 1979 cut-based and BDT analyses from this method are summarized in table 5.12. They are quite consistent with
 1980 one another within the statistical uncertainties. In the cut-based analysis, the same cut efficiency correction
 1981 factors shown in table 5.11 are also applied (in product with the $\beta_{\ell\ell}$) in the same flavor channels to the Z/γ^*
 1982 background.

	β_t
BDT Bin 1	1.01 ± 0.15
BDT Bin 2	0.89 ± 0.28
Cut-based	0.81 ± 0.21

Table 5.12: $Z/\gamma^* \rightarrow \ell\ell$ normalization factors for cut-based and BDT analyses. Uncertainties are statistical only.

1983 5.5.5 WW AND OTHER DIBOSON BACKGROUNDS

1984 The Standard Model WW and other diboson backgrounds have both their shape and normalization
1985 taken from MC simulation. They are validated in dedicated control regions and found to agree with data
1986 well.

1987 As SM WW production is the largest of these backgrounds and is irreducible, validating the estimate
1988 is of particular importance. A validation region is constructed by requiring the pre-selection requirements
1989 on leptons and $m_{\ell\ell}, n_b = 0$, and $m_T > 100$ GeV. The m_{T2} variable [86] is an additional discriminant
1990 that will isolate the SM WW background, and a requirement of $m_{T2} > 160$ GeV is placed to define
1991 the WW validation region. This requirement gives a 60% purity for the validation region. The derived
1992 normalization factor in the region is 1.15 ± 0.19 and is thus consistent with unity. Figure 5.ii shows the
1993 m_{T2} distribution and how it distinguishes the WW background.

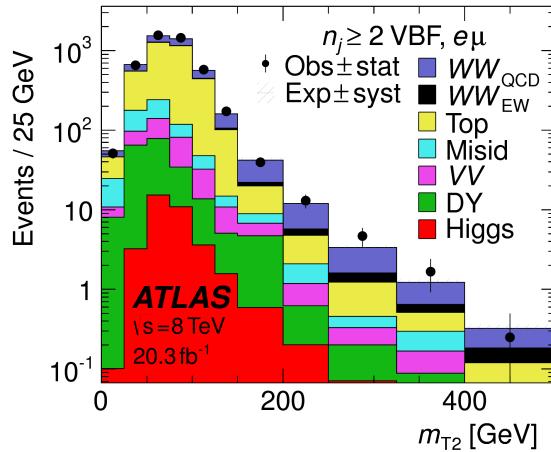


Figure 5.ii: Distribution of m_{T2} in the WW validation region of the VBF analysis [63].

1994 5.5.6 HIGGS PRODUCTION VIA GLUON-GLUON FUSION

1995 Because this analysis is dedicated to measuring the VBF contribution to Higgs production, the compo-
1996 nent of Higgs production from gluon-gluon fusion is treated as a background. The shape is taken directly
1997 from simulation, using the generators described in table 5.4. In the final combined fit of all different signal
1998 regions, the normalization is controlled by either a combined signal strength parameter μ , which controls

1999 the normalization of both ggF and VBF production, or a separate parameter μ_{ggF} depending on the in-
2000 terpretation being presented in the final results.

2001 **5.5.7 BACKGROUNDS WITH MISIDENTIFIED LEPTONS**

2002 As discussed previously, the $W + \text{jets}$ and QCD multijet backgrounds are derived with fully data-driven
2003 methods. These backgrounds do not make a large contribution to the final VBF signal region but their
2004 estimation methods are discussed briefly here.

2005 **$W + \text{jets}$ BACKGROUND**

2006 The $W + \text{jets}$ background enters the signal region by having one of the jets mis-reconstructed as a lepton.
2007 The background is estimated by constructing a control sample with two leptons, where one lepton passes
2008 the usual lepton quality requirements but the second lepton fails one of those requirements (also known
2009 as the “anti-identified” lepton). This control region is rich in the $W + \text{jets}$ contribution because if a second
2010 lepton is reconstructed in a $W + \text{jets}$ event it is likely to be poor quality. The purity of this $W + \text{jets}$ control
2011 sample is 85% to 90% depending on the exact configuration of leptons in the final state.

2012 The signal region estimate of $W + \text{jets}$ is estimated by extrapolation from the control sample to the sig-
2013 nal region using extrapolation factors derived in a $Z + \text{jets}$ control sample in data. The extrapolation factor
2014 is the ratio of the number of lepton candidates satisfying all quality criteria to the number of lepton can-
2015 didates anti-identified. This ratio is measured in bins of p_T and η . Thus, the final signal region estimate
2016 (binned as the extrapolation factor is binned) is simply the number of events in the anti-identified lepton
2017 control sample multiplied by the extrapolation factor derived from the $Z + \text{jets}$ control sample. Figure 5.12
2018 shows the extrapolation factors derived for electrons and muons.

2019 **QCD MULTIJET BACKGROUND**

2020 The method for estimating the multijet background is very similar to the $W + \text{jets}$ estimation method.
2021 The control sample in this case has two anti-identified leptons but otherwise satisfies all signal region re-
2022 quirements. The extrapolation factor is estimated from a multijet sample and applied twice to the control
2023 sample.

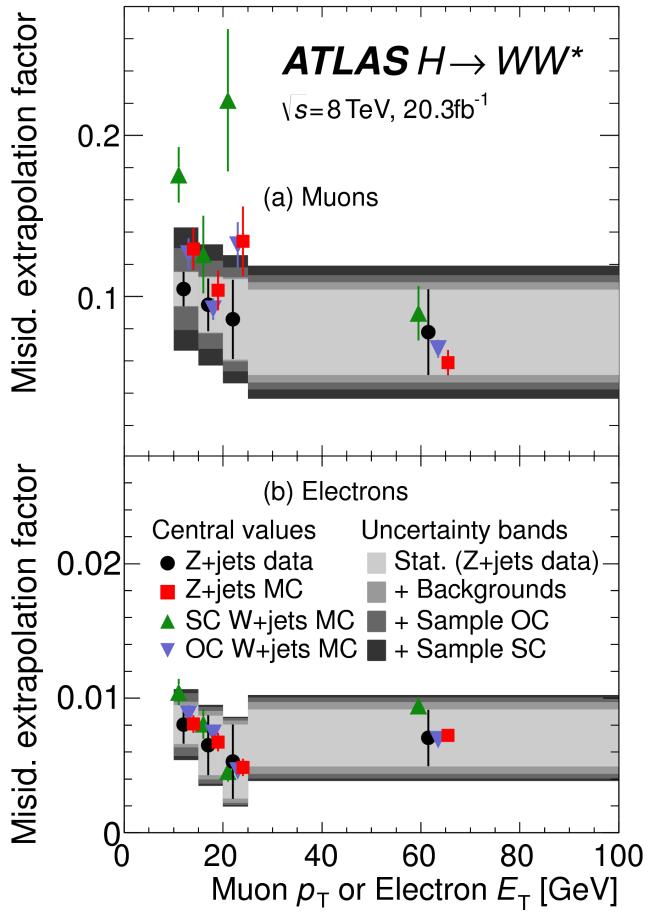


Figure 5.12: Extrapolation factors for the $W + \text{jets}$ estimate derived for muons (a) and electrons (b) as a function of lepton p_T [63].

2024 5.5.8 BACKGROUND COMPOSITION IN SIGNAL REGION

2025 After all of these estimation procedures, the signal region background composition can be calculated.
 2026 The estimated yields are all shown in table 5.8. Figure 5.13 shows the relative percentages of the differ-
 2027 ent background for the different flavor and same flavor final states. In $e\mu$, the leading backgrounds are
 2028 top backgrounds, ggF Higgs, and SM WW production. In $ee/\mu\mu$, the leading background is Drell-Yan,
 2029 followed by top and ggF Higgs.

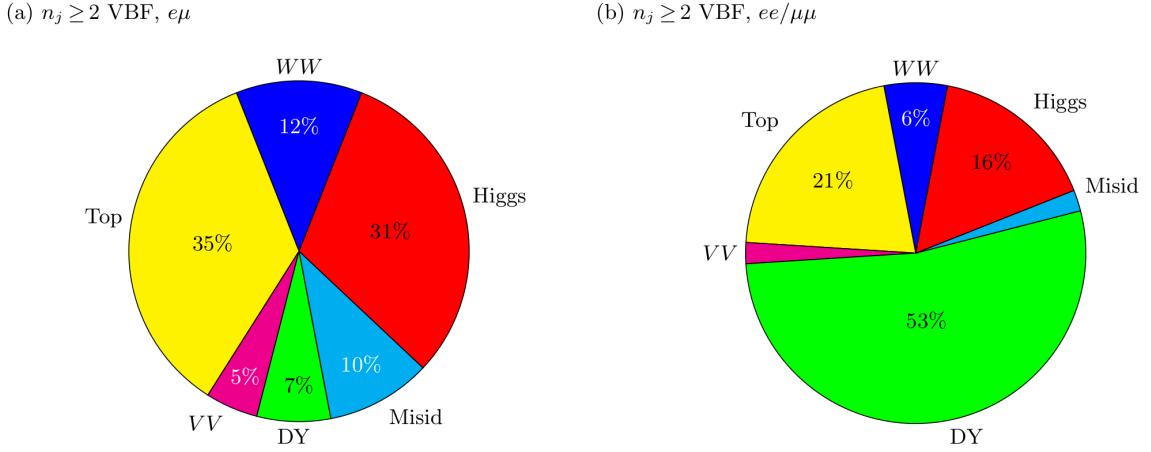


Figure 5.13: Background composition in final VBF signal region [63].

2030 5.6 SYSTEMATIC UNCERTAINTIES

2031 There are two main types of systematic uncertainties that are assessed for the analysis. First, theoretical
 2032 uncertainties associated with the signal and background yield estimates are discussed. Then, experimental
 2033 uncertainties due to detector effects are shown. Normalization uncertainties refer to uncertainties that
 2034 affect the cross section of the process in question in the signal region being probed. Shape uncertainties
 2035 refer to systematic uncertainties that affect the shape of the final discriminating variable (either m_T or
 2036 O_{BDT}).

2037 5.6.1 THEORETICAL UNCERTAINTIES

2038 There are four main components to theoretical uncertainties assigned to signal and background pro-
 2039 cesses taken from Monte Carlo. Each one is a different source of variation in the overall acceptance for
 2040 that process. The first involves variation of the QCD renormalization and factorization scales used in the
 2041 calculation. In this case, the two scales are varied both independently and simultaneously by factors of
 2042 two high or low. The resulting variation in normalization and shape for the process is taken as a systematic
 2043 uncertainty (referred to as scale uncertainty). This uncertainty approximates the level of the correction
 2044 to the cross section that would come from including the next order of the QCD calculation. Next, there
 2045 is an uncertainty associated with the PDF set used in generating the events. The uncertainty eigenvec-

2046 tors for the given PDF set are inspected, and the envelope of maximal variation is taken as an uncertainty
 2047 (referred to as PDF uncertainty). Finally, there are two uncertainties associated with the choice of MC
 2048 software. An uncertainty associated with the generator chosen for the hard scattering process is evaluated
 2049 by keeping the parton showering software constant but varying the matrix element generator and taking
 2050 the maximal variation as an uncertainty (referred to as the generator uncertainty). The converse variation
 2051 can also be done, where the matrix element generator remains constant and the generator used for the un-
 2052 derlying event/parton shower modeling is varied (referred to as the UE/PS uncertainty). In cases where
 2053 the background is normalized in a control region, the systematic uncertainty arises from variations of the
 2054 extrapolation factor α between the CR and the SR, which can affect the normalization of the background
 2055 in the SR.

2056 There are two additional uncertainties that are applied to the Higgs processes as well. First, there are
 2057 uncertainties assigned to the Higgs total production cross section. Then, there are uncertainties assigned
 2058 based on the fact that the analysis is done in exclusive jet bins and it is possible for signal events to migrate
 2059 from one bin to the next depending on the presence or absence of jets. These are assigned using the Jet Veto
 2060 Efficiency (JVE) procedure [18, 87] for ggF events and the Stewart-Tackmann (ST) method [88] for VBF
 2061 production. Table 5.13 shows the total theory uncertainties on the backgrounds in the cut-based analysis.
 2062 These are the sum in quadrature of the uncertainties from each of the variations described above.

Process	Theory syst. (%)
ggF H	48
Top	26
QCD WW	37
$Z/\gamma^* \rightarrow \tau\tau$	6.1

Table 5.13: Systematic uncertainties for various processes in the cut-based VBF analysis, given in units of % change
in yield. Values are given for the low m_{jj} signal region.

2063 Figures 5.14 and 5.15 show the variations in the extrapolation factor from the PDF and QCD uncertain-
 2064 ties on the top background estimate, binned in m_T , for the cut-based analysis. In both cases, there was
 2065 no significant shape uncertainty but normalization uncertainties were assigned according to the maximal
 2066 variation. These uncertainties enter into the 26% total uncertainty on top quark production quoted in
 2067 table 5.13

2068 While the estimate for the same-flavor $Z/\gamma^* \rightarrow \ell\ell$ background is data-driven, there is still a systematic
 2069 uncertainty taken for the non-closure of the method in Monte Carlo. This is taken as the maximum of the
 2070 deviation of the non-closure factor f_{corr} from unity and its uncertainty, or $\max(|1 - f_{\text{corr}}|, \delta f_{\text{corr}})$. For
 2071 the cut-based analysis this non-closure uncertainty 23%, while for the BDT analysis it is 17%.

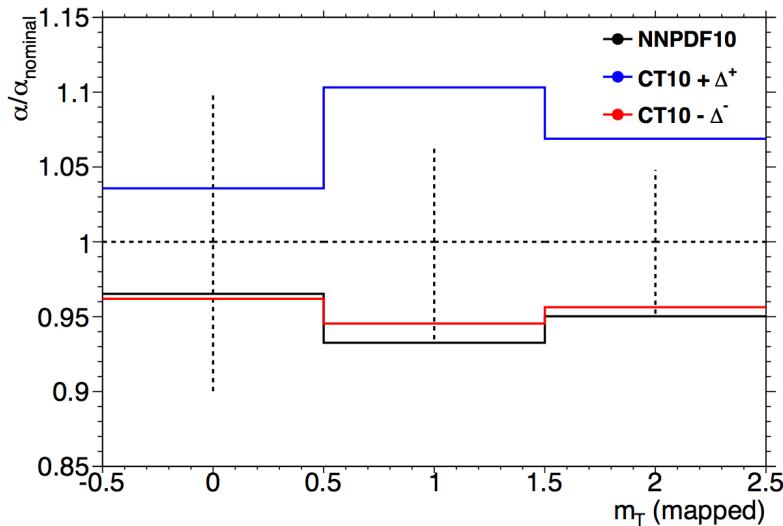


Figure 5.14: Variations in the top background extrapolation factor in the cut-based analysis due to PDF uncertainties. The uncertainties are shown in the three bins of m_T used in the final cut-based statistical fit. Variations from the eigenvector of the nominal PDF, CT10, as well as the result from an alternate PDF (NNPDF10), are compared.

2072 5.6.2 EXPERIMENTAL UNCERTAINTIES

2073 In this analysis, the theoretical uncertainties are the most dominant after statistical, but there are some
 2074 experimental uncertainties that make a contribution as well. The first is the uncertainty on the measured
 2075 integrated luminosity, which affects backgrounds whose normalizations are taken from MC and is mea-
 2076 sured to be 2.8% in the 8 TeV dataset [89]. The dominant sources of uncertainty overall are uncertain-
 2077 ties on the jet energy scale and resolution and the b -tagging efficiency. Additional sources include lepton
 2078 uncertainties on identification, resolution, and trigger efficiency, as well as uncertainties on the missing
 2079 transverse momentum.

2080 The jet energy scale uncertainty is split into several independent components, including jet-flavor de-
 2081 pendent calorimeter response uncertainties, uncertainties on modeling of pile-up interactions, uncertain-

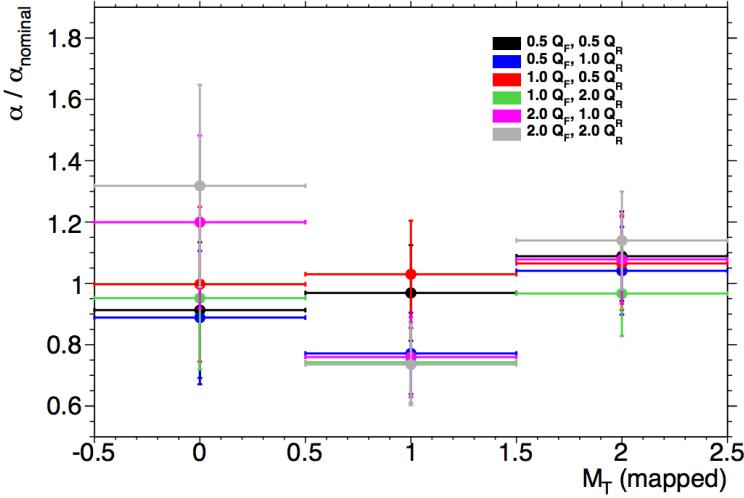


Figure 5.15: Variations in the top background extrapolation factor in the cut-based analysis due to QCD scale uncertainties. The uncertainties are shown in the three bins of m_T used in the final cut-based statistical fit. Q_F is the QCD factorization scale, while Q_R is the QCD renormalization scale.

ties on extrapolation from the central to forward detector regions, and MC non-closure [90]. The uncertainty on energy scale for jets used in this analysis ranges from 1% to 7% depending on the jet p_T and η . The jet energy resolution varies from 5% to 20%, with uncertainties ranging from 2% to 40% (the largest uncertainties occurring at the selection threshold).

The b-tagging efficiency is independently measured in data samples enriched in dileptonic decays of $t\bar{t}$ events or in events where a muon is reconstructed in the vicinity of a jet [91, 92]. The efficiencies and their uncertainties are binned in p_T and decomposed into uncorrelated components using an eigenvector method [93]. Uncertainties on the efficiency range from 1% to 7.8%. The uncertainty on the rate of misidentification of c -jets as b -jets ranges from 6-14%, while the uncertainty on the rate of light jet mistagging ranges from 9-19% depending on p_T and η .

The total experimental uncertainties on different signal and background components are summarized in table 5.14. They are compared to the level of other statistical and systematic uncertainties as well. Overall, the experimental uncertainties are sub-dominant compared to the statistical and theoretical uncertainties.

Sample	Total error	Stat. error	Expt. syst. err.	Theo. syst. err.
$n_j \geq 2$ VBF-enriched				
N_{sig}	13	—	6.8	12
N_{bkg}	9.2	4.7	6.4	4.5
N_{WW}	32	—	14	28
N_{top}	15	9.6	7.6	8.5
N_{misid}	22	—	12	19
N_{VV}	20	—	12	15
$N_{\tau\tau}$ (DY)	40	25	31	2.9
$N_{ee/\mu\mu}$ (DY)	19	11	15	—

Table 5.14: Composition of the post-fit uncertainties (in %) on the total signal (N_{sig}), total background (N_{bkg}), and individual background yields in the VBF analysis [63].

2095 5.7 RESULTS

2096 While the combined results of all the $H \rightarrow WW^*$ sub-analyses will be discussed in the next chapter,
 2097 this section presents the results of the VBF specific analysis and interpretations. As table 5.7 shows, the
 2098 final cut-based signal region contains 20 events in data with $m_T < 150$ GeV, 14 coming from the $e\mu$
 2099 channel and 6 coming from the $ee + \mu\mu$ channel. The BDT analysis has many more candidates due to its
 2100 looser selection, and the yields in each bin of O_{BDT} are shown in table 5.15.

2101 Figure 5.16(a) shows the final distribution of data candidates compared to the expected m_T distribution
 2102 for signal and background. The data are very consistent with a VBF Higgs hypothesis. Figure 5.16(b) shows
 2103 where the data candidates fall in the two-dimensional binning of m_T and m_{jj} used in the fit for the cut-
 2104 based analysis. Figure 5.17 shows the distributions of O_{BDT} and m_T in the VBF BDT analysis. Again the
 2105 data are quite consistent with a VBF Higgs hypothesis.

2106 Because the cut-based result is used as a validation for the BDT analysis and the two signal regions are
 2107 not fully orthogonal, it is interesting to explore which events overlap between the two analyses. Of the
 2108 twenty events in the cut-based signal region, only seven were not selected by the BDT analysis, while the
 2109 other thirteen also enter the BDT signal region. Figure 5.18 shows where the different analysis candidates
 2110 lie in the m_{jj} - m_T plane. This shows clearly that the advantage of the BDT analysis is that it can extract
 2111 signal candidates from the lower m_{jj} region due to its ability to recognize correlations with other variables.

(a) Before the BDT classification

Selection	Summary						Composition of N_{bkg}									
	$N_{\text{obs}}/N_{\text{bkg}}$	N_{obs}	N_{bkg}	N_{signal}			N_{WW}^{QCD}	N_{WW}^{EW}	$N_{t\bar{t}}$	N_t	N_{Wj}	N_{jj}	N_{VV}	$N_{\text{Drell-Yan}}$	$N_{ee/\mu\mu}^{\text{QCD}}$	$N_{\tau\tau}^{\text{EW}}$
				N_{ggF}	N_{VBF}	N_{VH}										
$e\mu$ sample	1.04 ± 0.04	718	689	13	15	2.0	90	11	327	42	29	23	31	2.2	130	2
$ee/\mu\mu$ sample	1.18 ± 0.08	469	397	6.0	7.7	0.9	37	3	132	17	5.2	1.2	10.1	168	23	1

(b) Bins in O_{BDT}

$e\mu$ sample																
Bin 0 (not used)	1.02 ± 0.04	661	650	8.8	3.0	1.9	83	9	313	40	26	21	28	2.2	126	1
Bin 1	0.99 ± 0.16	37	37	3.0	4.2	0.1	5.0	1.0	17	3.1	3.3	1.8	2.6	—	4.0	0.2
Bin 2	2.26 ± 0.63	14	6.2	1.2	4.2	—	1.5	0.5	1.8	0.3	0.4	0.3	0.8	—	0.3	0.3
Bin 3	5.41 ± 2.32	6	1.1	0.4	3.1	—	0.3	0.2	0.3	0.1	—	—	0.1	—	0.1	0.1
$ee/\mu\mu$ sample																
Bin 0 (not used)	1.91 ± 0.08	396	345	3.8	1.3	0.8	33	2	123	16	4.1	1.1	8.8	137	20.5	0.5
Bin 1	0.82 ± 0.14	53	45	1.5	2.2	0.1	3.0	0.5	10.4	1.8	0.8	0.2	0.9	26	1.7	0.1
Bin 2	1.77 ± 0.49	14	7.9	0.6	2.5	—	0.8	0.3	1.1	0.2	0.2	—	0.3	4.4	0.3	0.1
Bin 3	6.52 ± 2.87	6	0.9	0.2	1.7	—	0.1	0.2	0.2	—	—	—	0.7	—	—	—

Table 5.15: Event selection for the VBF BDT analysis. The event yields in (a) are shown after the pre-selection and the additional requirements applied before the BDT classification (see text). The event yields in (b) are given in bins in O_{BDT} after the classification [63].

While the context of these results in the broader $H \rightarrow WW^*$ statistical analysis will be presented in the next chapter, the statistical significance of the VBF Higgs result is shown here. In the BDT analysis, the expected signal significance was 2.7σ , while the observed significance was 3.1σ . In the cut-based analysis, the expected significance was 2.1σ and the observed significance was 3.0σ . The compatibility between these two results can be evaluated by computing the probability of observing a larger difference in Z_0 values than the one measured. Using toy Monte Carlo with the ggF signal strength fixed to unity and considering only statistical uncertainties, this probability is computed to be 79%, indicating good agreement between the analyses. This result represents the first evidence of the vector boson fusion production of a Higgs boson.

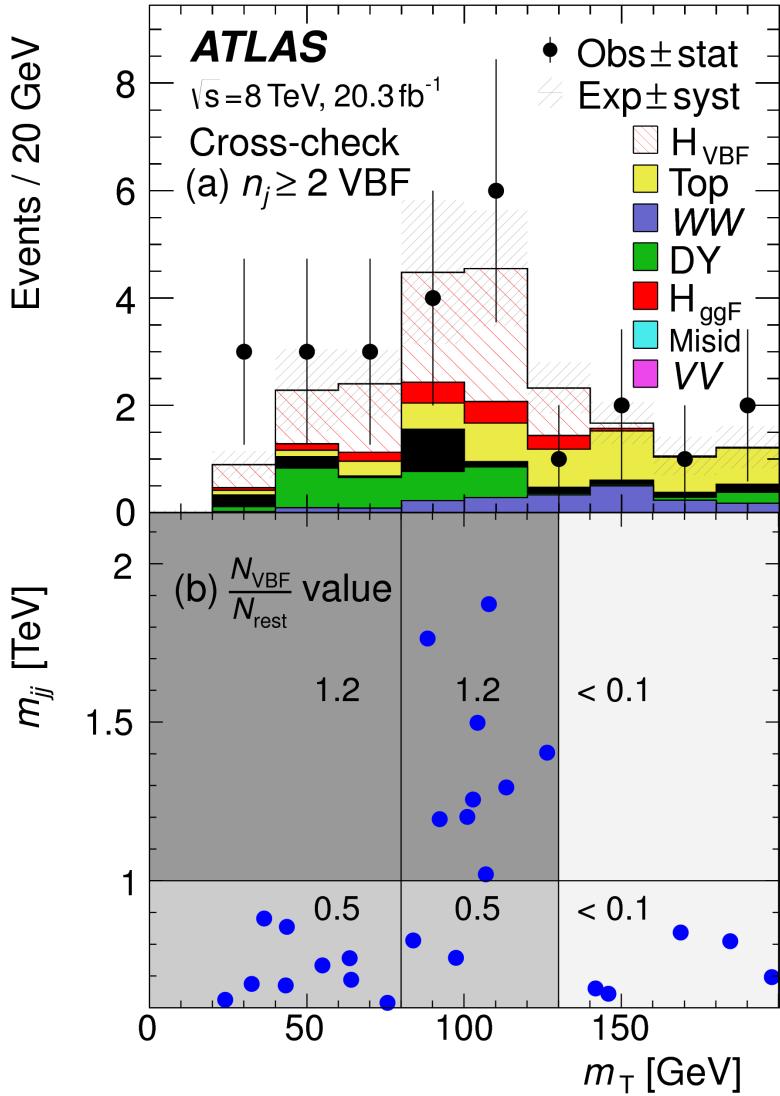


Figure 5.16: Post-fit distributions in the cut-based VBF analysis. Panel (a) shows the one-dimensional m_T distribution, while (b) shows the data candidates split into the bins of m_T and m_{jj} used in the final fit [63].

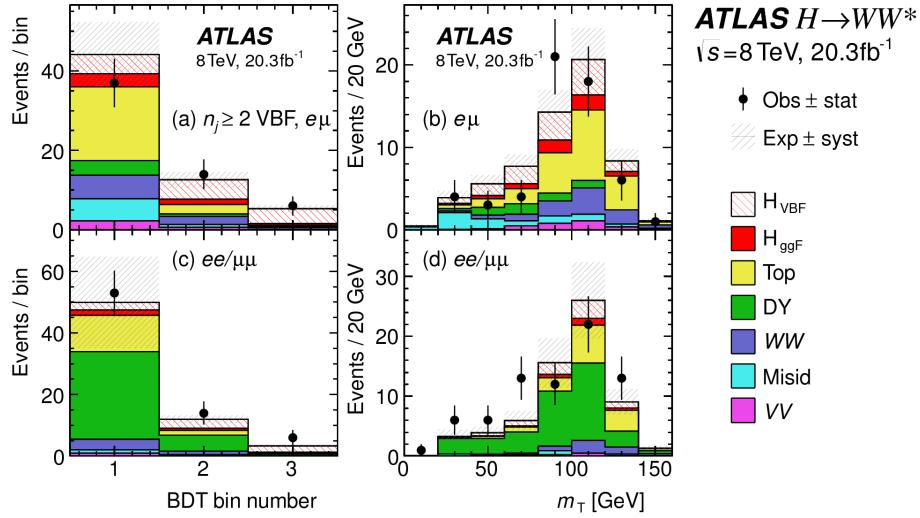


Figure 5.17: Postfit distributions in the BDT VBF analysis [63].

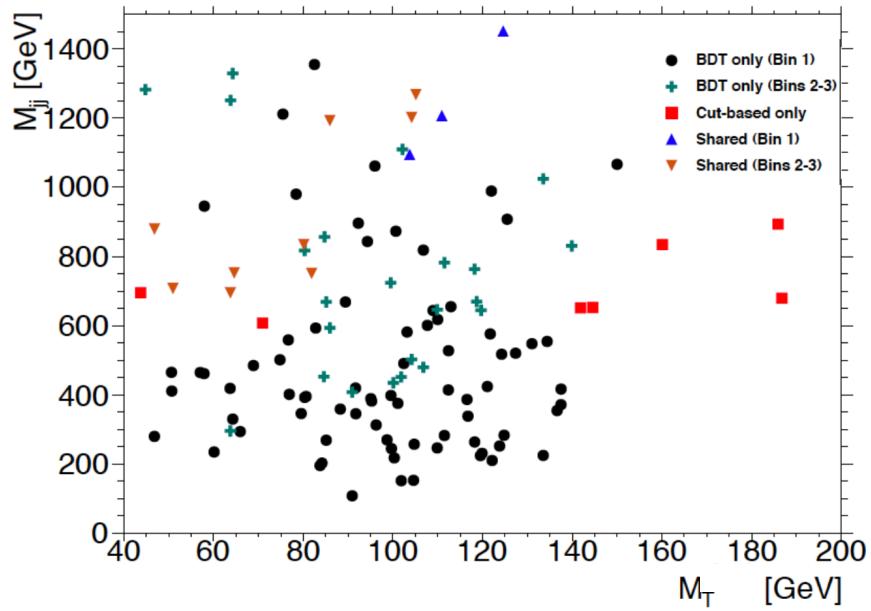


Figure 5.18: Overlap between cut-based and BDT VBF signal region candidates in the m_{jj} - m_T plane.

*The feeling is less like an ending than just another starting
point.*

Chuck Palahniuk

6

2121

2122

Combined Run I $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$

2123

results

2124

6.1 INTRODUCTION

2125

In the final statistical analysis of $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$, the dedicated gluon-gluon fusion and vector boson fusion sensitive signal regions are all combined into a single fit to determine the main parameters of interest, the Higgs signal strength μ and mass m_H . Therefore, while the specific requirements applied for the VBF sensitive analysis are discussed in chapter 5, the final measurement of these parameters can only be discussed in combination with the results of the ggF dedicated analysis. For example, because ggF Higgs production is considered a background in the VBF analysis, the ggF dedicated signal regions can actually constrain the normalization of this background in the VBF dedicated region.

2132

This chapter presents the combined interpretation of results in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis

2133 for gluon fusion and vector boson fusion Higgs production. First, the results of the dedicated gluon fu-
 2134 sion search are presented. Then, a comparison of the individual production mode signal strengths (μ_{ggF}
 2135 and μ_{VBF} and a measurement of the combined signal strength (μ) are shown. Subsequently, the mea-
 2136 sured values of the Higgs couplings to fermions and vector bosons is presented. Finally, the cross section
 2137 measurement for ggF and VBF production are shown.

2138 6.2 RESULTS OF DEDICATION GLUON FUSION $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ ANALYSIS

2139 The details of the dedicated gluon fusion $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis are not discussed in this thesis
 2140 and instead left to more comprehensive sources [63]. However, a brief summary of the results is essential
 2141 for describing the measurements of Higgs properties and interpreting the results of the dedicated VBF
 2142 search in a broader context. Additionally, the final Run 1 results on gluon fusion production make use of
 2143 the dedicated variables for same flavor final states developed in section 3.5. The results in the same flavor
 2144 final states will be shown here as well.

2145 6.2.1 RESULTS IN SAME FLAVOR ($ee/\mu\mu$) FINAL STATES

2146 Final states of the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel where both leptons have the same flavor ($ee/\mu\mu$)
 2147 were not included in the discovery result due to increased pileup conditions in the $\sqrt{s} = 8$ TeV data.
 2148 Dedicated techniques for background reduction in the same flavor final states were developed, as described
 2149 in section 3.5. The results shown in this section are the first published results using the same flavor channels
 2150 in the $H \rightarrow WW^*$ analysis.

2151 Table 6.1 shows the background estimate, expected signal yield, and event count in data for the same
 2152 flavor channels in the $n_j \leq 1$ signal regions. The dedicated same flavor background techniques allow this
 2153 channel to preserve a signal to background ratio similar to that of the different flavor channels.

	N_{obs}	N_{bkg}	N_{ggF}	N_{VBF}
$n_j = 0$	1108	1040 ± 40	77 ± 15	2.4 ± 1.7
$n_j = 1$	467	427 ± 21	22 ± 6	3.6 ± 1.8

Table 6.1: Post-fit yields in ggF dedicated signal regions for the $ee/\mu\mu$ final states [63].

2154 Figure 6.1 shows the final m_T distribution in data for the $n_j \leq 1$ channels. The data is very consistent
 2155 with the Higgs hypothesis and it can be seen that the same flavor channels are indeed sensitive to gluon
 2156 fusion production of the Higgs.

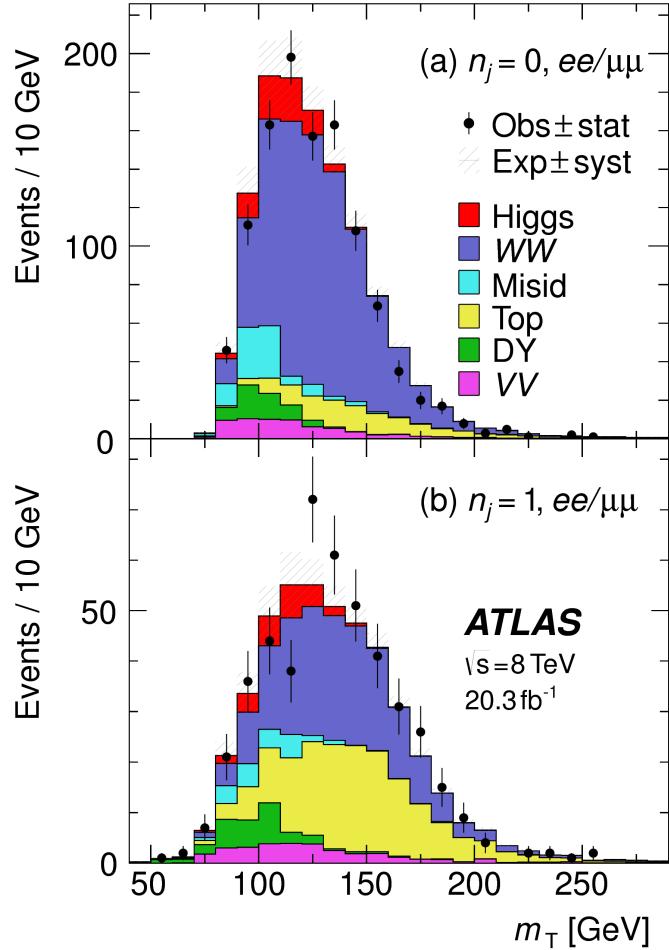


Figure 6.1: Post-fit m_T distribution in the $n_j \leq 1$ regions for the same flavor ($ee/\mu\mu$) final states [63].

2157 **6.2.2 COMBINED GLUON FUSION RESULTS**

2158 Table 6.2 shows the individual signal regions that were input into the final statistical fit. The ggF dedi-
 2159 cated bins use m_T as their discriminating variable and are separated into bins of p_T of the subleading
 2160 lepton as well. The VBF dedicated bin uses the O_{BDT} distribution as its final discriminant.

2161 Table 6.3 shows the yields in the various signal regions in both data and expected signal and back-
 2162 grounds. The yields for signal and background are all scaled according to the final normalizations cal-

SR category i			Fit var.
n_j , flavor	$\otimes m_{\ell\ell}$	$\otimes p_T^{\ell 2}$	
$n_j = 0$	$e\mu$	$\otimes [10, 30, 55]$	$\otimes [10, 15, 20, \infty]$
	$ee/\mu\mu$	$\otimes [12, 55]$	$\otimes [10, \infty]$
			m_T
$n_j = 1$	$e\mu$	$\otimes [10, 30, 55]$	$\otimes [10, 15, 20, \infty]$
	$ee/\mu\mu$	$\otimes [12, 55]$	$\otimes [10, \infty]$
			m_T
$n_j \geq 2$ ggF	$e\mu$	$\otimes [10, 55]$	$\otimes [10, \infty]$
$n_j \geq 2$ VBF	$e\mu$	$\otimes [10, 50]$	$\otimes [10, \infty]$
	$ee/\mu\mu$	$\otimes [12, 50]$	$\otimes [10, \infty]$
			O_{BDT}
			O_{BDT}

Table 6.2: All signal regions definitions input into final statistical fit [63].

culated in the fit.

	N_{obs}	N_{bkg}	N_{ggF}	N_{VBF}
$n_j = 0$	3750	3430 ± 90	300 ± 50	8 ± 4
$n_j = 1$	1596	1470 ± 40	102 ± 26	17 ± 5
$n_j \geq 2$, ggF $e\mu$	1017	960 ± 40	37 ± 11	13 ± 1.4
$n_j \geq 2$, VBF	130	99 ± 9	7.7 ± 2.6	21 ± 3

Table 6.3: Post-fit yields in the both ggF and VBF dedicated signal regions with all lepton flavor final states combined [63].

Figure 6.2 shows the final post-fit m_T distribution in the $n_j \leq 1$ regions. The data are very consistent with the hypothesis of ggF Higgs production. These yields are used as input, along with the VBF results in chapter 5, for the physical interpretation of results presented in subsequent sections.

6.3 SIGNAL STRENGTH MEASUREMENTS IN GGF AND VBF PRODUCTION

When all of the signal regions are combined in the fit, there can be a combined measurement of the signal strength as well as the individual ggF and VBF signal strengths. The combined signal strength is the ratio of the sum of the gluon fusion and VBF cross sections to the theory prediction, or a signal strength

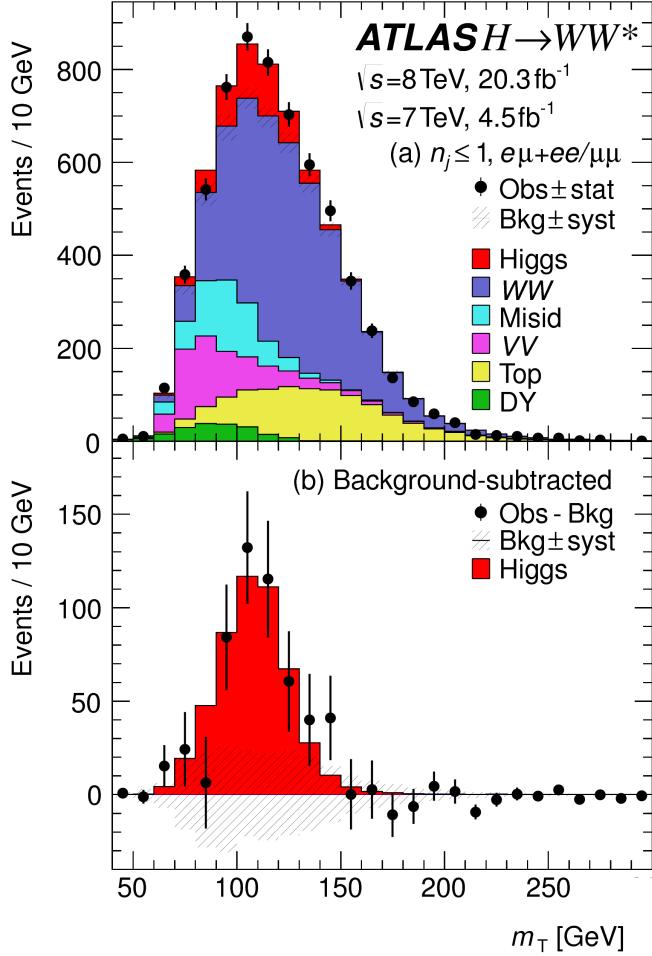


Figure 6.2: Post-fit m_T distribution in the $n_j \leq 1$ regions [63].

for the total Higgs production cross section that this analysis is sensitive to. The final measured combined signal strength μ is measured shown in equation 6.1.

$$\begin{aligned}
\mu &= 1.09 \quad {}^{+0.16}_{-0.15} \text{ (stat.)} \quad {}^{+0.08}_{-0.07} \left(\frac{\text{expt}}{\text{syst}} \right) \quad {}^{+0.15}_{-0.12} \left(\frac{\text{theo}}{\text{syst}} \right) \quad \pm 0.03 \left(\frac{\text{lumi}}{\text{syst}} \right) \\
&= 1.09 \quad {}^{+0.16}_{-0.15} \text{ (stat)} \quad {}^{+0.17}_{-0.14} \text{ (syst)} \\
&= 1.09 \quad {}^{+0.23}_{-0.21}.
\end{aligned} \tag{6.1}$$

Figure 6.3 gives the best fit signal strength $\hat{\mu}$ as a function of the hypothesized Higgs mass. The value at 125.36 GeV corresponds to the μ quoted in equation 6.1. This value of the Higgs mass is used because it is the most precise mass measurement from ATLAS, a result of the combined $\gamma\gamma$ and ZZ mass measure-

2176 ments [94].

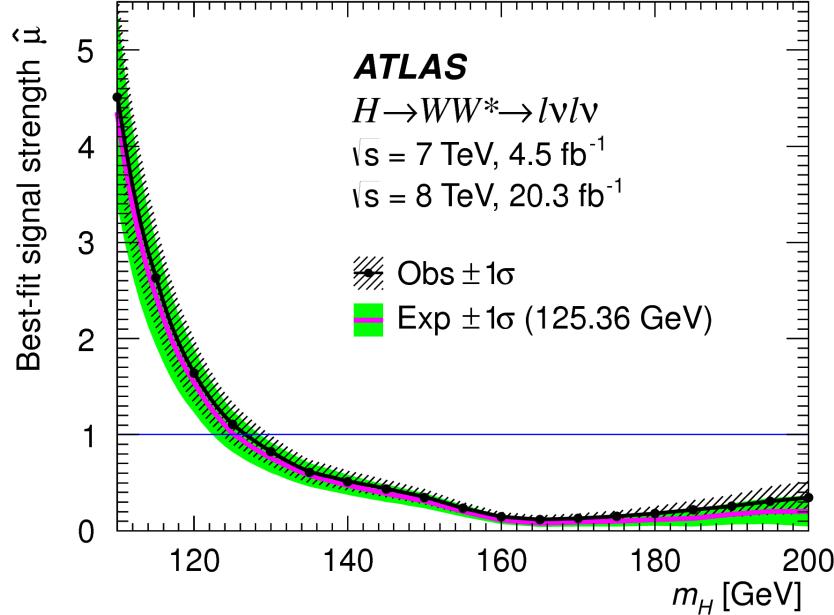


Figure 6.3: Best fit signal strength $\hat{\mu}$ as a function of hypothesized m_H [63].

2177 As explained in chapter 3, a probability p_0 can be computed using the test statistic q_0 to quantify the
 2178 probability that the background could fluctuate to produce an excess at least as large as the one observed
 2179 in the data. The local p_0 value is shown in figure 6.4 as a function of m_H . The minimum p_0 value is at
 2180 $m_H = 130$ GeV and corresponds to a significance of 6.1σ . The curve is relatively flat and the significance
 2181 is the same at 125.36 GeV within the quoted precision. The expected significance for a signal with strength
 2182 $\mu = 1.0$ is 5.8σ . This represents the first discovery level significance measurement in the $H \rightarrow WW^* \rightarrow$
 2183 $\ell\nu\ell\nu$ analysis.

2184 All the results presented so far in this section have been for the combined gluon fusion and VBF pro-
 2185 duction modes. However, each signal strength can be calculated separately in the likelihood as well. There
 2186 are two ways to do this. First, the likelihood can be parameterized in terms of a single parameter, the ratio
 2187 of the VBF and gluon fusion signal strengths. With this method, the statistical significance of the VBF
 2188 Higgs result can be evaluated. Figure 6.5 shows the likelihood as a function of the ratio $\mu_{\text{VBF}}/\mu_{\text{ggF}}$.

2189 The best fit value of the ratio of signal strengths is shown in equation 6.2. Within the quoted uncer-

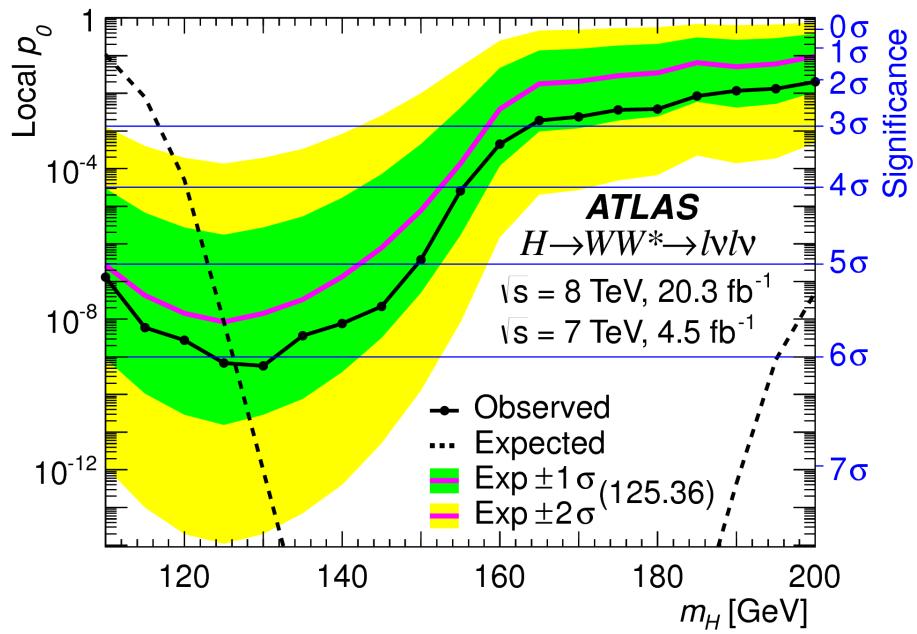


Figure 6.4: Local p_0 as a function of m_H [63].

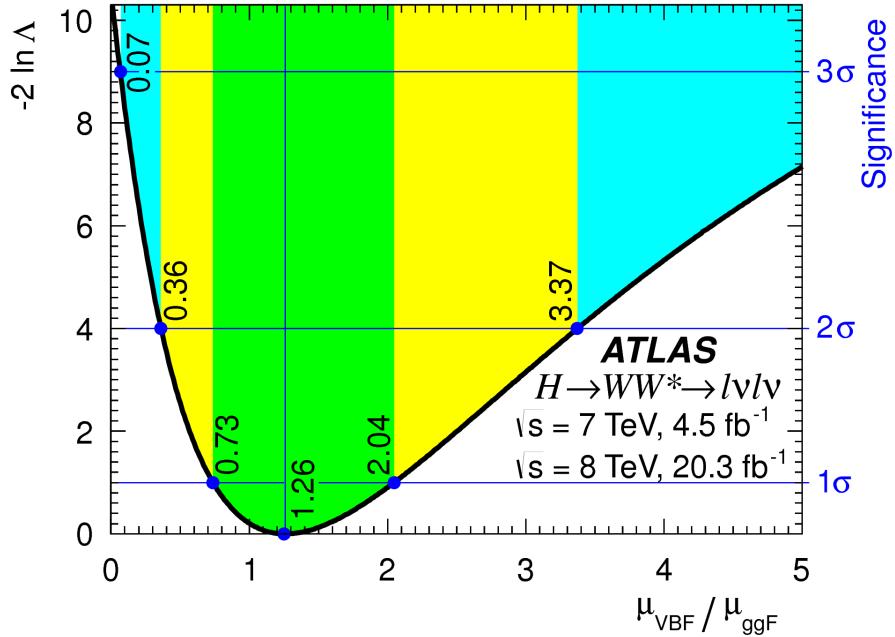


Figure 6.5: Likelihood as a function of $\mu_{\text{VBF}} / \mu_{\text{ggF}}$ [63].

tainties, it is consistent with a ratio of unity.

$$\frac{\mu_{\text{VBF}}}{\mu_{\text{ggF}}} = 1.26^{+0.61} (\text{stat.})^{+0.50} (\text{syst.}) = 1.26^{+0.79}_{-0.53} \quad (6.2)$$

2191 The null hypothesis for VBF production corresponds to a ratio of $\mu_{\text{VBF}}/\mu_{\text{ggF}} = 0$. The likelihood in
2192 figure 6.5 gives a significance of 3.2σ at $\mu_{\text{VBF}}/\mu_{\text{ggF}} = 0$, as quoted in chapter 5.

2193 In addition to the ratio of signal strengths, each signal strength can be varied independently in the like-
2194 lihood as well. Figure 6.6 shows the two dimensional likelihood scan in the $\mu_{\text{ggF}}-\mu_{\text{VBF}}$ plane. The best fit
2195 values of the two signal strengths are shown in equation 6.3. Both are consistent with unity within their
2196 uncertainties.

$$\begin{aligned} \mu_{\text{ggF}} &= 1.02 \pm 0.19^{+0.22}_{-0.18} = 1.02^{+0.29}_{-0.26} \\ \mu_{\text{VBF}} &= 1.27^{+0.44}_{-0.40}^{+0.29}_{-0.21} = 1.27^{+0.53}_{-0.45} \end{aligned} \quad (6.3)$$

(stat.) (syst.)

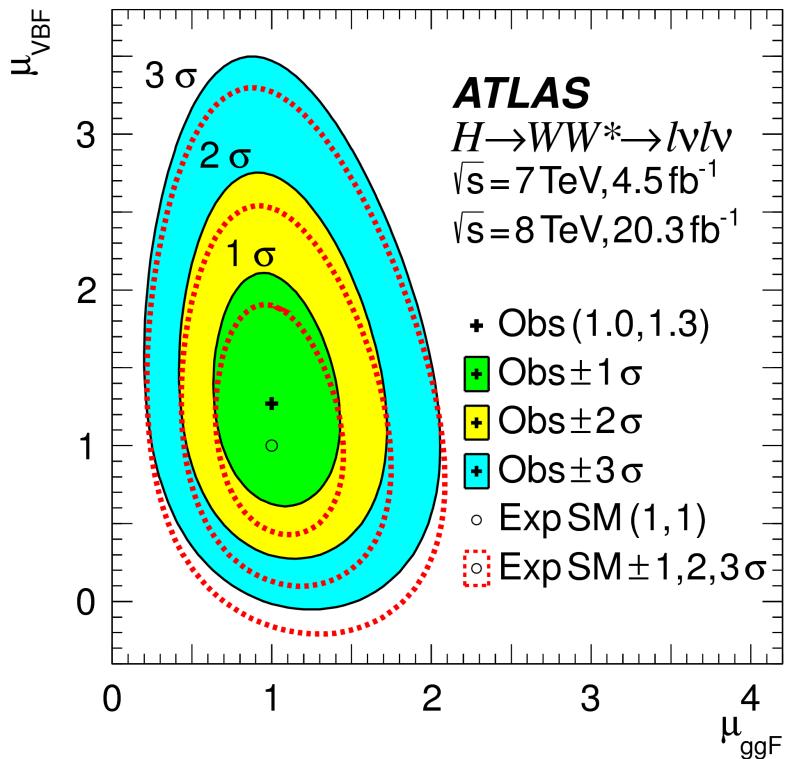


Figure 6.6: Likelihood scan as a function of μ_{VBF} and μ_{ggF} [63].

2197 **6.4 MEASUREMENT OF HIGGS COUPLINGS TO VECTOR BOSONS AND FERMIONS**

2198 Similar to the parameterization of signal strength, the couplings of the Higgs to fermions and bosons
 2199 can also be parameterized. The parameter of interest in this case is κ , or the ratio of the measured coupling
 2200 to the standard model expectation. Both the fermion and boson couplings have these so-called scale factors,
 2201 κ_F for fermions and κ_V for bosons. Gluon fusion production is sensitive to the fermion couplings through
 2202 the top quark loops in its production, while VBF production is sensitive to the vector boson couplings in
 2203 its production. Both modes are sensitive to the vector boson couplings in their decays. The signal strengths
 2204 will have dependence on the coupling scale factors as described in equation 6.4 [18].

$$\begin{aligned}\mu_{\text{ggF}} &\propto \frac{\kappa_F^2 \cdot \kappa_V^2}{(\mathcal{B}_{H \rightarrow f\bar{f}} + \mathcal{B}_{H \rightarrow gg}) \kappa_F^2 + (\mathcal{B}_{H \rightarrow VV}) \kappa_V^2} \\ \mu_{\text{VBF}} &\propto \frac{\kappa_V^4}{(\mathcal{B}_{H \rightarrow f\bar{f}} + \mathcal{B}_{H \rightarrow gg}) \kappa_F^2 + (\mathcal{B}_{H \rightarrow VV}) \kappa_V^2}.\end{aligned}\quad (6.4)$$

2205 Figure 6.7 shows the two-dimensional likelihood scan of κ_F and κ_V . The best-fit values are given in equa-
 2206 tion 6.5. The best-fit values are consistent with unity within their uncertainties.

$$\begin{aligned}\kappa_F &= 0.93 & {}^{+0.24}_{-0.18} & {}^{+0.21}_{-0.14} &= 0.93 & {}^{+0.32}_{-0.23} \\ \kappa_V &= 1.04 & {}^{+0.07}_{-0.08} & {}^{+0.07}_{-0.08} &= 1.04 & \pm 0.11.\end{aligned}\quad (6.5)$$

(stat.) (syst.)

2207

2208 **6.5 HIGGS PRODUCTION CROSS SECTION MEASUREMENT**

2209 Another measurement that comes naturally from the signal strength numbers quoted earlier is the pro-
 2210 duction cross section at 7 and 8 TeV for both gluon fusion and VBF production. The general equation

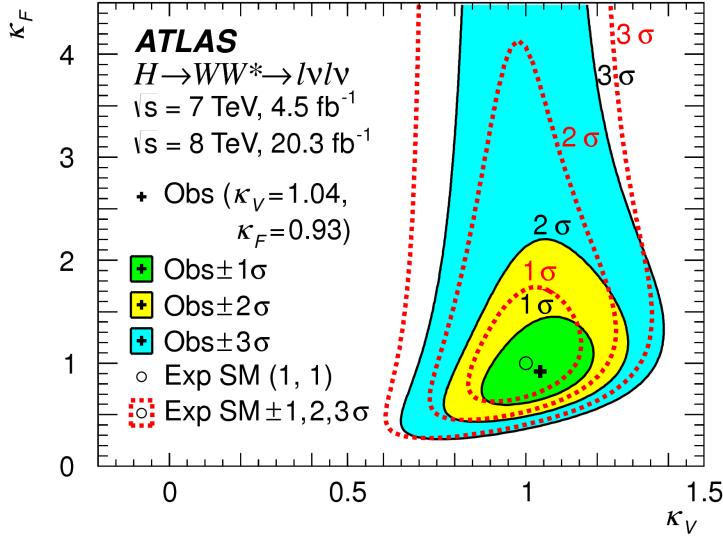


Figure 6.7: Likelihood scan as a function of κ_F and κ_V [63].

for calculating the cross section is given in equation 6.6.

$$\begin{aligned}
 (\sigma \cdot \mathcal{B}_{H \rightarrow WW^*})_{\text{obs}} &= \frac{(N_{\text{sig}})_{\text{obs}}}{\mathcal{A} \cdot \mathcal{C} \cdot \mathcal{B}_{WW \rightarrow l\nu l\nu}} \cdot \frac{1}{\int L dt} \\
 &= \hat{\mu} \cdot (\sigma \cdot \mathcal{B}_{H \rightarrow WW^*})_{\text{exp}}
 \end{aligned} \tag{6.6}$$

$(N_{\text{sig}})_{\text{obs}}$ is the number of events observed in data. \mathcal{A} is the geometric and kinematic acceptance of the detector, while \mathcal{C} is the efficiency of the signal region selection for events that are reconstructed in the detector. The branching ratio of a WW system to leptons must also be divided out. The production cross section depends on the center of mass energy and the production mode desired (gluon fusion or VBF), and so three separate cross section measurements are quoted in equation 6.7.

$$\begin{aligned}
 \sigma_{\text{ggF}}^{7\text{TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 2.0 \pm 1.7 \quad {}^{+1.2}_{-1.1} = 2.0 \quad {}^{+2.1}_{-2.0} \text{ pb} \\
 \sigma_{\text{ggF}}^{8\text{TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 4.6 \pm 0.9 \quad {}^{+0.8}_{-0.7} = 4.6 \quad {}^{+1.2}_{-1.1} \text{ pb} \\
 \sigma_{\text{VBF}}^{8\text{TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 0.51 \quad {}^{+0.17}_{-0.15} \quad {}^{+0.13}_{-0.08} = 0.51 \quad {}^{+0.22}_{-0.17} \text{ pb.}
 \end{aligned} \tag{6.7}$$

(stat.) (syst.)

2217 The predicted cross section values for gluon fusion are 3.3 ± 0.4 pb at 7 TeV and 4.2 ± 0.5 pb at 8 TeV,
2218 consistent with the measured values within their uncertainties. For vector boson fusion, the predicted
2219 cross section is 0.35 ± 0.02 pb, again consistent with the measured value.

2220 **6.6 CONCLUSION**

2221 The combined analysis of the gluon fusion and vector boson fusion processes in $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
2222 in the 7 and 8 TeV datasets has yielded the first discovery level significance for Higgs production in this
2223 decay channel. Additionally, precise measurements of the couplings to vector bosons and fermions are
2224 given. Finally, signal strengths and cross sections for each production mode are measured. Figure 6.8 shows
2225 the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ measurements in comparison with other Higgs decay channels in ATLAS. The
2226 measurement of signal strength from this channel remains the most sensitive in both the gluon fusion and
2227 VBF production modes for the Run 1 dataset.

ATLAS

Individual analysis

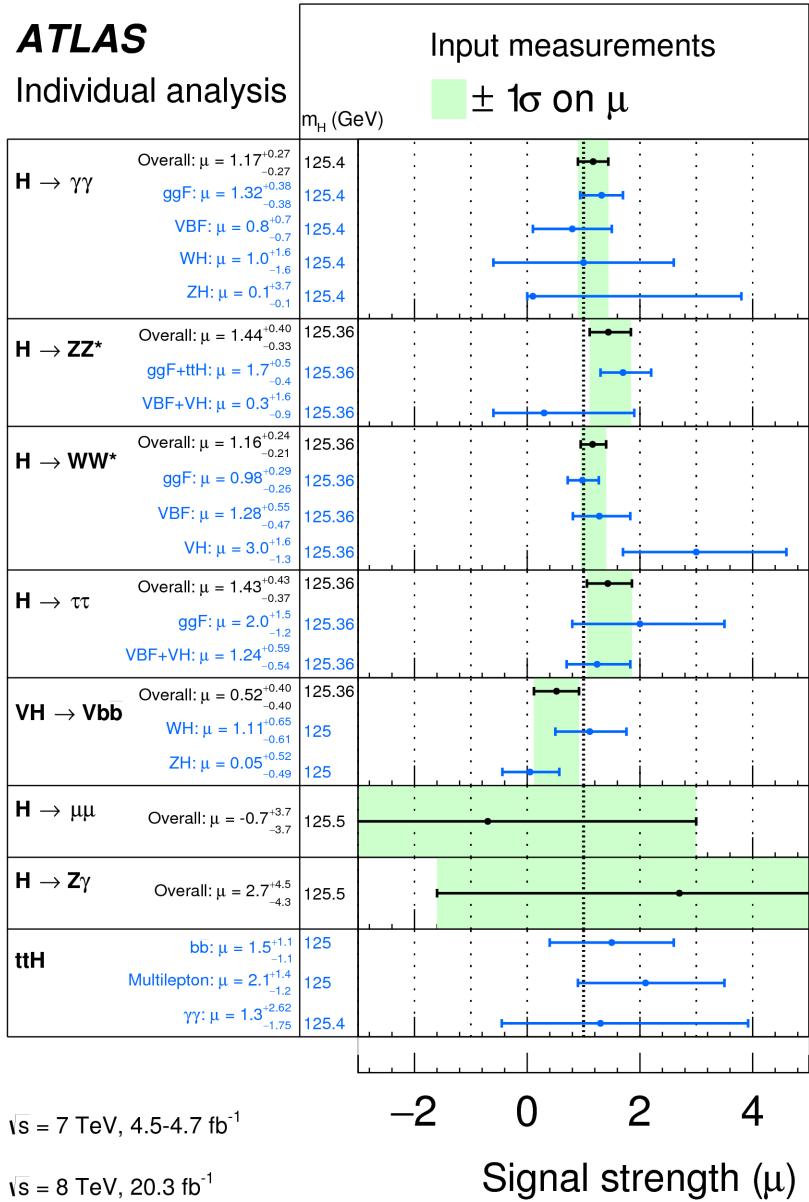


Figure 6.8: Comparison of signal strength measurements in different Higgs decay channels on ATLAS [95].

2228

Part III

2229

Search for Higgs pair production in the

2230

$HH \rightarrow b\bar{b}b\bar{b}$ channel in LHC Run 2 at $\sqrt{s} =$

2231

13 TeV

Passion is in all great searches and is necessary to all creative endeavors.

W. Eugene Smith

7

2232

2233 Search for Higgs pair production in boosted 2234 $b\bar{b}b\bar{b}$ final states

2235 7.1 INTRODUCTION

2236 After the discovery of the Higgs boson in the ATLAS Run 1 dataset and the subsequent measurements
2237 of its properties, the Higgs transformed into a potential tool in searches for physics beyond the Standard
2238 Model. The pair production cross section of the Higgs can be enhanced through BSM physics. Studying
2239 di-Higgs production also probes the Higgs self-coupling, shedding light on the structure of the Higgs po-
2240 tential. This chapter presents a search for resonant production of a Higgs pair in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$
2241 final state in 3.2 fb^{-1} of data collected at $\sqrt{s} = 13 \text{ TeV}$. In particular, this chapter focuses on a search for
2242 this final state in the regime where m_X is large ($\gtrsim 1 \text{ TeV}$) and the Higgs bosons in the decay are signifi-
2243 cantly boosted. A tailored selection for this boosted selection, using novel techniques in jet substructure

and b -tagging, is discussed. Then, the data-driven background estimate is presented. Finally, the results of the search are shown. The signal models used as benchmarks are a spin-2 Randall Sundrum graviton (RSG) and a narrow width spin-0 resonance. These models are described in more detail in Chapter 1. Limits on signal models are reserved for the next chapter where the results of this chapter are combined with the results of a separate selection dedicated to the lower m_X regime.

7.2 MOTIVATION

With the center of mass energy increase from $\sqrt{s} = 8$ TeV to $\sqrt{s} = 13$ TeV, the LHC and ATLAS are able to probe new resonances at higher mass scales than previously accessible in Run 1. This is a powerful motivator for searching for a new resonance in the early 13 TeV data. Figure 7.1 shows the ratios of parton luminosities between 8 and 13 TeV for different resonance masses. For a resonance of $M_X = 2$ TeV, the cross section at $\sqrt{s} = 13$ TeV is roughly a factor of 10 larger than at $\sqrt{s} = 8$ TeV.

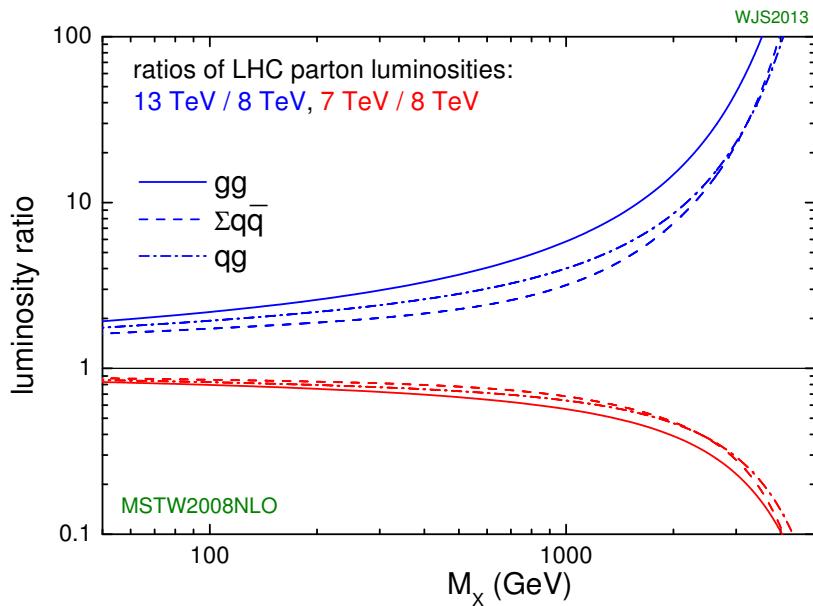


Figure 7.1: Parton luminosity ratios as a function of resonance mass M_X for 13/8 TeV and 7/8 TeV [96].

Higgs pair production offers a vast array of unprobed regions of phase space where searches for BSM physics can be made. Chapter 1 discusses some possibilities for both resonant and non-resonant enhance-

ment of the di-Higgs production cross section. Given the increased mass reach of the LHC in Run 2, it is particularly important to focus on resonant searches at high m_X . One consideration when conducting a search in the HH final state is which decay modes of the Higgs to consider. Figure 7.2 shows the branching ratio of the HH final state for different combinations of decays of each individual Higgs. As the largest branching ratio for the 125 GeV Higgs is $H \rightarrow b\bar{b}$, the $HH \rightarrow b\bar{b}b\bar{b}$ branching ratio is also the largest at 33%.

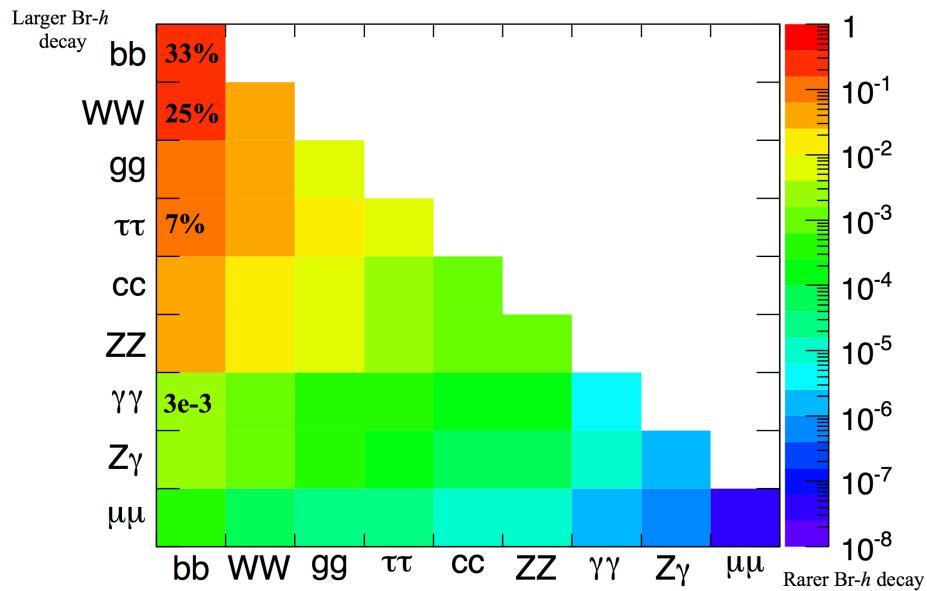


Figure 7.2: Summary of HH branching ratios [97].

At high m_X , the Higgs bosons resulting from the decay of a heavy resonance will have large p_T ¹. The ΔR between the decay products of the Higgs is inversely proportional to the Higgs p_T , as shown in equation 7.1.

$$\Delta R \approx \frac{2m}{p_T} \quad (7.1)$$

Figure 7.3 shows the minimum ΔR between truth level B decay vertices in simulation samples for Randall-Sundrum gravitons of different masses. The figure shows that as the mass of the graviton increases, the ΔR distribution between the b quarks in the Higgs decay tends to shift to lower values. Because of this effect, it is necessary to tailor a selection to target these merged b -jets.

¹In the limit that $m_H \ll m_X$, the Higgs p_T is roughly $m_X/2$.

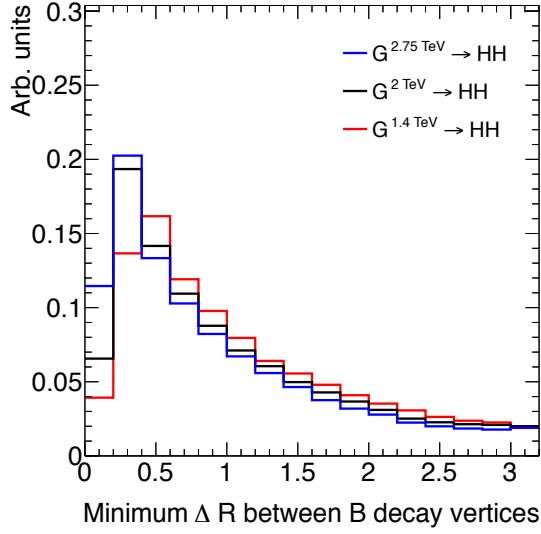


Figure 7.3: Minimum ΔR between B decay vertices for different RSG masses in a $G_{KK}^* \rightarrow HH \rightarrow 4b$ sample with $c = 1$

2270 7.3 DATA AND SIMULATION SAMPLES

2271 7.3.1 SIGNAL MODELS

2272 While the resonance search is by its nature generic (as it is a simple search for a peak in the $4b$ invariant
 2273 mass spectrum), there are two signal models that the selection requirements have been optimized for.
 2274 The first is Randall-Sundrum (RSG) model, where a tower of massive spin-2 Kaluza-Klein gravitons is
 2275 predicted. The second is a heavy narrow spin-0 resonance, the so-called “heavy Higgs”. This type of res-
 2276 onance arises, for example, in the two Higgs doublet model (2HDM). More details about the physics of
 2277 these models and their motivation is given in chapter 1.

2278 Signal graviton (G_{KK}^*) events are generated at leading order (LO) with **MADGRAPH5 v2.2.2** [98]. The
 2279 PDF set used is the **NNPDF2.3 LO** set [99]. For modeling parton shower and hadronization in jets, **PYTHIA**
 2280 8.186 is used with the A14 tune [76, 100]. The free parameters in the RSG model are the graviton mass
 2281 and the coupling constant $c \equiv k/\bar{M}_{\text{Pl}}^2$. Both the production cross section and width of the graviton are
 2282 proportional to c^2 . Samples are generated at both $c = 1$ and $c = 2$ for a variety of mass points between

² k is the curvature constant for the warped extra dimension and \bar{M}_{Pl} is the Planck mass divided by 8π

2283 300 GeV and 3 TeV.

2284 The second signal sample is a heavy spin-0 resonance H with a fixed width of $\Gamma_H = 1$ GeV. This is
2285 generated with **MADGRAPH5** and uses the **CT10** PDF set [79]. The parton shower and hadronization are
2286 handled by **HERWIG ++** with the **CTEQ6L1** PDF set and the **UEEE5** event tune [80, 101, 102]. Because
2287 the width and branching ratios depend on 2HDM parameters, each mass point generated with this fixed
2288 width corresponds to a different point in the 2HDM parameter phase space. Mass points are generated
2289 between 300 GeV and 1 TeV as with the RSG signal samples.

2290 **7.3.2 BACKGROUND SAMPLES**

2291 While the dominant **QCD** multijet background is estimated with a fully data-driven method, the sub-
2292 dominant backgrounds $t\bar{t}$ and Z +jets are modeled with some input from simulation.

2293 $t\bar{t}$ events are simulated at next-to-leading order (NLO) with the **POWHEG-BOX** version 1 generator us-
2294 ing the **CT10** PDF set [103]. The parton shower, hadronization, and underlying event are simulated with
2295 **PYTHIA 6.428** with the **CTEQ6L1** PDF set [75]. The Perugia 2012 tune is used [104]. NNLO QCD cor-
2296 rections to the cross sections are computed in **Top++ 2.0** [105]. The top quark mass is set to 172.5 GeV.
2297 The shapes of distributions in $t\bar{t}$ are taken from MC while the normalization is taken from data.

2298 Finally, the Z +jets background is simulated with **PYTHIA 8.186** and the **NNPDF2.3** LO PDF set. This
2299 background is negligible compared to the others and is taken fully from MC.

2300 **7.3.3 DATA SAMPLE AND TRIGGER**

2301 This analysis is done on 3.2 fb^{-1} of data taken in 2015 at $\sqrt{s} = 13$ TeV. The details of the machine
2302 conditions during this time can be found in Chapter 2. Only data which was taken during stable beam con-
2303 ditions with all detectors functioning is used. Events must pass a trigger which requires a single 360 GeV
2304 large radius ($R = 1.0$) jet to be reconstructed in the HLT. Figure 7.4 shows the trigger efficiency for vari-
2305 ous trigger options as a function of graviton mass. Above $m_G > 1$ TeV, the single large radius jet trigger
2306 is 99% efficient for events passing the signal selection.

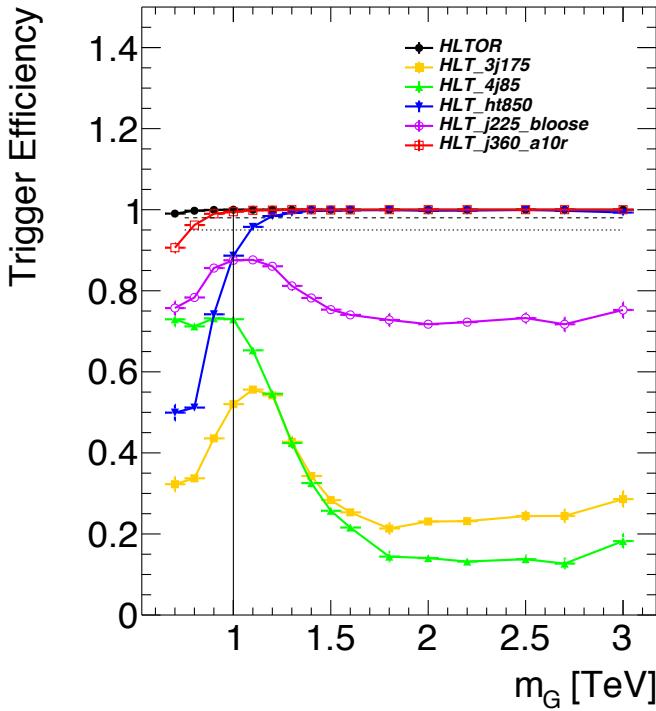


Figure 7.4: Trigger efficiency for events passing all signal region selections as a function of mass in $G_{\text{KK}}^* \rightarrow HH \rightarrow 4b$ samples with $c = 1$ [106]. In the trigger names, “j” refers to a jet or jets. “ht” refers to H_T , the scalar sum of transverse momenta in the event. “bloose” refers to a loose b -tagging requirement applied to the jet. “a10r” refers to anti- k_T jets with $R = 1.0$. The numbers at the end are the thresholds on the given quantity in GeV.

2307 7.4 EVENT RECONSTRUCTION AND OBJECT SELECTION

2308 The boosted selection first begins by defining a unique set of objects that can be exploited to increase
 2309 signal efficiency in the kinematic regime where the final state b -jets are very merged.

2310 7.4.1 LARGE RADIUS ($R = 1.0$) JETS

2311 The first step towards reconstructing the final state is to define objects that can be used to measure the
 2312 kinematics of the Higgs bosons. In the boosted selection anti- k_T jets with a radius parameter of 1.0 are
 2313 used. These jets are much larger in angular size than the typical $R = 0.4$ jets and are intended to encompass
 2314 both jets resulting from the Higgs decay³. The jets are built from clusters in the calorimeter calibrated with

³This is in contrast to the resolved selection, which uses two $R = 0.4$ anti- k_T jets for each Higgs

local calibration weighting [58].

Because of the large extent of these jets, great care must be taken to remove potential contributions of calorimeter clusters from pile-up. This is done using a technique called jet trimming [107]. With trimming, the constituents of the large radius jet are re-clustered with a smaller radius with the k_T algorithm. Then, these so-called subjets are removed from the larger jet if $p_T^{\text{subjett}} / p_T^{\text{jet}} < f_{\text{cut}}$. In this analysis, the subjet radius is $R = 0.2$ and $f_{\text{cut}} = 0.05$. Trimming has been shown to improve the mass resolution of large radius jets. Figure 7.5 shows the effect of trimming on the large radius jet mass (M_J). Because the large radius jet fully contains the higgs decay products, its invariant mass should correspond to the 125 GeV mass of the Higgs. The trimming algorithm brings the jet mass much closer to the expected Higgs mass and improves the mass resolution.

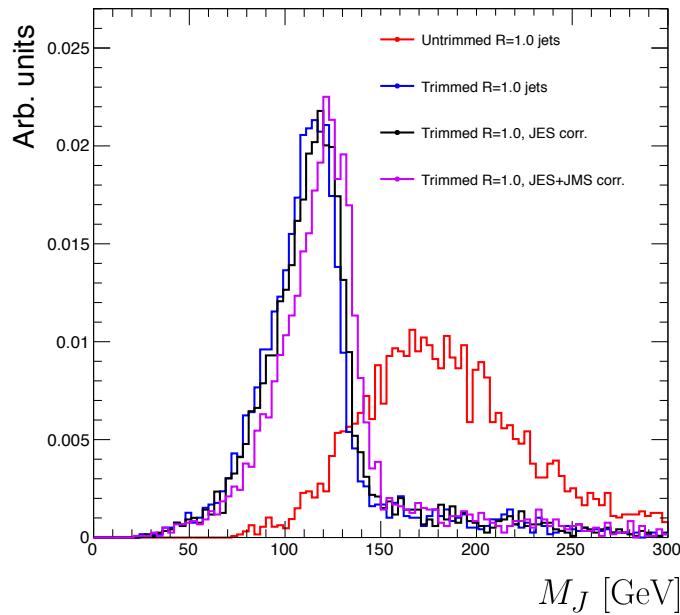


Figure 7.5: Comparison of untrimmed and trimmed jet masses for large radius jets in a RSG sample with $m_{G_{KK}^*} = 1$ TeV. JES (JMS) refers to the standard jet energy (mass) scale calibration for ATLAS [58].

The large radius jets are required to satisfy $250 < p_T < 1500$ GeV. They must also be within $|\eta| < 2.0$ in order to ensure that the full jet is within the inner detector tracking volume. Finally, they are required to have $M_J > 50$ GeV. The upper p_T cut and lower threshold on mass are applied to correspond to the kinematic range where uncertainties are available in ATLAS calibrations [108, 109].

2329 7.4.2 TRACK JETS AND b -TAGGING

2330 Because the b -jets from boosted Higgs decays are so close together (as illustrated in figure 7.3), narrow ra-
 2331 dius jets are required to fully resolve both b -jets. The minimum radius feasible for jets based on calorimeter
 2332 deposits is determined by the calorimeter granularity. However, because b -tagging relies on information
 2333 from the inner detector, it is possible to define another type of jet that can have a smaller radius and better
 2334 b -tagging resolution. These jets are called “track jets” [109, 110].

2335 Track jets are formed by applying the usual anti- k_T clustering algorithm to tracks that are required to be
 2336 consistent with the primary vertex. After the jet axis has been determined using these tracks, a second step
 2337 of track association is also performed to add tracks that can be useful for b -tagging [110]. In this analysis,
 2338 the tracks are clustered with a radius parameter of $R = 0.2$. This radius has been shown to give good
 2339 performance in boosted Higgs tagging [109, 110]. Figure 7.6 shows a comparison among different track jet
 2340 radii of the efficiency for reconstructing two b -jets from each Higgs in a RSG sample as a function of mass.
 Track jets with radius of 0.2 give the best performance, especially at high mass. In this analysis, track jets

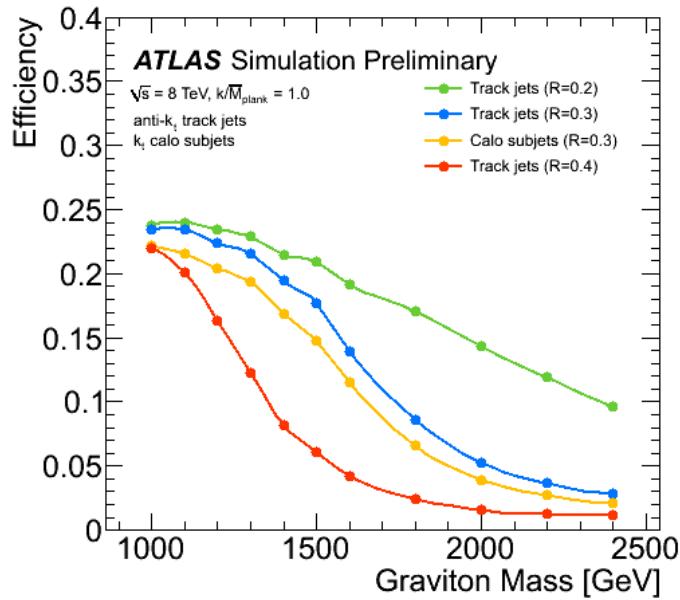


Figure 7.6: Efficiency of finding two b -jets from each Higgs in an RSG event using calorimeter jets with $R = 0.3$ or
 2341 different track jet radii [110]

2342 are required to have $p_T > 10 \text{ GeV}$ and $|\eta| < 2.5$. They must also have at least two tracks.

2343 **7.4.3 MUONS**

2344 Muons are used in this study to correct the four-momenta of calorimeter jets by accounting for semi-
2345 leptonic b decays. The muons used are combined ID and MS muons which must satisfy tight identification
2346 requirements [54]. The muons must have $p_T > 4 \text{ GeV}$ and $|\eta| < 2.5$. Table 7.1 summarizes the object
2347 requirements described in this section.

	R	p_T	$ \eta $	M
Calorimeter jets	1.0	$250 < p_T < 1500 \text{ GeV}$	< 2.0	$> 50 \text{ GeV}$
Track jets	0.2	$> 10 \text{ GeV}$	< 2.5	-
Muons	-	4 GeV	< 2.5	-

Table 7.1: Summary of requirements on objects used in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search

2348 **7.5 EVENT SELECTION**

2349 The first requirement in the boosted selection is for ≥ 2 large radius jets satisfying the selections out-
2350 lined above. The two highest momentum large-R jets in the event are referred to as “Higgs candidates”.
2351 The leading jet is required to have $p_T > 350 \text{ GeV}$.

2352 Track jets satisfying the object selections are matched to Higgs candidate jets via ghost association [33].
2353 Each Higgs candidate must have at least 2 track jets associated with it. These basic requirements are illus-
2354 trated in figure 7.7.

2355 The QCD multijet background produces less central jets than high mass resonances, so there is an ad-
2356 dditional requirement that the two Higgs candidates be close together in η . The large-R jets are required to
2357 satisfy $|\Delta\eta(JJ)| < 1.7$.

2358 **7.5.1 MASS REQUIREMENTS**

2359 The final set of requirements ensures that the Higgs candidates are consistent with expected properties
2360 of the 125.0 GeV Higgs. First, a variable (X_{hh}) is defined to measure the consistency of both of the Higgs

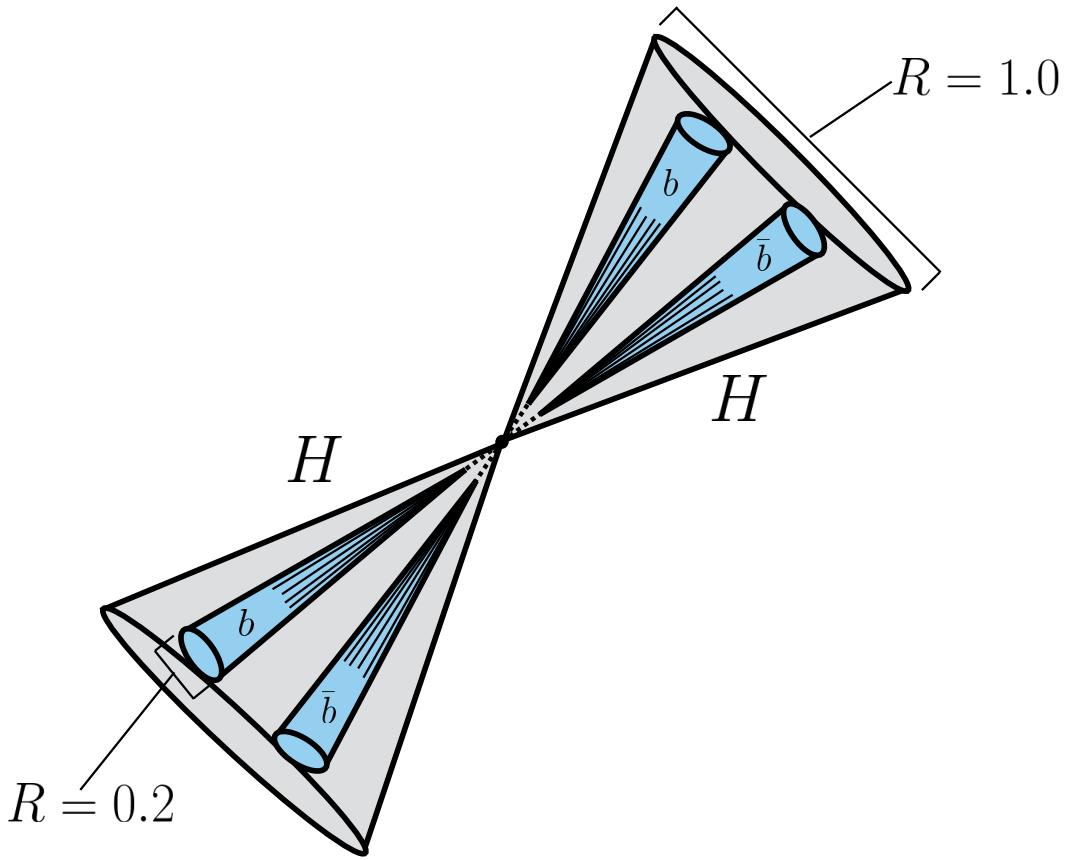


Figure 7.7: Illustration of the boosted selection requirements on Higgs candidates. Each large-radius calorimeter jet (Higgs candidate) must contain two track jets

²³⁶¹ candidate jets with the SM Higgs mass. This is shown in equation 7.2.

$$X_{hh} = \sqrt{\left(\frac{M_J^{\text{lead}} - 124 \text{ GeV}}{0.1 M_J^{\text{lead}}}\right)^2 + \left(\frac{M_J^{\text{sublead}} - 115 \text{ GeV}}{0.1 M_J^{\text{sublead}}}\right)^2} \quad (7.2)$$

²³⁶² The mass values in the X_{hh} formula are optimized to maximize signal efficiency. The sub-leading jet
²³⁶³ typically has a lower mass due to semi-leptonic b decays and final state radiation. X_{hh} effectively acts as
²³⁶⁴ a χ^2 measurement of the consistency of the two Higgs candidate masses with the signal hypothesis. The
²³⁶⁵ denominators of each term ($0.1M$) give the uncertainty on the mass measurement for the large radius jets.
²³⁶⁶ Events are required to satisfy $X_{hh} < 1.6$.

²³⁶⁷ Before making the requirement on X_{hh} , the masses of the Higgs candidates are corrected for semi-

leptonic b decays using muons with the criteria outlined in the previous section. Any muons within a $\Delta R < 0.2$ of a b -tagged track jet (as described in the next section) have their four-momenta added to the four-momentum of the Higgs candidate. This correction does not affect the pre-selection requirements but does affect the X_{hh} requirement and the final invariant mass distribution used.

7.5.2 b -TAGGING REQUIREMENTS

The last requirement applied is on the number of b -tagged track jets. There are two signal regions defined. The first requires exactly four b -tagged track jets, two in each Higgs candidate (known as the $4b$ signal region). At high resonance masses, this requirement is inefficient, so an additional signal region requiring only three b -tagged track jets is also defined (known as the $3b$ signal region). While this has a larger background it is also more efficient for high resonance masses. For both signal regions, the MV2c20 algorithm, where the training sample for the algorithm has 20% charm events is used. More details for this algorithm can be found in Chapter 2.

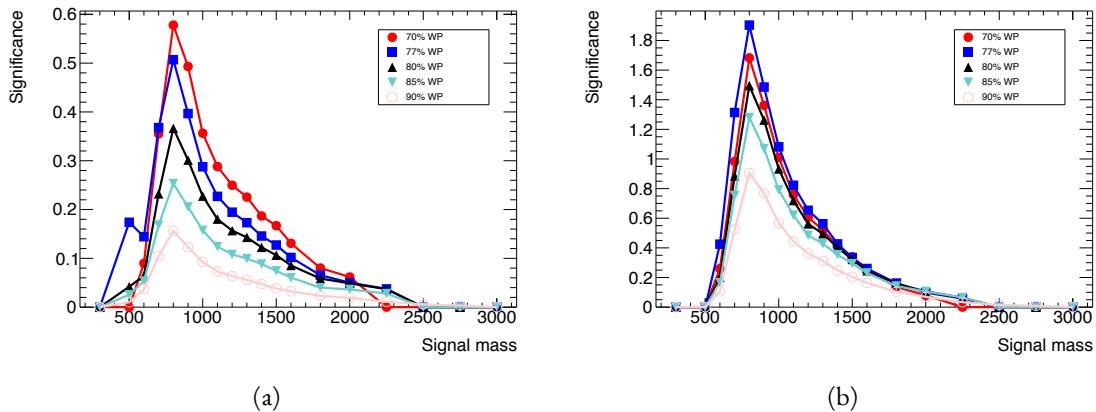


Figure 7.8: Estimated significance as a function of signal mass for RSG $c = 1$ models in the $3b$ (a) and $4b$ (b) regions for different b -tagging efficiency working points

Once the algorithm is selected, an efficiency working point must also be chosen. This working point defines the efficiency with which true b -jets are tagged and also fixes the overall background rejection of the algorithm. Higher efficiency working points accept more true b -jets but also allow for more background. Five different working points (70%, 77%, 80%, 85%, 90%) are tested. With each working point, the full data driven background estimation method is run to quantify the amount of background that will be

present in the final signal region. The significance is quantified using the median discovery significance for signal and background with Poisson errors, given in equation 7.3 [112].

$$Z = \sqrt{2 \left((s + b) \ln \left(1 + \frac{s}{b} \right) - s \right)} \quad (7.3)$$

Here, s is the expected number of signal events and b is the expected number of background events. This formula is derived using Poisson statistics with errors on both the signal and background. It is used because it is valid in the regime where s and b are of the same order. Note that in the limit where s is much smaller than b , this equation reduces to the more well known s/\sqrt{b} . Figure 7.8 shows the estimated significance as a function of signal mass in RSG $c = 1$ models for the $3b$ and $4b$ signal regions. The 77% working point gives the best performance over a wide range of masses in the $4b$ signal region. As this is the region which contributes the most to the total discovery significance, the 77% efficiency working point is chosen for the analysis.

7.5.3 SELECTION EFFICIENCY

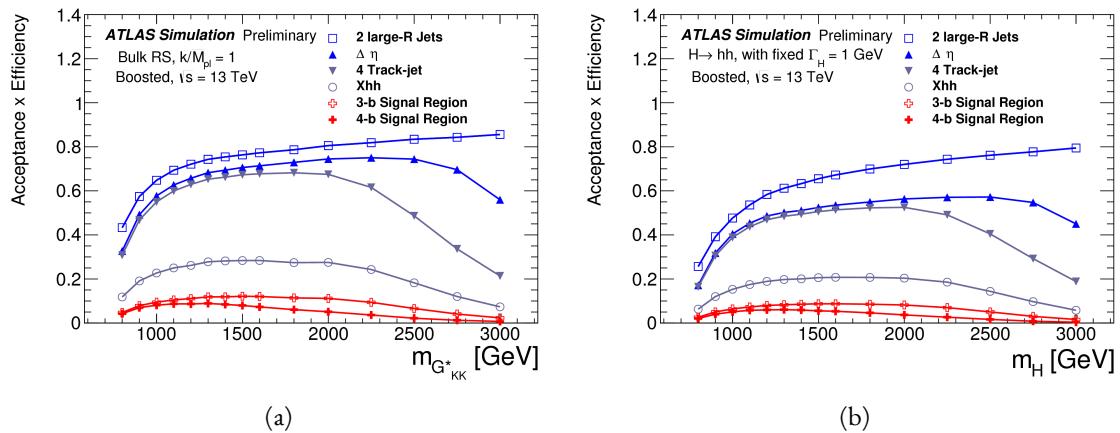


Figure 7.9: Acceptance \times efficiency as a function of mass for (a) RSG and (b) narrow heavy scalar signal models [113].

Figure 7.9 shows the product of acceptance and efficiency as a function of mass for both the RSG and narrow heavy scalar resonance signal models. After $m_X > 1$ TeV, the efficiency of the $4b$ requirement begins to decline. After $m_X > 2$ TeV, the efficiency of requiring two track jets in each Higgs candidate

begins to decline as well. Both of these behaviors illustrate the difficulty of resolving the merged decay products at high mass. Figure 7.10 shows a more detailed comparison of the signal efficiency in the $3b$ vs $4b$ signal regions for the RSG model. The efficiencies shown here are relative to all prior selection requirements. It can be seen there that at high masses the $3b$ signal region is more efficient for signal.

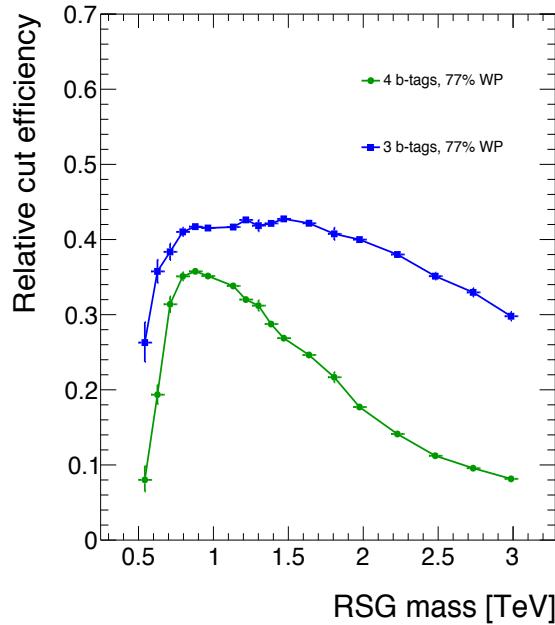


Figure 7.10: Efficiency of requiring 3 or 4 b -tagged track jets vs. RSG mass. The efficiency quoted is relative to the previous selection requirements (rather than an absolute efficiency).

To investigate the degradation of b -tagging efficiency at high p_T , the individual jet tagging efficiencies can be compared as a function of signal mass. This is shown in figure 7.11. The figure shows that the leading jet tagging efficiency in both calorimeter jets degrades heavily, while the sub-lead jet tagging efficiency remains relatively constant. More details on the cause of this degradation are shown in appendix A.

The final discriminating variable used in the boosted analysis is M_{2J} , the invariant mass of the two Higgs candidates. In order to improve the mass resolution, the four-momenta of each Higgs candidate are scaled by m_h/M_J . The effect of this correction is small in the boosted analysis but is done for consistency with the resolved selection. Table 7.2 shows the effect of the selection requirements on signal and background simulations as well as data.

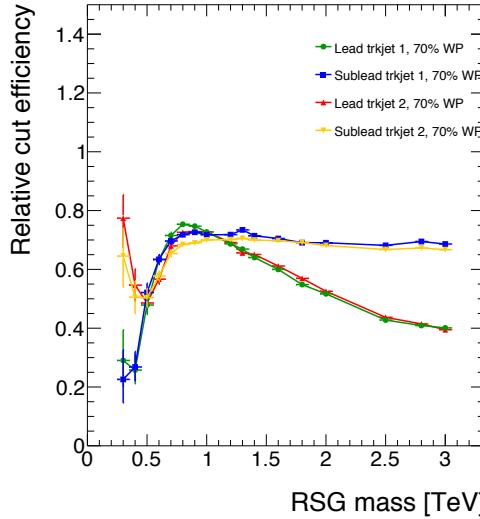


Figure 7.11: MV2c20 b -tagging efficiency for each of the four track jets in the boosted $4b$ selection as a function of RSG mass for $c = 1$ models.

Selection	Data	$m_{G_{KK}^*} = 1\text{TeV}$	$m_{G_{KK}^*} = 2\text{TeV}$	$t\bar{t}$	$Z + \text{jets}$
$N(\text{fiducial large-R jets}) \geq 2$	2202396	23.3	0.48	32345.2	4255.7
leading large-R jet $p_T > 350\text{ GeV}$	1873741	22.9	0.48	26511.7	3649.9
Both large-R jet $m > 50\text{ GeV}$	1854625	21.2	0.47	24369.8	3575.8
Both large-R jet $p_T < 1500\text{ GeV}$	1853601	21.2	0.46	24346.5	3572.9
$ \Delta\eta(JJ) < 1.7$	1435273	20.8	0.44	20751.0	3265.8
≥ 2 track-jets per large-R jet	1224727	19.8	0.40	18234.5	2692.6
3 b -tags, $X_{hh} < 1.6$	316	3.4	0.067	46.7	2.0
4 b -tags, $X_{hh} < 1.6$	20	2.9	0.030	1.4	0.0

Table 7.2: Effect of boosted selection on data, RSG signal models, $t\bar{t}$, and $Z + \text{jets}$. The numbers from simulation are normalized with the MC generator cross section and do not take into account the data driven estimates described in section 7.6 [114].

7.6 DATA-DRIVEN BACKGROUND ESTIMATION

The largest background to this final state is QCD multijet production, constituting 80-90% of the total background. Because of the difficulties in modeling higher order QCD processes, this background is estimated with a fully data-driven method. The only other non-negligible background is $t\bar{t}$, constituting the other 10-20%. Due to the presence of $t\bar{t}$ in the sideband region where the QCD background will be

⁴The $Z + \text{jets}$ background is a sub-percent level contribution

2417 estimated, the normalization of the QCD and $t\bar{t}$ backgrounds are simultaneously estimated.

2418 **7.6.1 MASS REGION DEFINITIONS**

2419 The first step in the data-driven background estimate is to define a sideband mass region where the
2420 background normalization can be derived. Additionally, a control region is defined where the background
2421 estimate can be validated. The control (CR) and sideband (SB) regions are defined using a radial distance
2422 in the two-dimensional large-R jet mass plane, R_{hh} , which is defined in equation 7.4.

$$R_{hh} = \sqrt{(M_J^{\text{lead}} - 124 \text{ GeV})^2 + (M_J^{\text{sublead}} - 115 \text{ GeV})^2} \quad (7.4)$$

2423 Events in the sideband region are required to fail the signal region $X_{hh} < 1.6$ requirement and have
2424 $R_{hh} > 35.8 \text{ GeV}$. The control region consists of those events which are not in the signal or sideband
regions. Figure 7.12 shows the definition of the signal, control, and sideband mass regions. Table 7.3 sum-

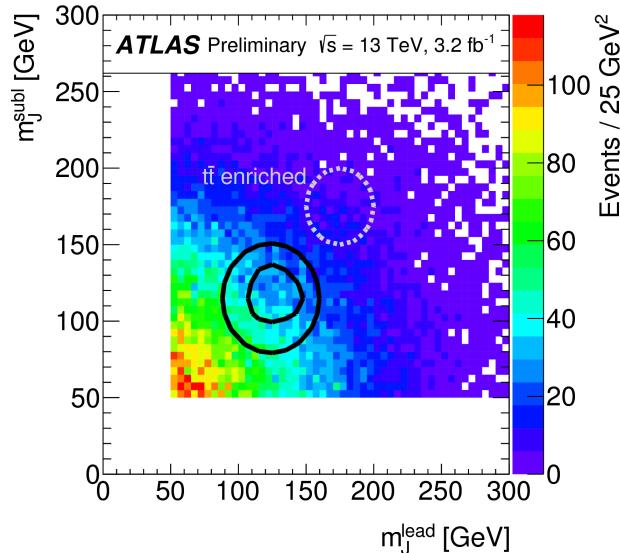


Figure 7.12: M_J^{sublead} vs. M_J^{lead} in a 2 b -tag data sample. The signal region is defined by the inner black contour ($X_{hh} < 1.6$) and the sideband region is defined by the outer contour ($R_{hh} > 35.8 \text{ GeV}$). The region between the black contours is the control region. The mass region which is enriched in $t\bar{t}$ background is also shown for illustration. [[113](#)]

2425

2426 marizes the mass region selections for the three different regions used in the analysis.

Region	Requirement	Notes
Signal Region (SR)	$X_{hh} < 1.6$	-
Control Region (CR)	$R_{hh} < 35.8 \text{ GeV}$ and $X_{hh} > 1.6$	Used for validation of background estimates
Sideband Region (SB)	$R_{hh} > 35.8 \text{ GeV}$	Used to derive background normalization

Table 7.3: Mass region definitions used for background estimation

2427 7.6.2 BACKGROUND ESTIMATION

2428 The method for estimating the background in this analysis is similar to the ABCD method presented in
 2429 Chapter 5. In this case, the two handles used to define different regions for the estimate are the number of
 2430 b -tagged track jets and the mass regions. A region requiring exactly two b -tagged track jets in one large-R
 2431 jet (referred to as the 2-tag or $2b$ region) is defined for use in the background estimate. The number of
 2432 expected background events in the $3b$ and $4b$ signal regions is then given by equation 7.5.

$$N_{\text{bkg}}^{3(4)-\text{tag},\text{SR}} = \mu_{\text{Multijet}} N_{\text{Multijet}}^{2-\text{tag},\text{SR}} + \beta_{t\bar{t}} N_{t\bar{t}}^{3(4)-\text{tag},\text{SR}} + N_{Z+\text{jets}}^{3(4)-\text{tag},\text{SR}} \quad (7.5)$$

2433 In this equation, $N_{\text{bkg}}^{3(4)-\text{tag}}$ is the expected number of background events in the $3b$ or $4b$ signal regions.
 2434 $N_{\text{Multijet}}^{2-\text{tag}}$ is the number of multijet events in the 2-tag region. $N_{t\bar{t}}^{3(4)-\text{tag}}$ is the number of $t\bar{t}$ events pre-
 2435 dicted in the MC for the $3b$ or $4b$ signal region, and the variable is similarly defined for the $Z+\text{jets}$ back-
 2436 ground. The $\beta_{t\bar{t}}$ parameter is a scale factor used to correct the normalization of the $t\bar{t}$ estimate in the signal
 2437 region. μ_{Multijet} is an extrapolation factor that is derived in the sideband region and used to estimate the
 2438 ratio of 2-tag events to 3(4)-tag events in the signal region. It is defined in equation 7.6.

$$\mu_{\text{Multijet}} = \frac{N_{\text{Multijet}}^{3(4)-\text{tag},\text{SB}}}{N_{\text{Multijet}}^{2-\text{tag},\text{SB}}} = \frac{N_{\text{data}}^{3(4)-\text{tag},\text{SB}} - \beta_{t\bar{t}} N_{t\bar{t}}^{3(4)-\text{tag},\text{SB}} - N_{Z+\text{jets}}^{3(4)-\text{tag},\text{SB}}}{N_{\text{data}}^{2-\text{tag},\text{SB}} - \beta_{t\bar{t}} N_{t\bar{t}}^{2-\text{tag},\text{SB}} - N_{Z+\text{jets}}^{2-\text{tag},\text{SB}}} \quad (7.6)$$

2439 The $t\bar{t}$ scale factor ($\beta_{t\bar{t}}$) and the QCD multijet extrapolation factor (μ_{Multijet}) are estimated together in
 2440 a simultaneous fit in the sideband region. Then, the number of events in the 2-tag signal region is used,
 2441 along with the $t\bar{t}$ estimate in the $3b$ and $4b$ signal regions and μ_{Multijet} , to estimate the total number
 2442 of background events in the two final signal regions. The shape of the final discriminant M_{2J} is also

2443 taken from the 2-tag signal region where there are more statistics. This method is illustrated graphically in
2444 figure 7.13.

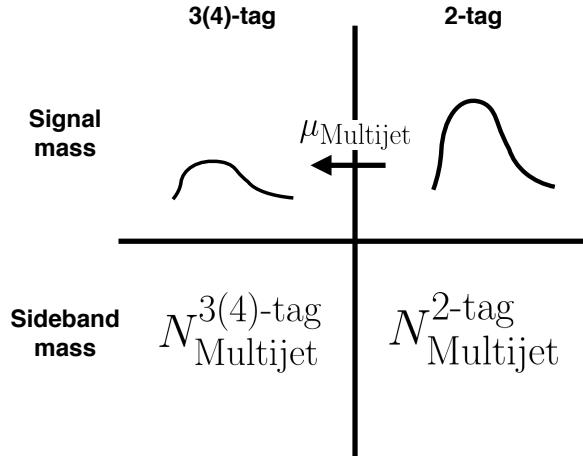


Figure 7.13: An illustration of the data-driven background estimation technique for the boosted analysis

2445 In the $3b$ region, the fit yields values of $\mu_{\text{Multijet}} = 0.160 \pm 0.03$ and $\beta_{t\bar{t}} = 1.02 \pm 0.09$. In the $4b$
2446 region, the fit gives $\mu_{\text{Multijet}} = 0.0091 \pm 0.0007$ and $\beta_{t\bar{t}} = 0.82 \pm 0.39$. The uncertainties quoted are
2447 statistical only. The larger uncertainties in the $4b$ values indicate the lower statistics available in that region.

2448 Figure 7.14 shows the distributions of data and background estimates in the $3b$ and $4b$ sideband regions
2449 after the background fit has been done. The normalizations are constrained from the fit to match that of
2450 the data, but good modeling of the shape of the mass of the leading large-R jet is seen as well. The shapes
2451 of the kinematic distributions in the $4b$ region are taken from the $3b$ region due to the better MC statistics
2452 in that region.

2453 7.6.3 BACKGROUND SHAPE FIT

2454 As mentioned in the previous section, the background shape in the 3-tag and 4-tag signal regions is
2455 taken from the 2-tag signal mass region. Due to the limited statistics available, the background shapes are
2456 additionally smoothed after being extrapolated to the 3-tag and 4-tag signal regions. Only the data in the
2457 range $900 < M_{2J} < 2000$ GeV is included in the fit due to the limited statistics available above 2 TeV.
2458 Both the $t\bar{t}$ and QCD multijet background are independently fit with an exponential shape, $y = e^{ax+b}$.

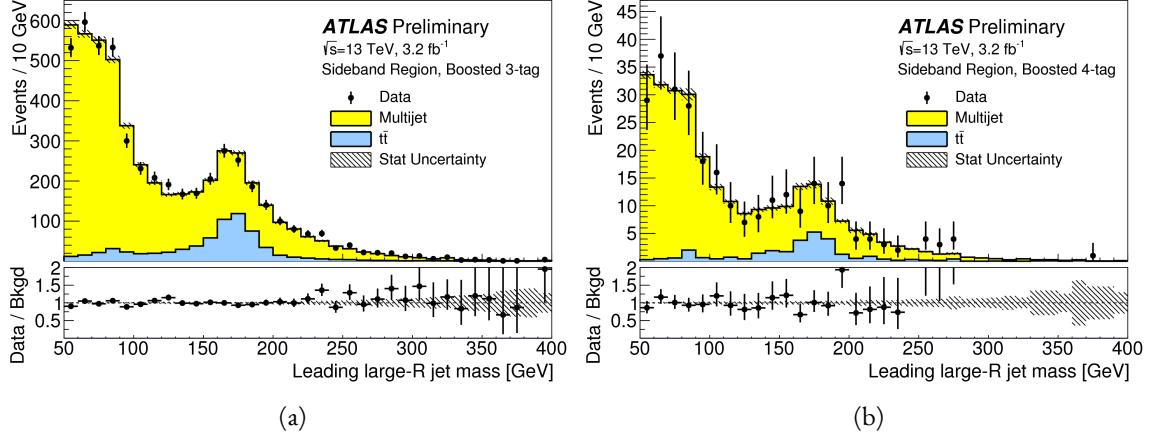


Figure 7.14: Leading large-R jet mass in the 3b (a) and 4b (b) sideband regions. The multijet and $t\bar{t}$ backgrounds are estimated using the data-driven methods described above. Because their normalizations are derived in the sideband region, the total background normalization is constrained by default to match the normalization of the data [113].

2459 Other shapes are considered and used for the systematic uncertainties. Table 7.4 shows the fit values for
 2460 the parameters. Because both the 3b and 4b QCD shapes come from the 2-tag region, the slopes derived
 2461 are very similar.

	a	b
QCD (4b)	0.00545 ± 0.00021	5.44 ± 0.24
$t\bar{t}$ (4b)	0.00746 ± 0.00021	4.88 ± 0.36
QCD (3b)	0.00545 ± 0.00021	8.30 ± 0.24
$t\bar{t}$ (3b)	0.00746 ± 0.00021	8.58 ± 0.36

Table 7.4: Parameters derived for exponential fit to background M_{2J} shape in the 3b and 4b signal regions [114]

2462 7.6.4 VALIDATION OF BACKGROUND ESTIMATE

2463 The background estimate can be validated by using the method to estimate the number of events in the
 2464 control mass region rather than the signal mass region. Figure 7.15 shows the M_{2J} distribution in the 3b
 2465 and 4b control regions, comparing data and background estimates. In both cases, both the background
 2466 shape and normalization are consistent with the data, indicating good agreement. The ratio of data to the
 2467 background estimates is also fit to a line in the figure to test for any shape difference. The slope of the
 2468 line is within 1σ (from the fit uncertainties) of flat, further indicating that the data is consistent with the
 2469 background estimate in the control region.

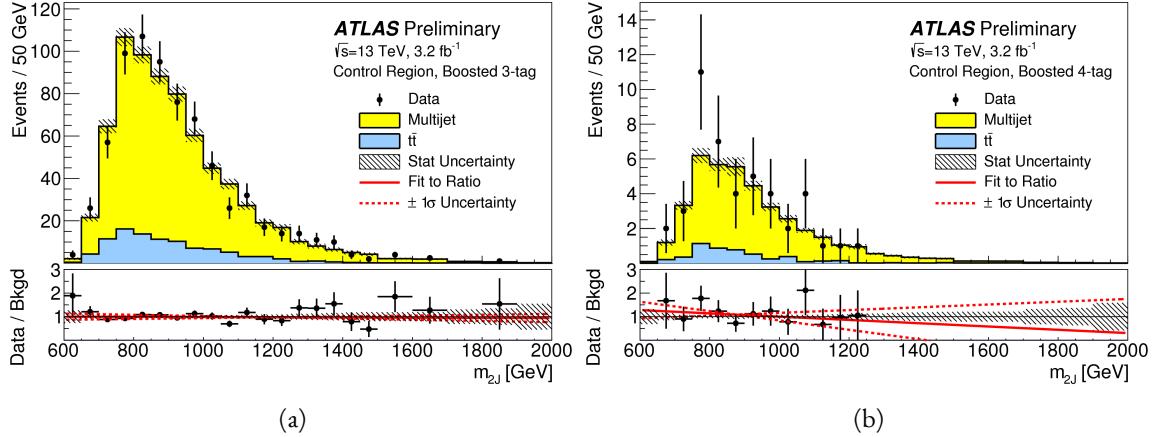


Figure 7.15: Di-jet invariant mass (M_{2J}) in the 3b (a) and 4b (b) control regions. The multijet and $t\bar{t}$ backgrounds are estimated using the data-driven methods described above [113].

Table 7.5 shows the yields in data and background estimates in the 3-tag and 4-tag sideband and control regions. Again, here, it can be seen that the total number of predicted background events from the data driven method is consistent with the number of data events in the region.

Sample (3-tag)	Sideband Region	Control Region
Multijet	4328 ± 27	607 ± 10
$t\bar{t}$	683.5 ± 8.1	99.6 ± 3.1
Z+jets	31.8 ± 3.7	7.7 ± 1.8
Total	5043 ± 28	715 ± 11
Data	5043	724
Sample (4-tag)	Sideband Region	Control Region
Multijet	247.4 ± 1.5	34.7 ± 0.6
$t\bar{t}$	28.4 ± 1.5	5.1 ± 0.7
Z+jets	3.4 ± 1.2	0.6 ± 0.5
Total	279.2 ± 2.5	40.3 ± 1.0
Data	279	45

Table 7.5: The number of events in data and predicted background events in the boosted 3-tag and 4-tag sideband and control regions. The uncertainties shown are statistical only. [113]

2473 7.7 SYSTEMATIC UNCERTAINTIES

2474 The systematic uncertainties in this analysis can be divided into two broad categories. The first type
2475 is uncertainties associated with the modeling of the signal processes. The second type of uncertainty is
2476 associated with both the shape and normalization of the background prediction.

2477 7.7.1 SIGNAL MODELING UNCERTAINTIES

2478 The signal modeling uncertainty has three main components: theoretical uncertainty on the acceptance,
2479 experimental uncertainties on the large-R jets, and experimental uncertainties on the track jets related to
2480 b -tagging. In this analysis the experimental uncertainties are the most significant.

2481 The first uncertainty on signal modeling is the theoretical uncertainty on the acceptance. As explained
2482 in section 5.6.1, there are four components to this uncertainty. The first is related to missing higher order
2483 terms from the matrix element calculations which is estimated by varying the QCD renormalization and
2484 factorization scales. The second is uncertainty due to the PDF set used. The third is a generator uncer-
2485 tainty which is estimated by modifying the generator used to model the underlying event and hadroniza-
2486 tion. Finally, there is an uncertainty associated with the modeling of the initial state and final state radia-
2487 tion (ISR/FSR). The total theoretical uncertainty on the signal yield is 3%, and this is dominated by the
2488 ISR/FSR modeling.

2489 There are uncertainties on the large-R jets in both the jet energy scale (JES) and jet energy resolution
2490 (JER) as well as the jet mass scale (JMS) and jet mass resolution (JMR). These are evaluated using $\sqrt{s} =$
2491 8 TeV data from Run 1 of ATLAS and extrapolated to the Run 2 beam and detector conditions using
2492 MC⁵. The details of these uncertainties can be found in reference [115].

2493 Uncertainties on the track jets are related to the b -tagging efficiency. The total uncertainty on the signal
2494 yield due to b -tagging is evaluated by propagating variations of the b -tagging efficiency through the boosted
2495 selection requirements. The uncertainties are calculated jet-by-jet and parameterized as a function of b -jet
2496 p_T and η [93]. For high p_T b -jets (with $p_T > 300$ GeV), the uncertainties are extrapolated using MC
2497 simulation from the lower p_T b -jets [116].

⁵The uncertainties are correspondingly larger due to the uncertainty of this extrapolation.

2498 Table 7.6 shows the systematic uncertainties on the signal normalization for models with $m_{G_{\text{KK}}^*} =$
 2499 1.5 TeV and both $c = 1$ and $c = 2$ as well as a narrow width heavy scalar. The dominant uncertainty
 2500 comes from b -tagging and this uncertainty is larger in the 4-tag region than the 3-tag region.

Source	Background	G_{KK}^*		H
		$c = 1$	$c = 2$	
Luminosity	-	5.0	5.0	5.0
3-tag				
JER	< 1	< 1	< 1	< 1
JES	2	< 1	< 1	< 1
JMR	1	12	12	11
JMS	5	14	13	17
b -tagging	1	23	22	23
Theoretical	-	3	3	3
Multijet Normalization	3	-	-	-
Statistical	2	1	1	1
Total	7	31	30	33
4-tag				
JER	< 1	< 1	< 1	< 1
JES	< 1	< 1	< 1	< 1
JMR	4	12	13	13
JMS	5	13	13	14
b -tagging	2	36	36	36
Theoretical	-	3	3	3
Multijet Normalization	14	-	-	-
Statistical	3	1	1	1
Total	15	42	42	43

Table 7.6: Summary of systematic uncertainties in the total background and signal event yields (expressed in %) in the boosted 3-tag and 4-tag signal regions. Systematic uncertainties on the signal normalization are shown for models with $m_{G_{\text{KK}}^*} = 1.5 \text{ TeV}$ and both $c = 1$ and $c = 2$ as well as a narrow width heavy scalar.

2501 **7.7.2 BACKGROUND UNCERTAINTIES**

2502 Uncertainties on the QCD multijet background normalization and shape are estimated using the con-
 2503 trol mass region. As shown previously, the background predictions in the control region match with the

2504 data yields within the statistical uncertainty in both the 3-tag and 4-tag control regions. As an additional
 2505 protection, the statistical uncertainty on the background prediction in the control region is assigned as a
 2506 systematic uncertainty on the normalization of the QCD background.

2507 Additional robustness tests are done by varying the definition of the control mass region and the b -
 2508 tagging requirements used to define the 2-tag sample. In all cases, the effect of the variations is found to be
 2509 within the statistical uncertainties on the background normalization in the control region.

2510 Shape uncertainties on the background are evaluated using two techniques. First, as shown in fig-
 2511 ure 7.15, the ratio between the data and background prediction is fit with a linear function. The uncer-
 2512 tainties on the slope of this fit are assigned as shape uncertainties. An additional uncertainty is assigned by
 2513 using alternate power law fit functions for the smoothing of the background shape. Table 7.7 shows the
 2514 alternate shapes used. The largest difference between the nominal fit function and the alternates, taking
 2515 into account the 1σ uncertainty band on each fit as well, is taken as a shape uncertainty.

Functional Form
$f_1(x) = p_0(1 - x)^{p_1}x^{p_2}$
$f_2(x) = p_0(1 - x)^{p_1}e^{p_2 x^2}$
$f_3(x) = p_0(1 - x)^{p_1}x^{p_2}x$
$f_4(x) = p_0(1 - x)^{p_1}x^{p_2} \ln x$
$f_5(x) = p_0(1 - x)^{p_1}(1 + x)^{p_2}x$
$f_6(x) = p_0(1 - x)^{p_1}(1 + x)^{p_2} \ln x$
$f_7(x) = \frac{p_0}{x}(1 - x)^{p_1-p_2} \ln x$
$f_8(x) = \frac{p_0}{x^2}(1 - x)^{p_1-p_2} \ln x$

Table 7.7: Alternate fit functions used to model the M_{2J} distribution in the QCD multijet background. In the equations, $x = M_{2J}/\sqrt{s}$.

2516 The uncertainties on the $t\bar{t}$ background are obtained by propagating the various experimental variations
 2517 (JES, JER, JMS, JMR, b -tagging) through the analysis selection requirements. Table 7.6 summarizes the
 2518 background uncertainties in the 3-tag and 4-tag regions.

2519 7.8 RESULTS

2520 Table 7.8 shows the observed yields in the 3-tag and 4-tag signal regions for the boosted analysis com-
2521 pared to the predicted number of background events. In the 3-tag region, 316 events are observed with
2522 a predicted background of 285 ± 19 . In the 4-tag region, 20 events are observed with a predicted back-
2523 ground of 14.6 ± 2.4 . Figure 7.16 shows the M_{2J} distribution in the 3-tag and 4-tag regions. There are
2524 some small excesses in the data, in particular in the 3-tag region around $M_{2J} \approx 900$ GeV and in the region
2525 of $1.6 < M_{2J} < 2.0$ TeV. The significance of these excesses will be evaluated in the next chapter in the
2526 statistical combination with the resolved results.

Sample	Signal Region (3-tag)	Signal Region (4-tag)
Multijet	235 ± 14	13.5 ± 2.4
$t\bar{t}$	48 ± 22	1.2 ± 1.0
$Z + \text{jets}$	2.0 ± 2.2	-
Total	285 ± 19	14.6 ± 2.4
Data	316	20
G_{KK}^* (1000 GeV), $c = 1$	3.4 ± 0.9	2.9 ± 1.1

Table 7.8: Observed yields in the 3-tag and 4-tag signal regions for the boosted analysis compared to the predicted number of background events Errors correspond to the total uncertainties in the predicted event yields. The yields for a graviton with $m_{G_{\text{KK}}^*} = 1$ TeV and $c = 1$ are also shown. [113]

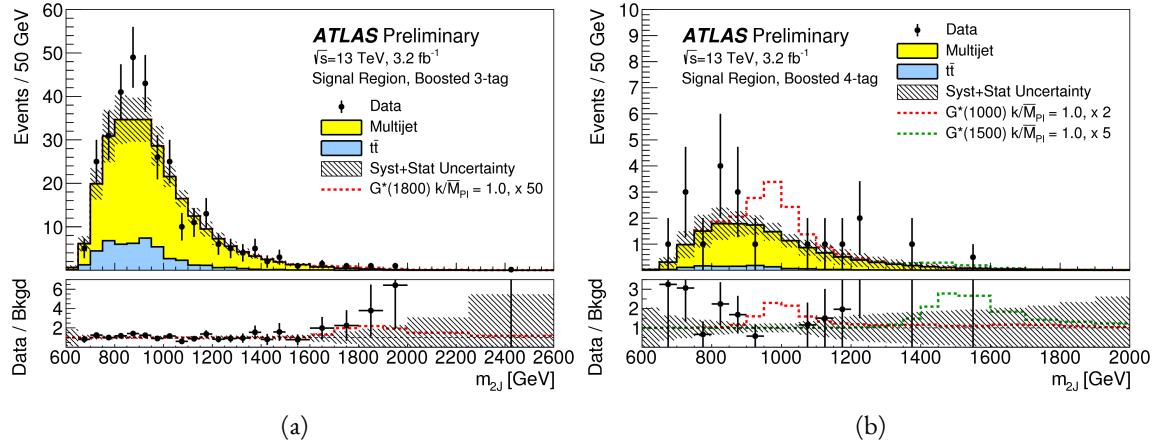


Figure 7.16: Di-jet invariant mass (M_{2J}) in the $3b$ (a) and $4b$ (b) signal regions. The multijet and $t\bar{t}$ backgrounds are estimated using the data-driven methods described above. In the $3b$ region, a graviton signal with $m_{G_{KK}^*} = 1.8$ TeV and $c = 1$ is overlaid, with the cross section multiplied by a factor of 50 so that the signal is visible. In the $4b$ region, signals with $m_{G_{KK}^*} = 1.0$ TeV and $m_{G_{KK}^*} = 1.5$ TeV are overlaid, both with $c = 1$ and the yields multiplied by factors of 2 and 5 respectively [113].

*There is no real ending. It's just the place where you stop
the story.*

Frank Herbert

8

2527

2528

Combined limits from boosted and resolved searches

2529

2530

8.1 INTRODUCTION

2531 In order to cover the full mass range of possible resonances decaying to di-Higgs final states, two distinct
2532 tailored selections were produced. The resolved selection is more sensitive in the mass range of $400 < m_X < 1100$ GeV while the boosted selection is more sensitive to masses in the range $1100 < m_X < 3000$ GeV. Chapter 7 presents the details of the boosted selection and results. In setting limits on spin-2
2533 Randall-Sundrum graviton (RSG) and narrow width heavy scalar (H) models, the results of the boosted
2534 selection are combined with the results of the resolved selection to cover the full mass range.
2535

2536 This chapter presents limits on signal models resulting from the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search in both
2537 the resolved and boosted selections. It first presents a brief overview of the resolved results that go into

2539 the limit setting. Then, an overview of the statistical methods used for the search and limit setting is given.
 2540 Finally, limits on the RSG and heavy scalar models are presented.

2541 **8.2 RESOLVED RESULTS**

2542 The details of the resolved selection will not be presented here and can be found in reference [113]. In
 2543 basic terms, the selection searches for four $R = 0.4$ b-tagged calorimeter jets (where each pair of jets is
 2544 one Higgs candidate). This is distinct from the boosted methodology which searches for merged decay
 2545 products. The backgrounds to the resolved selection are the same as those presented in Chapter 7 for the
 2546 boosted analysis.

2547 Table 8.1 shows the results for data yields and expected background in the resolved signal region. Fig-
 2548 ure 8.1 shows the M_{2J} distribution in the resolved signal region. The total number of events is consistent
 2549 with the prediction and no significant excess is seen. One event in the boosted 4-tag signal is shared with
 2550 the resolved signal region and has a mass of 852 GeV.

Sample	Signal Region Yield
Multijet	43.3 ± 2.3
$t\bar{t}$	4.3 ± 3.0
$Z + \text{jets}$	-
Total	47.6 ± 3.8
Data	46
SM hh	0.25 ± 0.07
$G_{\text{KK}}^*(800 \text{ GeV}), c = 1$	5.7 ± 1.5

Table 8.1: Observed yields in the resolve selection 4-tag signal region compared to the predicted number of background events Errors correspond to the total uncertainties in the predicted event yields. The yields for a graviton with $m_{G_{\text{KK}}^*} = 800$ GeV and $c = 1$ are also shown. [113]

2551 **8.3 SEARCH TECHNIQUE AND RESULTS**

2552 The statistical technique used for the search in this analysis is the same as that used in the $H \rightarrow WW^*$
 2553 analysis presented in section 3.6.2. The test statistic q_0 is used to define the p -values which measure the

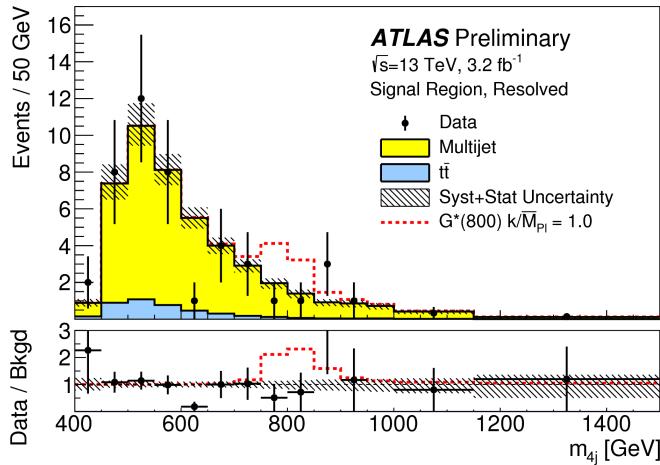


Figure 8.1: Di-jet invariant mass (M_{2J}) in the resolved signal region. A graviton signal with $m_{G_{KK}^*} = 800$ GeV and $c = 1$ is overlaid. [113].

compatibility of the data with the background-only hypothesis corresponding to a signal strength $\mu = 0$.

Local p_0 values are computed to quantify the probability that the background could produce a fluctuation greater than or equal to the one observed in the data. In the resolved analysis, no significant excesses are observed. The largest discrepancy with respect to the background only hypothesis occurs near a resonance mass of 900 GeV and is found to be less than 2σ in significance.

In the boosted selection, the largest local excess is a broad excess in the $3b$ signal region that begins near $M_{2J} \approx 1.7$ GeV. Assuming a G_{KK}^* with this mass and $c = 1.0$, the local significance of this excess is 2.0σ .

8.4 LIMIT SETTING

In the absence of any significant excess observed in the data, limits on different signal models can be set. This section describes the limit setting procedure and presents combined results of the resolved and boosted analyses.

8.4.1 LIMIT SETTING PROCEDURE

The procedure used for setting exclusion limits in this analysis is the CL_s method [117]. The first step in setting the limits is to define a test statistic which will be used. For limit setting, the test statistic is shown

2569 in equation 8.1.

$$\tilde{q}_\mu = \begin{cases} -2 \ln \frac{L(\mu, \hat{\theta}(\mu))}{L(0, \hat{\theta}(0))} & \hat{\mu} < 0 \\ -2 \ln \frac{L(\mu, \hat{\theta}(\mu))}{L(\hat{\mu}, \hat{\theta})} & 0 \leq \hat{\mu} < \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad (8.1)$$

2570 In the above equation, μ is the value of the signal strength under test, $\hat{\mu}$ is the best fit μ , $\hat{\theta}$ is the
2571 best fit value of the nuisance parameters, $\hat{\theta}$ is the best fit value of the nuisance parameters under the fixed
2572 μ value, and L is the Poisson likelihood of the data (as described in section 3.6.2).

2573 The test statistic \tilde{q}_μ is constructed to protect against two interesting corner cases when setting the upper
2574 limit on the cross section. First, it protects against negative signal strengths μ which are unphysical. Second,
2575 it does not count excesses in the data larger than those expected by a signal strength μ as evidence against
2576 the μ hypothesis.

2577 The CL_s statistic is constructed by taking a ratio of two probabilities. CL_{s+b} is the probability that the
2578 signal+background hypothesis would produce a value of the test statistic that is less than or equal to the
2579 observed value¹. CL_b is the probability that the background only hypothesis will produce a value
2580 of the test statistics less than or equal to the observed. The CL_s statistic is then the ratio CL_{s+b}/CL_b . A
2581 95% upper limit on the cross section is set at the value of μ that makes the CL_s statistic less than 5%.

2582 In practice, the limits are computed numerically within an asymptotic approximation for the distribu-
2583 tion of the test statistic \tilde{q}_μ . The details of this approximation can be found in reference [65].

2584 The resolved and boosted analyses are combined using a very simple procedure rather than a full statis-
2585 tical combination. For each mass point tested, the limit which gives the most stringent constraint is used.
2586 This means that for mass points below 1.1 TeV the resolved signal region is used, while at and above this
2587 point the combination of the orthogonal 3b and 4b boosted signal regions is used.

2588 8.4.2 LIMIT SETTING RESULTS

2589 Figure 8.2 shows the combined 95% upper bounds as a function of mass for three different models:
2590 G_{KK}^* with $c = 1$, G_{KK}^* with $c = 2$, and a narrow heavy scalar H .

¹Lower values of \tilde{q}_μ mean better compatibility

2591 The cross section of $\sigma(pp \rightarrow G_{\text{KK}}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $c = 1$ is constrained to be less than 70 fb
2592 for masses in the range $600 < m_{G_{\text{KK}}^*} < 3000$ GeV. For the RSG model with $c = 2$, cross sections limits
2593 between 40 fb and 200 fb are set for the mass range of $500 < m_{G_{\text{KK}}^*} < 3000$ GeV. Masses in the range
2594 of $475 < m_{G_{\text{KK}}^*} < 785$ GeV are excluded with $c = 1$ (with an exclusion of the range 465 to 745 GeV
2595 expected). Masses less than 980 GeV are excluded with $c = 2$ (with an exclusion for masses less than
2596 1 TeV expected).

2597 In the heavy Higgs model, the cross section upper limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges from
2598 30 to 300 fb in the mass range of $500 < m_H < 3000$ GeV. The resolved analysis can also set an upper
2599 limit on the Standard Model di-Higgs production cross section discussed in chapter 3. The upper limit on
2600 $\sigma(pp \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ in the Standard Model is constrained to be less than 1.22 pb.

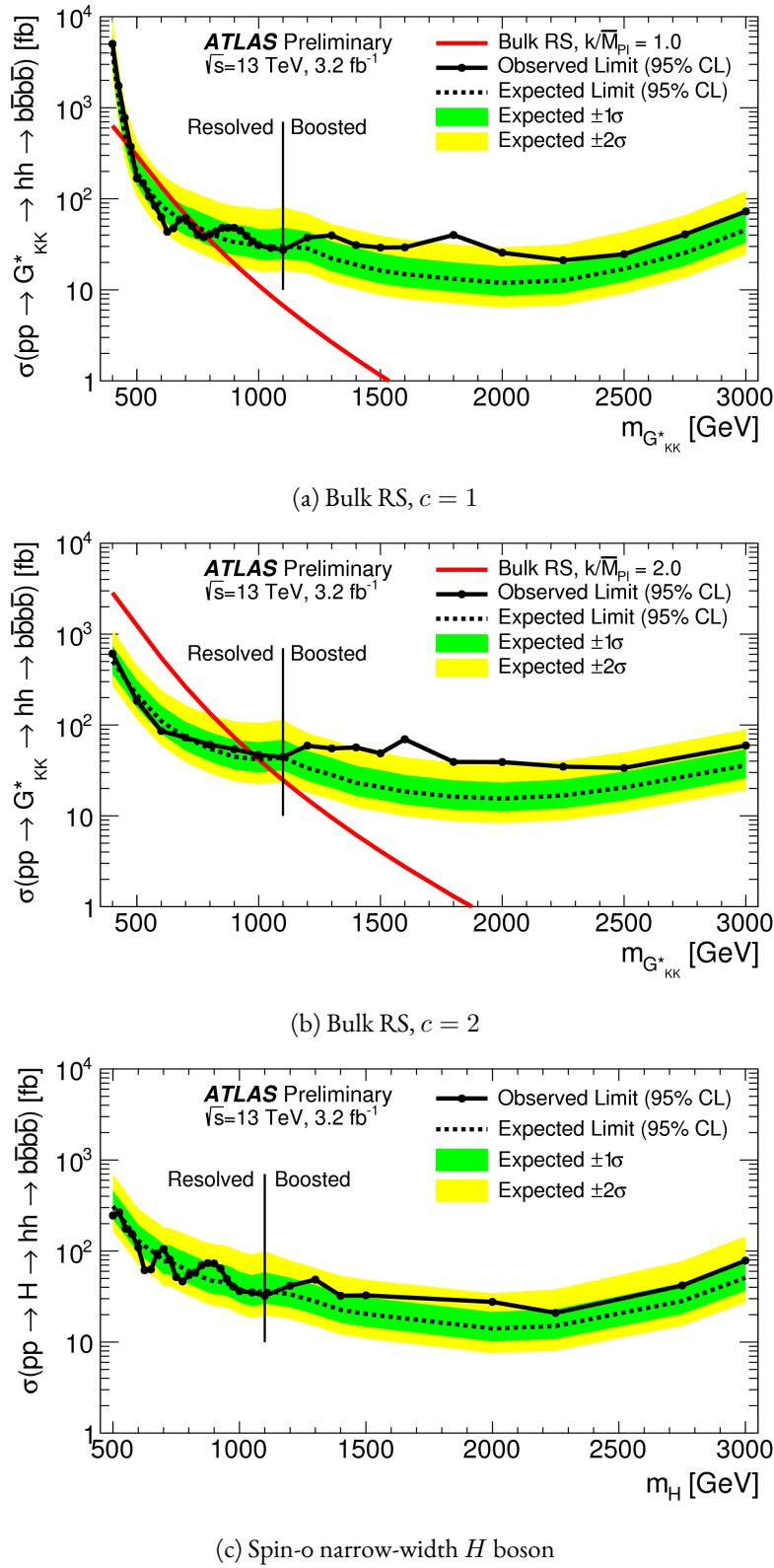


Figure 8.2: Expected and observed upper limit as a function of mass for G_{KK}^* in the RSG model with (a) $c = 1$ and (b) $c = 2$, as well as (c) H with fixed $\Gamma_H = 1$ GeV, at the 95% confidence level in the CL_s method. [13]

2601

Part IV

2602

Looking ahead

9

2603

2604

Conclusion

2605 After being sought for many years at different collider experiments, the Higgs boson was discovered by
2606 the ATLAS and CMS experiments in 2012, confirming the leading theory for the source of electroweak
2607 symmetry breaking and filling in the last missing piece of the Standard Model. After its discovery, mea-
2608 surements of the particle's detailed properties and searches for new particles decaying to Higgs final states
2609 were both extremely important in constraining physics beyond the Standard Model. This dissertation
2610 presented this evolution through two results: the observation and measurement of the Higgs boson in the
2611 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 8$ TeV and a search for Higgs pair production
2612 in the $HH \rightarrow b\bar{b}b\bar{b}$ channel at $\sqrt{s} = 13$ TeV with the ATLAS detector in pp collisions at the Large
2613 Hadron Collider.

2614 In the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$, results from both the discovery of the Higgs boson and the full ATLAS
2615 Run 1 dataset were presented. The Higgs boson was discovered with a 6.1σ significance in a combination
2616 of the $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ4\ell$, $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ with 4.2 fb^{-1} at $\sqrt{s} = 7$ TeV and 5.2 fb^{-1} at

$\sqrt{s} = 8$ TeV. With the full 20.3 fb^{-1} at $\sqrt{s} = 8$ TeV and 4.2 fb^{-1} at $\sqrt{s} = 7$ TeV, ATLAS achieved discovery level significance in the $H \rightarrow WW^*$ channel alone and obtained the first evidence of vector boson fusion production in that channel. The combined signal strength is measured to be $\mu = 1.09^{+0.23}_{-0.21}$. The total observed significance of the $H \rightarrow WW^*$ process is observed to be 6.1σ (with 5.8σ expected). Advanced methods for background reduction and estimation, particularly in same-flavor lepton final states, are shown. The VBF signal strength is measured to be $\mu_{\text{VBF}} = 1.27^{+0.53}_{-0.45}$ with an observed significance of 3.2σ (with 2.7σ expected).

These results required many novel innovations. The increase of pileup interactions in the higher instantaneous luminosity LHC conditions of 2012 led to a degradation of missing transverse momentum resolution. As a result, the prominent $Z/\gamma^* + \text{jets}$ background of the same flavor $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ final states increased greatly. New variables, including a track-based missing transverse momentum and a measurement of the balance between the dilepton system and recoiling jets, allowed for significant reduction of this background. In the VBF channel, selections were optimized to exploit the unique VBF final state topology. Incorporating these variables into a boosted decision tree technique allowed the analysis to exceed the 3σ statistical significance threshold.

After the end of Run 1, the results of Higgs measurements from ATLAS were combined with those from CMS to produce the most precise measurements of the Higgs boson so far [118]. Figure 9.1 shows the combination of ATLAS and CMS data for the Higgs signal strength in and coupling measurements. In the signal strength measurements of gluon fusion and vector boson fusion, the $H \rightarrow WW^*$ channel provides the tightest constraints. Additionally, the Higgs coupling to W bosons is the most precisely measured with a relative uncertainty of 10%.

With the discovery of the Higgs firmly established and its properties measured, a natural next step was to search for new physics with Higgs final states. At $\sqrt{s} = 13$ TeV, a search for Higgs pair production in the $b\bar{b}b\bar{b}$ final state with 3.2 fb^{-1} was conducted. A signal region optimized for the boosted final states arising from high mass resonances was constructed. This signal region utilized large-radius calorimeter jets and b -tagging with small radius track jets to maximize the signal acceptance. No significant excesses were observed, and upper limits on cross sections are placed for spin-2 Randall Sundrum gravitons (RSG) and

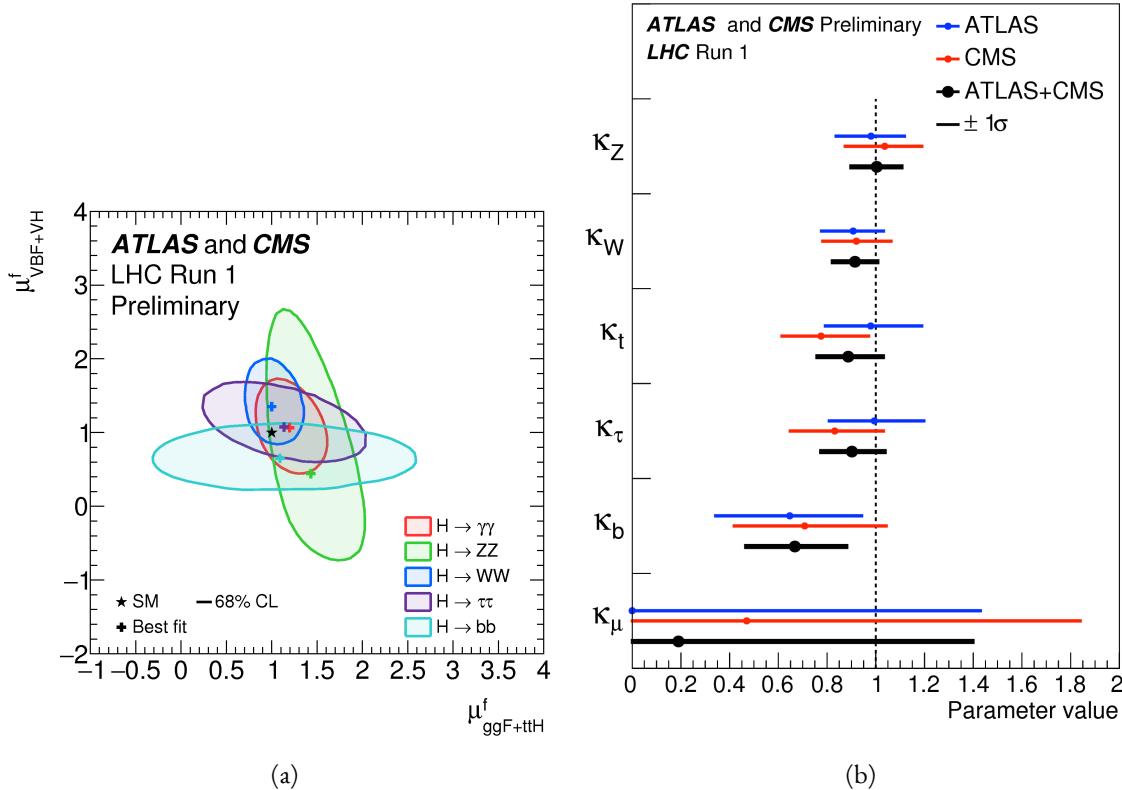


Figure 9.1: Combined ATLAS and CMS measurements in Run 1 for (a) Higgs signal strength in gluon fusion and VBF and (b) Higgs couplings normalized to their SM predictions

narrow spin-0 resonances. The increase in center of mass energy in Run 2 allowed this analysis to extend upper limits up to 3 TeV, while previous results from ATLAS in Run 1 only quotes limits up to 2 TeV. The cross section of $\sigma(pp \rightarrow G_{KK}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $k/\bar{M}_{\text{Pl}} = 1$ is constrained to be less than 70 fb for masses in the range $600 < m_{G_{KK}^*} < 3000$ GeV. For the RSG model with $k/\bar{M}_{\text{Pl}} = 2$, cross sections limits between 40 fb and 200 fb are set for the mass range of $500 < m_{G_{KK}^*} < 3000$ GeV. The cross section upper limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges from 30 to 300 fb in the mass range of $500 < m_H < 3000$ GeV.

While there has been a rigorous program of measurements and searches involving the Higgs, there is still much room for improvement at the High Luminosity LHC (HL-LHC) and beyond. The measured signal strength for VBF production in $H \rightarrow WW^*$ still has a relative error at the level of 40%, largely dominated by statistical uncertainty. Projections for the HL-LHC show that the uncertainty on the VBF signal strength can be reduced to approximately 15% with 3000 fb^{-1} [119, 120]. This uncertainty also

2656 assumes that theoretical uncertainties on the signal, which would be the largest contribution in this dataset,
 2657 remain as they are now. Improvements in the theoretical understanding of the Higgs signal would also
 2658 reduce the signal strength uncertainty dramatically. Such precision measurements allow for measurements
 2659 of the Higgs coupling to vector bosons precise to the few percent level, therefore giving much power to
 2660 constrain or discover new physics.

2661 The prospects for detection of beyond the Standard Model resonant di-Higgs production at the HL-
 2662 LHC are also quite promising. Figure 9.2 shows projections for the discovery significance of RSG signals
 2663 at the HL-LHC in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search [120]. In all detector budget scenarios, a 1.5 TeV
 2664 resonance is above or near 5σ significance, while a 2 TeV resonance is between $4-5\sigma$ except for the lowest
 budget.

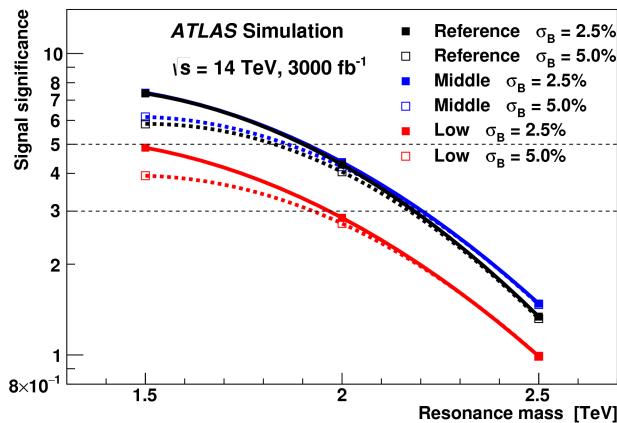


Figure 9.2: Discovery significance for RSG models at the HL-LHC in three different budget scenarios [120].
 Systematic uncertainties on the background prediction (σ_B) of 2.5% and 5.0% are both tested.

2665
 2666 The Higgs will continue to be an incredibly powerful tool in the understanding of nature at the HL-
 2667 LHC and beyond. Through both precision measurements and searches, the nature of electroweak symme-
 2668 try breaking will be better understood and the potential for the discovery of physics beyond the Standard
 2669 Model has never been greater.

A

2670

2671

b-tagging performance at high p_T

2672 One of the limiting factors of the signal acceptance in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search at high resonance
2673 masses is the degradation of the *b*-tagging efficiency for high p_T jets. This appendix presents a study of the
2674 underlying causes of this degradation.

2675 A.I CHANGES IN MV₂ SCORE AT HIGH p_T

2676 The degradation of *b*-tagging at high p_T was studied in particular in the context of RSG models at high
2677 mass. Figure A.I shows the p_T of the leading track jet inside of the leading calorimeter jet in RSG events.
2678 At high $m_{G_{KK}^*}$, the p_T spectrum of track jets is much harder than at lower masses due to the increased
2679 Higgs p_T .

2680 Figure A.2 shows the MV_{2c2o} algorithm score for the leading and subleading track jets inside of the
2681 leading calorimeter jet. In both cases, it can be seen that at higher RSG masses the MV₂ score shifts towards
2682 more background like (negative) values. Additionally, this effect is more pronounced in the leading track

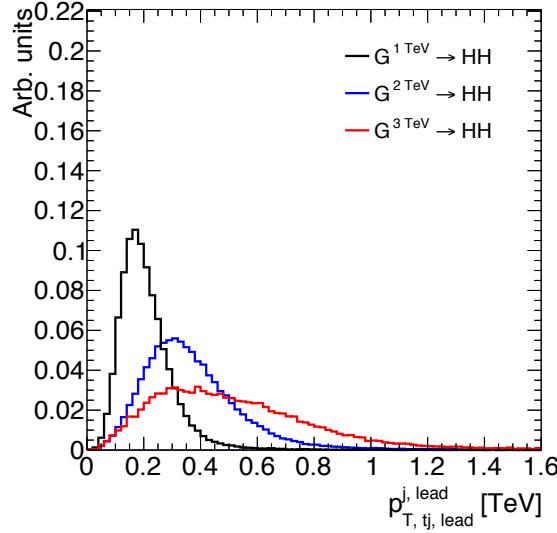


Figure A.1: p_T of the leading track jet in the leading calorimeter jet for different signal masses in RSG $c = 1$ models

2683 jet than the subleading.

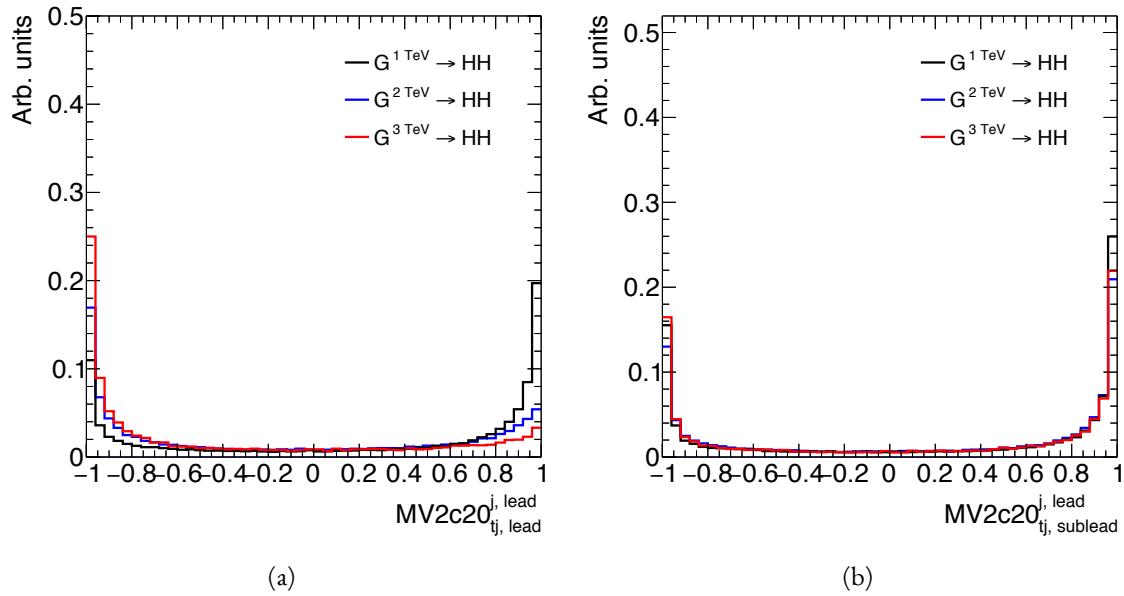


Figure A.2: MV2c20 score for the leading track jet (a) and subleading track jet (b) of the leading calorimeter jet for different signal masses in RSG $c = 1$ models

2684 To understand what is causing this change in the MV2c20 score, the same comparisons can be made for
 2685 the input variables of MV2c20. The focus in these comparisons will be on the leading track jet as this is the
 2686 one seen to have the largest difference in MV2 score. Figure A.3 shows the log likelihood ratio $\log(p_b/p_u)$

from the IP₃D (three dimensional impact parameter) algorithm. At higher masses, the IP₃D likelihood ratio distribution does become more background-like.

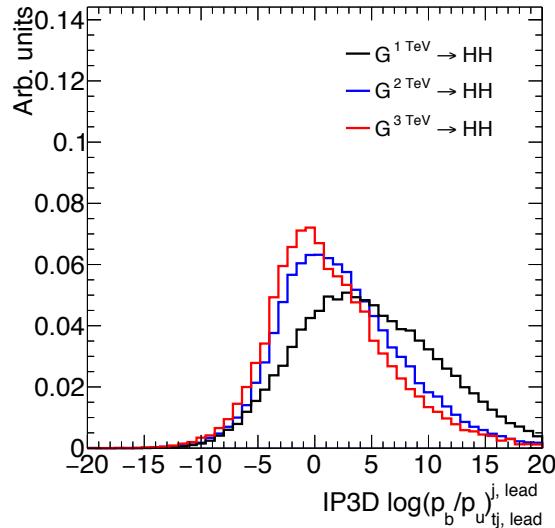


Figure A.3: IP₃D log-likelihood ratio ($\log(p_b/p_u)$) of the leading track jet in the leading calorimeter jet for different signal masses in RSG $c = 1$ models

Figure A.4 shows the mass and number of tracks at the secondary vertex computed by the SV1 algorithm. When there is no secondary vertex found, the algorithm assigns a default negative value for these quantities. Both of these distributions show that there is a significantly larger fraction of jets where no secondary vertex is found in the high mass samples compared to the $m_{G_{KK}^*} = 1$ TeV sample. The SV1 algorithm's inability to find a secondary vertex could be an important factor in the overall MV₂ score shift, as this eliminates eight of the input variables that would normally contribute information to the algorithm.

Figure A.5 shows the same quantities for the JetFitter algorithm. In this case, there is also a change in the fraction of jets which have their secondary vertices successfully reconstructed, but this change is not as drastic as that seen in SV1. There is also an increase in the number of jets which have high values of mass.

A.2 EFFECT OF MULTIPLE b -QUARKS INSIDE ONE JET

One hypothesis for why the efficiency of b -tagging the leading track jet degrades is that at high masses, the b quarks get close enough together that both of them are inside of the leading track jet. Because MV₂

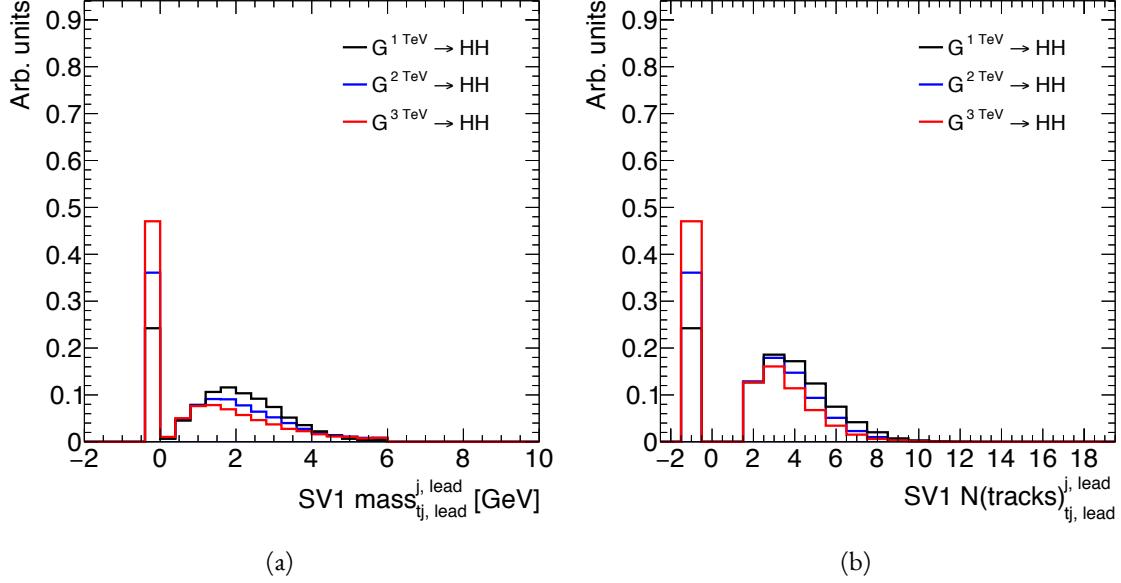


Figure A.4: Mass (a) and number of tracks (b) for the secondary vertices computed with the SV1 algorithm. When no secondary vertex is found, the quantities are assigned to default negative values.

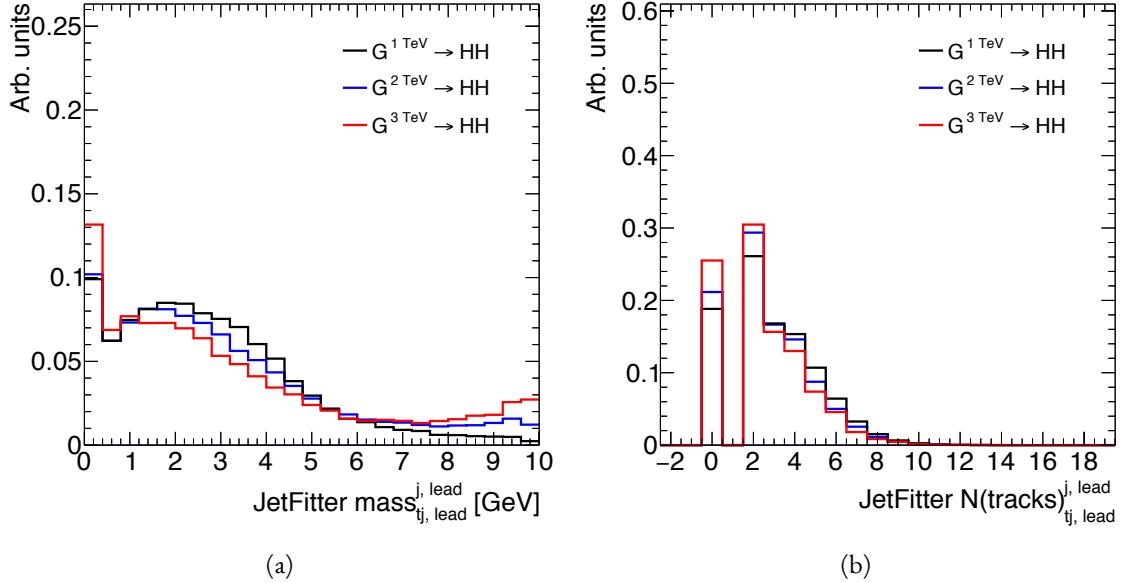


Figure A.5: Mass (a) and number of tracks (b) for vertices computed with the JetFitter algorithm. When no vertices are found, the quantities are assigned to default negative values.

²⁷⁰¹ is not tuned for tagging multiple b quarks inside one jet, the tagging efficiency could degrade. Figure A.6
²⁷⁰² shows MV₂ scores and SV1 mass for cases where there are two b quarks at truth level within the radius of

the leading track jet compared to cases where there is only one true b ¹. This figure suggests that the presence

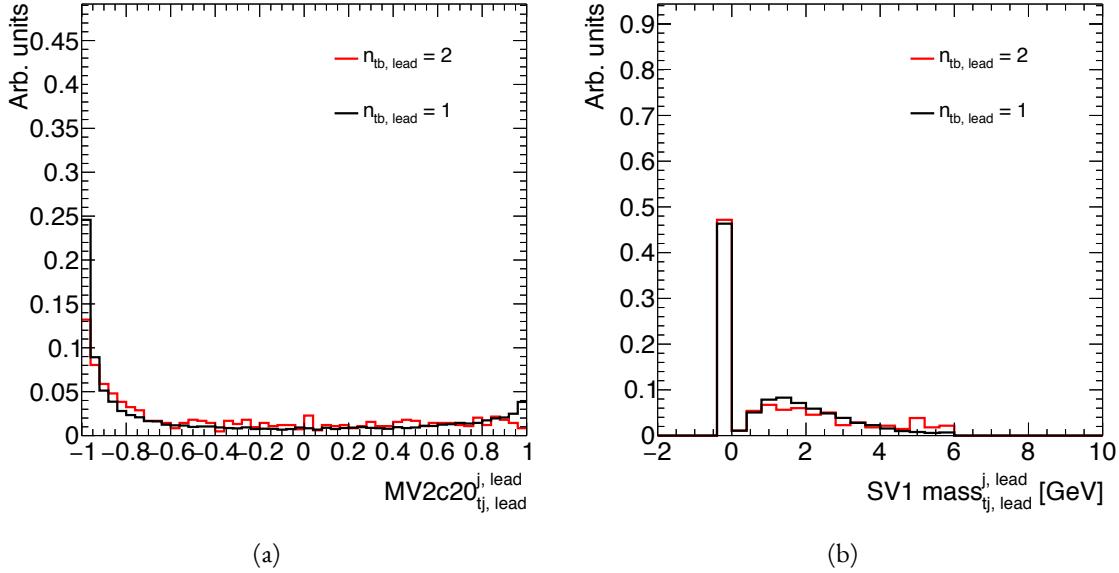


Figure A.6: MV₂c20 score (a) and SV1 mass (b) for leading track jets with two truth b quarks ($n_{tb,\text{lead}} = 2$) compared to those with only one truth b ($n_{tb,\text{lead}} = 1$).

of two b -quarks inside the leading jet is not the cause of the degradation in efficiency. There is a change in the shape of the MV₂ score distribution, but it is not nearly as pronounced as that seen in A.2 at higher masses. Additionally, the fraction of jets with no secondary vertex found is nearly identical in the track jets with two truth b -quarks.

A.3 CHANGES IN TRACK QUALITY AT HIGH p_T

Another hypothesis for the degradation of the b -tagging efficiency is a decrease in track quality for high p_T b jets. One way to check the overall quality of the tracking inside the jet is to investigate quantities related to the leading track inside of the track jet. Figure A.7 shows the fit χ^2/n_{DOF} and number of hits in the pixel detector for the leading track of the leading track jet. In both cases, the figure shows that in higher mass samples, the quality of the leading track inside of the track jet degrades substantially. The fit quality is lessened and the tracks have less hits in the pixel detector. This is likely due to the fact that at higher p_T ,

¹When two truth b quarks are required in the leading jet, the subleading jet is required to have zero. When one is required for the leading, one is also required for the subleading.

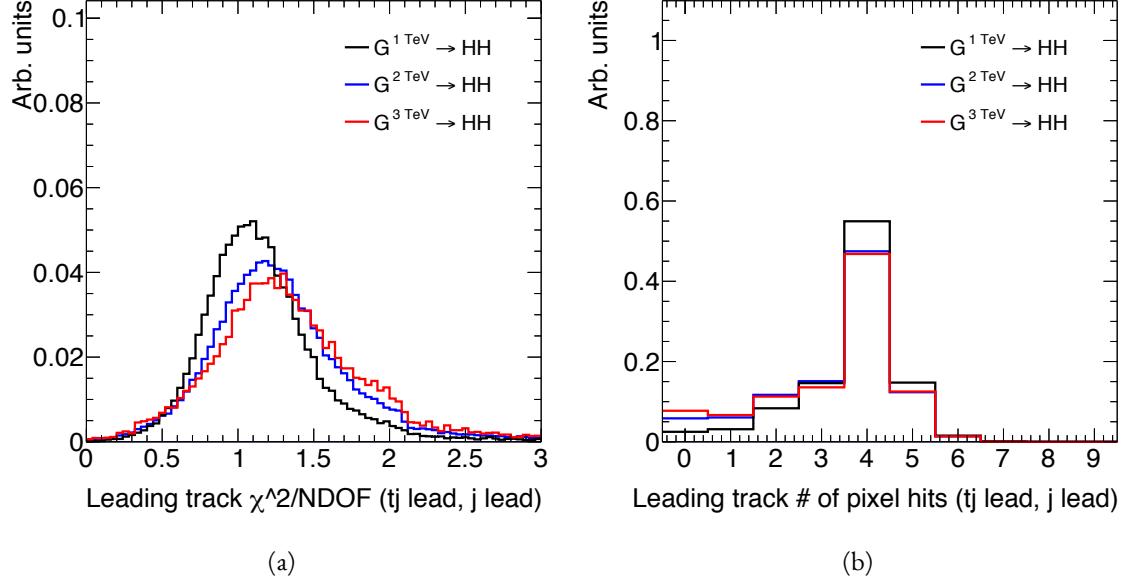


Figure A.7: Track fit χ^2/nDOF (a) and number of pixel detector hits (b) for the leading track of the leading track jet in different mass RSG $c = 1$ samples

the B -hadron will sometimes live long enough to miss the IBL and first pixel layer, thus decreasing the number of hits on the track.

To check whether this is the cause for the shift in the MV_2 score and the higher difficulty in reconstructing secondary vertices, jets whose leading track have at least four pixel hits are compared with those whose tracks have less than four pixel hits. The results for the MV_2 score and SV_1 mass are shown in figure A.8. Track jets where the leading track does not have at least four pixel hits are more likely to not have a secondary vertex reconstructed. Additionally, their $\text{MV}_{2\text{c}20}$ score is shifted more significantly to background-like values. This seems to confirm the hypothesis that degrading track quality is responsible for the lowered b -tagging efficiency at high p_{T} .

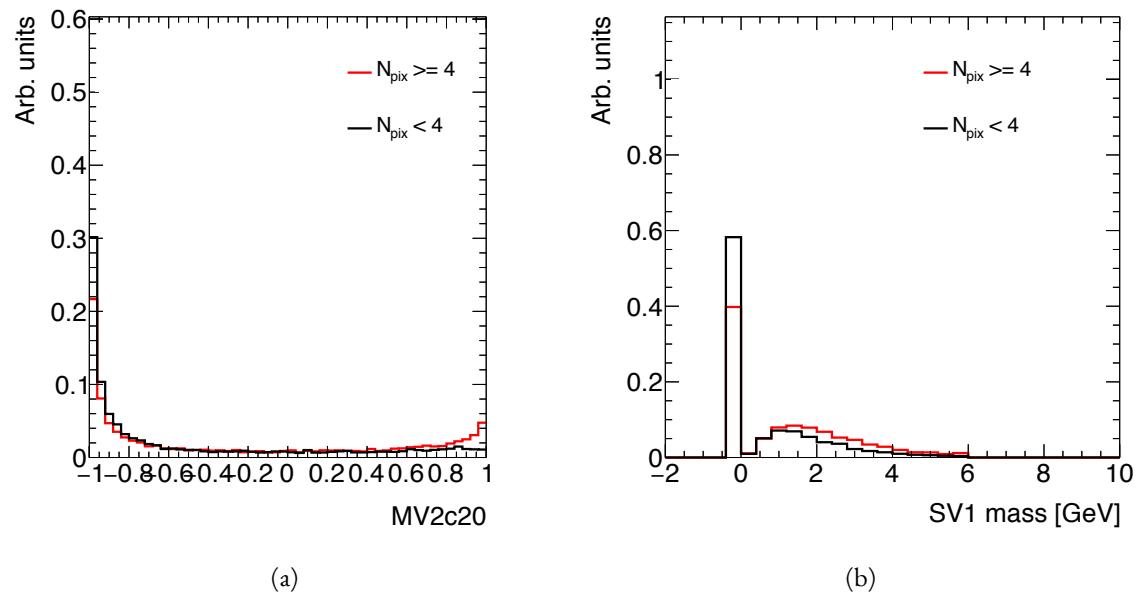


Figure A.8: MV₂c₂₀ score (a) and SV₁ mass (b) for leading track jets whose leading track jet has at least four pixel hits ($N_{\text{pix}} \geq 4$) compared to those which do not ($N_{\text{pix}} < 4$).

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