

¹ Observation of the Higgs boson in the WW^*
² channel and search for Higgs boson pair
³ production in the $b\bar{b}b\bar{b}$ channel with the
⁴ ATLAS detector

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20 **Observation of the Higgs boson in the WW^* channel and search
 21 for Higgs boson pair production in the $b\bar{b}b\bar{b}$ channel with the
 22 ATLAS detector**

23 ABSTRACT

24 This dissertation presents two studies: the observation and measurement of the Higgs boson in the
 25 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 8$ TeV and a search for Higgs pair pro-
 26 duction in the $HH \rightarrow b\bar{b}b\bar{b}$ channel at $\sqrt{s} = 13$ TeV with the ATLAS detector in pp collisions at the
 27 Large Hadron Collider.

28 First, the discovery of a particle consistent with the Higgs in 4.8 fb^{-1} at $\sqrt{s} = 7$ TeV and 5.8 fb^{-1} at
 29 $\sqrt{s} = 8$ TeV is discussed. Then, the measurement of the Higgs boson signal strength and cross section
 30 in both the gluon fusion and vector boson fusion (VBF) production modes using 20.3 fb^{-1} of $\sqrt{s} =$
 31 8 TeV data combined with 4.8 fb^{-1} of 7 TeV data is shown. The combined signal strength is measured
 32 to be $\mu = 1.09^{+0.23}_{-0.21}$. The total observed significance of the $H \rightarrow WW^*$ process is observed to be
 33 6.1σ (with 5.8σ expected). Advanced methods for background reduction and estimation, particularly in
 34 same-flavor lepton final states, are shown. The VBF signal strength is measured to be $\mu_{\text{VBF}} = 1.27^{+0.53}_{-0.45}$
 35 with an observed significance of 3.2σ (with 2.7σ expected). In the VBF channel, a selection requirement
 36 based method, the precursor to the final multivariate technique used for the result, is detailed.

37 Finally, a search for Higgs pair production in the $b\bar{b}b\bar{b}$ final state with 3.2 fb^{-1} at $\sqrt{s} = 13$ TeV
 38 is presented. A particular focus is placed on a tailored signal region for resonant production of Higgs
 39 pairs at high masses, utilizing novel techniques in object reconstruction to increase signal acceptance in
 40 boosted final state topologies. No significant excesses are observed, and upper limits on cross sections are
 41 placed for spin-2 Randall Sundrum gravitons (RSG) and narrow spin-0 resonances. The cross section of
 42 $\sigma(pp \rightarrow G_{\text{KK}}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $k/\bar{M}_{\text{Pl}} = 1$ is constrained to be less than 70 fb for masses in the
 43 range $600 < m_{G_{\text{KK}}^*} < 3000 \text{ GeV}$. For the RSG model with $k/\bar{M}_{\text{Pl}} = 2$, cross sections limits between
 44 40 fb and 200 fb are set for the mass range of $500 < m_{G_{\text{KK}}^*} < 3000 \text{ GeV}$. The cross section upper
 45 limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges from 30 to 300 fb in the mass range of $500 < m_H <$
 46 3000 GeV .

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0

481

482

Introduction

483 The Higgs boson is often described as one of the cornerstones of the Standard Model. Since the con-
484 ception of the Higgs mechanism as the source of electroweak symmetry breaking in the early 1960s,
485 countless collider experiments have searched for this elusive particle. This dissertation presents multi-
486 ple studies of the Higgs boson with the ATLAS detector at the Large Hadron Collider (LHC).

487 One of the first priorities of the early LHC was the search for the Higgs boson. This search was first
488 tackled in three main channels: $H \rightarrow \gamma\gamma$, $H \rightarrow ZZ^*$, and $H \rightarrow WW^*$. Each channel has its own
489 merits, but the WW^* is particularly suited to searching over a wide range of masses. The $H \rightarrow WW$
490 branching ratio is large and it is the primary decay channel above the $2m_W$ mass threshold.

491 While the rate of events produced in $H \rightarrow WW^*$ is large, the channel poses some challenges. First,
492 the most common mode of study for this channel is $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$. With neutrinos in the
493 final state, it is not possible to fully reconstruct the invariant mass of the parent Higgs like the $\gamma\gamma$ and
494 $ZZ \rightarrow 4\ell$ channels. Second, the final state topology is mimicked by a wide variety of backgrounds that

495 need to be properly estimated. This means tailored selection requirements for background reduction
496 and robust background estimation techniques must both be developed.

497 In 2012, the ATLAS and CMS experiments announced the discovery of a new particle consistent with
498 the Higgs boson [22, 34]. In ATLAS, this discovery was made with 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$
499 and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. The $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis played an important role in this
500 discovery. After the discovery, measurement of the properties of the newly discovered particle and con-
501 firmation of its consistency with the Standard Model Higgs were the main priorities. The WW^* chan-
502 nel is also uniquely suited to these types of measurements. Because of its good rate, it offers some of the
503 best cross section measurements available among the various Higgs decay modes. It is also suited for
504 measurement of multiple Higgs production modes, like the vector boson fusion (VBF) mode, where in-
505 coming quarks radiate W/Z bosons which fuse to make a Higgs. In VBF production with the WW^*
506 decay channel, the coupling of the Higgs to the W boson is present in both the production and decay
507 which allows for more precise measurements of this coupling than other channels which rely on gluon
508 fusion production (where gluons couple to the Higgs through a top loop in the production). The mea-
509 surement of VBF carries the additional challenge that its cross section is an order of magnitude smaller
510 than that of gluon fusion, meaning that the large branching ratio to WW^* offers an additional advan-
511 tage in isolating this production mode. In the final ATLAS Run 1 results, combining 20.3 fb^{-1} taken at
512 $\sqrt{s} = 8 \text{ TeV}$ with the 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$, the WW^* channel achieved its first observa-
513 tion of VBF production of the Higgs.

514 After Run 1 of the LHC, with the existence of the Higgs now firmly established, the focus shifted to
515 searches for physics beyond the Standard Model. In particular, searches for high mass resonances benefit
516 from the LHC's increase to $\sqrt{s} = 13 \text{ TeV}$ in Run 2. The newly discovered Higgs can be used as a tool
517 in these searches. Higgs pair production in the Standard Model has a low cross section that requires large
518 datasets (on the order of the LHC's lifetime) for full measurement. However, new physics can modify
519 this cross section, especially new resonances which decay to two Higgs bosons. A search for Higgs pair
520 production in the $HH \rightarrow b\bar{b}b\bar{b}$ final state was performed with 3.2 fb^{-1} collected with ATLAS at $\sqrt{s} =$
521 13 TeV in 2015.

522 This dissertation begins by discussing the discovery of the Higgs and the role of the $H \rightarrow WW^* \rightarrow$
523 $\ell\nu\ell\nu$ channel. It then discusses the first observation of the VBF production mode in $H \rightarrow WW^* \rightarrow$
524 $\ell\nu\ell\nu$ with the full ATLAS Run 1 dataset, as well as the final combined Run 1 measurements from this
525 channel. Finally, it presents a search for Higgs pair production in the $HH \rightarrow b\bar{b}b\bar{b}$ channel. It is orga-
526 nized into four parts.

527 Part 1 presents the theoretical and experimental background required for the subsequent parts. Chap-
528 ter 1 gives an overview of Higgs physics, particularly single and double Higgs production in the Standard
529 Model and beyond. Chapter 2 presents details regarding the Large Hadron Collider and the ATLAS
530 experiment. The evolution of machine conditions, descriptions of the ATLAS sub-detectors, and an
531 overview of object reconstruction in ATLAS are all shown. A brief interlude on the ATLAS Muon New
532 Small Wheel upgrade is also given, as this upgrade has been a focus of my graduate work and will have
533 important impact on ATLAS' ability to study the Higgs at the High Luminosity LHC.

534 Part 2 discusses the observation and measurement of the Higgs in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
535 channel in the ATLAS Run 1 dataset at $\sqrt{s} = 7$ and 8 TeV. Because I worked in this channel from
536 before the discovery through to the final analysis of the Run 1 dataset, Part 2 is organized in such a way
537 to allow easy presentation of multiple analyses on different subsets of the full Run 1 dataset. Chapter 3
538 presents a general overview of the $H \rightarrow WW^*$ analysis strategy and defines many of the variables and
539 common elements used in the rest of Part 2. Chapter 4 presents the discovery of the Higgs boson, fo-
540 cusing on the role of the WW^* channel in this discovery. Chapter 5 presents the first observation of the
541 VBF production mode of the Higgs in the WW^* channel, a study which was done on the full Run 1
542 ATLAS dataset. In this chapter, the focus is mainly on the selection cut-based VBF analysis. The cut-
543 based analysis was an important first step to the final VBF result which used a Boosted Decision Tree
544 (BDT). Where appropriate, connections between the cut-based and BDT analyses are shown and their
545 compatibility is discussed. Finally, the VBF analysis was an important input into the combined Run 1
546 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ result, which used both the gluon fusion and VBF channels in a combined
547 fit to infer properties of the Higgs, including its couplings to the gauge bosons and its production cross
548 section. This is the topic of Chapter 6.

549 Part 3 presents a search for Higgs pair production in the $HH \rightarrow b\bar{b}b\bar{b}$ channel. Chapter 7 presents
550 an overview of this search in the boosted regime, where the Higgs pairs are the result of the decay of a
551 heavy resonance. Chapter 8 shows the combined results between the boosted regime and the resolved
552 regime, which is sensitive to lower mass resonances and non-resonant Higgs pair production. Finally,
553 Part 4 presents a conclusion and brief outlook of future Higgs physics with ATLAS.

Part I

Theoretical and Experimental Background

1

556

557

The Physics of the Higgs Boson

558 This chapter presents an overview of the Standard Model of Particle Physics and in particular the physics
559 of the Higgs boson. First, a brief overview of the Standard Model and its history are presented. Then, a
560 description of the Higgs mechanism of electroweak symmetry breaking is given. Next, the physics of sin-
561 ggle Higgs boson production and decay is described. The Standard Model also allows for production of
562 two Higgs bosons and this is detailed as well. Finally, di-Higgs production in two beyond the Standard
563 Model (BSM) theories - Randall-Sundrum gravitons (RSG) and Two Higgs Doublet Models (2HDM) -
564 is shown.

565 **I.I THE STANDARD MODEL OF PARTICLE PHYSICS**

566 The Standard Model (SM) of Particle Physics is a quantum field theory describing the fundamental
567 particles of nature and the forces that govern their interactions. Several comprehensive treatments of
568 the SM already exist in the literature[1, 29, 35–38] and this section will not rehash those. Rather, this

569 section presents a brief overview of the SM particles and forces in order to define them for subsequent
 570 discussions.

571 The Standard Model consists of two primary categories of fundamental particles: fermions (spin 1/2
 572 particles) and bosons (integer spin particles). The SM also describes three forces: electromagnetism, the
 573 weak nuclear force, and the strong nuclear force. Gravity is not included in the theory and is largely irrele-
 574 vant at the scales currently probed by collider experiments. Within the fermions, there are both quarks
 575 (which interact via all three forces) and the leptons. The charged leptons interact via electromagnetic and
 576 weak interactions, while neutrinos (neutral leptons) interact only via the weak force. Within the bosons,
 577 there are the W^\pm and Z bosons (the mediators of the weak force), the gluon (g , the mediator of the
 578 strong force), and the photon (γ), the mediator of the electromagnetic force. Finally, there is the Higgs
 579 boson, a fundamental spin-0 particle resulting from the Higgs mechanism of electroweak symmetry
 580 breaking. Figure 1.1 summarizes the fermions and bosons of the SM.

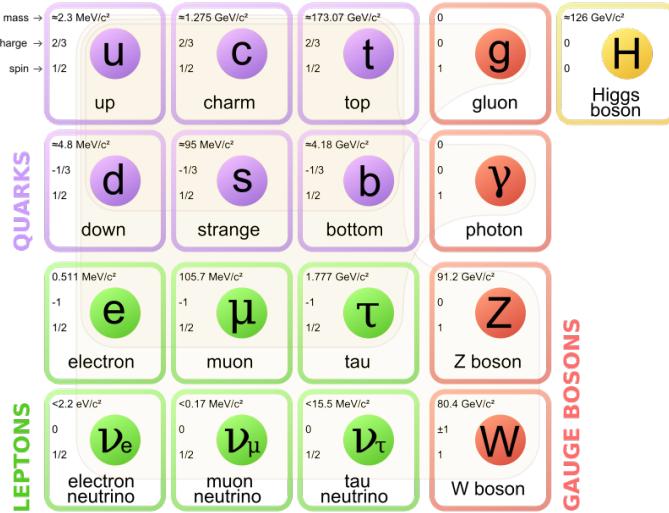


Figure 1.1: The particles of the Standard Model and their properties[1].

581 The Standard Model coalesced into a unified theoretical framework in the 1960s through the work of
 582 Glashow, Weinberg, Salam, and others on the theory of electroweak interactions[39–42]. This theory
 583 characterized both the electromagnetic and weak interactions as unified under a single gauge symme-
 584 try group, namely $SU(2) \times U(1)$. At low enough energy scales (on the order of the W and Z masses,
 585 the electroweak symmetry is broken, as evidenced by the fact that the weak bosons have mass while the

586 photon does not. The discovery of the Higgs boson in 2012 confirmed the Higgs mechanism as the most
 587 likely candidate for this electroweak symmetry breaking[22, 34]. The electroweak theory is then com-
 588 bined with the theory of quantum chromodynamics (which models the strong sector as a non-Abelian
 589 $SU(3)$ gauge group) to form the complete SM[43].

590 1.2 ELECTROWEAK SYMMETRY BREAKING AND THE HIGGS

591 In the Standard Model Lagrangian, it is difficult to include mass terms for the W and Z bosons without
 592 breaking the fundamental gauge symmetry of the Lagrangian. A traditional mass term does not preserve
 593 the $SU(2) \times U(1)$ symmetry. Additionally, scattering of massive W and Z bosons violate unitarity
 594 and these diagrams diverge at high energy scales. In the 1960s, Higgs, Brout, Englert, Guralnik, Kibble,
 595 and Hagen developed a mechanism for spontaneous symmetry breaking via the addition of a complex
 596 scalar doublet to the SM. Three of the four real degrees of freedom of this complex field would go to the
 597 longitudinal modes of the W^\pm and Z , thus allowing them to have mass[44–47]. The remaining degree
 598 of freedom would manifest as an additional scalar, known now as the Higgs boson.

599 The mechanism works by introducing a Lagrangian for the newly introduced field that still respects
 600 the symmetry of the Standard Model inherently, but with a minimum at a non-zero vacuum expecta-
 601 tion value for the field. In this minimum of the potential, the electroweak symmetry is broken. Specifi-
 602 cally, consider a complex scalar doublet Φ with four degrees of freedom, as shown in equation 1.1.

$$\Phi = \begin{pmatrix} \phi^+ \\ \phi^0 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi_1^+ + i\phi_2^+ \\ \phi_1^0 + i\phi_2^0 \end{pmatrix} \quad (1.1)$$

603 The minimal potential of a self-interacting Higgs that still respects the SM symmetry is given in equa-
 604 tion 1.2.

$$V(\Phi) = \mu^2 \Phi^\dagger \Phi + \lambda (\Phi^\dagger \Phi)^2 \quad (1.2)$$

605 If the μ^2 term of this potential is positive, then the potential has a minimum at $\Phi = 0$ and the SM

⁶⁰⁶ symmetry is preserved. However, if instead $\mu^2 < 0$, then the minimum is at a finite value of Φ , namely

$$\Phi_{\min} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v \end{pmatrix} \quad (1.3)$$

⁶⁰⁷ where $v = \sqrt{\mu^2/\lambda}$. Because this is the location of the minimum, it corresponds to the vacuum expecta-
⁶⁰⁸ tion value for the field ($\langle \Phi \rangle = \Phi_{\min}$). The excitations of the Higgs can then be parameterized as

$$\Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 \\ v + H \end{pmatrix} \quad (1.4)$$

⁶⁰⁹ The full scalar Lagrangian, including the kinetic term, is then given as

$$\mathcal{L}_s = (D^\mu \Phi)^\dagger (D_\mu \Phi) - V(\Phi) \quad (1.5)$$

⁶¹⁰ where the covariant derivative is defined as

$$D_\mu = \partial_\mu + \frac{ig}{2} \tau^a W_\mu^a + ig' Y B_\mu \quad (1.6)$$

⁶¹¹ and W^1, W^2, W^3 and B are the $SU(2)$ and $U(1)$ gauge fields of the electroweak theory, respectively. g
⁶¹² and g' are the corresponding coupling constants. With the scalar Lagrangian in place, the physical gauge
⁶¹³ fields can then be written as

$$W_\mu^\pm = \frac{1}{\sqrt{2}} (W_\mu^1 \mp i W_\mu^2) \quad (1.7)$$

⁶¹⁴

$$Z_\mu = \frac{-g' B_\mu + g W_\mu^3}{\sqrt{g^2 + g'^2}} \quad (1.8)$$

⁶¹⁵

$$A_\mu = \frac{g B_\mu + g' W_\mu^3}{\sqrt{g^2 + g'^2}} \quad (1.9)$$

⁶¹⁶ Equation 1.7 corresponds to the charged W^+ and W^- bosons, equation 1.8 corresponds to the neutral

₆₁₇ Z boson, and equation 1.9 corresponds to the neutral photon. The masses of the particles also arise from
₆₁₈ the Lagrangian. The photon has zero mass, while the masses of the W and Z bosons are given in equa-
₆₁₉ tion 1.10.

$$\begin{aligned} M_W^2 &= \frac{1}{4}g^2v^2 \\ M_Z^2 &= \frac{1}{4}(g^2 + g'^2)v^2 \end{aligned} \quad (1.10)$$

₆₂₀ The fermion masses also arise through a coupling with the Higgs via the Yukawa interaction (for a de-
₆₂₁ tailed description, see [38]). In this case the coupling between the Higgs and the fermions goes as

$$g_{Hf\bar{f}} = \frac{m_f}{v} \quad (1.11)$$

₆₂₂ The full Lagrangian of Higgs interactions can be written as

$$\mathcal{L}_{\text{Higgs}} = -g_{Hf\bar{f}}\bar{f}fH + \frac{g_{HHH}}{6}H^3 + \frac{g_{HHHH}}{24}H^4 + \delta_V V_\mu V^\mu \left(g_{HVV}H + \frac{g_{HHVV}}{2}H^2 \right) \quad (1.12)$$

₆₂₃ with

$$\begin{aligned} g_{HVV} &= \frac{2m_V^2}{v} & g_{HHVV} &= \frac{2m_V^2}{v^2} \\ g_{HHH} &= \frac{3m_H^2}{v} & g_{HHHH} &= \frac{3m_H^2}{v^2} \end{aligned} \quad (1.13)$$

₆₂₄ Here, V refers to the W^\pm and Z , and $\delta_W = 1$ while $\delta_Z = 1/2$. Phenomenologically, there are a few
₆₂₅ features of this Lagrangian that are useful to note. First, note that the Higgs mass is a free parameter
₆₂₆ of the theory that must be determined experimentally. Second, note that the coupling of the Higgs to
₆₂₇ the vector bosons and fermions scales with the masses of these particles, a fact that is important when
₆₂₈ considering both the production and decays of the particle. Also note that the branching ratio of the
₆₂₉ Higgs to W bosons will be twice that of the branching ratio to Z if the Higgs mass is large enough to
₆₃₀ produce the particles on shell because of the extra symmetry factor associated with the W coupling.
₆₃₁ Finally, note the presence of the cubic and quartic Higgs self interaction terms, which can lead to final
₆₃₂ states with multiple Higgs bosons produced.

633 1.3 HIGGS BOSON PRODUCTION AND DECAY

634 This section discusses the properties of Higgs production and decay mechanisms. The details presented
635 here will focus on the properties of a 125 GeV Higgs boson, as this is the mass closest to that of the
636 newly discovered Higgs.

637 1.3.1 HIGGS PRODUCTION

638 The Higgs is produced by four main production modes at the Large Hadron Collider - gluon-gluon
639 fusion (ggF), vector boson fusion (VBF), associated production with a W or Z boson, or associated
640 production with top quarks ($t\bar{t}H$). Figure 1.2 shows the Feynman diagrams for these four modes.

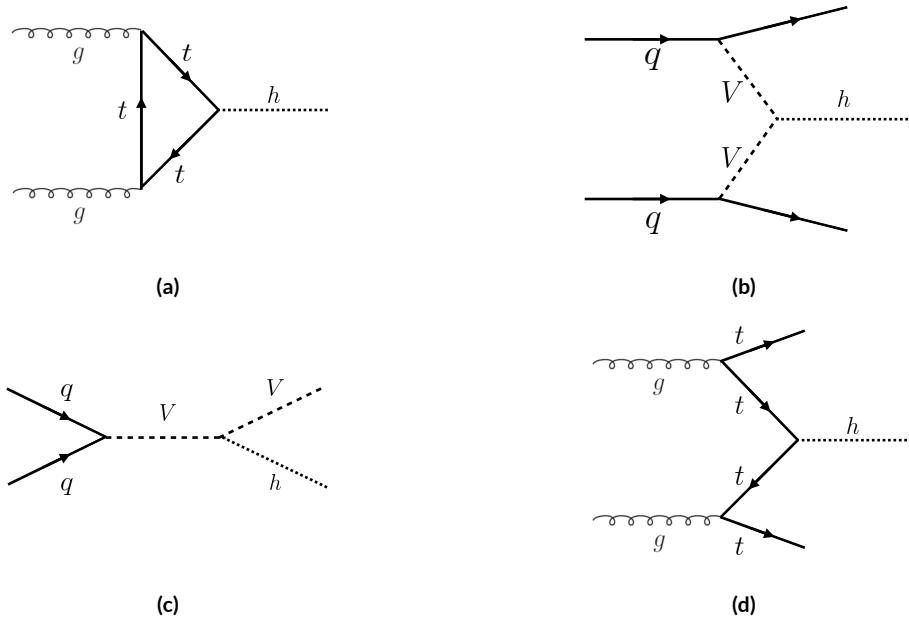


Figure 1.2: The four most common Higgs boson production modes at the LHC: (a) gluon-gluon fusion, (b) vector boson fusion, (c) $W/Z + H$ production, (d) $t\bar{t}H$ production

641 In gluon-gluon fusion, gluons from the incoming protons fuse via a top-quark loop to produce a
642 Higgs. The top quark is the dominant contribution in the loop due to its heavy mass and the fact that
643 the Higgs-fermion coupling constant scales with fermion mass. In vector boson fusion, the incoming
644 quarks each radiate a W or Z boson which fuse to produce the Higgs. This production mode results in
645 a final state with a Higgs boson and two additional jets which tend to be forward because they carry the

646 longitudinal momentum of the incoming partons. The Higgs can also be produced in association with a
 647 W or Z boson. The W/Z is produced normally and then radiates a Higgs (this mode is also sometimes
 648 known as “Higgs-strahlung”). Finally, the Higgs can be produced in association with two top quarks.
 649 Each incoming gluon splits into a $t\bar{t}$ pair, and one of the top pairs combines to create a Higgs.

650 Figure 1.3 shows the production cross section for a 125 GeV Higgs boson in each of these modes at a
 651 pp collider as a function of center of mass energy.

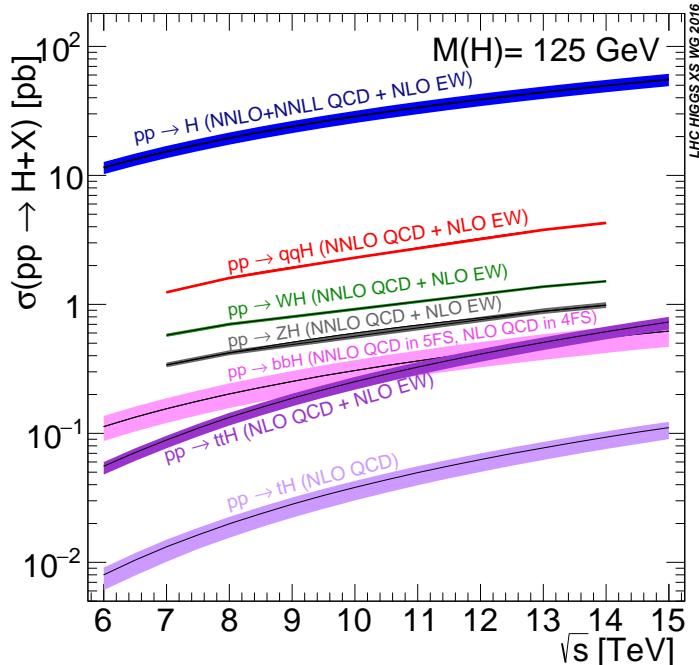


Figure 1.3: Higgs production cross sections as a function of center of mass energy (\sqrt{s}) at a pp collider[2].

652 In figure 1.3, note that gluon fusion has the largest cross section, while VBF is the second largest at
 653 approximately a factor of 10 smaller. The figure also includes the less commonly studied $b\bar{b}H$ and tH
 654 modes. The $b\bar{b}H$ and tH modes are not studied as commonly as $t\bar{t}H$ due to the larger background con-
 655 tributions and lower cross sections, respectively. At $\sqrt{s} = 8$ TeV, ggF production of a 125 GeV Higgs
 656 has a cross section of 19.47 pb, while VBF has a cross section of 1.601 pb[2]. The cross sections of all
 657 of the main Higgs production modes at this center of mass energy, as well as their uncertainties from
 658 varying the renormalization and factorization scales and PDFs, are summarized in table 1.1 for a 125 GeV
 659 Higgs.

Production mode	σ (pb)	QCD scale uncert. (%)	PDF + α_s uncert. (%)
Gluon fusion	19.47	+7.3 / - 8.0	3.1
Vector boson fusion	1.601	+0.3 / - 0.2	2.2
WH	0.7026	+0.6 / - 0.9	2.0
ZH	0.4208	+2.9 / - 2.4	1.7
$t\bar{t}H$	0.1330	+4.1 / - 9.2	4.3
$b\bar{b}H$	0.2021	+20.7 / - 22.3	
tH (t -channel)	0.01869	+7.3 / - 16.5	4.6
tH (s -channel)	1.214×10^{-3}	+2.8 / - 2.4	2.8

Table 1.1: Production cross sections for a 125 GeV Higgs boson at $\sqrt{s} = 8$ TeV with scale and PDF uncertainties [2].

660 1.3.2 HIGGS BRANCHING RATIOS

661 The fact that the Higgs couples more strongly to more massive particles is crucial for understanding its
 662 branching ratios. The width for Higgs decays to fermions is given in equation 1.14 [29].

$$\Gamma(H \rightarrow f\bar{f}) = \frac{N_c \sqrt{2} G_F m_f^2 m_H}{8\pi} \quad (1.14)$$

663 In this case, N_c is the number of colors, G_F is the Fermi constant, m_f is the mass of the fermion, and
 664 m_H is the mass of the Higgs. Note that the width scales with the square of the fermion mass. (This also
 665 assumes that the Higgs mass is large enough to decay with both the fermions on shell.)

666 The decay width to WW is given in equation 1.15 [29].

$$\Gamma(H \rightarrow W^+W^-) = \frac{\sqrt{2} G_F M_W^2 m_H}{16\pi} \frac{\sqrt{1-x_W}}{x_W} (3x_W^2 - 4x_W + 4) \quad (1.15)$$

667 where m_W is the mass of the W and $x_W = 4M_W^2/m_H^2$. To get the branching ratio to ZZ , the
 668 equation is divided by 2 to account for identical particles in the final state, and x_W is replaced with
 669 $x_Z = 4M_Z^2/m_H^2$. This is shown in equation 1.16 [29].

$$\Gamma(H \rightarrow ZZ) = \frac{\sqrt{2} G_F M_Z^2 m_H}{32\pi} \frac{\sqrt{1-x_Z}}{x_Z} (3x_Z^2 - 4x_Z + 4) \quad (1.16)$$

670 These formulas can also be visualized as a function of Higgs mass. Figure 1.4 shows the branching

671 ratios as a function of the Higgs mass.

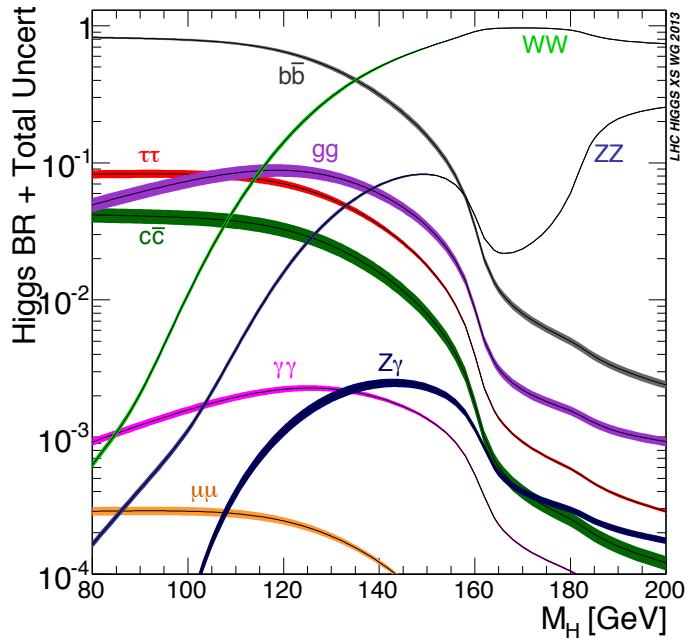


Figure 1.4: Higgs boson branching ratios as a function of m_H [2].

672 There are a few interesting features to note in this figure. First, note that at high Higgs masses, once
 673 on-shell production of both W and Z bosons is possible, these two decays are the dominant ones due to
 674 the large masses of the W/Z . Also note that the branching ratio to W s is twice that of Z s at these large
 675 masses due to the δ_V symmetry factor noted previously. At 125 GeV, the Higgs is accessible through
 676 many different decay modes. The largest branching ratio is the decay $H \rightarrow b\bar{b}$ at 58.24% [2]. This
 677 branching is larger than the WW/ZZ decays because one of the two bosons must be produced off-
 678 shell for $m_H = 125$ GeV. The second largest branching ratio is to WW^* at 21.37 % (before taking
 679 into account the branching ratios of the W). Table 1.2 summarizes the branching ratios for a 125 GeV
 680 Higgs. Note that there is in fact a Higgs branching ratio to $\gamma\gamma$ even though photons are massless. This
 681 decay happens through a loop (the largest contributions to the loop are top and W) which suppresses
 682 the branching ratio.

683 Note that the branching ratios alone do not tell the full story of which Higgs channels are the most
 684 sensitive. For example, a $H \rightarrow b\bar{b}$ search in gluon fusion production is incredibly difficult due to the

Decay	Branching ratio (%)
$b\bar{b}$	58.24
WW^*	21.37
gg	8.187
$\tau\tau$	6.272
$c\bar{c}$	2.891
ZZ^*	2.619
$\gamma\gamma$	0.2270
$Z\gamma$	0.1533
$\mu\mu$	0.02176

Table 1.2: Branching ratios for a 125 GeV Higgs boson[2].

large QCD dijet background at the LHC. However, in associated production of the Higgs, where a W or Z gives additional final state particles that can be used to reduce background, a search for $H \rightarrow b\bar{b}$ can be sensitive. The combinations of production and decay modes that are most commonly studied are summarized in table 1.3 [29].

Decay	Inclusive (incl. ggF)	VBF	WH/ZH	$t\bar{t}H$
$H \rightarrow \gamma\gamma$	✓	✓	✓	✓
$H \rightarrow b\bar{b}$			✓	✓
$H \rightarrow \tau^+\tau^-$		✓		
$H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$	✓	✓	✓	
$H \rightarrow ZZ \rightarrow 4\ell$	✓			
$H \rightarrow Z\gamma \rightarrow \ell\ell\gamma$	very low			

Table 1.3: Possible channels for Higgs searches. Checkmarks denote the most sensitive production modes [29].

1.4 HIGGS PAIR PRODUCTION IN THE STANDARD MODEL

The Standard Model also allows for processes that produce two Higgs bosons in the final state, known as Higgs pair production or di-Higgs production. The two main production mechanisms are shown in figure 1.5.

The two diagrams in figure 1.5 interfere destructively with one another, resulting in a low overall cross section for di-Higgs production at the LHC. Nevertheless, Higgs pair production is quite interesting to study because it gives direct access to the λ parameter of the Higgs potential, also known as the Higgs self

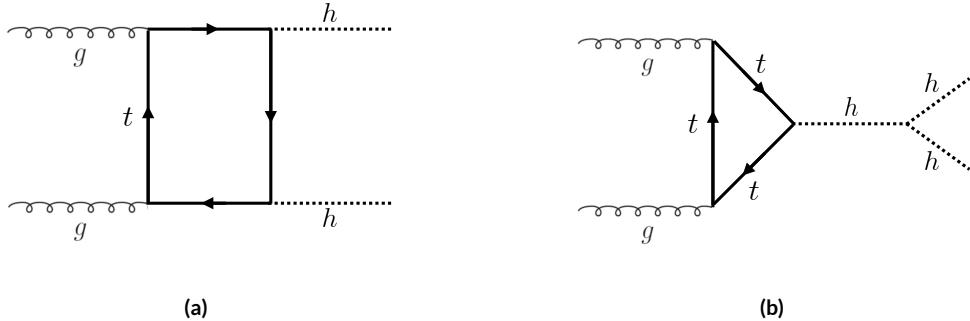


Figure 1.5: The two leading diagrams for Standard Model di-Higgs production at the LHC: (a) box diagram, (b) Higgs self coupling

coupling. The diagram in figure 1.5(b) is sensitive to this coupling through the triple Higgs vertex.

One can substitute the gluon fusion production of diagram 1.5(b) with any of the other production modes previously discussed. These other modes do not suffer from interference with the box diagram in figure 1.5(a) due to the presence of additional particles in the final state. They still have a lower cross section than the gluon fusion mode, however. The cross sections for di-Higgs production in the different modes, as well as their uncertainties, are shown in table 1.4 [30]. These are shown for $\sqrt{s} = 14$ TeV as the higher center of mass energy is more sensitive to this process. Note that the scale of cross section quoted is now in fb rather than pb.

Production mode	σ (fb)	Total uncert. (%)
Gluon fusion	33.89	+37.2 / - 27.8
Vector boson fusion	2.01	+7.6 / - 5.1
$W H H$	0.57	+3.7 / - 3.3
$Z H H$	0.42	+7.0 / - 5.5
$t\bar{t}H$	1.02	-

Table 1.4: Production cross sections for pair production of a 125 GeV Higgs boson at $\sqrt{s} = 14$ TeV with total uncertainty [30]. The uncertainties include QCD scale and PDF variations as well as uncertainties on α_S .

703

1.5 HIGGS PAIR PRODUCTION IN THEORIES BEYOND THE STANDARD MODEL

The Standard Model Higgs pair production cross section is rather small, and datasets on the scale of the full lifetime of the LHC will be required to obtain sensitive measurements of the Higgs self-coupling.

707 However, the discovery of the Higgs also gives particle physicists a new tool that can be exploited in the
 708 search for new physics beyond the Standard Model. In particular, Higgs pair production is a promising
 709 channel in the search for new physics. The cross section for di-Higgs production can be altered through
 710 both resonant and non-resonant production of Higgs pairs. In non-resonant production, di-Higgs pro-
 711 duction vertices can arise from the presence of a new strong sector and additional colored particles[48–
 712 50]. Figure 1.6 shows examples of the types of vertices that can arise. In the resonant case, new heavy
 713 particle can decay to Higgs pairs. Such new particles can include heavy Higgs bosons arising in two
 714 Higgs doublet models (2HDM) or Higgs portal models as well as heavy gravitons in Randall-Sundrum
 715 theories[3, 5, 48, 51–55]. Figure 1.7 shows a generic diagram for a heavy resonance decaying to two Higgs
 716 bosons. In the 2HDM, X corresponds to the heavy CP-even scalar H . In the Randall-Sundrum model,
 717 X corresponds to a heavy spin-2 graviton G .

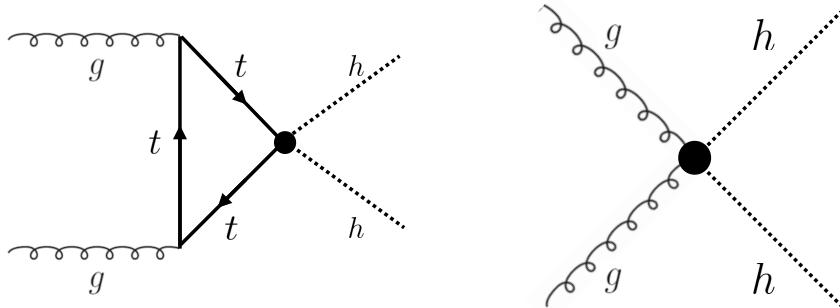


Figure 1.6: Diagrams with new vertices for non-resonant Higgs pair production arising in composite Higgs models

718 The next sections provide more detail on the phenomenology of resonant Higgs production in
 719 Randall-Sundrum and 2HDM models, as these models will later be tested in a dedicated search for reso-
 720 nant production of boosted Higgs pairs.

721 1.5.1 RANDALL-SUNDRUM GRAVITONS

722 The Randall-Sundrum model is a proposed solution to the hierarchy problem that posits a five-dimensional
 723 warped spacetime that contains two branes: one where the force of gravity is very strong and a second
 724 brane at the TeV scale corresponding to the known Standard Model sector [51]. In the theory, the
 725 branes are weakly coupled and the graviton probability function drops exponentially going from the

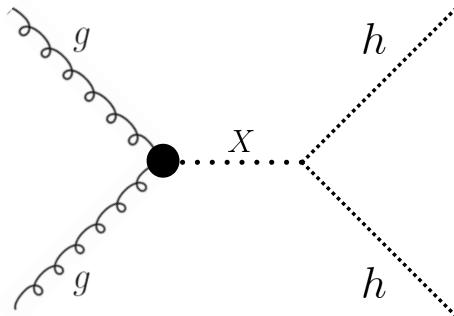


Figure 1.7: Generic Feynman diagram for resonant Higgs pair production in BSM theories

726 gravity brane to the SM brane, rendering gravity weak on the SM brane. The experimental consequence
 727 of this theory is a tower of widely spaced (in mass) Kaluza-Klein graviton resonances. In theories where
 728 the fermions are localized to the SM brane, production of gravitons from fermion pairs is suppressed
 729 and the primary mode of production is gluon fusion[3]. These gravitons have a substantial branching
 730 fraction to Higgs pairs, ranging from 6.43% for gravitons with a mass of 500 GeV to 7.66% at 3 TeV.
 731 Figure 1.8 shows the branching ratios of the spin-2 Randall Sundrum graviton (RSG) as a function of its
 732 mass. The predominant decays are to $t\bar{t}$ above the mass threshold for that channel.

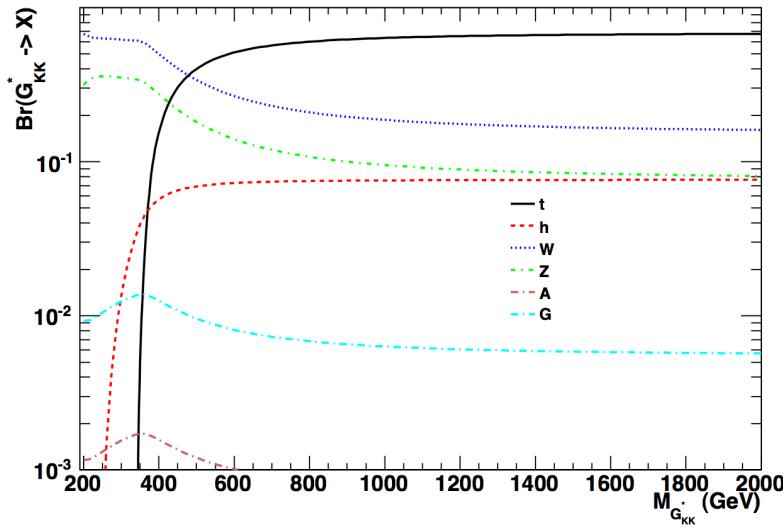


Figure 1.8: Branching ratios for a spin-2 Randall-Sundrum graviton as a function of mass computed in MadGraph with the CP3-Origins implementation [3, 4]

733 These models have two free parameters - the mass of the graviton and a curvature parameter k . Typ-

734 ically, rather than k , the theory is parameterized using $c \equiv k/\bar{M}_{\text{pl}}$, where \bar{M}_{pl} is the reduced Planck
 735 mass. The cross section for production of the RSG decreases as a function of mass and is strongly depen-
 736 dent on the gluon PDF. The increase in center of mass energy from 8 to 13 TeV in LHC Run 2 greatly
 737 increases the cross section at higher mass. Figure 1.9 shows the cross section as a function of graviton
 738 mass at $\sqrt{s} = 13$ TeV for RSG models with $c = 1.0$ and $c = 2.0$.

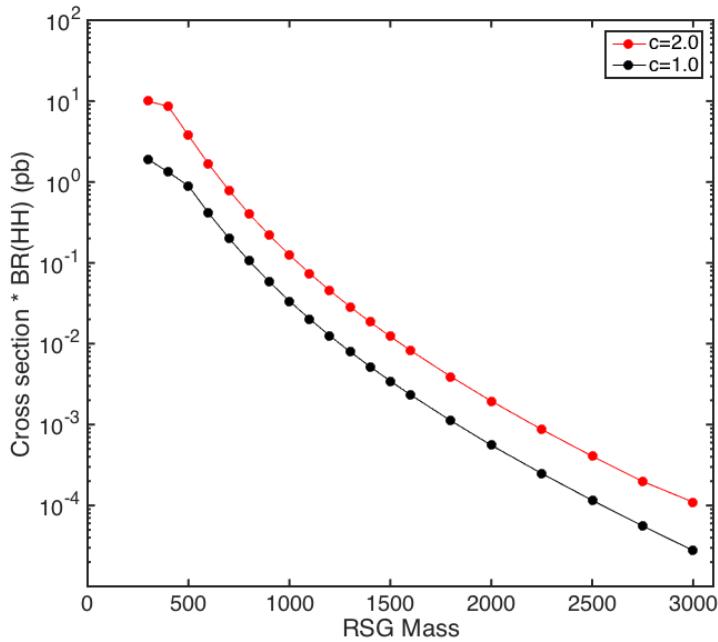


Figure 1.9: $\sigma \times \text{BR}(HH)$ for RSG as a function of mass computed in MadGraph with the CP3-Origins implementation [3, 4]

739 Another interesting feature of the theory is that the width of the graviton increases with both c and
 740 m_G . Figure 1.10 shows the graviton width for both $c = 1.0$ and $c = 2.0$ as a function of mass. In
 741 $c = 1.0$, the width starts at 8.365 GeV for a mass of 300 GeV and increases to 187.2 GeV at a mass of
 742 3 TeV. Similarly, with $c = 2.0$, the width starts at 33.46 GeV for $m_G = 300$ GeV and increases to
 743 748.8 GeV at a mass of 3 TeV.

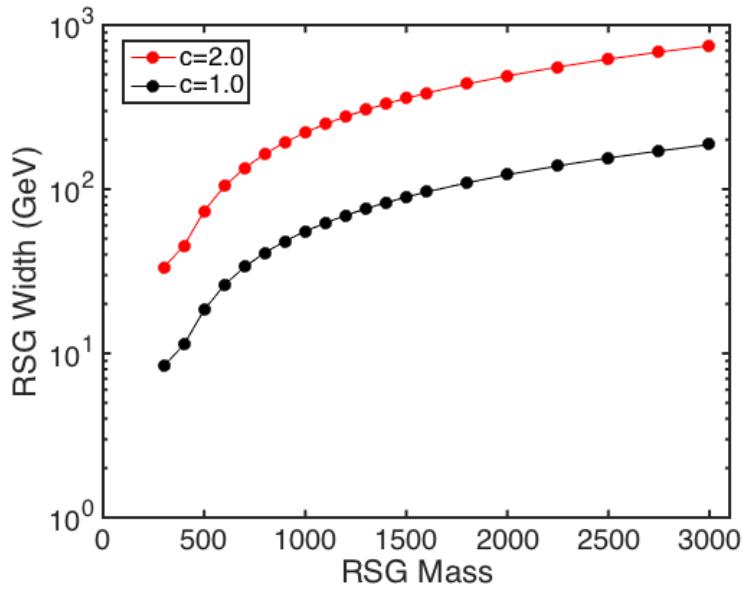


Figure 1.10: RSG width as a function of mass computed in MadGraph with the CP3-Origins implementation [3, 4]

744 1.5.2 TWO HIGGS DOUBLET MODELS

745 In Two Higgs Doublet Models (2HDM), a second complex scalar doublet is added to SM[5, 53, 54]. In
 746 this case, all four degrees of freedom in the second doublet correspond to new particles, meaning that
 747 there are five total scalars from the two Higgs doublets - h (light CP-even Higgs), H (heavy CP-even
 748 Higgs), A (heavy CP-odd Higgs), and H^\pm (charged Higgs). The model is parameterized by two main
 749 parameters. The first, $\tan \beta \equiv \frac{v_2}{v_1}$, is the ratio of the vacuum expectation values of the two Higgs dou-
 750 blets (where v_1 corresponds to the v in the SM Higgs model described above). The second parameter
 751 is α , a mixing angle between the heavy and light Higgs fields. Models are also often parameterized with
 752 $\cos(\beta - \alpha)$ rather than α directly. The limit where $\cos(\beta - \alpha) = 0$ is called the alignment limit, and
 753 it is in this limit that the light Higgs h has the same couplings as a Standard Model Higgs.

754 2HDM models are usually separated into two main types - Type I and Type II. In Type I models, the
 755 charged fermions only couple to the second Higgs doublet, leading to a fermiophobic light Higgs. In
 756 Type II models, up-type quarks couple to the first doublet while down-type quarks couple to the second
 757 doublet. One specific realization of a Type II 2HDM is the Minimal Supersymmetric Standard Model
 758 (MSSM).

759 Resonant di-Higgs production in this model can proceed through decays of the heavy CP-even Higgs
 760 $H \rightarrow hh$. The branching ratio for $H \rightarrow hh$ depends on the model type as well as the values of $\tan \beta$
 761 and $\cos(\beta - \alpha)$. Figure 1.11 shows the branching ratios as a function of the mass of the heavy scalar H for
 762 both Type I and Type II models. Depending on the type of model hh can be a substantial fraction of the
 763 decays of H .

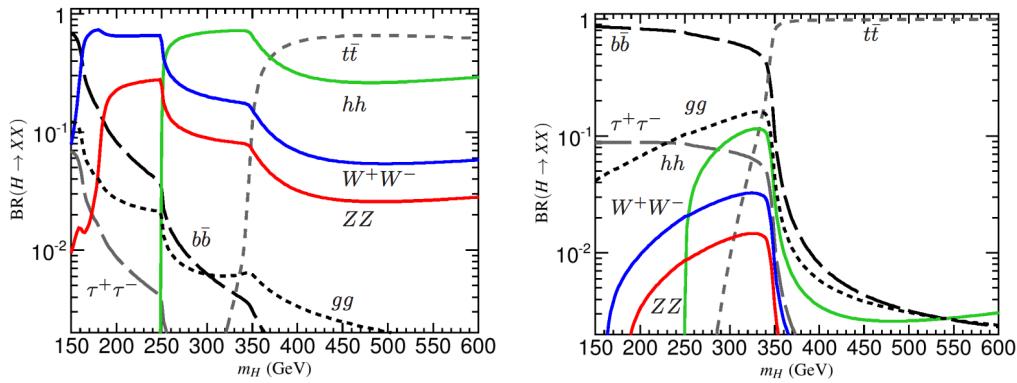


Figure 1.11: Branching ratios for heavy Higgs H in Type I (left) and Type II (right) 2HDM models with $\tan \beta = 1.5$ and $\cos(\beta - \alpha) = 0.1$ (Type I) (0.01 for Type II). [5]

764 1.6 CONCLUSION

765 Studying the Higgs sector is essential for understanding the details of how mass arises in the Standard
 766 Model and how the electroweak symmetry is broken. The discovery of the Higgs boson also opens the
 767 door for its use as a tool to search for new physics, and Higgs pair production is an ideal candidate for
 768 this study. Even if no BSM physics is found in Higgs pair production, searches for Higgs pairs will put
 769 constraints on the Higgs self coupling and thus further knowledge of the Standard Model and the details
 770 of the Higgs potential.

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The ATLAS detector and the Large Hadron Collider

774 This chapter presents an overview of the experimental systems used to conduct the measurements pre-
775 sented in this thesis. First, a brief overview of the accelerator, the Large Hadron Collider, will be given.
776 In this section, the accelerator conditions relevant to data-taking are presented as well. Next, an overview
777 of the ATLAS experiment is given. The basics of each sub-detector's role are summarized, as well as the
778 details of the datasets accumulated. Then, a brief interlude on the ATLAS Muon New Small Wheel up-
779 grade is presented. While this new detector does not have a direct impact on any of the datasets taken
780 so far, it will have an impact on future analyses and the work done on it is briefly summarized here. Fi-
781 nally, an overview of object reconstruction in ATLAS is given. While the details of all of the algorithms
782 will not be presented in detail, aspects of the reconstruction performance such as object resolutions are
783 shown as these are relevant to the two studies presented later in this thesis.

784 2.1 THE LARGE HADRON COLLIDER

785 The Large Hadron Collider (LHC) is a proton-proton collider at the CERN laboratory in Geneva,
786 Switzerland[56]. It is designed for a maximum collision center of mass energy of $\sqrt{s} = 14 \text{ TeV}$ and
787 has a circumference of 26.7 kilometers. Four main experiments are located at the interaction points (IP)
788 of the accelerator: ATLAS (A Toroidal LHC ApparatuS), CMS (the Compact Muon Solenoid), ALICE
789 (A Large Ion Collider Experiment), and LHCb [7, 57–59]. The studies performed in this thesis were all
790 completed with the ATLAS detector.

791 Figure 2.1 shows a schematic of the LHC ring and the various experiments.

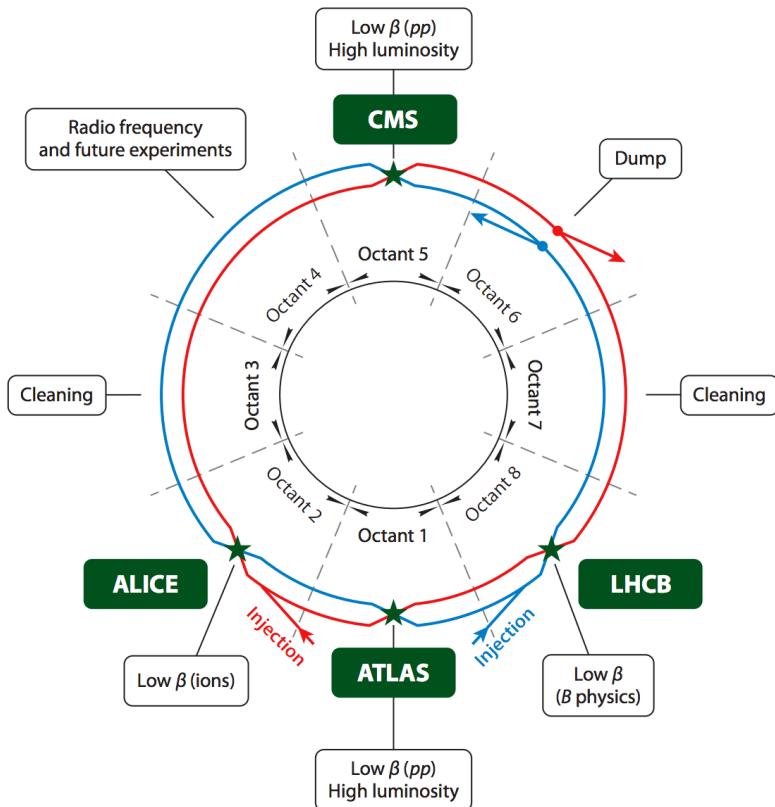


Figure 2.1: A schematic view of the LHC ring [6]

792 One of the most interesting features of the LHC is in its magnet design. Because the tunnel does not
793 have room for separate superconducting magnets for each of the beam pipes, the LHC employs a twin-
794 bore magnet design. Each magnet must hold an 8.3 Tesla magnetic field in order to bend the proton

795 beams at $\sqrt{s} = 14$ TeV. The superconducting magnets are cooled to a temperature of 1.9 Kelvin with
 796 superfluid helium.

797 2.1.1 INSTANTANEOUS LUMINOSITY

798 The rate of physics events expected from the accelerator is dependent on the instantaneous luminosity
 799 of the machine and the cross section of the physics process, $R_{\text{events}} = L\sigma$. Here, R_{events} is the number
 800 of events per second, L is the instantaneous luminosity of the machine, and σ is the cross section for the
 801 physics process being measured. The instantaneous luminosity of the LHC is determined by numerous
 802 factors related to machine conditions. Equation 2.1 gives the equation for instantaneous luminosity of
 803 Gaussian beam profile [6].

$$L = \frac{N_b^2 n_b f_{\text{rev}} \gamma_r}{4\pi \epsilon_n \beta^*} F \quad (2.1)$$

804 The LHC collides protons in bunches, and in the above equation N_b is the number of protons per
 805 bunch while n_b is the number of bunches per beam. Nominally, the LHC can hold up to 2808 pro-
 806 ton bunches. f_{rev} is the revolution frequency. ϵ_n is the normalized transverse beam emittance, a mea-
 807 surement of the average spread of the particles position-momentum space which has the dimension of
 808 length. β^* is the value of the *beta* function for the beam at the interaction point. It relates the emmi-
 809 tance to the Gaussian width of the beam with $\sigma_{\text{beam}} = \sqrt{\epsilon \cdot \beta}$. F is a reduction factor that corrects for
 810 the fact that the beams are colliding at an angle at the IP.

811 Another way of writing the instantaneous luminosity is shown in equation 2.2. In this case, the in-
 812 stantaneous luminosity is written as the ratio of the rate of inelastic collisions with the inelastic cross
 813 section[60].

$$L = \frac{R_{\text{inel}}}{\sigma_{\text{inel}}} = \frac{\mu n_b f_{\text{rev}}}{\sigma_{\text{inel}}} \quad (2.2)$$

814 In this case, μ is the average number of interactions per bunch crossing in the accelerator. μ is a useful
 815 parameter for characterizing the amount of activity recorded in an experiment. As the instantaneous

816 luminosity and thus μ increase, there are more interactions per bunch crossing and more activity in the
 817 detector. This is often characterized with $\langle \mu \rangle$, the measured per bunch crossing μ value averaged over
 818 all bunch crossings. The interactions inside each bunch crossing that are not the main physics process of
 819 interest are often referred to as “pileup” interactions, and $\langle \mu \rangle$ is a measurement of the level of pileup in
 820 the detector.

821 2.1.2 EVOLUTION OF MACHINE CONDITIONS

822 This thesis uses datasets taken at three different center of mass energies: $\sqrt{s} = 7$ TeV data taken in the
 823 year 2011, $\sqrt{s} = 8$ TeV data taken in the year 2012, and $\sqrt{s} = 13$ TeV dataa taken in the year 2015. In
 824 addition to increasing center of mass energy, the instananeous luminosity and parameters that determine
 825 it were evolving. Table 2.1 summarizes that machine conditions in each of these datasets.

	2011	2012	2015	Design
\sqrt{s} [TeV]	7	8	13	14
Number of bunches	1380	1380	1825	2808
Max. protons per bunch	1.45×10^{11}	1.7×10^{11}		1.15×10^{11}
Bunch spacing [ns]	50	50	25	25
Max. instantaneous luminosity [$\text{cm}^{-2}\text{s}^{-1}$]	3.7×10^{33}	7.7×10^{33}	5×10^{33}	10^{34}
β^* [m]	1.0	0.6	0.8	0.55
$\langle \mu \rangle$	11.6	20.7	13.7	-

Table 2.1: Evolution of LHC machine conditions [31, 32]

826 2.2 THE ATLAS DETECTOR

827 The ATLAS detector is a multi-purpose particle detector experiment at the LHC’s Point 1 [7]. It has
 828 nearly 4π coverage in solid angle around the interaction point. It consists of an inner detector for mea-
 829 suring charged particles, electromagnetic and hadronic calorimeters, and a muon spectrometer. Fig-
 830 ure 2.2 gives an overview of the detector.

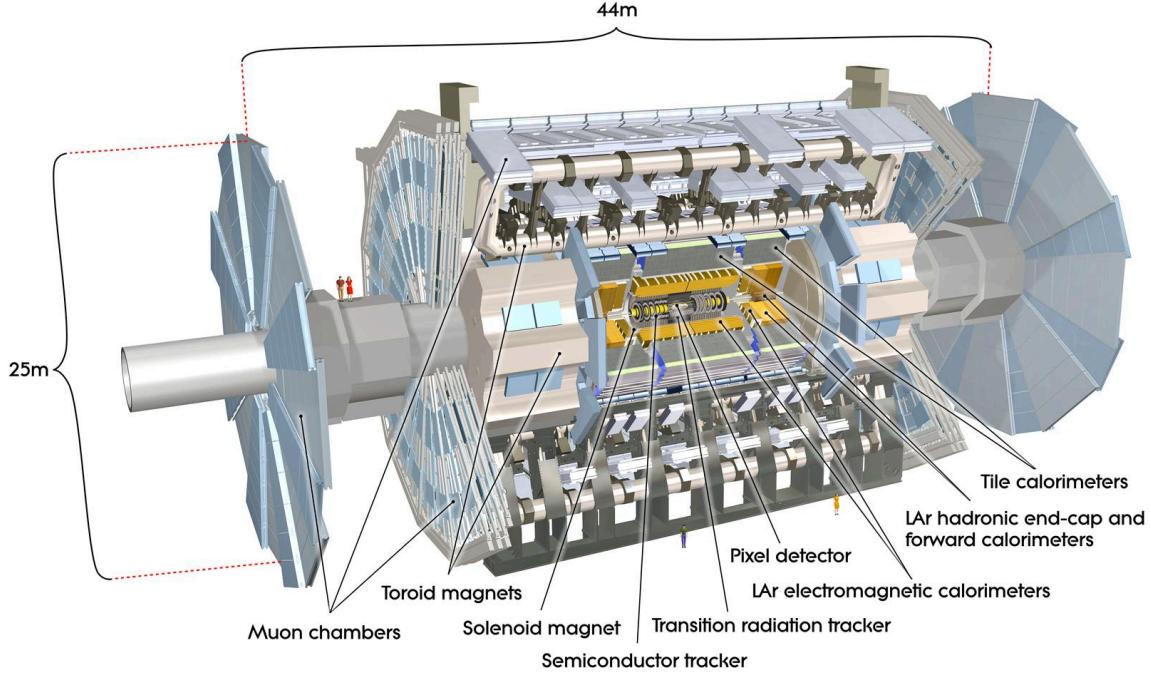


Figure 2.2: A full diagram of the ATLAS detector [7]

831 2.2.1 COORDINATE SYSTEM

832 Before defining the properties of the individual detectors, it is important to establish the coordinate
 833 system used. Figure 2.3 shows a schematic of the coordinate system. The azimuthal plane (perpendicular
 834 to the beam line) is defined as the x - y plane. The angle in this plane is referred to as ϕ . The angle relative
 835 to the beam axis is referred to as θ . Rather than using θ directly as a coordinate, the experiment often
 836 uses the pseudorapidity η . η is defined in equation 2.3.

$$\eta = \ln \left(\tan \left(\frac{\theta}{2} \right) \right) \quad (2.3)$$

837 Pseudorapidity is the massless approximation of rapidity, the angle used to parameterize boosts in
 838 special relativity. This is important for two reasons. First, it means that differences in η are Lorentz in-
 839 variant. Second, particle production is roughly constant in pseudorapidity. Particles with η close to zero
 840 are referred to as “central”, while those at high $|\eta|$ are called “forward”. In general, two main detector

841 topologies can be seen in figure 2.2. There are “barrel” elements, which surround the beam line cylind-
 842 drically and are in the central region of the detector. In the forward region, there are “endcap” regions
 843 which are arranged as disks perpendicular to the beam line.

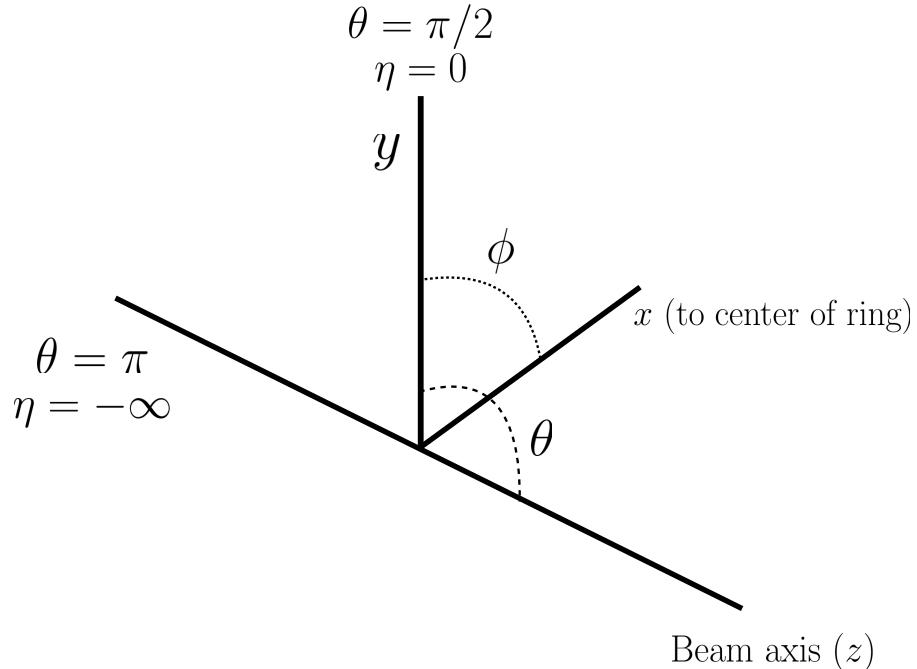


Figure 2.3: The ATLAS coordinate system

844 2.2.2 INNER DETECTOR

845 The ATLAS Inner Detector (ID) system is built for precision tracking of charged particles. It covers the
 846 range $|\eta| < 2.5$. In this range, approximately 1000 particles are generated every bunch crossing in the
 847 detector. This requires having fine granularity to achieve the resolutions required for good momentum
 848 measurement and vertex reconstruction.

849 The ID consists of three sub-components: the pixel detector, semiconductor tracker (SCT), and trans-
 850 sition radiation tracker (TRT). It is surrounded by a solenoid providing a 2 T axial magnetic field which
 851 bends particles in the transverse plane to allow for momentum measurement. Figure 2.4 shows the layout
 852 of each of these components.

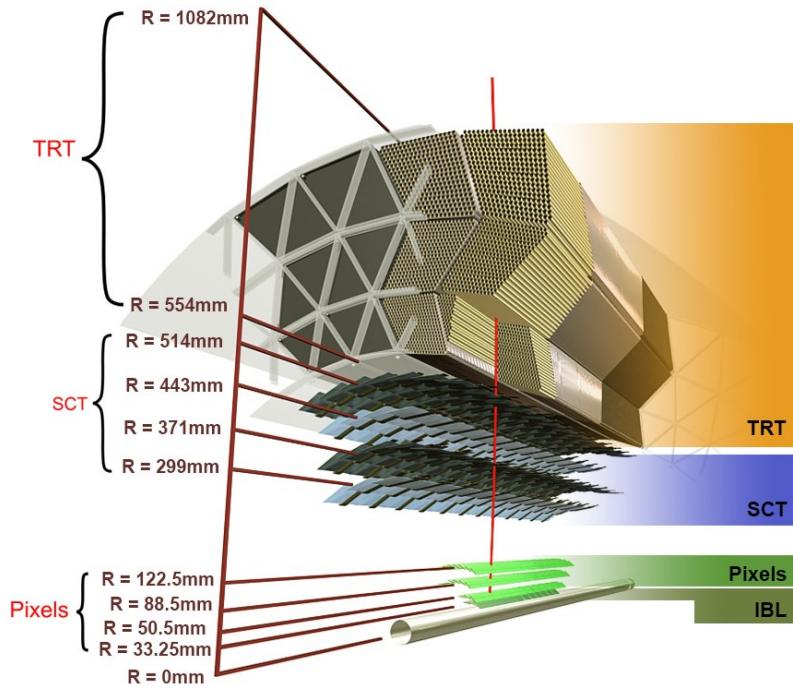


Figure 2.4: Layout of the ATLAS Inner Detector system [8]

853 PIXEL DETECTOR

854 The pixel detector is the first detector particles traverse after being generated in proton collisions and is
 855 the most granular detector. Its operation is crucial for precision tracking and vertex reconstruction as
 856 well as higher level object reconstruction like tagging of jets from b -quarks. The basic sensing element in
 857 this subdetector is a silicon pixel detector. The operating principle for the silicon pixels is that of a p - n
 858 junction. When a charged particle passes through, it creates electron-hole pairs that are then separated
 859 by the electric field. The sensors are $250 \mu\text{m}$ thick and use oxygenated n -type wafers with readout pixels
 860 on the n^+ side of the detector [7]. Overall, the pixel detector has 1744 sensors and 80.4 million readout
 861 channels.

862 In the barrel region, the pixel detector has three concentric layers of sensors surrounding the beam-
 863 line. In the endcap region, it consists of disks perpendicular to the beam axis. The detector is segmented
 864 in the R - ϕ plane and in z . Usually, three pixel layers are crossed by a charged particle track. The intrinsic
 865 accuracies of the sensors are $10 \mu\text{m}$ in R - ϕ and $115 \mu\text{m}$ in z (or R for the endcap).

866 **INSERTABLE B-LAYER**

867 In Run 2, a new innermost pixel layer, known as the insertable B-layer (IBL), was added to the Inner
868 Detector [61]. This layer was added to cope with the higher luminosities planned in LHC Run 2 and
869 at the high luminosity HL-LHC. Additionally it improves tracking position resolution which in turn
870 improves the vertexing and b -tagging capabilities in ATLAS. The detector sits directly on a new beam
871 pipe, only 33.25 mm away from the collision points in the azimuthal plane.

872 **SEMICONDUCTOR TRACKER (SCT)**

873 The semiconductor tracker (SCT) consists of silicon microstrips and comprises the next four layers of
874 the ID. This sub-detector has 6.4cm long sensors that are daisy-chained into strips with a strip pitch of
875 $80 \mu\text{m}$ [7]. Some of the strips have a small stereo angle to allow for measurement of both angular coor-
876 dinates. In total there are 6.3 million readout channels. The intrinsic accuracies are $17 \mu\text{m}$ in $R\text{-}\phi$ and
877 $580 \mu\text{m}$ in z (or R in the endcap).

878 **TRANSITION RADIATION TRACKER (TRT)**

879 The transition radiation tracker (TRT) serves two purposes. First, it consists of 4mm diameter straw
880 tubes filled with a 70/27/3% gas mixture of xenon, carbon dioxide, and oxygen to provide tracking of
881 charged particles. Particles typically have 36 TRT straw tube hits per track. The material in between
882 the straws is designed to induce transition radiation which can be useful for particle identification. As
883 particles pass between media with different dielectric constants, they emit transition radiation that can
884 cause additional showers in the TRT. In particular it is useful for discrimination between electrons and
885 pions or other charged hadrons, as the amount of transition radiation is proportional to the Lorentz
886 factor of the particle.

887 **2.2.3 CALORIMETERS**

888 The calorimeter system consists of two main sub-components: a fine granularity electromagnetic calorime-
889 ter tailored for the measurement of photons and electrons and multiple coarser hadronic calorimeters

890 dedicated to the measurement of hadronic showers [7]. The calorimeter system has broader coverage
891 than the inner detector, covering the region out to $|\eta| < 4.9$. It is also designed to deliver good contain-
892 ment of showers so as to limit leakage into the muon system. Figure 2.5 shows the layout of the calorime-
893 ter system.

894 Both the electromagnetic and hadronic calorimeters are sampling calorimeters. They alternate active
895 material for energy measurement with passive material for energy absorption. The materials used for
896 each purpose vary based on the type of calorimeter and its location in the detector.

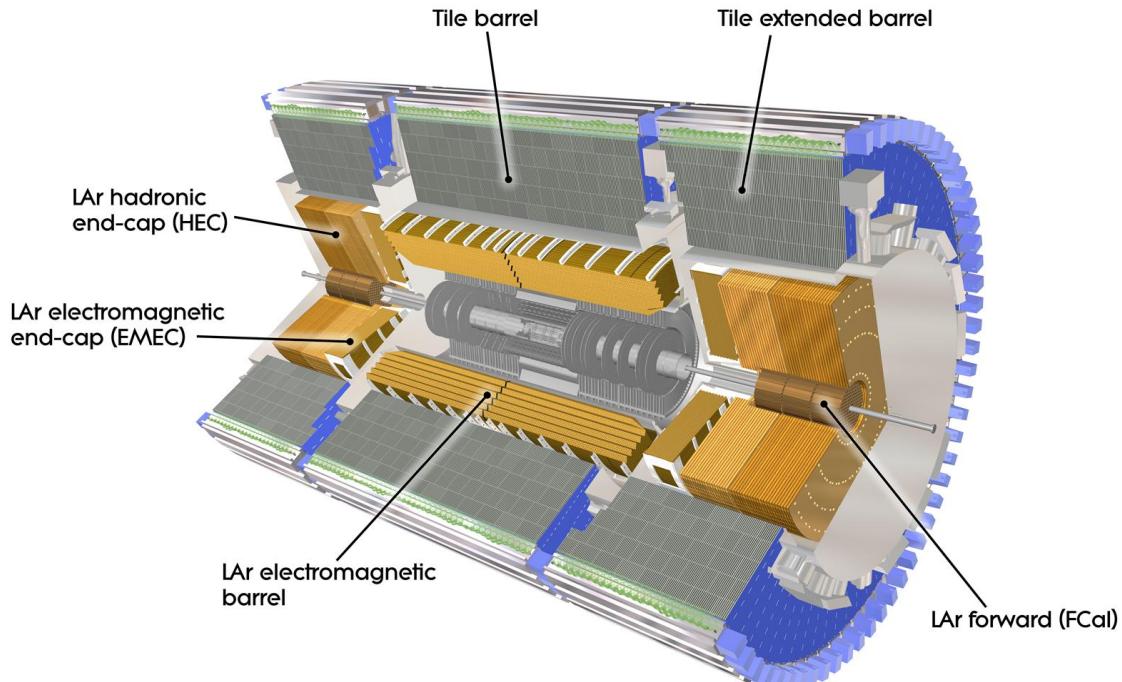


Figure 2.5: Layout of the ATLAS calorimeter system [7]

897 ELECTROMAGNETIC CALORIMETER

898 The electromagnetic calorimeter (EM calorimeter) use liquid Argon (LAr) as its active material and lead
899 as its passive material. It is arrange in an accordion geometry to increase the absorption area while still
900 allowing it to have no azimuthal cracks (complete symmetry in ϕ). The EM calorimeter is divided into a

901 barrel portion that extends to $|\eta| < 1.475$ and an endcap portion going from $1.375 < |\eta| < 3.2$. The
902 region where these two units overlap is called the “transition region”.

903 In order to provide good containment the calorimeter depth must be optimized. Typically, for elec-
904 tronagnetic calorimeters the depth is measured in radiation lengths. In general, the intensity of a par-
905 ticle beam attenuates exponentially in distance with a constant equal to the radiation length. That is,
906 $I(x) = I_0 e^{-x/X_0}$, where I is the intensity, x is the distance traveled, and X_0 is the radiation length.
907 The ATLAS EM calorimeter is designed to have > 22 radiation lengths in the barrel and > 24 in the
908 endcap [7].

909 **HADRONIC CALORIMETERS**

910 There are three types of hadronic calorimeters present in ATLAS: the tile calorimeter (TileCal), hadronic
911 endcap (HEC), and forward calorimeter (FCal). Each one is optimized for stopping of hadronic showers
912 and the materials chosen are specific to their placement in the detector.

913 The TileCal is a scintillating tile calorimeter placed directly outside the EM calorimeter. It uses steel as
914 the absorber and plastic scintillator tiles as the active material. It has coverage in the barrel at $|\eta| < 1.0$
915 and in the “extended barrel” region of $0.8 < |\eta| < 1.7$.

916 The HEC had two wheels perpendicular to the beam line per endcap and is located directly behind
917 the EM calorimeter endcap modules. The HEC covers the region from $1.5 < |\eta| < 3.2$, overlapping
918 slightly with both the tile calorimeter and the forward calorimeter. Like the EM calorimeter, it uses liq-
919 uid Argon as the active material, but it uses copper as the absorber.

920 The FCal covers the most forward regions of the calorimeter system, extending to the region of $3.1 <$
921 $|\eta| < 4.9$. It again uses liquid argon as its active material. For absorber, it consists of an innermost
922 module made of copper followed by a module made of tungsten.

923 The hadronic equivalent of radiation length is called the interaction length and is denoted as λ . In the
924 barrel, the hadronic calorimeter depth is approximately 9.7λ , while in the endcap is is 10λ . The outer
925 supports contribute an additional 1.3λ . This is been shown to be sufficient to limit punch-through of
926 showers to the muon system [7].

927 2.2.4 MUON SPECTROMETER

928 The muon spectrometer is dedicated to measuring the momentum and position of muons. It consists
929 of tracking and trigger chambers which are unique in the barrel and endcap regions. The magnetic field
930 for bending of muons is provided by a system of three large air-core toroid magnets (from which ATLAS
931 derives its name.) These magnets provide 1.5 to 5.5 Tm of bending power at $0 < |\eta| < 1.4$ and approx-
932 imately 1 to 7.5 Tm in the endcap region of $1.6 < |\eta| < 2.7$. The entire muon system covers the range
933 $0 < |\eta| < 2.7$. Monitored drift tubes (MDTs) are used for tracking in the barrel and the two outer
934 layers of the endcap, while cathode strip chambers (CSCs) are used to provide tracking in the innermost
935 endcap wheel. In the barrel, resistive plate chambers (RPCs) are used as trigger chambers while thin gap
936 chambers (TGCs) are used in the endcap. Figure 2.6 shows the layout of the ATLAS muon system. The
937 entire muon system is designed with the specification of providing a 10% momentum resolution for a
938 1 TeV muon.

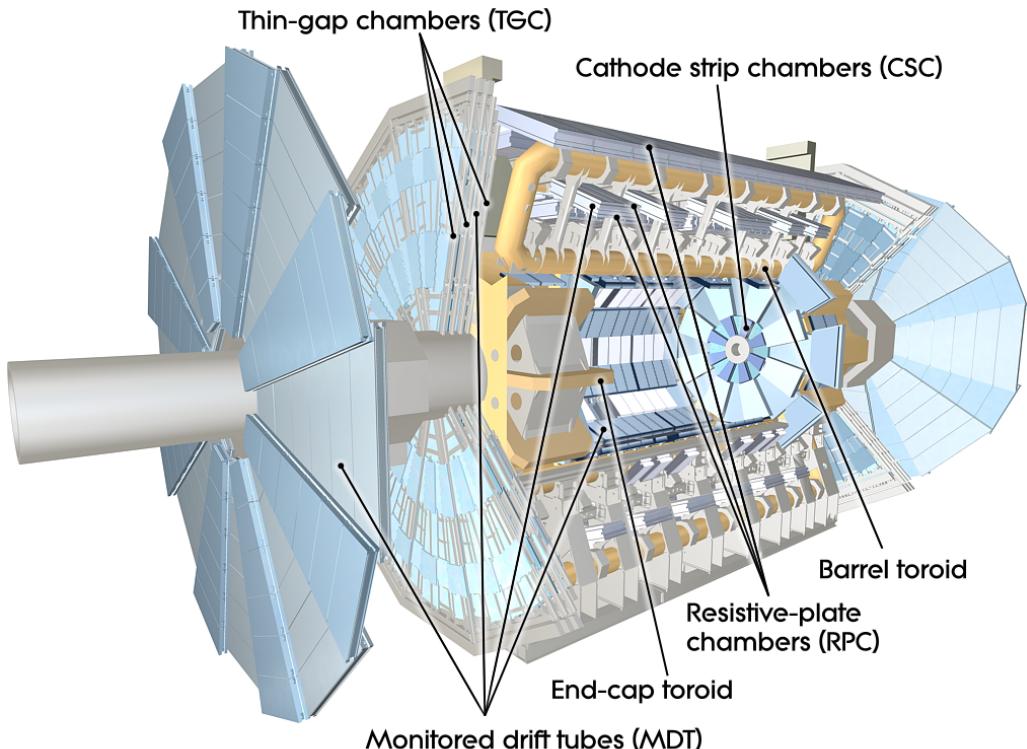


Figure 2.6: Layout of the ATLAS muon system [7]

939 MONITORED DRIFT TUBES (MDTs)

940 The monitored drift tubes (MDTs) are aluminum 3cm diameter tubes filled with a 93/7 % mixture of
941 Argon and CO₂, with trace amounts of water. As a charged particle traverses the tube, it ionizes the gas
942 and the ions drift to a wire at the center of the tube. The radial distance of traversal of the particle in the
943 tube is determined by the drift time of the electrons, allowing for fine position resolution. The tubes
944 have an average resolution of 80 μm per tube and a maximum drift time of approximately 700ns. The
945 tubes are oriented so that they give precision measurement in η and run along ϕ . They cover $|\eta| < 2.7$,
946 except in the innermost layer of the endcap where they only go to $|\eta| < 2.0$ [7].

947 CATHODE STRIP CHAMBERS (CSCs)

948 The cathode strip chambers cover a narrow window of the innermost endcap region at $2.0 < |\eta| <$
949 2.7. In this region the background rates in the cavern are particularly high and the CSCs are designed
950 to handle these higher rates. The CSCs are multiwire proportional chambers with wires pointing in
951 the radial direction (away from the beam pipe). The wire serves as an anode and there are two types of
952 segmented cathode strip, one perpendicular to the wires which gives the precision measurement and one
953 parallel which provides the transverse coordinate. It has an 80/20 gas mixture of Argon and CO₂ [7].

954 RESISTIVE PLATE CHAMBERS (RPCs)

955 The resistive plate chambers (RPCs) are gaseous electrode-plate detectors covering the region $|\eta| <$
956 1.05. They consist of two resistive plates separated by a distance of 2 mm. The gas mixture used is a
957 94.7/5/0.3% mixture of C₂H₂F₄, Iso-C₄H₁₀, and SF₆. It has readout strips with a pitch of 23-35 mm
958 for both η and ϕ measurement and thus provides measurement of the azimuthal coordinate in the barrel
959 that the MDTs do not. The thin gas gap allows for a quick response time which makes it ideal for use in
960 the trigger. There are three layers of RPCs which are referred to as the three trigger stations. They allwo
961 for both a low p_T and high p_T trigger. The coincidence of hits in the innermost chambers allows for
962 triggering of muons between 6 and 9 GeV, while the outermost layer allows the trigger to select high
963 momentum tracks in the range of 9 to 35 GeV [7].

964 THIN GAP CHAMBERS (TGCs)

965 The thin gap chambers (TGCs) are multiwire proportional chambers where the wire to cathode dis-
966 tance (1.4mm) is smaller than the wire-to-wire distance (1.8 mm). They contain a gas mixture of CO₂
967 and *n*-pentane and use a hih electric field to gain good time resolution. They serve two functions in the
968 end-cap system. First, they serve as the trigger chambers. Second, they also provide azimuthal coordi-
969 nate measurement which the MDTs do not. They sit on the inner and middle layers of the endcap. The
970 outermost layer's azimuthal coordinate is determined by extrapolation [7].

971 2.2.5 MAGNET SYSTEM

972 As mentioned previously, there are two independent magnet systems in ATLAS. The first is a 2 T solenoid
973 field in the inner detector which provides bending in the azimuthal plane. The second is an approxi-
974 mately 0.5 T toroidal field in the muon system which provides bending in $|\eta|$. Figure 2.7 shows the pre-
975 dicted field integral as a function of $|\eta|$ [7].

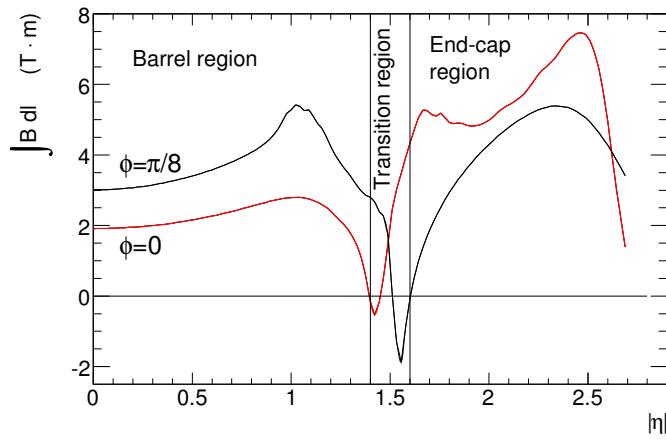


Figure 2.7: Predicted field integral as a function of $|\eta|$ for the ATLAS magnet system [7]

976 2.2.6 TRIGGER SYSTEM

977 The ATLAS trigger system searches for signatures of muons, electrons, photons, hadronically decay-
978 ing τ leptons, and jets in order to save these events for further analysis. The trigger system in ATLAS

979 is designed to reduce the maximum LHC event rate of 40 MHz to a more reasonable rate that can be
 980 recorded. The trigger first consists of a fast, hardware based system called the Level-1 (L1) trigger. The
 981 L1 trigger consists of independent dedicated detector sub-components that can seed regions of in-
 982 terest (RoIs) for further analysis downstream. For muons, the RPCs and TGCs are used, while in the
 983 calorimeter coarsely grained sections of calorimeter cells called towers are used. Once regions of interest
 984 are seeded, a software based system called the High Level Trigger (HLT) is used to reconstruct objects
 985 and integrate information from different parts of the detector. In Run 1 of ATLAS, the HLT consisted
 986 of two separate stages: the level 2 (L2) trigger and the event filter (EF).

987 The maximum trigger rate that the L1 trigger can handle is 75 kHz. In the HLT, the rate of events
 988 written to disk is approximately 200 Hz. Figure 2.8 shows the trigger rates for different L1 triggers in
 989 2012 and 2015 for ATLAS [9].

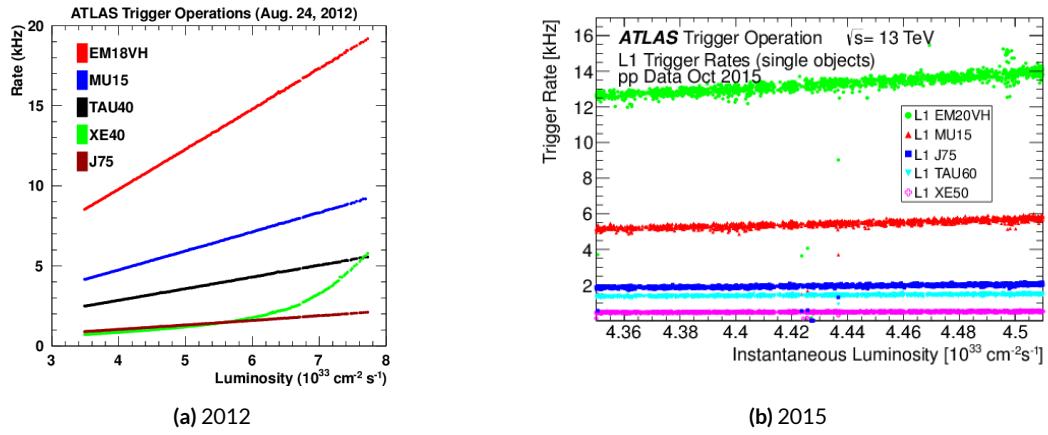


Figure 2.8: ATLAS trigger rates for Level-1 triggers as a function of instantaneous luminosity in 2012 and 2015 operation. These are single object triggers for electromagnetic clusters (EM), muons (MU), jets (J), missing energy (XE), and τ leptons (TAU). The threshold of the trigger is given in the name in GeV. [9]

990 2.2.7 ATLAS DATASETS

991 ATLAS has collected data at center of mass energies of 7, 8, and 13 TeV. Figure 2.9 shows the integrated
 992 luminosity as a function of time for each of the three collected datasets. At $\sqrt{s} = 7$ TeV, ATLAS
 993 recorded 5.08 fb^{-1} . Increased instantaneous luminosity in 2012 led to a larger dataset of 21.3 fb^{-1}
 994 recorded at $\sqrt{s} = 8$ TeV. After Long Shutdown 1 (LS1) of the LHC and a restart in 2015, ATLAS

995 recorded 3.9 fb^{-1} of data at $\sqrt{s} = 13 \text{ TeV}$. [10, 11]

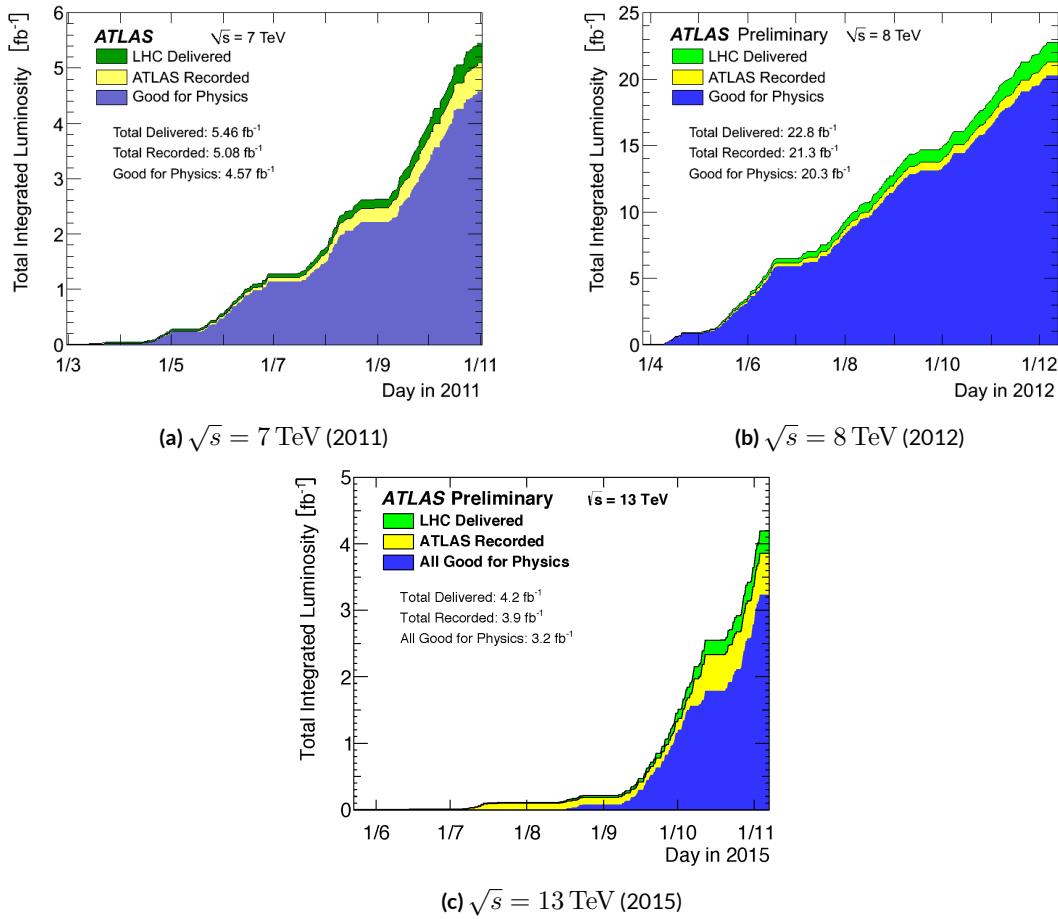


Figure 2.9: Instantaneous luminosity as a function of time for data recorded by ATLAS at different center of mass energies [10, 11]

996 2.2.8 DETECTOR PERFORMANCE

997 Table 2.2 summarizes the design requirements for each of the different sub-detectors. This table shows
 998 the energy and momentum resolution of each tracking, calorimetry, and muon measurements.

999 2.3 THE ATLAS MUON NEW SMALL WHEEL UPGRADE

1000 As the LHC continues operation, it is scheduled to be upgraded in several phases to allow it to reach
 1001 higher instantaneous luminosities and thus collect larger datasets. These conditions will open new doors

	Required resolution
Tracking	$\sigma_{p_T}/p_T = 0.05\% p_T \oplus 1\%$
EM calorimetry	$\sigma_E/E = 10\%/\sqrt{E} \oplus 0.7\%$
Hadronic calorimetry	
Barrel and end-cap	$\sigma_E/E = 50\%/\sqrt{E} \oplus 3\%$
Forward	$\sigma_E/E = 100\%/\sqrt{E} \oplus 10\%$
Muon spectrometer	σ_{p_T}/p_T at $p_T = 1$ TeV

Table 2.2: Performance requirements for the ATLAS detector [7].

for study of rare physics processes but will also present interesting challenges that must be faced. ATLAS will require new detector technologies to cope with the increased background rates in the cavern in these high luminosity conditions. One such upgrade, scheduled to be installed during Long Shutdown 2 (LS2) of the LHC in 2018, is the ATLAS Muon New Small Wheel (NSW) upgrade [12]. The NSW will replace the innermost end-cap wheel of the muon system with new technologies, as this is the part of the muon detector closest to the beam and thus suffers from the highest rates.

2.3.1 MOTIVATION

The motivation of the NSW is two-fold. First, the objective is to alleviate the decreased tracking efficiency that comes in a high rate environment. As figure 2.10, at the LHC design luminosity both the efficiency of recording hits and reconstructing track segments in the MDTs decreases at the LHC design luminosity.

Second, the NSW will work to alleviate the rate of fake triggers arising in the endcap. Figure 2.11 shows the extrapolated trigger rates as a function of the p_T threshold with and without the NSW upgrade. As the figure shows, the NSW upgrade will reduce the trigger rate by an order of magnitude compared to the current endcap trigger system.

2.3.2 NSW DETECTOR TECHNOLOGIES

The NSW will use two new detector technologies - micromesh gaseous structure detectors (micromegas) and small-strip thin gap chambers (sTGCs) [12, 62]. Unlike the previous detectors, both of these detector technologies can be used for tracking or trigger. However, the micromegas is more suited to tracking

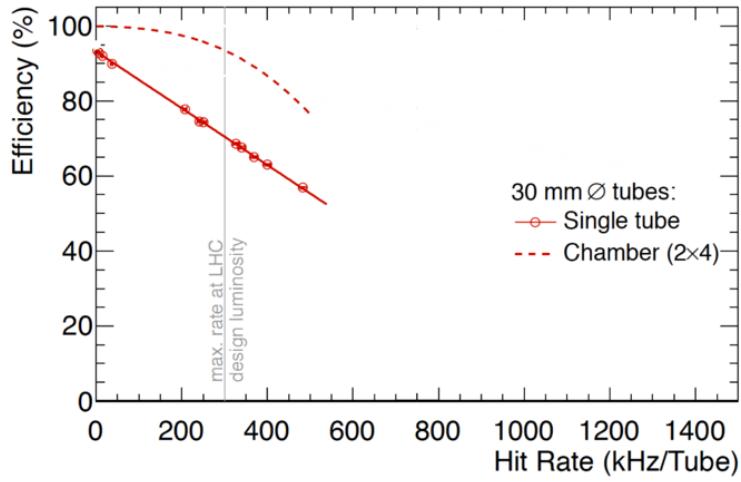


Figure 2.10: MDT tube hit (solid) and segment (dashed) efficiency as a function of hit rate per tube [12]

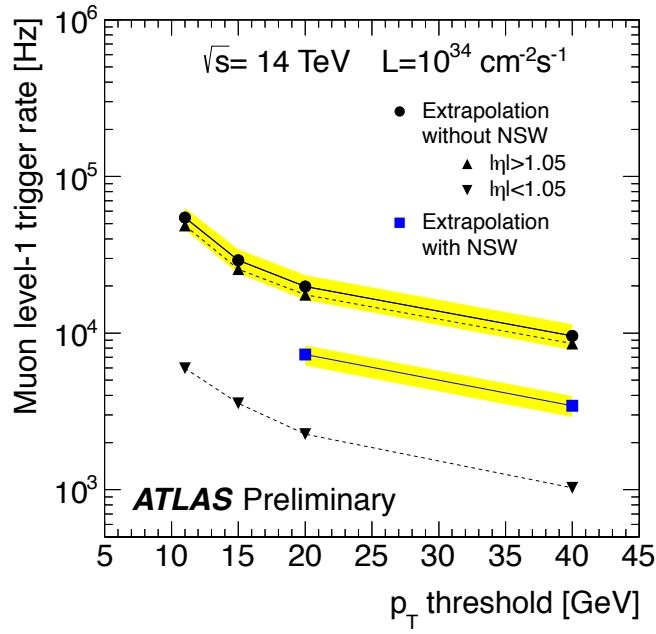


Figure 2.11: Trigger rate as a function of p_T threshold with and without the NSW upgrade [12]

¹⁰²¹ because of its good spatial resolution, while the sTGCs have better time resolution and are more suited
¹⁰²² for the trigger. To maintain a fully redundant system, both technologies are used for both purposes.

1023 **MICROMEGAS**

1024 Micromegas detectors operate using a thin metallic mesh that sits approximately $100 \mu\text{m}$ away from
1025 the readout electrodes to create the amplification region. Above this mesh, there is a drift region on the
1026 order of a few mm in length capped by a drift electrode. As a charged particle traverses the detector, it
1027 ionizes gas and the electrons drift down towards readout strips. The timing of the drift can be used to
1028 reconstruct the angle of traversal of the particle. This is illustrated in figure 2.12. The micromegas used in
1029 ATLAS will be resistive micromegas, where the readout electrodes are topped with resistive strips [63].
1030 This alleviates the risk of sparking in the large area detectors that ATLAS will use.

1031 In ATLAS, the micromegas drift gap will be 5 mm and the amplification gap will be $128 \mu\text{m}$. They
1032 are filled with the same gas mixture as the MDTs. They will be stacked in an octuplet in an XXUV-
1033 UVXX geometry, where X refers to straight strips and U and V refer to stero strips at an angle of $\pm 1.5^\circ$.
1034 This arrangement allows for measurement of the azimuthal coordinate and gives a large lever arm be-
1035 tween the straight strips for triggering purposes. Figure 2.12 shows the geometry of a single micromegas
1036 detector as well as its operating principle [12].

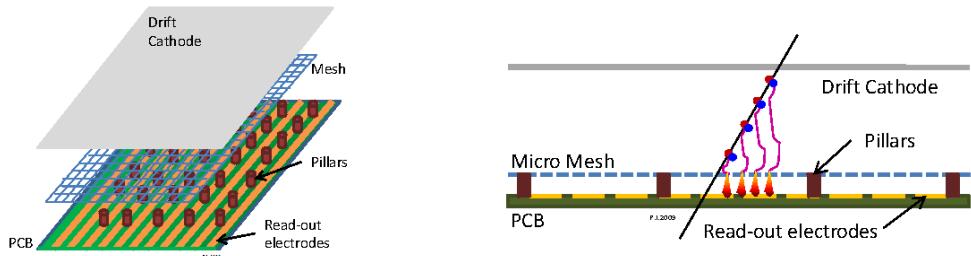


Figure 2.12: Illustrations of the geometry (left) and operating principle (right) of the micromegas detector [12]

1037 **sTGCs**

1038 The sTGCs are similar to the TGCs already described. They consist of gold-plated tungsten wires with
1039 a 1.8 mm pitch between two cathode planes 1.4 mm away from the wire plane. One cathode plane
1040 consists of strips with a 3.2 mm pitch (much smaller pitch than the TGCs), while the other consists
1041 of coarser pads that are used for defining regions of interest in the sTGC trigger algorithm. Figure 2.13
1042 shows the basic detector geometry.

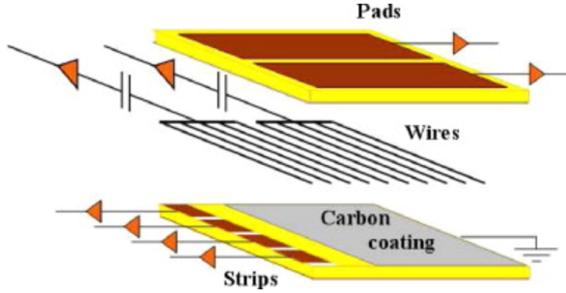


Figure 2.13: Geometry of the sTGC detector [12]

1043 **2.3.3 PHYSICS IMPACT**

1044 Maintaining low p_T thresholds for muons while still staying within the trigger rate budget at Level 1
 1045 (20 kHz) for the muon system is crucial for physics analyses to be successful in high luminosity condi-
 1046 tions. One realm where the lepton trigger threshold is especially important is in Higgs physics. In the
 1047 $H \rightarrow WW^*$ analysis, one of the W bosons is off shell and tends to decay to soft leptons. In associated
 1048 production of a Higgs with a W , the lepton is also important because the lepton provides the main han-
 1049 dle which allows the event to be triggered. Table 2.3 shows the impact of increasing the trigger thresholds
 1050 on these analyses. It shows that either raising the threshold or using only the barrel both have signifi-
 1051 cant impacts on the signal efficiency. With the NSW, the signal efficiency is largely maintained and the
 1052 triggers can be unprescaled.

Threshold	$H \rightarrow b\bar{b}$ (%)	$H \rightarrow WW^*$ (%)
$p_T > 20$ GeV	93	94
$p_T > 40$ GeV	61	75
$p_T > 20$ GeV (barrel only)	43	72
$p_T > 20$ GeV (with NSW)	90	92

Table 2.3: Signal efficiencies for WH production with $H \rightarrow b\bar{b}$ and $H \rightarrow WW^* \rightarrow \mu\nu qq$ under different trigger configurations [12].

1053 **2.4 OBJECT RECONSTRUCTION IN ATLAS**

1054 ATLAS analyses first start by requiring the presence of certain reconstructed physics objects in the event.
 1055 This section will present a brief overview of the algorithms used to reconstruct electrons, muons, jets

1056 (including b -jets), and missing energy^{*}. The performance of object reconstruction and measurement will
 1057 also be discussed as these are relevant to the analyses presented later. Figure 2.14 gives an overview of the
 1058 different sub-detectors that each type of particle will interact with in ATLAS.

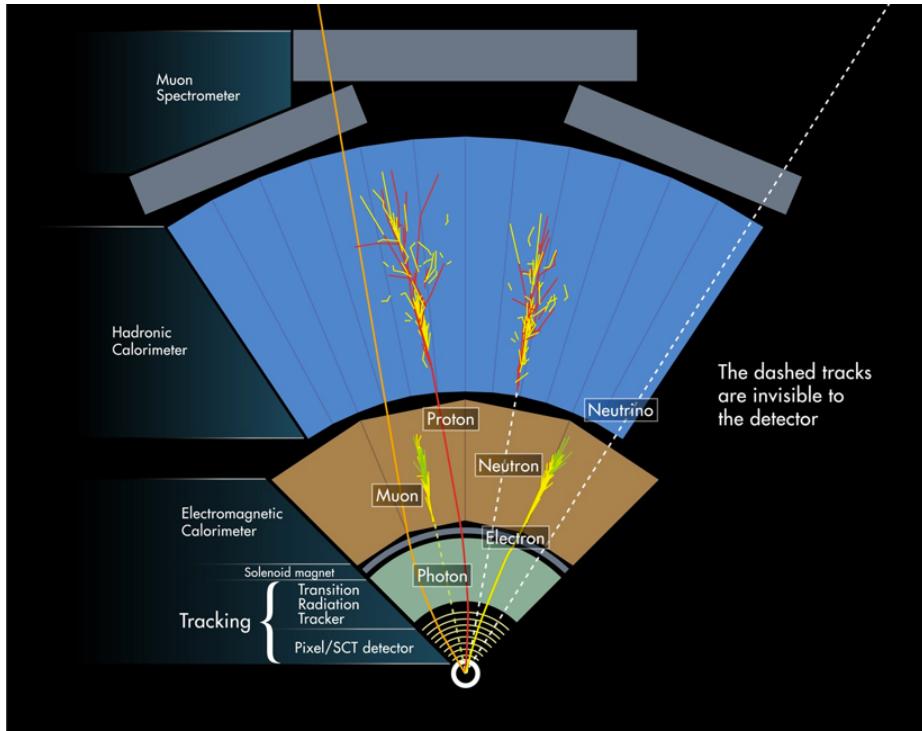


Figure 2.14: Illustration of particle interactions in ATLAS [13]

1059 2.4.I ELECTRONS

1060 Electrons in ATLAS will leave tracks in the inner detector and energy deposits in the electromagnetic
 1061 calorimeter. The algorithm for recognizing the signature of electrons proceeds in two steps: reconstruc-
 1062 tion and identification.

1063 In reconstruction, an electron candidate is formed by matching EM calorimeter deposits with ID
 1064 tracks. The algorithm first chooses seed clusters in the EM calorimeter by using a sliding window algo-
 1065 rithm that searches for towers with transverse energy larger than 2.5 GeV. In addition to seed clusters,
 1066 track candidates must be identified in the ID. The algorithm selects seed tracks with $p_T > 1$ GeV that

*Reconstruction algorithms for other objects, such as photons and τ leptons, are not detailed here as these objects are not used in the presented studies.

1067 do not fit well with a pion hypothesis. Once candidate tracks are selected, they are re-fit with a Gaussian Sum Filter (GSF) algorithm to estimate electron parameters [64]. Finally, an electron candidate is
 1068 formed if at least one track matches to a seed cluster in the calorimeter. The full details of the reconstruction
 1069 algorithm can be found in reference [14].
 1070

1071 Once an electron candidate is present, identification criteria must be applied in order to reject fake
 1072 electrons from background. Many different variables are used for this identification, most of them re-
 1073 lated to the shower shape in the EM calorimeter and the amount of leakage into the hadronic calorime-
 1074 ter, as well as information from the ID and in particular the TRT. There are both cut-based and likelihood-
 1075 based criteria that range from “loose” to “very tight”. For details, see reference [14].

1076 Figure 2.15 shows the algorithm’s reconstruction efficiency of true electrons for different identification
 1077 criteria as well as the electron energy resolution in simulation [14, 15]. The reconstruction efficiency is
 1078 measured using both Z and J/ψ tag and probe techniques.

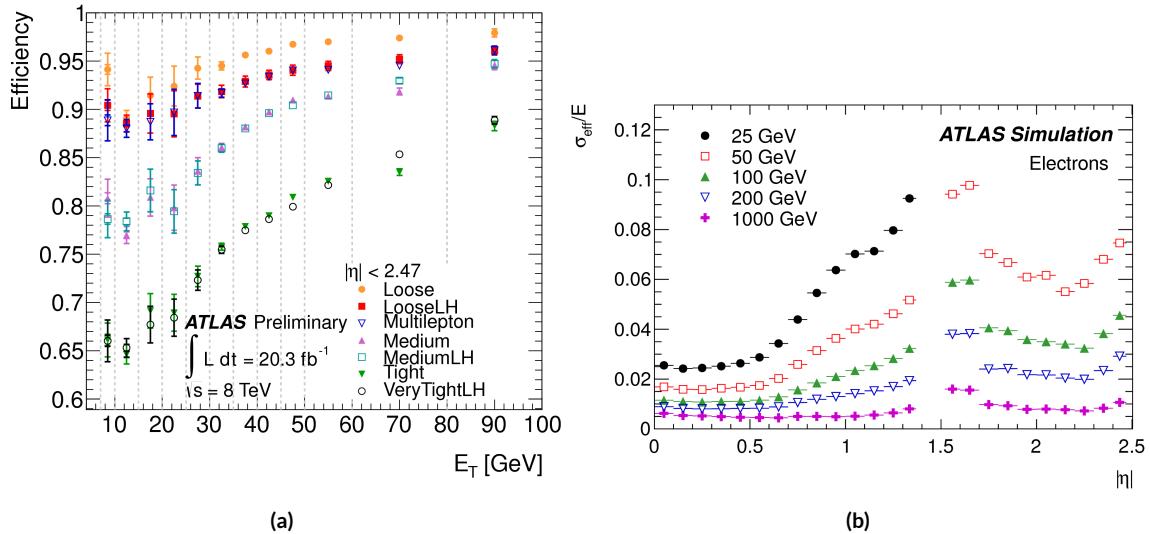


Figure 2.15: Electron performance: (a) reconstruction efficiency as a function of electron E_T [14] (b) energy resolution in simulation as a function of $|\eta|$ for different energy electrons [15]

1079 2.4.2 MUONS

1080 The ATLAS detector is designed to stop most particles before they reach the muon spectrometer. Muons,
 1081 however, are minimum ionizing particles, meaning that they will not lose a significant amount of energy

1082 through interactions with the detector and will thus pass through. Therefore, the muon reconstruction
1083 works to match tracks in the muon spectrometer with tracks in the inner detector.

1084 The first step of reconstruction is to reconstruct local straight line tracks, called segments, in each
1085 muon chamber. Segments are then fit to larger tracks that traverse the entire muon spectrometer. Such
1086 muon tracks are referred to as “standalone” tracks (SA) as they only use information from the muon
1087 spectrometer. The standalone tracks are then matched to tracks in the inner detector to form “com-
1088 bined” (CB) muons, where the combined ID and MS fit are used to determine the momentum and di-
1089 rection of the muon. To improve acceptance, segment-tagged and calorimeter-tagged muons are also
1090 reconstructed. In these cases, ID tracks are matched to segments in the MS and calorimeter deposits con-
1091 sistent with a minimum ionizing particle, respectively. The details of the reconstruction can be found in
1092 reference [16].

1093 As with electrons, once muon candidates are reconstructed they have identification criteria applied to
1094 reduce background. These criteria include the χ^2 match between the ID and MS tracks, the number of
1095 hits in the ID, overall ID and MS track fit quality, and additional variables [16]. The criteria range from
1096 “loose” to “tight” as with electrons.

1097 Figure 2.16 shows the muon reconstruction efficiency (measured with Z and J/ψ tag and probe) and
1098 invariant mass resolution [16].

1099 2.4.3 JETS

1100 When a quark or gluon is produced in collisions, it is not measured directly in ATLAS. Rather, due to
1101 QCD effects, it produces a collimated spray of hadrons in the direction of the original parton, which is
1102 known as a jet. Jets are reconstructed in ATLAS using energy deposits in the hadronic calorimeter. The
1103 first step is build “topological clusters” out of energy deposits in calorimeter cells [65, 66]. This is done
1104 using strategy where seed cells are chosen by picking cells whose energy measurements are four times the
1105 amount of noise expected for that cell. Adjacent cells with at least 2σ energy measurements are added
1106 to the cluster, then a final layer of clusters with energy above 0σ are added. Once calorimeter clusters
1107 are formed, they are clustered further into jet candidates using the anti- k_T jet clustering algorithm [67].

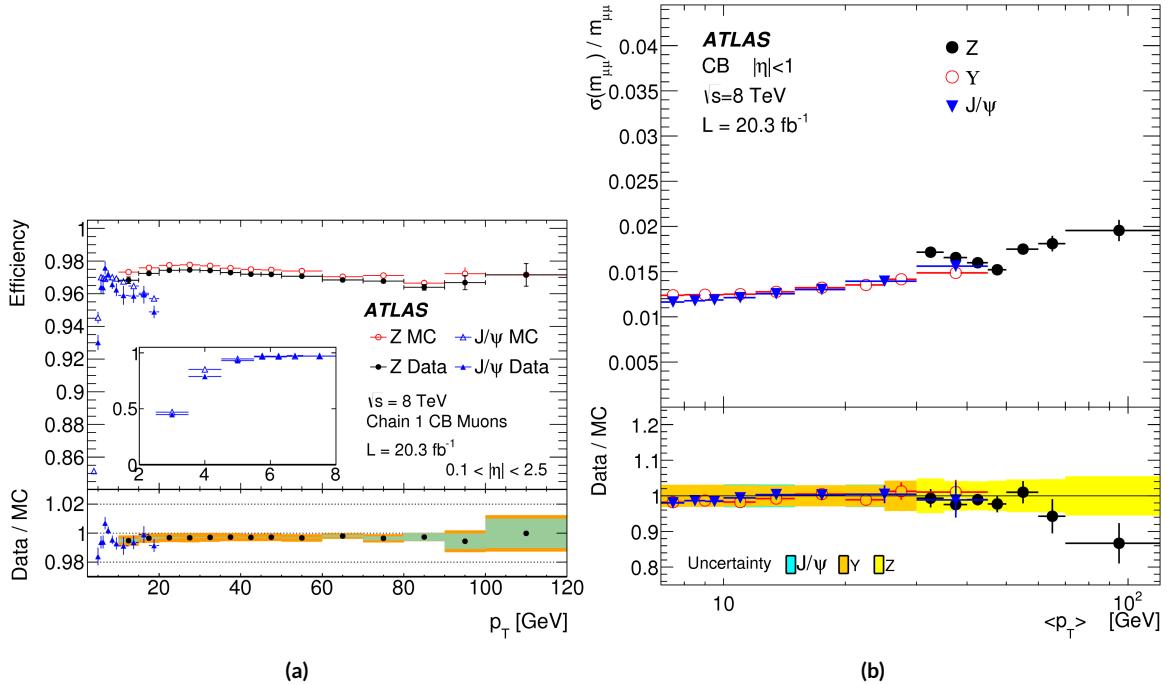


Figure 2.16: Muon performance: (a) reconstruction efficiency as a function of muon p_T (b) dimuon mass resolution as a function of average p_T [16]

This algorithm uses a parameter R that appears in the denominator of the clustering distance metric and defines the radial size of the jet in η - ϕ space.

The energy response of the calorimeter must be properly characterized in order to reconstruct jet energy. Calorimeter clusters can be calibrated either with the EM calibration, where each cluster is assumed to have come from the energy deposit of an electron or photon, or the LCW calibration, where local cluster weights are computed to allow for local calibration of clusters as hadronic or electromagnetic.

The details of the jet energy calibration are not detailed here and are discussed in reference [17].

Figure 2.17 shows the jet energy response after calibration in Monte Carlo as a function of the true p_T of the jet [17].

2.4.4 b -TAGGING

One important aspect of jet physics is the task of identifying the flavor of parton that produced the measured jet. While in general this is very difficult, jets from b -quarks offer an interesting case where such identification is possible. B mesons have a lifetime on the order of 10^{-12} seconds, which makes a $c\tau$

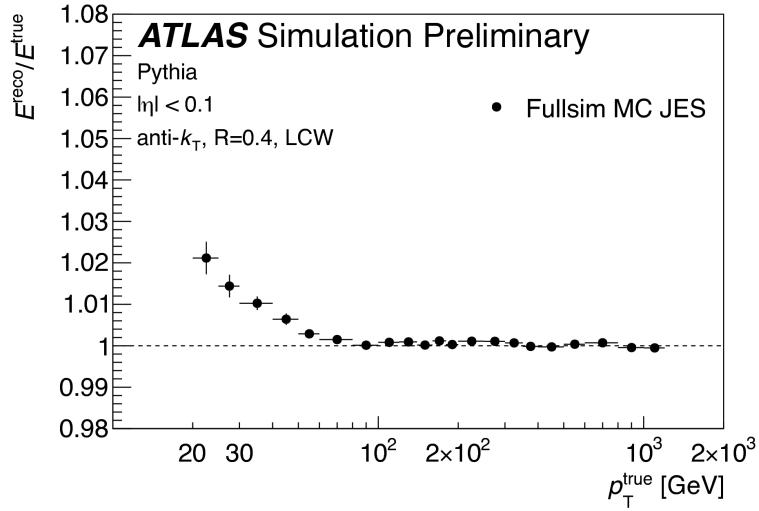


Figure 2.17: Jet energy response after calibration as a function of true p_T in simulation [17]

on the order of millimeters [1]. This type of displaced decay vertex can be identified in detectors like ATLAS and allows b -jets to be distinguished from other flavors of jets[†].

ATLAS uses a multivariate machine learning algorithm to identify jets from b -quarks. The inputs to this algorithm are determined from lower level reconstruction algorithms. There are three distinct algorithms that reconstruct variables which are used as input to the multivariate technique.

The first family is referred to as IPxD (where the x can either be 2 or 3). These algorithms use the transverse and longitudinal impact parameters d_0 and z_0 of the tracks inside a jet to determine their consistency with the primary vertex. They two or three dimensional (hence the x) templates for light flavor, charm, and bottom jets and then evaluate the likelihood of the jet coming from each of these types. The likelihood ratios are used as inputs to the multivariate algorithm.

The next two algorithms used as input are referred to as the secondary vertex (SV) and JetFitter (JF) algorithms. The SV algorithm uses tracks inside the jet to fit for vertices that are displaced from the primary vertex. The JF algorithm attempts to reconstruct the full flight path of the b by looking for multiple displaced vertices along the same line (as B decays often result in subsequent c decays).

In Run 1, the multivariate b -tagging algorithm used a neural network and was referred to as MV1.

[†]Jets from charm quarks can also be detected in this way but they do not live quite as long so the displacement of the vertex is harder to distinguish

1136 The details of this algorithm and its inputs are given in reference [18]. In Run 2, the number of inputs
 1137 was simplified and a boosted decision tree with 24 input variables was used, referred to as MV2. The
 1138 details of this algorithm are in reference [19]. Figure 2.18 shows the performance of each of these algo-
 1139 rithms.

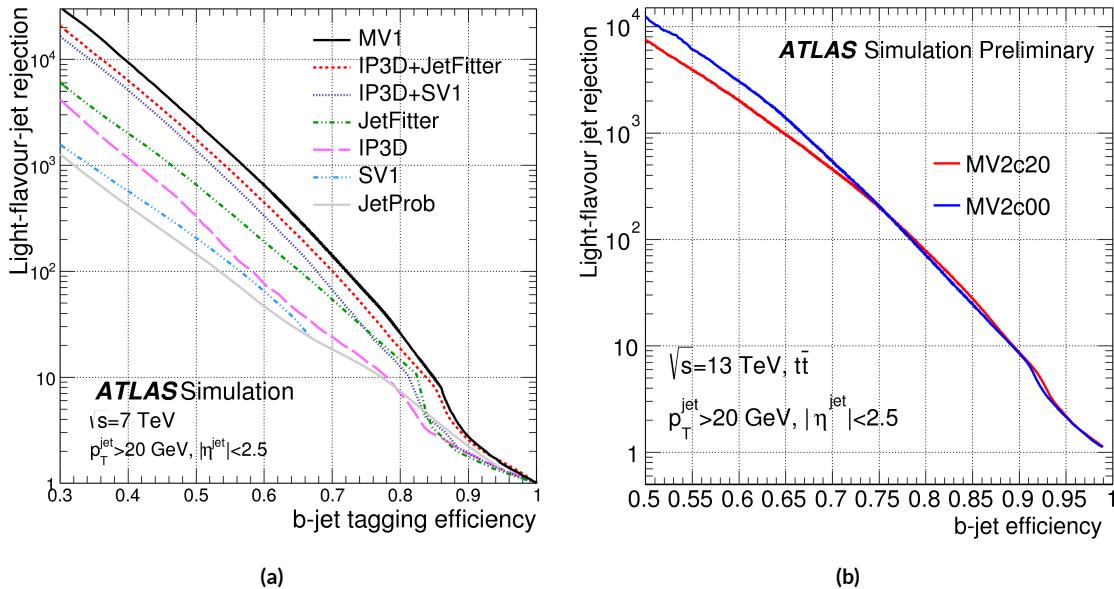


Figure 2.18: Light jet rejection (1/efficiency) vs. b -jet efficiency for MV1 and its input algorithms (a) [18] and MV2 (b) [19] in simulated $t\bar{t}$ events. The numbers in the algorithm names in (b) refer to the fraction of charm events used in the MV2 training.

1140 2.4.5 MISSING TRANSVERSE ENERGY

1141 As noted in figure 2.14, neutrinos produced in ATLAS will pass through the detector without inter-
 1142 acting. The only way of detecting the presence of particles like neutrinos (or BSM particles that are
 1143 long-lived) is to use missing transverse momentum. The basic principle of missing transverse energy is
 1144 to use the momentum balance of the incoming protons to infer the presence of missing particles. The
 1145 net longitudinal momentum of the incoming partons that collide is not known (since each carries an un-
 1146 known fraction of the proton's momentum). However, the protons (and thus incoming partons) have
 1147 no net momentum in the plane transverse to the beam line (the x - y plane). Therefore, if there are no
 1148 un-measured particles in the final state, the transverse momenta of all of the final state particles should

balance. The magnitude of this imbalance is known as missing transverse momentum (E_T^{miss}).

The basic calculation of missing transverse momentum from calorimeter cells is given in equation 2.4 [68].

$$\begin{aligned} E_x^{\text{miss}} &= -\sum_{i=1}^{N_{\text{cell}}} E_i \sin \theta_i \cos \phi_i \\ E_y^{\text{miss}} &= -\sum_{i=1}^{N_{\text{cell}}} E_i \sin \theta_i \sin \phi_i \end{aligned} \quad (2.4)$$

The E_T^{miss} calculation is separated into different terms based on the objects that the calorimeter clusters are associated with. This way, each cell's contribution is calibrated appropriately according to the object. This separation of terms is shown in equation 2.5 [68].

$$\begin{aligned} E_{x(y)}^{\text{miss,calo}} &= E_{x(y)}^{\text{miss},e} + E_{x(y)}^{\text{miss},\gamma} + E_{x(y)}^{\text{miss},\tau} + E_{x(y)}^{\text{miss,jets}} \\ &\quad + E_{x(y)}^{\text{miss,softjets}} + E_{x(y)}^{\text{miss},\mu} + E_{x(y)}^{\text{miss,CellOut}} \end{aligned} \quad (2.5)$$

The CellOut term of the above equation corresponds to calorimeter cells with energy deposits that are not associated with other objects. The soft jets term comes from cells associated to jets with p_T between 7 and 20 GeV, while the jets term comes from jets with $p_T > 20$ GeV. Because muons do not deposit significant energy in the calorimeter, the muon momentum is used for the muon term [68]. The final E_T^{miss} is calculated using equation 2.6.

$$E_T^{\text{miss}} = \sqrt{(E_x^{\text{miss}})^2 + (E_y^{\text{miss}})^2} \quad (2.6)$$

Figure 2.19 shows the resolution of the components of the E_T^{miss} under different pileup suppression techniques [20].

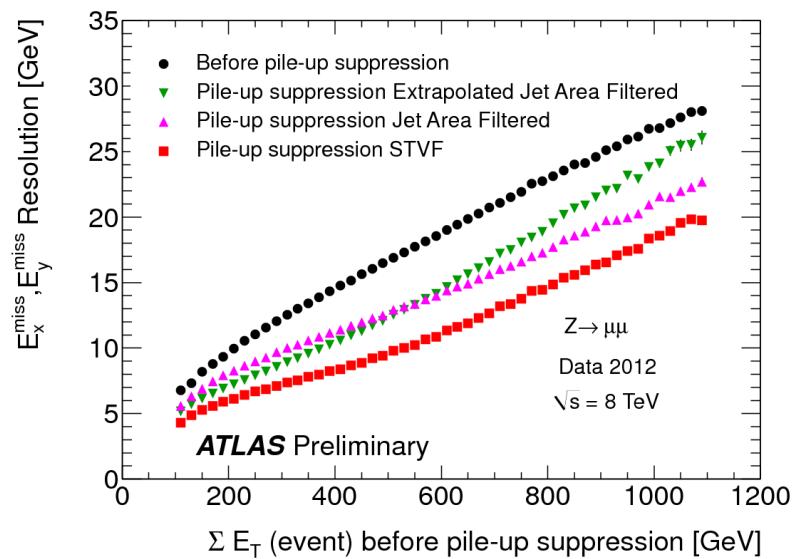


Figure 2.19: Resolution of E_T^{miss} components as a function of $\sum E_T$ before pileup suppression with different pileup techniques [20]

1161

Part II

1162

Observation and measurement of Higgs

1163

boson decays to WW^* in LHC Run I at

1164

$\sqrt{s} = 7$ and 8 TeV

*Basic research is what I am doing when I don't know
what I am doing.*

Wernher von Braun

3

1165

1166

$H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ Analysis Strategy

1167 3.1 INTRODUCTION

1168 This chapter presents an overview of the strategy for searching for a Higgs boson in the $H \rightarrow WW^* \rightarrow$
1169 $\ell\nu\ell\nu$ decay topology. Its purpose is to define in broad terms how the search and measurement are un-
1170 dertaken, before going into details on the specific sub-categories within the larger analysis. First, the
1171 properties of the Higgs signal are discussed and the associated backgrounds are presented. Next, the ob-
1172 servables used to enhance the signal to background ratio are defined. Finally, the parameters of interest
1173 in the search and measurement will be shown, along with a brief overview of the statistical treatment of
1174 the final Higgs candidates.

1175 Following this chapter, the results of three different studies within the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel
1176 are shown. Chapter 4 presents a search for Higgs boson production in gluon fusion mode and the role
1177 of the $H \rightarrow WW^*$ channel in its discovery. Chapter 5 shows the search and first observation in ATLAS
1178 of the Vector Boson Fusion (VBF) production mode of the Higgs in the $H \rightarrow WW^*$ decay channel.

1179 Finally, chapter 6 shows the combined Run 1 $H \rightarrow WW^*$ results for the measurement of the Higgs
1180 cross section and relative coupling strengths to other SM particles.

1181 3.2 THE $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ SIGNAL IN ATLAS

1182 The signal studied in this and subsequent chapters is the Higgs boson in the WW^* final state, where
1183 each W boson subsequently decays into a charged lepton and a neutrino. In its simplest decay path, the
1184 final state consists of two neutrinos and two charged leptons, each of which can be either an electron or
1185 a muon. If one or both of the W s decay to τ leptons, only leptonic decays of the τ are considered. This
1186 decay path produces additional neutrinos in the final state but still gives two charged leptons as before.
1187 Neutrinos are not detected in ATLAS, so the final state ultimately consists of two reconstructed leptons
1188 and missing transverse momentum (denoted as E_T^{miss}). Final states where both of the charged leptons
1189 are electrons or muons are referred to as the “same flavor” ($ee/\mu\mu$) final states, while those with one
1190 electron and one muon are referred to as “different flavor” ($e\mu$ or μe).

1191 While the basic final state consists of two leptons and E_T^{miss} , there can be additional objects depend-
1192 ing on the production mode of the Higgs. As described in detail in Chapter 1, if the Higgs is produced
1193 via vector boson fusion production, there will be two additional forward jets in the event. Even in gluon
1194 fusion, one or more jets can be produced through initial state radiation from the incoming gluons. Be-
1195 cause of the varying background composition as a function of jet multiplicity, each bin in this variable
1196 has its own dedicated requirements applied in the search and measurement. The $n_j = 0$ and $n_j = 1$
1197 bins are dedicated to gluon fusion production, while the $n_j \geq 2$ bin has separate dedicated searches for
1198 ggF and VBF production.

1199 Figure 3.1 shows the relative branching fractions for the $H \rightarrow WW^*$ process, calculated from the
1200 Particle Data Group values for the W and τ branching ratios[?]. The largest branching ratio is both
1201 W bosons decaying to quark pairs at 45.44%. The next largest is one W decaying leptonically and the
1202 other decaying to quarks, a branching ratio of 34.18%. In all cases, ℓ denotes either an electron or muon,
1203 and the leptonic branching ratios of the τ are included. For example, the $\ell\nu qq$ final state includes one W
1204 decaying to $e\nu$, $\mu\nu$, or $\tau\nu$. In the case of the $W \rightarrow \tau\nu$ decay, the τ lepton then decays to an electron or

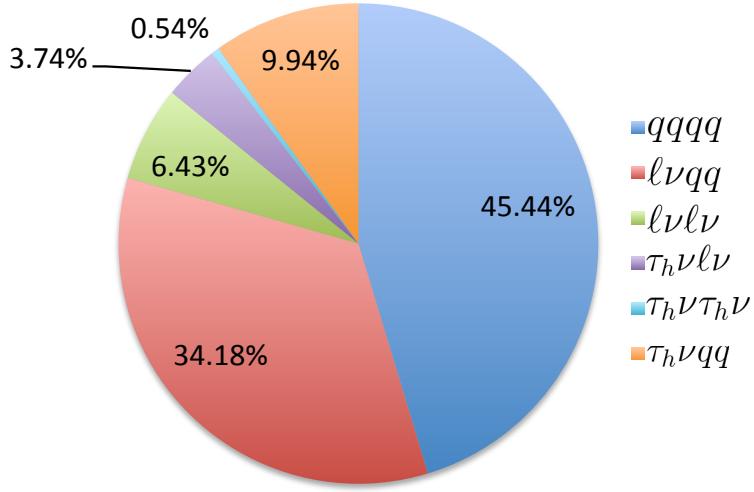


Figure 3.1: Branching ratios for a WW system. q refers to quarks. ℓ can be either an electron or muon, and the leptonic branching ratios of the τ are included. For example, the $\ell\nu qq$ final state includes one W decaying to $e\nu$, $\mu\nu$, or $\tau\nu$. τ_h refer to hadronic decays of the τ .

1205 muon via $\tau \rightarrow \nu_\tau \ell \nu_\ell$. Final states with a τ_h refer to hadronic decays of the τ . The branching ratio to the
 1206 $\ell\nu\ell\nu$ final state is 6.43%.

1207 While the $\ell\nu\ell\nu$ final state is not a large fraction of the branching ratio, there are significant advantages
 1208 in this channel. First, both the $qqqq$ and $\ell\nu qq$ channels suffer from a large QCD multijet background,
 1209 which is often difficult to model. Second, events in the the $\ell\nu\ell\nu$ channel in data can be triggered more
 1210 efficiently due to the presence of two leptons.

1211 3.3 BACKGROUND PROCESSES

1212 Many processes from the Standard Model can also produce a final state with two leptons and missing
 1213 transverse momentum . This section lists the dominant backgrounds to Higgs production. It gives gen-
 1214 eral descriptions of how the backgrounds mimic Higgs production and how they can be reduced. Ta-
 1215 ble3.1 summarizes the different processes.

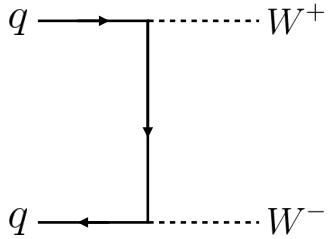


Figure 3.2: Feynman diagram for Standard Model WW production

1216 3.3.1 STANDARD MODEL WW PRODUCTION

1217 Non-resonant Standard Model diboson production, as shown in figure 3.2, is an irreducible background
 1218 to Higgs boson production in the WW final state. It produces the same exact final state objects, namely
 1219 leptonically decaying W bosons. There are no additional objects in the final state that allow for back-
 1220 ground reduction. Therefore the analysis solely relies on the correlations between the leptons to reduce
 1221 this background.

1222 3.3.2 TOP QUARK PRODUCTION

1223 Production of top quarks, either in pairs ($t\bar{t}$ production) or singly (e.g. Wt production), can also mimic
 1224 Higgs production. Because top quarks decay via $t \rightarrow Wb$, top pair production can produce a final state
 1225 with two W bosons that then decay leptonically. In this case, however, there are two additional jets from
 1226 the bottom quarks in the final state. This allows the analysis to veto on the presence of jets identified as
 1227 originating from a b in order to reduce the size of the background.

1228 Single top production can occur via s -channel, t -channel, or associated production (Wt). The mode
 1229 which most closely resembles the Higgs final state is Wt . In this case, there are two real W bosons pro-
 1230 duced, as with $t\bar{t}$. However, the decay of the single top quark will still also produce one b -jet, meaning a
 1231 b veto will reduce this background as well.

1232 Figure 3.3 shows the Feynman diagrams for $t\bar{t}$ and Wt production.

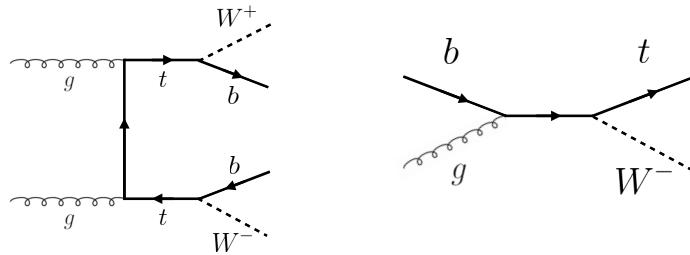


Figure 3.3: Feynman diagrams for top pair production (left) and Wt production (right)

1233 3.3.3 W +JETS BACKGROUND

1234 Single W boson production, in association with jets, is a unique background. The other background
 1235 considered so far have all included real leptons in the final state. In this case, however, only one real lep-
 1236 ton from the decay of a W exists in the final state. The second reconstructed lepton can arise from two
 1237 different cases. First, the lepton may truly be an algorithm “fake”, or a jet misidentified as a lepton by
 1238 either the electron or muon reconstruction algorithms. Second, the lepton may be a real lepton but
 1239 coming from semi-leptonic decays of particles inside the shower of the jet. This background can be re-
 1240 duced by requiring that the reconstructed lepton have little activity surrounding it in the calorimeter
 1241 (also known as an “isolated” lepton). Figure 3.4 shows the Feynman diagram for W +jets production.

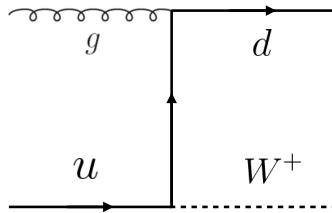


Figure 3.4: An example Feynman diagram of W +jets production

1242 3.3.4 Z/γ^* +JETS BACKGROUND

1243 Production of a Z/γ^* in association with jets (also known as Drell-Yan) is also a background to Higgs
 1244 production. In particular, the same flavor final states have a large Z +jets background, as the Z decays
 1245 into two leptons of the same flavor. (This background also enters the different flavor final state through

¹²⁴⁶ the leptonic decays of $Z \rightarrow \tau\tau$). Figure 3.5 shows the production of a Z in association with one jet. Be-
¹²⁴⁷ cause there are no neutrinos in this final state, variables like E_T^{miss} can be used to reduce the background.

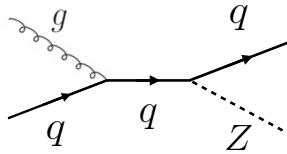


Figure 3.5: An example Feynman diagram of $Z + \text{jets}$ production

¹²⁴⁸ 3.3.5 OTHER (SUBDOMINANT) BACKGROUNDS

¹²⁴⁹ There are additional processes which contribute to the background composition but are not produced
¹²⁵⁰ as frequently as those listed already. The first of these are referred to as VV or “Other diboson” pro-
¹²⁵¹ cesses and include multiple Standard Model diboson processes, including WZ , ZZ , $W\gamma$, $W\gamma^*$, and
¹²⁵² $Z\gamma$ production. Additionally, there is background from QCD multijet production, where two jets are
¹²⁵³ misidentified as leptons.

Category	Process	Description
SM WW	$WW \rightarrow \ell\nu\ell\nu$	Real leptons and neutrinos
Top quark production	$t\bar{t} \rightarrow WbWb \rightarrow \ell\nu b\ell\nu b$ $tW \rightarrow WbW \rightarrow \ell\nu\ell\nu b$ $t\bar{b}, t\bar{q}\bar{b}$	Real leptons, untagged b s Real leptons, untagged b Untagged b , jet misidentified as lepton
Drell-Yan	$Z/\gamma^* \rightarrow ee, \mu\mu$ $Z/\gamma^* \rightarrow \tau\tau \rightarrow \ell\nu\ell\nu\nu$	“Fake” E_T^{miss} Real leptons and neutrinos
Other dibosons	$ZZ \rightarrow \ell\ell\nu\nu$ $W\gamma^*, WZ \rightarrow \ell\nu\ell\ell, ZZ \rightarrow \ell\ell\ell\ell$ $W\gamma, Z\gamma$	Real leptons and neutrinos Unreconstructed leptons γ reconstructed as e , unreconstructed lepton
$W + \text{jets}$	$Wj \rightarrow \ell\nu j$	Jet reconstructed as lepton
QCD multijet	jj	Jets reconstructed as leptons

Table 3.1: A summary of backgrounds to the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ signal

1254 **3.4 SHARED SIGNAL REGION SELECTION REQUIREMENTS**

1255 As presented in section 3.2, there are many different combinations of objects that can define a $H \rightarrow$
 1256 $WW^* \rightarrow \ell\nu\ell\nu$ final state. The multiplicity of jets and the flavor combinations of the leptons both lead
 1257 to many potential signal regions. Additionally, signal regions can be optimized separately to be sensi-
 1258 tive to the distinct production modes of the Higgs. Gluon fusion, vector boson fusion, and associated
 1259 production of a Higgs all lead to unique final state topologies. Figure 3.6 delineates the different signal
 1260 regions used in the gluon fusion and vector boson fusion $H \rightarrow WW^*$ analyses. While there are different
 1261 optimizations possible in each signal region, there are also some commonly shared selections that will be
 1262 described here.

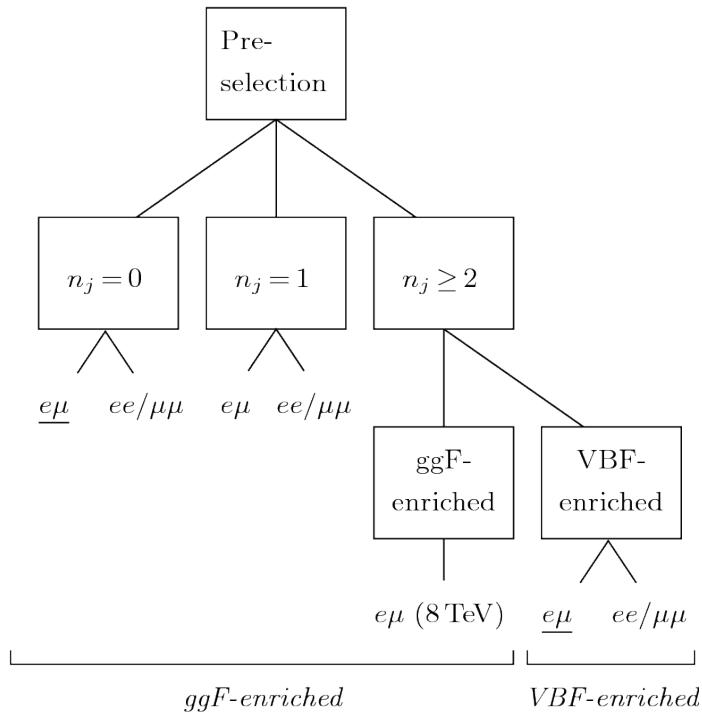


Figure 3.6: An illustration of the unique analysis signal regions[21]

1263 **3.4.1 EVENT PRE-SELECTION**

1264 Before being sorted into the distinct signal regions, basic requirements are applied on the reconstructed
 1265 objects in the event to select Higgs-like event candidates. First, two oppositely charged leptons are re-

1266 quired.

1267 Once the leptons are selected, the last requirement for event pre-selection is the presence of neutrinos.
1268 As neutrinos cannot be detected directly in ATLAS, E_T^{miss} can be used as a proxy for the combined neu-
1269 trino momentum in the transverse plane. In general, it is expected that the signal should have a harder
1270 E_T^{miss} spectrum than backgrounds, especially if those backgrounds did not contain neutrinos. One ad-
1271 ditional consideration when using E_T^{miss} is the fact that mis-measurements of objects in the detector
1272 can lead to imbalances in the transverse plane that are not due to real particles escaping the detector.
1273 One indicator that this is the case is that the E_T^{miss} vector in the transverse plane will be pointing in the
1274 same direction as the mis-measured object. Therefore, a new variable, $E_{T,\text{rel}}^{\text{miss}}$, is used in the pre-selection.
1275 $E_{T,\text{rel}}^{\text{miss}}$ is defined in equation 3.1.

$$E_{T,\text{rel}}^{\text{miss}} = \begin{cases} E_T^{\text{miss}} \sin \Delta\phi_{\text{near}} & \text{if } \Delta\phi_{\text{near}} < \pi/2 \\ E_T^{\text{miss}} & \text{otherwise,} \end{cases} \quad (3.1)$$

1276 If the closest object to the E_T^{miss} vector is within $\pi/2$ radians in the transverse plane, the E_T^{miss} is pro-
1277 jected away from this object. Otherwise, the normal E_T^{miss} vector is used. Figure 3.7 shows a graphical
1278 illustration of this concept.

1279 Once both the lepton and E_T^{miss} pre-selections are made, the analysis can be divided into different
1280 regions according to jet multiplicity.

1281 3.4.2 JET MULTIPLICITY

1282 Jet multiplicity, denoted as n_j , is used to sub-divide the analysis into its distinct signal regions. The rea-
1283 son for this is twofold. First, different jet multiplicity bins will be more or less sensitive to different Higgs
1284 production modes. For example, the $n_j \geq 2$ region is more sensitive to VBF production because of
1285 the two high momentum jets produced at matrix element level. For gluon fusion production to enter
1286 this bin, two initial state radiation jets must be emitted. Second, background composition varies greatly
1287 in different bins of n_j . Figure 3.8 shows the jet multiplicity in both the different flavor and same flavor
1288 regions. It also shows the background composition in the bins of n_b . There are a few clear trends from

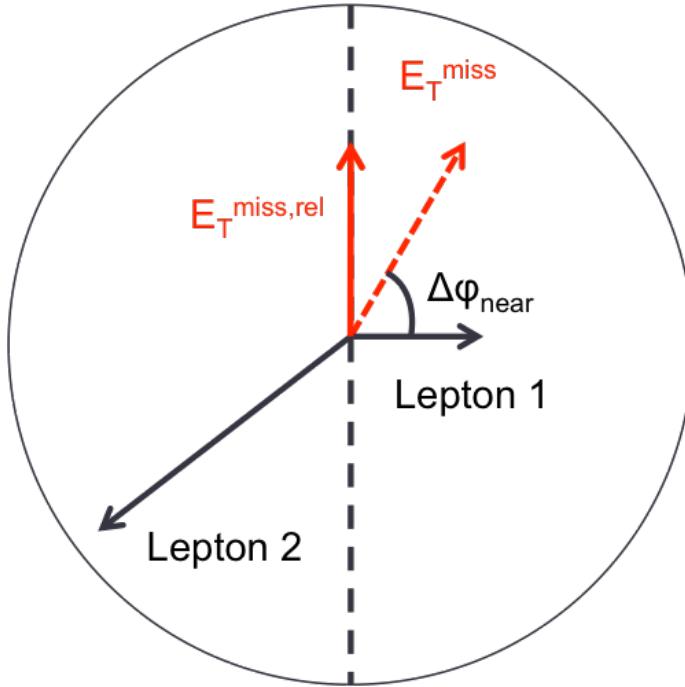


Figure 3.7: A graphical illustration of the $E_{T,\text{rel}}^{\text{miss}}$ calculation

1289 this distribution. The first is that the Drell-Yan background dominates in the same flavor channels for
 1290 $n_j \leq 1$. Second, the top background becomes a clear contributor to the total background for $n_j \geq 1$.
 1291 Lastly, the SM WW production dominates in the $n_j = 0$ bin, as it is an irreducible background to
 1292 $H \rightarrow WW^*$ production. Because of these distinct features, each jet multiplicity bin is treated separately.

1293 3.5 BACKGROUND REDUCTION IN SAME-FLAVOR FINAL STATES

1294 As described in section 3.4.2, the background composition of the same flavor final states is unique to that
 1295 of the different flavor states. In particular, Drell Yan processes play a much larger role because the Z/γ^*
 1296 decays to same flavor leptons. Because real neutrinos are absent in the Z/γ^* decays to ee and $\mu\mu$, a re-
 1297 quirement on E_T^{miss} should largely reduce the background. However, as this section will demonstrate,
 1298 with increasing pileup conditions the resolution of the calorimeter-based E_T^{miss} degrades greatly. There-
 1299 fore, two new variables for Z/γ^* background reduction are constructed and described in this section.

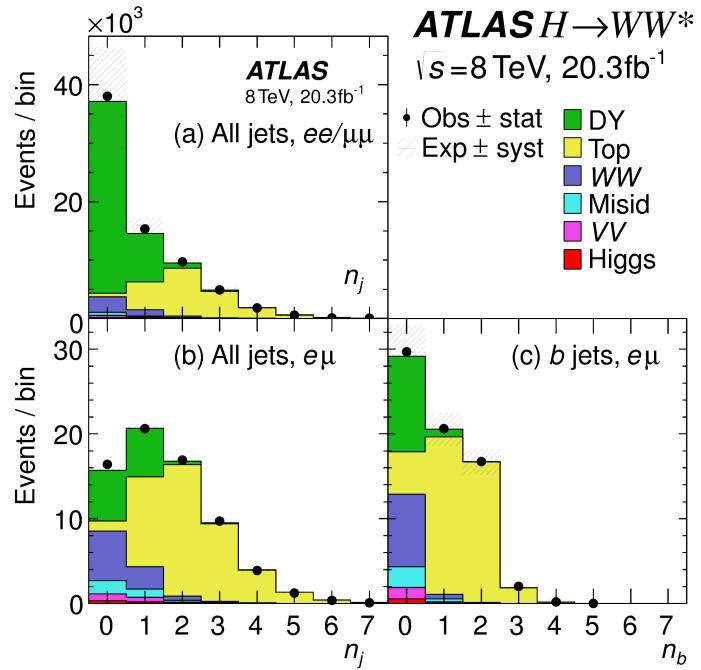


Figure 3.8: Predicted backgrounds (compared with data) as a function of n_j (a and b) and n_b (c)

3.5.1 PILEUP AND E_T^{miss} RESOLUTION

Secondary interactions of protons in the colliding bunches of the LHC (known as pileup interactions, described in detail in Chapter 2) deposit energy into the ATLAS calorimeter on top of the energy that comes from the hard scatter process that is being searched for or analyzed. The calculation of E_T^{miss} is fundamentally Poissonian, as summing up all of the energy deposits in individual calorimeter cells or clusters is similar to a counting experiment. Thus, the energy resolution scales as \sqrt{E} , just as the error on a mean of N in a Poisson distribution is \sqrt{N} . As more energy is deposited in the calorimeter, the E_T^{miss} resolution degrades, meaning that the E_T^{miss} resolution is particularly sensitive to LHC instantaneous luminosity conditions.

Figure 3.9 shows an event display of a $Z/\gamma^* + \text{jets}$ event candidate with the twenty-five reconstructed primary vertices. This display illustrates that while the interaction of interest only has tracks coming from the hardest primary vertex, all of the secondary interactions will deposit energy in the calorimeter as well.

Figure 3.10 shows the RMS of the E_T^{miss} distribution in $Z \rightarrow \mu\mu$ events (where there are no real neu-

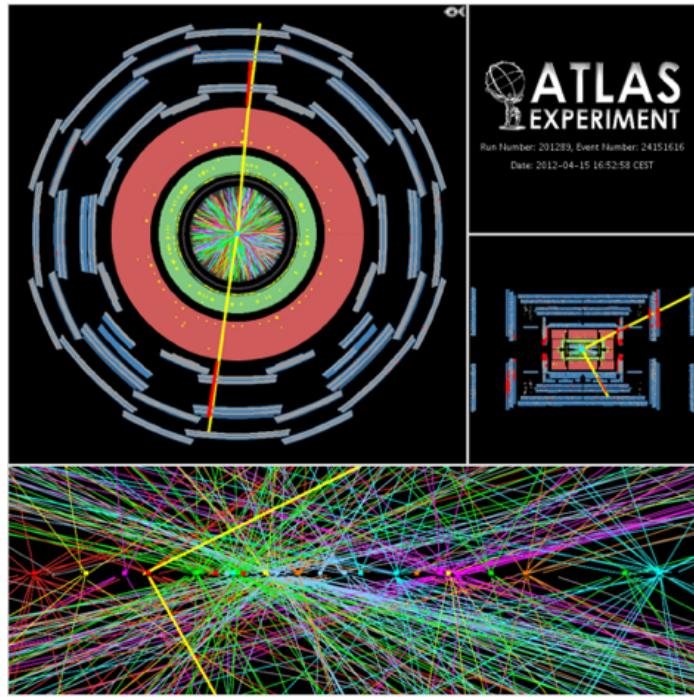


Figure 3.9: An event display of a Z/γ^* + jets event illustrating the effect of pileup interactions

trinos) as a function of the number of the average number of interactions. Under 2011 LHC conditions, this RMS was approximately 9 GeV, while under 2012 running conditions the resolution worsened to 12 GeV. This worsening dilutes the E_T^{miss} variable's ability to reduce the Z/γ^* background.

3.5.2 TRACK-BASED DEFINITIONS OF MISSING TRANSVERSE MOMENTUM

Because the increasing number of secondary proton-proton interactions degrades calorimeter-based E_T^{miss} resolution, a new variable using only contributions from the primary interaction vertex is necessary to further reduce the Z/γ^* background. While it is not possible to associate calorimeter energy deposits with a particular vertex, individual charged particle tracks in the Inner Detector are associated to unique vertices. Thus, two track-based definitions of missing transverse momentum , using only tracks coming from the primary vertex in the event, are used in the analysis. The simplest variable, $p_T^{\text{miss}(\text{trk})}$, is the vectorial sum of the p_T of all of the tracks from the primary vertex and the selected leptons (excluding the tracks associated with the selected leptons to avoid double counting). This is defined in equa-

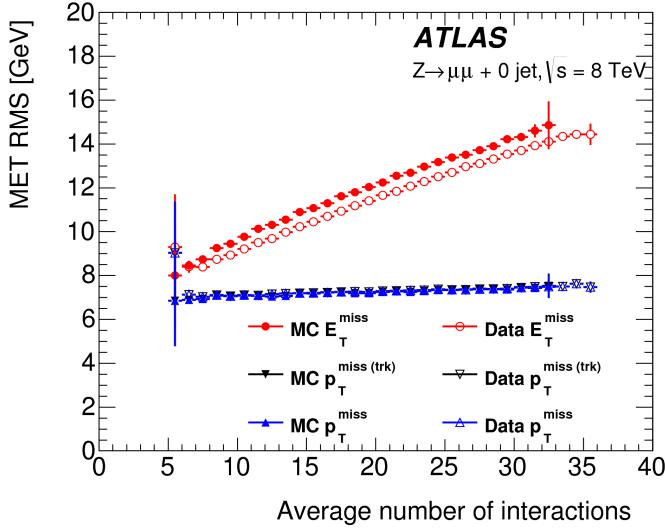


Figure 3.10: The RMS of different missing transverse momentum definitions as a function of the average number of interactions per bunch crossing

tion 3.2.

$$p_T^{\text{miss}(\text{trk})} = - \left(\sum_{\text{selected leptons}} p_T + \sum_{\text{other tracks}} p_T \right), \quad (3.2)$$

In events with hard jets, a better resolution on the missing transverse momentum is obtained by including the calorimeter based measurement of the hard jets rather than the track based measurements. Thus, another variable, p_T^{miss} , is defined, using the nominal measurements of p_T for the selected leptons and jets and using tracks rather than calorimeter clusters for the soft component of the missing transverse momentum. This is defined in equation 3.3.

$$p_T^{\text{miss}} = - \left(\sum_{\text{selected leptons}} p_T + \sum_{\text{selected jets}} p_T + \sum_{\text{other tracks}} p_T \right), \quad (3.3)$$

Figure 3.10 illustrates that these two new variables accomplish their intended purpose. The resolution as a function of mean number of interactions for both $p_T^{\text{miss}(\text{trk})}$ and p_T^{miss} is much flatter compared to the dependence for E_T^{miss} .

Figure 3.11a shows the difference between the true and reconstructed values of missing transverse mo-

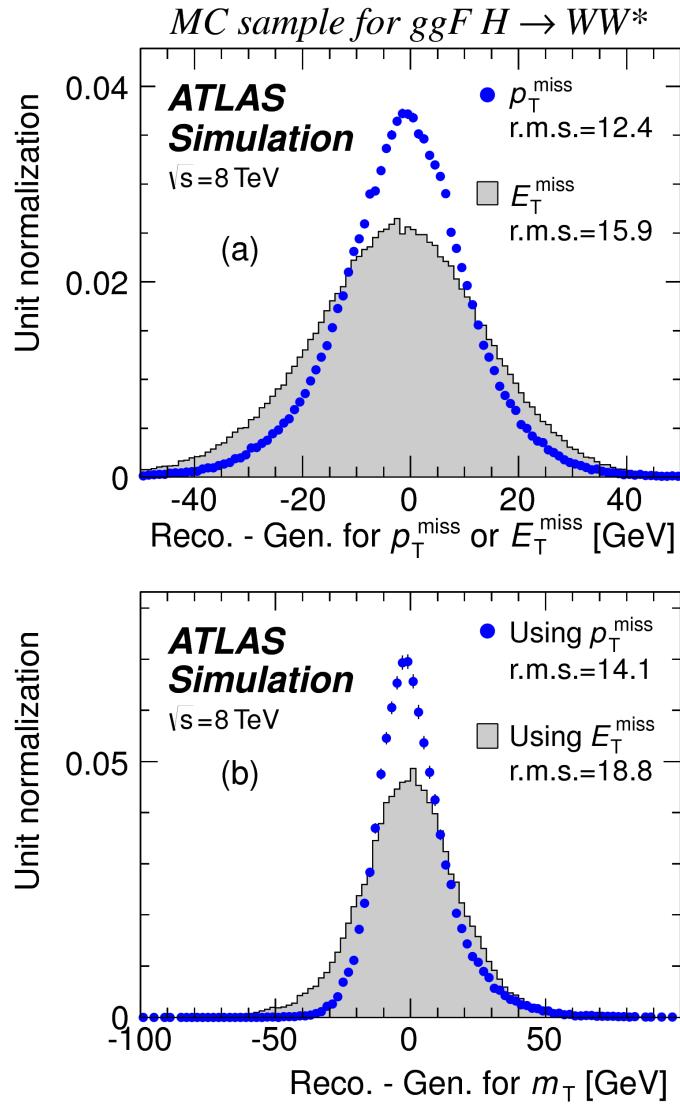


Figure 3.11: The difference between the true and reconstructed values of the missing transverse momentum (a) and m_T (b) in a gluon fusion signal sample

¹³³⁶ momentum using both the track-based p_T^{miss} and calorimeter based E_T^{miss} . The RMS of the distribution
¹³³⁷ improves by 3.5 GeV when using p_T^{miss} .

¹³³⁸ 3.5.3 DISTINGUISHING Z/γ^* +JETS AND $H \rightarrow WW^*$ TOPOLOGIES

¹³³⁹ The track-based definitions of missing transverse momentum were constructed to mitigate degrading
¹³⁴⁰ performance as a function of pileup. However, an additional variable can be constructed to exploit kine-

matic and topological differences between the Z/γ^* background and $H \rightarrow WW^*$ signal. Because there
 1341 are no real neutrinos in the final state (in the case of $Z/\gamma^* \rightarrow ee, \mu\mu$ decays), the dilepton system of a
 1342 Z/γ^* will be balanced with the jets produced in the hard scatter. A new variable, f_{recoil} , is constructed
 1343 to estimate the balance between the dilepton system and the jets in the quadrant opposite the dilepton
 1344 vector in the transverse plane. It is defined in equation 3.4. The numerator of f_{recoil} is the magnitude of
 1345 the vectorial sum of the p_T of jets in the quadrant opposite the dilepton system, weighted by each jet's Jet
 1346 Vertex Fraction (JVF, described in chapter 2). The denominator is the magnitude of the dilepton p_T .
 1347

$$f_{\text{recoil}} = \left| \sum_{\text{jets } j \text{ in } \wedge} \text{JVF}_j \cdot \mathbf{p}_T^j \right| / p_T^{\ell\ell}. \quad (3.4)$$

1348 Figure 3.12 shows a shape comparison of the distribution of f_{recoil} in a simulated $Z/\gamma^* + \text{jets}$ sample,
 1349 a $H \rightarrow WW^*$ signal sample, and other backgrounds that contain real neutrinos. The $Z/\gamma^* + \text{jets}$ events
 1350 tend to be more balanced between the dilepton system and recoiling jets, while the processes containing
 1351 real neutrinos are less balanced in the transverse plane. Thus, a requirement on f_{recoil} will also reduce
 1352 the $Z/\gamma^* + \text{jets}$ background while maintaining a good signal efficiency.

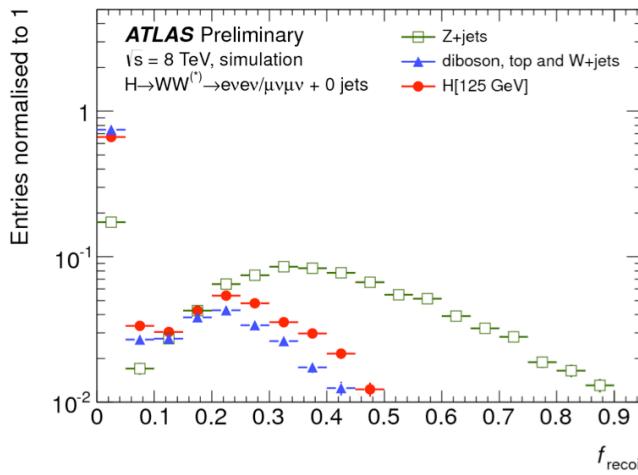


Figure 3.12: Comparison of f_{recoil} distributions for $Z/\gamma^* + \text{jets}$, $H \rightarrow WW^*$, and other backgrounds with real neutrinos.

1353 3.5.4 OPTIMIZING BACKGROUND REDUCTION SELECTION REQUIREMENTS

1354 The requirements on $p_T^{\text{miss(trk)}}$ and f_{recoil} used to reduce the Z+jets background must be optimized
 1355 to maximize their efficacy. Figure 3.13 shows an early attempt to optimize the combination of the two
 1356 requirements in the gluon fusion zero jet bin. Each bin shows the expected signal significance if the
 1357 $p_{T,\text{rel}}^{\text{miss(trk)}}$ is required to be greater than the left edge of the bin and the f_{recoil} is required to be less than
 1358 the top edge of the bin. The figure shows that the best signal significance comes from requiring low val-
 1359 ues of $f_{\text{recoil}} (< 0.05)$ and $p_{T,\text{rel}}^{\text{miss(trk)}}$ values greater than 45 GeV.

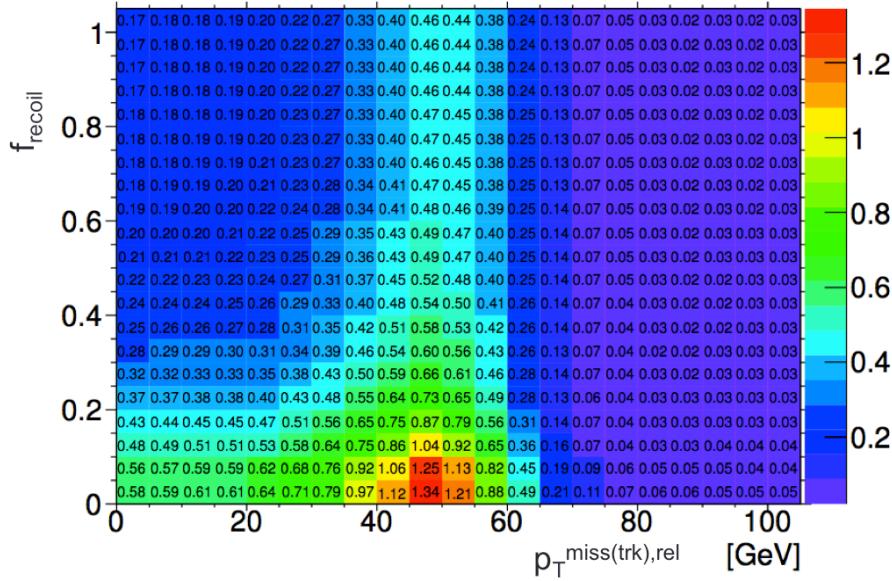


Figure 3.13: Signal significance as a function of required value for f_{recoil} and $p_{T,\text{rel}}^{\text{miss(trk)}}$ in the ggF $H \rightarrow WW^*$ with $n_j = 0$

1360 3.6 PARAMETERS OF INTEREST AND STATISTICAL TREATMENT

1361 As with any search or measurement, there are particular parameters of the Higgs that the $H \rightarrow WW^*$
 1362 analysis is interested in measuring. In this case, the parameters of interest are the mass of the Higgs bo-
 1363 son and its production cross section. Because the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ process does not have a closed
 1364 final state, it is not possible to measure the full invariant mass of the particle that may have produced
 1365 the final state. However, a proxy for the invariant mass using transverse plane information can be de-

1366 fined. This is described in more detail in section 3.6.1. The second parameter of interest is the ratio of the
 1367 measured cross section to that expected from the Standard Model Higgs, which is denoted a μ . This is
 1368 defined in equation 3.5.

$$\mu = \frac{\sigma}{\sigma_{\text{SM}}} \quad (3.5)$$

1369 All of the likelihoods used in the statistical analysis of the final signal region events are paramaterized
 1370 as a function of μ . μ is a natural variable for hypothesis testing, as $\mu = 0$ corresponds to a background
 1371 only hypothesis and $\mu = 1$ corresponds exactly to a Standard Model Higgs.

1372 3.6.1 TRANSVERSE MASS

1373 Because the longitudinal information about the neutrinos is not attainable, the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
 1374 analysis uses a mass variable, the transverse mass, that exploits information in the transverse plane as a
 1375 proxy for the full invariant mass. The transverse mass is defined in equation 3.6.

$$m_T = \sqrt{(E_T^{\ell\ell} + p_T^{\text{miss}})^2 - |\mathbf{p}_T^{\ell\ell} + \mathbf{p}_T^{\text{miss}}|^2}, \quad (3.6)$$

1376 Here the $E_T^{\ell\ell}$ and $p_T^{\ell\ell}$ are the transverse energy and momentum of the dilepton system, while p_T^{miss} is
 1377 a proxy for the transverse momentum of the di-neutrino system. The track-based p_T^{miss} is used in the m_T
 1378 rather than the calorimeter based E_T^{miss} because it has a better resolution on the true transverse mass.
 1379 Figure 3.11b shows the improvement in the RMS of the difference between the true and reconstructed
 1380 transverse mass in a ggF signal sample. The RMS improves by 4.7 GeV using p_T^{miss} in the m_T calcula-
 1381 tion.

1382 3.6.2 STATISTICAL TREATMENT^{*}

1383 LIKELIHOOD FUNCTION

1384 The statistical analysis of final event candidates is framed as a hypothesis test, where the null hypothe-
1385 sis is background-only (no Standard Model Higgs). The first step in the analysis is to form a likelihood
1386 function for the data. In its simplest form, this likelihood is the probability of observing the number
1387 of events seen in the final signal region given knowledge of the signal strength. Because observation of
1388 events is fundamentally a Poisson counting experiment, this simple likelihood can be expressed as a Pois-
1389 son probability of observing N events given a total number of predicted signal and background events.
1390 This basic likelihood is shown in equation 3.7.

$$\mathcal{L}(\mu) = P(N|\mu S + B) \quad (3.7)$$

1391 Here, P is the Poisson probability density function, N is the total number of observed events, μ is
1392 the signal strength, S is the predicted number of signal events, and B is the predicted number of back-
1393 ground events.

1394 In particle physics, certain background estimates are commonly normalized in so-called “control” re-
1395 gions and those predictions are scaled by the same normalization factor in the signal region. This leads
1396 to a slightly more complicated likelihood, which is a function of both the signal strength and the back-
1397 ground normalization. This is shown in equation 3.8.

$$\mathcal{L}(\mu, \theta) = P(N|\mu S + \theta B) P(N_{\text{CR}}|\theta B_{\text{CR}}) \quad (3.8)$$

1398 Here, θ is a so-called “nuisance parameter”, a parameter that is not a primary parameter of interest but
1399 still enters the likelihood. The second Poisson term adds an extra term to the likelihood, enforcing the
1400 fact that the background normalization must be consistent with the number of observed events in data
1401 in the control region, N_{CR} .

1402 So far, these two formulations of likelihoods have assumed a single signal region and do not take into

*Many thanks to Aaron Armbruster, whose thesis[69] inspired parts of this section.

1403 account any shape information of potential discriminating variables. The $H \rightarrow WW^*$ analysis is di-
 1404 vided into many different categories, and we can perform the same counting experiment described above
 1405 in each individual category. As mentioned in section 3.6.1, the transverse mass is used as the primary dis-
 1406 criminating variable in many of the $H \rightarrow WW^*$ sub-analyses, so additionally we can perform the same
 1407 counting experiment in each bin of the m_T distribution to incorporate some shape information. Thus,
 1408 the total likelihood becomes a product over signal regions and bins of the m_T distribution. Finally, there
 1409 are usually many backgrounds that are normalized in control regions, so the new formulation of the like-
 1410 lihood takes this into account as well by including a product over control regions in the second Poisson
 1411 term. All of these modifications are shown in equation 3.9.

$$\mathcal{L}(\mu, \theta) = \prod_{\substack{\text{SRs } i \\ \text{bins } b}} P \left(N_{ib} \middle| \mu S_{ib} + \sum_{\text{bkg } k} \theta_k B_{kib} \right) \prod_{\text{CRs } l} P \left(N_l \middle| \sum_{\text{bkg } k} \theta_k B_{kl} \right) \quad (3.9)$$

1412 The final step to get the full likelihood used in the analysis is to add nuisance parameters for the sys-
 1413 tematic uncertainties. In cases where the uncertainty does not affect the shape of m_T bin-by-bin, each
 1414 systematic uncertainty ϵ is allowed to affect the expected event yields through a linear response function
 1415 of the nuisance parameter, namely $\nu(\theta) = (1 + \epsilon)^\theta$. If instead the uncertainty does affect the shape,
 1416 the effect is instead parameterized by $\nu_b(\theta) = 1 + \epsilon_b \theta$. The value of the nuisance parameters for the
 1417 systematic uncertainty are constrained with a Gaussian term that is added to the likelihood as well. This
 1418 is of the form $g(\delta|\theta) = e^{-(\delta-\theta)^2/2}/\sqrt{2\pi}$, where δ is the central value and θ is a nuisance parameter.
 1419 Finally, a last term is added to account for the statistical uncertainty in the Monte Carlo samples used,
 1420 which adds an additional poisson term. The full likelihood used in the final statistical analysis is defined

¹⁴²¹ in equation 3.10.

$$\begin{aligned} \mathcal{L}(\mu, \boldsymbol{\theta}) = & \prod_{\substack{\text{SRs i} \\ \text{bins b}}} P\left(N_{ib} \middle| \mu S_{ib} \cdot \prod_{\substack{\text{sig.} \\ r}} \nu_{br}(\theta_r) + \sum_{\text{bkg k}} \theta_k B_{kib} \cdot \prod_{\substack{\text{bkg.} \\ \text{syst.} \\ s}} \nu_{bs}(\theta_s)\right) \\ & \cdot \prod_{\text{CRs l}} P\left(N_l \middle| \sum_{\text{bkg k}} \theta_k B_{kl}\right) \\ & \cdot \prod_{\substack{\text{syst.} \\ t}} g(\delta_t | \theta_t) \cdot \prod_{\text{bkg k}} P(\xi_k | \zeta_k \theta_k) \end{aligned} \quad (3.10)$$

¹⁴²² In the fourth term of the equation, quantifying uncertainty due to finite Monte Carlo sample size, ξ
¹⁴²³ represents the central value of the background prediction, θ is the associated nuisance parameter, $\zeta =$
¹⁴²⁴ $(B/\delta B)^2$, where δB is the statistical uncertainty of B .

¹⁴²⁵ The best fit value of the signal strength μ is determined by finding the values of μ and $\boldsymbol{\theta}$ that maxi-
¹⁴²⁶ mize the likelihood, while setting $\delta = 0$ and $\xi = \zeta$.

¹⁴²⁷ Once the likelihood is defined, a test statistic must be built for use in hypothesis testing.

¹⁴²⁸ TEST STATISTIC

¹⁴²⁹ To distinguish whether the data match a background only or background and signal hypothesis, a test
¹⁴³⁰ statistic must be used. The $H \rightarrow WW^*$ analysis used the profile likelihood technique[70]. The first step
¹⁴³¹ in formulating this test statistic is to define the profile likelihood ratio, shown in equation 3.11.

$$\lambda(\mu) = \frac{\mathcal{L}(\mu, \hat{\theta}_\mu)}{\mathcal{L}(\hat{\mu}, \hat{\theta})} \quad (3.11)$$

¹⁴³² Here $\hat{\theta}_\mu$ is the value of θ that maximizes the likelihood for the choice of μ being tested. Additionally,
¹⁴³³ $\hat{\theta}$ and $\hat{\mu}$ represent the values of θ and μ that gives the overall maximum value of the likelihood.

¹⁴³⁴ Once this is defined, a test statistic q_μ is constructed. This is shown in equation 3.12.

$$q_\mu = -2 \ln \lambda(\mu) \quad (3.12)$$

1435 A higher value of q_μ means that the data are more incompatible with the hypothesized value of μ , and
1436 q_0 then corresponds to the value of the test statistic for the background only hypothesis. A p_0 value is
1437 then defined to quantify the compatibility between the data and the null hypothesis. The p_0 value is the
1438 probability of obtaining a value of q_0 larger than the observed value, and this is shown in equation 3.13.

$$p_0 = \int_{q_0^{\text{obs}}}^{\infty} f(q_\mu | \mu = 0) dq_\mu \quad (3.13)$$

1439 Here $f(q_\mu)$ is the probability distribution function of the test statistic. Finally, the p_0 value can be
1440 converted into a signal significance, using the formula in equation 3.14, or the one-sided tail of the Gaus-
1441 sian distribution.

$$Z_0 = \sqrt{2} \operatorname{erf}^{-1}(1 - 2p_0) \quad (3.14)$$

1442 The threshold for discovery used in particle physics is $Z_0 \geq 5$, more commonly known as a value of
1443 5σ .

The real voyage of discovery consists not in seeking new landscapes, but in having new eyes.

Marcel Proust

4

1444

1445

The discovery of the Higgs boson and the 1446 role of the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel

1447 4.1 INTRODUCTION

1448 This chapter presents the results of the search for the Higgs boson in 4.8 fb^{-1} collected at $\sqrt{s} = 7 \text{ TeV}$
1449 and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. The results of three searches at $\sqrt{s} = 8 \text{ TeV}$ in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
1450, $H \rightarrow \gamma\gamma$, and $H \rightarrow ZZ \rightarrow 4\ell$ channels are combined with results of searches at $\sqrt{s} = 7 \text{ TeV}$
1451 in the same search channels (as well as the $H \rightarrow \tau\tau$ production and associated production searches for
1452 $H \rightarrow b\bar{b}$). The results of this combination are a 5.9σ detection of a new particle consistent with a Higgs
1453 boson. Rather than going into detail for all of the different Higgs decay searches, this chapter will discuss
1454 the three most sensitive channels and in particular focus on $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$. While the focus is
1455 on WW^* , some of the ZZ^* and $\gamma\gamma$ results are shown for completeness. The results not discussed here

1456 can be found in the ATLAS Higgs discovery publication[[22](#)].

1457 **4.2 DATA AND SIMULATION SAMPLES**

1458 The data sample used for the following results was taken in 2011 and 2012 at center of mass energies of 7
1459 and 8 TeV, respectively, with 4.8 fb^{-1} collected at 7 TeV and 5.8 fb^{-1} collected at 8 TeV. Higgs pro-
1460 duction in the gluon fusion and vector boson fusion modes is modeled with **POWHEG** for the hard scat-
1461 tering event and **PYTHIA** for the showing and hadronization. Associated production of a Higgs with a
1462 vector boson or top quarks is modeled via **PYTHIA**.

1463 Table 4.1 shows the Monte Carlo generators used for modeling the signal and background processes
1464 relevant for the three analyses to be discussed.

Process	Generator
$\text{ggF}, \text{VBF } H$	POWHEG + PYTHIA
$WH, ZH, t\bar{t}H$	PYTHIA
$W + \text{jets}, Z/\gamma^* + \text{jets}$	ALPGEN + HERWIG
$t\bar{t}, tW, tb$	MC@NLO + HERWIG
tqb	ACERMC + PYTHIA
$q\bar{q} \rightarrow WW$	MC@NLO + HERWIG
$gg \rightarrow WW$	GG2WW+ HERWIG
$q\bar{q} \rightarrow ZZ$	POWHEG + PYTHIA
$gg \rightarrow ZZ$	GG2ZZ+ HERWIG
WZ	MADGRAPH+ PYTHIA , HERWIG
$W\gamma + \text{jets}$	ALPGEN + HERWIG
$W\gamma^*$	MADGRAPH+ PYTHIA
$q\bar{q}/gg \rightarrow \gamma\gamma$	SHERPA

Table 4.1: Monte carlo generators used to model signal and background for the Higgs search[[22](#)].

1465 **4.3 $H \rightarrow WW \rightarrow e\nu\mu\nu$ SEARCH**

1466 The $H \rightarrow WW \rightarrow e\nu\mu\nu$ search is unique compared to the ZZ and $\gamma\gamma$ channels. The Higgs mass can-
1467 not be fully reconstructed due to the presence of neutrinos in the final state, so the transverse mass m_T
1468 is used as the final discriminating variable. Compared to the other channels, there are more backgrounds

1469 here as well, as discussed in chapter 3. The same flavor final states are excluded from this search due to
1470 high pileup in the 8 TeV dataset.

1471 4.3.1 EVENT SELECTION

1472 The analysis requires two opposite charge isolated leptons, with the leading (sub-leading) lepton required
1473 to have $p_T > 25(15)$ GeV. The events are separated into different signal regions depending on which
1474 flavor of lepton is leading ($e\mu$ for leading electron, μe for leading muon). Strict lepton quality cuts are
1475 applied to the sample to reduce backgrounds from fake leptons.

1476 Jets are reconstructed with the anti- k_T algorithm with a radius parameter $R = 0.4$. The jets are
1477 required to have $p_T > 25$ GeV and $|eta| < 4.5$, with jets in the tracking volume required to have a
1478 jet vertex fraction of 0.5 and jets in the forward region required to have $p_T > 30$ GeV. The analysis is
1479 separated into three different signal regions based on jet multiplicity: $n_j = 0, 1, \geq 2$.

1480 To indicate the presence of neutrinos in the event, a requirement of $E_{T,\text{rel}}^{\text{miss}} > 25$ GeV is made*. This
1481 requirement significantly reduces the QCD multijet and Z/γ^* + jets backgrounds. Figure 4.1 shows the
1482 distribution of n_j in data and simulation after applying these “pre-selection” requirements.

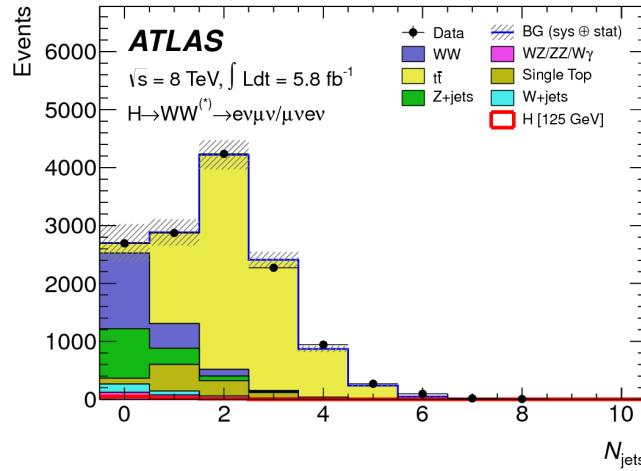


Figure 4.1: Jet multiplicity distribution in data and MC after applying lepton, jet, and $E_{T,\text{rel}}^{\text{miss}}$ selections. The WW and top backgrounds have been normalized using control samples, and the hashed band indicates the total uncertainty on the prediction. [22]

*For the definition of $E_{T,\text{rel}}^{\text{miss}}$, see chapter 3

1483 Additional selections are applied to require the dilepton topology to correspond to that of a SM
1484 Higgs. The requirements are presented here - more detailed discussion on the motivation for each re-
1485 quirement is saved for Chapter 5. In all of the jet multiplicity channels, the dilepton system is required to
1486 have a small gap in azimuthal angle, $\Delta\phi_{\ell\ell} < 1.8$. Similarly, the $m_{\ell\ell}$ is required to be less than 50 GeV
1487 in the lower jet multiplicity channels and less than 80 GeV in the $n_j \geq 2$ channel. In the $n_j = 0$
1488 channel, the magnitude of the dilepton p_T , $p_T^{\ell\ell}$, is required to be greater than 30 GeV.

1489 In the higher jet multiplicity channels ($n_j \geq 1$), the top background is a more important component
1490 and must be reduced. The total transverse momentum p_T^{sum} is thus required to be less than 30 GeV.
1491 Additionally, the di- τ invariant mass $m_{\tau\tau}$ (dilepton mass computed under the assumption that the neu-
1492 trinos from the τ decay are emitted collinear to the charged leptons) is used to reject $Z \rightarrow \tau\tau$ events by
1493 requiring $|m_{\tau\tau} - m_Z| > 25$ GeV. These variables are also discussed in more detail in Chapter 5.

1494 In the $n_j \geq 2$ channel, requirements are made to isolate the VBF contribution to Higgs production.
1495 The kinematics of the two leading jets are used to make these requirements. In particular, the event must
1496 have $\Delta y_{jj} > 3.8$ and $m_{jj} > 500$ GeV, along with a veto on having any additional jets with rapidity
1497 between the two leading jets.

1498 4.3.2 BACKGROUND ESTIMATION

1499 The details of the background estimation techniques used in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis are dis-
1500 cussed in section 5.5. As that section refers to a later iteration of the analysis, a general discussion is given
1501 here for completeness. The dominant backgrounds are SM WW production and top (both pair and
1502 single) production, and these backgrounds have their normalizations estimated from dedicated control
1503 regions while their shapes are taken from simulation.

1504 The control sample for the Standard Model WW background is defined by making the same require-
1505 ments as the signal region with the $m_{\ell\ell}$ requirement inverted (now requiring $m_{\ell\ell} > 80$ GeV) and
1506 removing the $\Delta\phi_{\ell\ell}$ requirement. This creates a control sample that is 70% (40%) pure in the 0(1)-jet
1507 region. The correction to the pure MC-based background estimate is quantified by defining a normal-
1508 ization factor β which is the ratio of the data yield to the MC yield ($N_{\text{data}}/N_{\text{MC}}$) in this control sample.

1509 Table 4.2 shows the WW normalization factors in the $n_j = 0$ and $n_j = 1$ bins (the $n_j \geq 2$ estimate is
 1510 taken directly from MC).

n_j	β_{WW}	β_t
= 0	1.06 ± 0.06	1.11 ± 0.06
= 1	0.99 ± 0.15	1.11 ± 0.05
≥ 2	-	1.01 ± 0.26

Table 4.2: Normalization factors (ratio of data and MC yields in a control sample) for the Standard Model WW and top backgrounds in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis [22]. Only statistical uncertainties are shown.

1511 The top background estimate is also computed separately in each jet multiplicity bin. In the $n_j = 0$
 1512 channel, the background is first normalized using data after pre-selection requirements with no selection
 1513 on n_j . Then, a dedicated b -tagged control sample is used to evaluate the ratio of one-jet to two-jet events
 1514 in data. The details of this technique are shown in reference [71]. In the $n_j = 1$ and the $n_j \geq 2$ regions,
 1515 the top background is normalized in a control sample where the signal region selections are applied,
 1516 but the b -jet veto is reversed and the Higgs topology requirements on $m_{\ell\ell}$ and $\Delta\phi_{\ell\ell}$ are removed. The
 1517 resulting normalization factors for these techniques are shown in table 4.2.

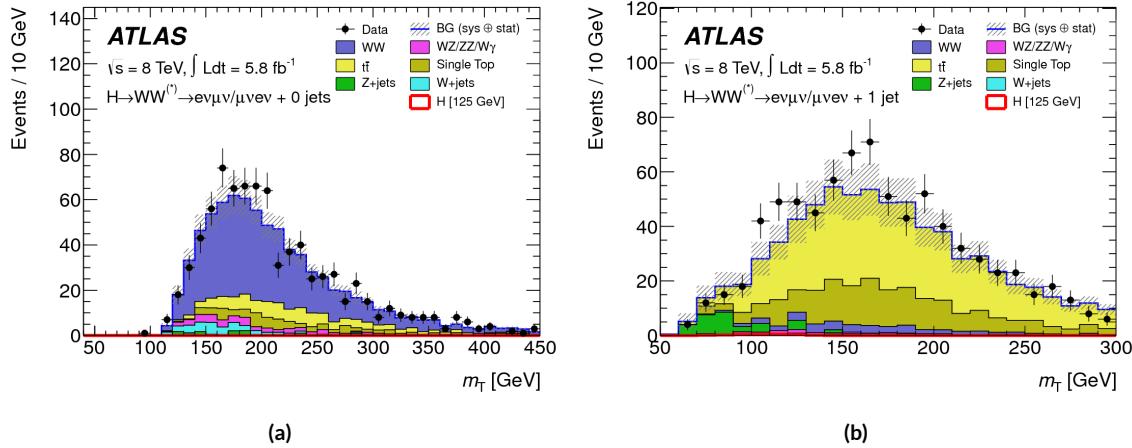


Figure 4.2: Comparison of m_T between data and simulation in the $n_j = 0$ WW (a) and $n_j = 1$ top (b) control samples [22]

1518 The control samples which are used for background normalization can also be used to validate the
 1519 modeling of the m_T distribution for each background. Figure 4.2 shows the comparison between data

1520 and MC for the m_T distribution after correcting the normalization of the backgrounds in the WW and
 1521 top control regions. Good agreement between data and simulation is seen in both cases.

1522 The $W + \text{jets}$ background estimate is taken entirely from data using a control sample with one well
 1523 reconstructed lepton and one anti-identified lepton. All other backgrounds are taken purely from simu-
 1524 lation.

1525 4.3.3 SYSTEMATIC UNCERTAINTIES

1526 The systematic uncertainties that have the largest impact on the analysis are the theoretical uncertainties
 1527 associated with the signal cross section, and these are shared with the ZZ^* and $\gamma\gamma$ channels. The un-
 1528 certainties resulting from variations of the QCD scale are $+7\% / -8\%$ on the final singal yield. Those
 1529 coming from variations of the parton distribution function (PDF) used in the simulation add a $\pm 8\%$
 1530 uncertainty on the yield. The uncertainties on the branching ratios of the Higgs are $\pm 5\%$.

1531 The main experimental uncertainties come from variations of the jet energy scale (JES), jet energy
 1532 resolution, pile-up, E_T^{miss} , b -tagging efficiency, $W + \text{jets}$ background estimate, and integrated luminosity.
 1533 For more details, see reference [22].

1534 4.3.4 RESULTS

1535 Table 4.3 shows the signal and background yields in the final signal region after normalizing the back-
 1536 grounds according to the methods described above.

	$n_j = 0$	$n_j = 1$	$n_j \geq 2$
Signal	20 ± 4	5 ± 2	0.34 ± 0.07
WW	101 ± 13	12 ± 5	0.10 ± 0.14
Other dibosons	12 ± 3	1.9 ± 1.1	0.10 ± 0.10
$t\bar{t}$	8 ± 2	6 ± 2	0.15 ± 0.10
Single top	3.4 ± 1.5	3.7 ± 1.6	-
$Z/\gamma^* + \text{jets}$	1.9 ± 1.3	0.10 ± 0.10	-
$W + \text{jets}$	15 ± 7	2 ± 1	-
Total background	142 ± 16	26 ± 6	0.35 ± 0.18
Observed in data	185	38	0

Table 4.3: Data and expected yields for signal and background in the final $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ signal region.
 Uncertainties shown are both statistical and systematic. [22]

1537 Figure 4.3 shows the m_T distribution in the $n_j \leq 1$ channels for 8 TeV data. (No events are observed
 1538 in data in the $n_j \geq 2$ channels in this dataset). The excess shown here relatively flat as a function of
 1539 hypothesized Higgs mass. The combined 7 and 8 TeV data gives an excess with local significance of 2.8σ
 1540 with an expected significance of 2.3σ , corresponding to a μ measurement of 1.3 ± 0.5 .

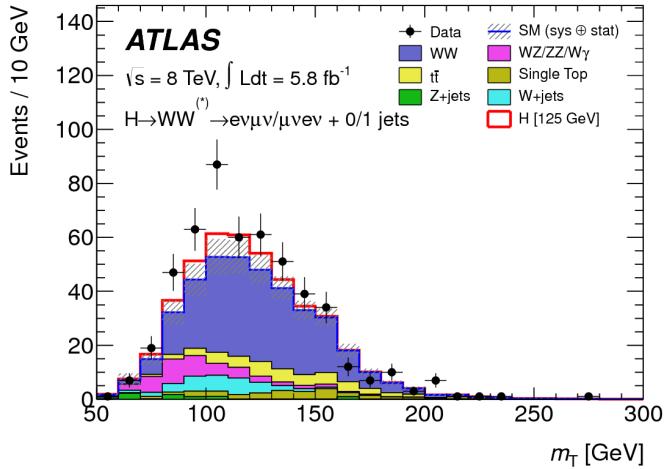


Figure 4.3: m_T distribution in the $H \rightarrow WW \rightarrow e\nu\mu\nu$ $n_j \leq 1$ channels for 8 TeV data[22].

1541 4.4 $H \rightarrow \gamma\gamma$ SEARCH

1542 The $H \rightarrow \gamma\gamma$ search is in essence a search for a peaked excess above the falling SM diphoton mass spec-
 1543 trum, with $m_{\gamma\gamma}$ as the ultimate discriminating variable. Events are selected by requiring two isolated
 1544 photons, with the leading (sub-leading) photon required to have $E_T > 40(30)$ GeV. In the 8 TeV
 1545 data, the photons are required to pass cut-based identification criteria consistent with a photon in the
 1546 electromagnetic calorimeter and little leakage in the hadronic calorimeter.

1547 The main challenges for this analysis are accurate mass reconstruction and background estimation.
 1548 In order to accurately reconstruct the invariant mass of the di-photon system, both the energy and di-
 1549 rection of the photons must be measured well. Therefore, the identification of the primary vertex of the
 1550 hard interaction is particularly important, and is done using a multivariate likelihood which combines
 1551 information about the photon direction and vertex position. The background is modeled with a falling
 1552 spectrum in $m_{\gamma\gamma}$ that is parameterized by different functions depending on the category of the event.

1553 4.4.1 RESULTS

1554 The resulting diphoton mass spectrum is shown in figure 4.4. The best fit mass value in the $\gamma\gamma$ channel
 1555 alone in the combined 7 and 8 TeV data is 126.5 GeV. The local significance at this point is 4.5σ , with
 1556 an expected significance of 2.5σ . Therefore, the measured signal strength μ is 1.8 ± 0.5 in this channel.

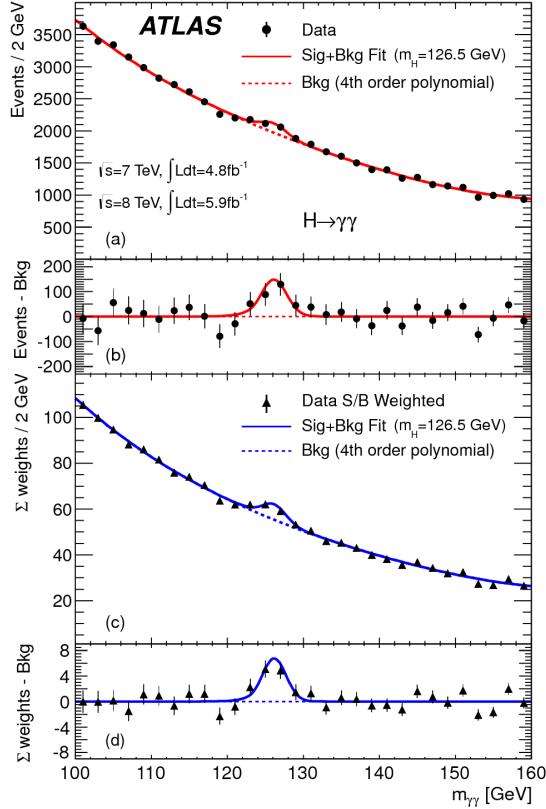


Figure 4.4: Diphoton mass spectrum in 7 and 8 TeV data. Panel a) shows the unweighted data distribution superimposed on the background fit, while panel c) shows the data where each event category is weighted by its signal to background ratio. Panels b) and d) show the respective distributions with background subtracted[22].

1557 4.5 $H \rightarrow ZZ \rightarrow 4\ell$ SEARCH

1558 The $H \rightarrow ZZ \rightarrow 4\ell$ analysis searches for a Standard Model Higgs boson decaying to two Z bosons,
 1559 each of which decays to a pair of same flavor, opposite charge isolated leptons. The ultimate discriminat-
 1560 ing variable is $m_{4\ell}$, or the invariant mass of the four selected leptons. The ℓ denotes an e or μ as with the
 1561 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis.

1562 Four distinct signal regions are constructed depending on the flavors of the final state, additionally
 1563 separated by the flavor of the leading lepton pair. These are referred to as $4e$, $2e2\mu$, $2\mu2e$, 4μ .
 1564 The main backgrounds in the $H \rightarrow ZZ \rightarrow 4\ell$ search are continuum ZZ^* production, $Z + \text{jets}$ pro-
 1565 duction, and $t\bar{t}$. The $m_{4\ell}$ distribution for background is estimated from simulation. The normalization
 1566 of the SM ZZ^* background is also taken from MC simulation, while the $Z + \text{jets}$ and $t\bar{t}$ normalizations
 1567 are taken from data-driven methods.

1568 4.5.1 RESULTS

1569 Figure 4.5 shows the $m_{4\ell}$ spectrum measured in the 7 and 8 TeV datasets. The total number of events
 1570 observed in the window between 120 and 130 GeV is 13, with 6 events in the 4μ channel, 2 events in
 1571 the $4e$ channel, and 5 events in the $2e2\mu/2\mu2e$. The best fit μ value in the combined 7 and 8 TeV data
 1572 occurs at 125 GeV and is measured to be 1.2 ± 0.6 . The observed significance at this mass is 3.6σ , with
 1573 an expected significance of 2.7σ .

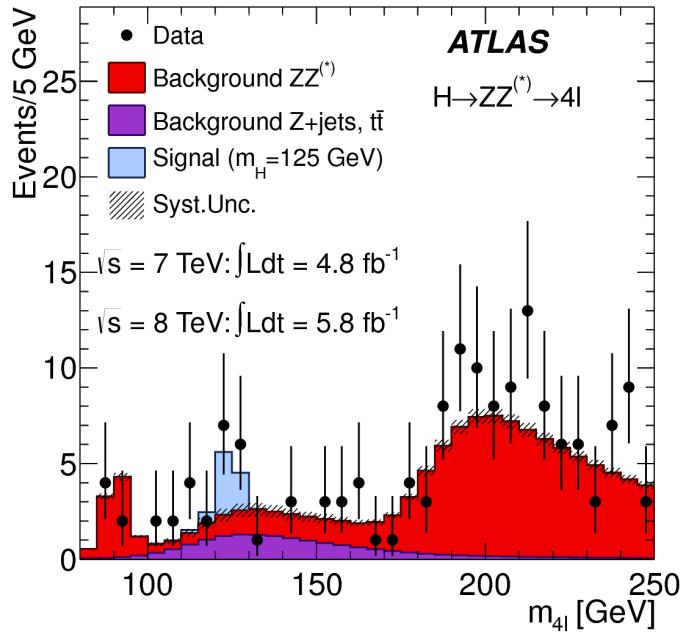


Figure 4.5: Four lepton invariant mass spectrum ($m_{4\ell}$) in 7 and 8 TeV data compared to background estimate. A 125 GeV SM Higgs signal is shown in blue[22].

1574 4.6 COMBINED RESULTS

1575 The statistical interpretation of the combined results is undertaken as described in section 3.6.2, with a
 1576 hypothesis test based on a likelihood ratio parameterized by the Higgs signal strength μ . The null hy-
 1577 pothesis corresponds to $\mu = 0$, while the SM Higgs corresponds to $\mu = 1$.

1578 Table 4.4 summarizes the properties of the individual channels as well as the significances of the ex-
 1579 cesses seen. The most significant observed local excess comes from the $\gamma\gamma$ channel. Figure 4.6 shows a
 1580 comparison of the observed local p_0 values as a function of hypothesized mass for the three different
 1581 search channels. Both the ZZ^* and $\gamma\gamma$ channels have very peaked excesses, while the WW^* excess can
 1582 be seen as very broad because the m_T distribution does not provide detailed information about the true
 1583 Higgs mass.

Channel	Fit var.	Observed Z_l	Expected Z_l	$\hat{\mu}$
$H \rightarrow ZZ^* \rightarrow 4\ell$	$m_{4\ell}$	3.6	2.7	1.2 ± 0.6
$H \rightarrow \gamma\gamma$	$m_{\gamma\gamma}$	4.5	2.5	1.8 ± 0.5
$H \rightarrow WW^* \rightarrow e\nu\mu\nu$	m_T	2.8	2.3	1.3 ± 0.5
Combined	-	6.0	4.9	1.4 ± 0.3

Table 4.4: Summary of the expected and observed significance and measured signal strengths in the combined 7 and 8 TeV datasets for the Higgs discovery analysis[22].

1584 Figure 4.7 shows the combined exclusion limit, p_0 , and signal strength. The highest local excess comes
 1585 at a value of 126.5 GeV and corresponds to a 6.0σ observed excess.

1586 Figure 4.8 shows a comparison of the measured signal strengths between the different Higgs search
 1587 channels. All measured μ are consistent with unity within their uncertainty, and the combined μ mea-
 1588 surement is 1.4 ± 0.3 .

1589 The likelihood can also be computed in a two-dimensional plane of m_H and μ , and this is shown in
 1590 figure 4.9. The figure shows that while the $\gamma\gamma$ and ZZ^* channels have very good mass resolution, the
 1591 excess in WW^* covers a broad mass range. The banana shape of the WW^* result is due to the fact that
 1592 the excess in this channel can either be explained by increasing the signal strength or by changing the
 1593 mass (and thus the cross section). The two parameters are correlated due to the lack of mass sensitivity in
 1594 this channel.

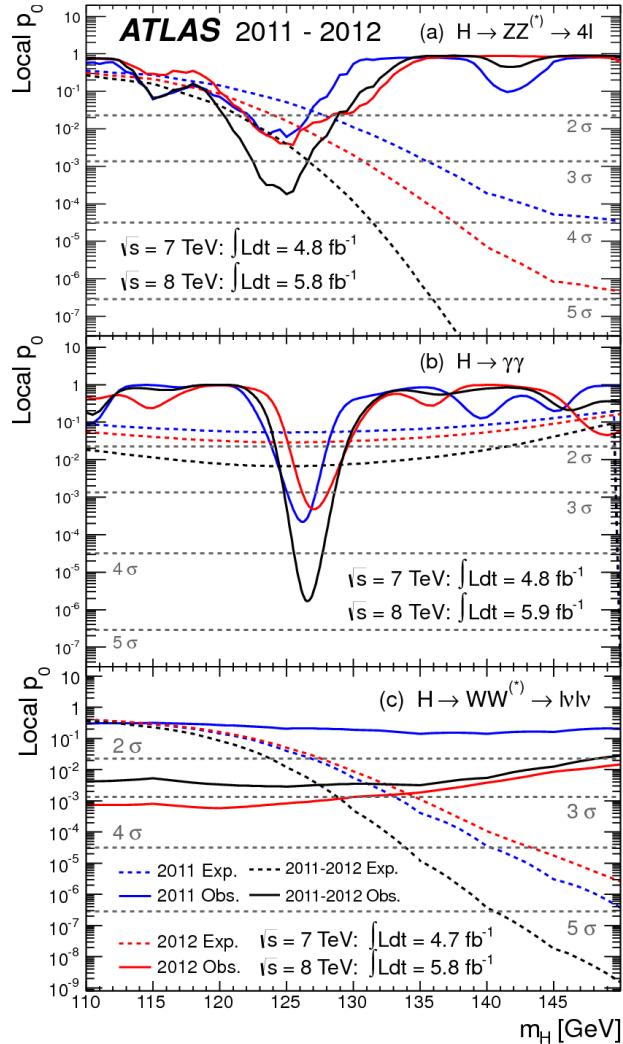


Figure 4.6: Local p_0 distribution as a function of hypothesized Higgs mass for the $H \rightarrow ZZ^* \rightarrow 4\ell$ (a), $H \rightarrow \gamma\gamma$ (b), and $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ (c) channels. Dashed curves show expected results, while solid curves show observed. Red curves are from 7 TeV data, blue curves from 8 TeV, and black curved combined[22].

Because multiple Higgs mass points are searched for, the local significance must be corrected for a look-elsewhere effect to compute a true global significance. The global significance for finding a Higgs anywhere in the mass range of 110 GeV to 600 GeV is 5.1σ . This increases slightly to 5.3σ if only mass range from 110 to 150 GeV.

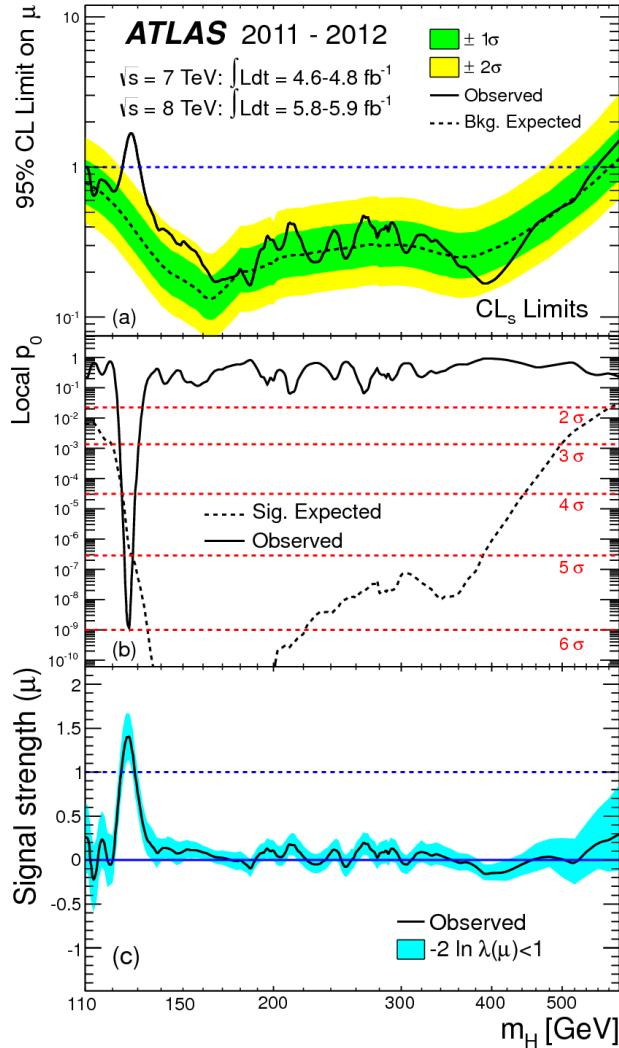


Figure 4.7: Combined 95% CL limits (a), local p_0 values (b), and signal strength measurement (c) as a function of Higgs mass[22].

1599 4.7 CONCLUSION

1600 A search for the production of a Standard Model Higgs boson was conducted in 4.8 fb^{-1} collected at
 1601 $\sqrt{s} = 7 \text{ TeV}$ and 5.8 fb^{-1} at $\sqrt{s} = 8 \text{ TeV}$. A new particle consistent with the Higgs boson was
 1602 observed, with a mass of 126.5 GeV and a global (local) significance of $5.1(6.0)\sigma$. This is the first dis-
 1603 covery level observation of a particle consistent with the Higgs.

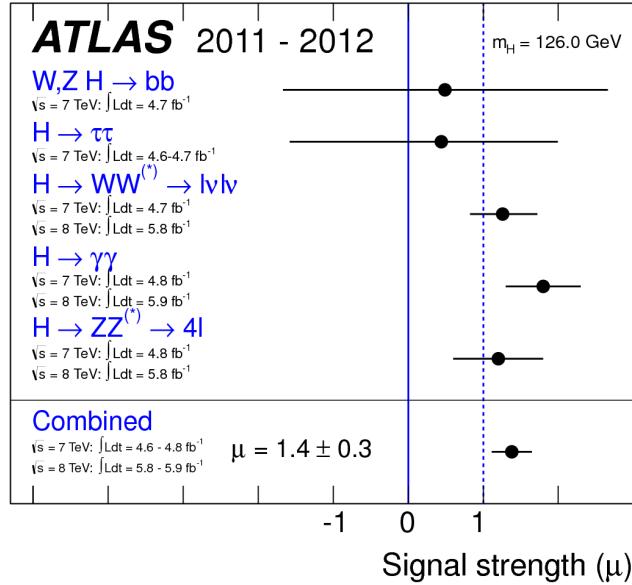


Figure 4.8: Comparison of measured signal strength μ for a 126 GeV Higgs in the 7 and 8 TeV datasets[22].

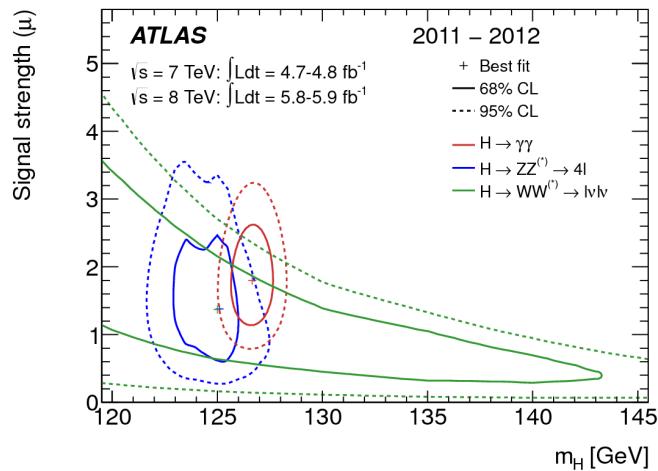


Figure 4.9: Two dimensional likelihood as a function of signal strength μ and Higgs mass m_H [22].

*The imagination of nature is far, far greater than the
imagination of man.*

Richard Feynman

5

1604

1605

Observation of Vector Boson Fusion

1606

production of $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$

1607

5.1 INTRODUCTION

1608 After the discovery of a particle consistent with the Higgs boson, the $H \rightarrow WW^*$ analysis had two main
1609 goals. The first goal was to increase the sensitivity of the analysis to fully confirm that the $H \rightarrow WW^*$
1610 process did indeed exist. The second goal was to characterize the particle as much as possible, including
1611 searching for the lower cross-section production modes, in order to confirm that it was indeed a Higgs
1612 boson. This chapter presents a dedicated search for Vector Boson Fusion (VBF) production of a Higgs
1613 boson decaying via the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ mode. First, basics of the topology of VBF production
1614 are presented. Then, the details of the analysis are shown, including signal region definition, background
1615 estimation techniques, and systematic uncertainties. Finally, the results of the analysis are presented. As

1616 will be shown, this analysis is the first and most sensitive observation of the VBF production mode of the
1617 Higgs on ATLAS.

1618 In the VBF channel, there are both a selection requirement based signal region analysis (known as the
1619 “cut-based”) and a multivariate analysis which uses a boosted decision tree (known as the BDT analysis).
1620 The focus of this chapter will be on the cut-based signal region, as this is an important component of the
1621 VBF analysis and in particular acts as strong validation for the final BDT result. Connections between
1622 the cut-based and BDT analyses will be discussed where appropriate.

1623 5.2 DATA AND SIMULATION SAMPLES

1624 The results presented here are with 20.3 fb^{-1} taken at $\sqrt{s} = 8 \text{ TeV}$ and 4.5 fb^{-1} taken at $\sqrt{s} = 7 \text{ TeV}$.
1625 The details of the LHC and detector conditions during this period are given in Chapter 2. The trigger
1626 selection defining the dataset is discussed in section 5.2.1. The simulation samples used for signal and
1627 background modeling are given in section 5.2.2.

1628 5.2.1 TRIGGERS

1629 The analysis uses a combination of single lepton and dilepton triggers to allow lowering of the p_T thresh-
1630 olds and increased signal acceptance. The p_T threshold on the leptons is a particularly important consid-
1631 eration for this signal. Because the second W produced in the decay can be off-shell, it tends to pro-
1632 duce lower momentum leptons. Thus, being able to lower the p_T threshold while still maintaining
1633 a low background rate is critical. Figure 5.1 shows an example of the subleading lepton p_T for a VBF
1634 $H \rightarrow WW^*$ signal compared to the corresponding $t\bar{t}$ background. Note that the lepton p_T spectrum is
1635 considerably softer in the signal sample.

1636 As discussed in Chapter 2, there are multiple levels in the ATLAS trigger system, and there are differ-
1637 ent p_T thresholds imposed for the leptons at each level. Additionally, some triggers have a loose selection
1638 on the isolation of the lepton (looser than that applied offline in the analysis object selection). Table 5.1
1639 shows the thresholds used for single lepton triggers, while table 5.2 shows the thresholds coming from
1640 di-lepton triggers. The single lepton trigger efficiency for muons that pass the analysis object selection is

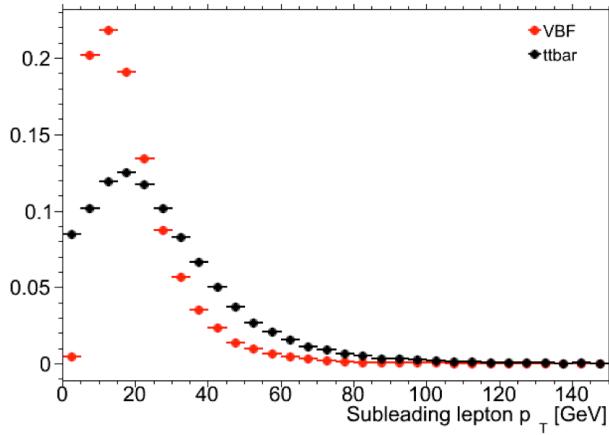


Figure 5.1: A comparison of the subleading lepton p_T spectrum between VBF $H \rightarrow WW^*$ production and $t\bar{t}$ background

1641 70% for muons in the barrel region ($|\eta| < 1.05$) and 90% in the endcap region. The electron trigger ef-
 1642 ficiency increases with electron p_T but the average is approximately 90%. These efficiencies are measured
 1643 by combined performance and trigger signature groups[72, 73].

	Level-1 threshold	High-level threshold
Electron	18	$24i$
	30	60
Muon	15	$24i$
		36

Table 5.1: Single lepton triggers used for electrons and muons. A logical “or” of the triggers listed for each lepton type is taken. Units are in GeV, and the i denotes an isolation requirement in the trigger.

	Level-1 threshold	High-level threshold
ee	10 and 10	12 and 12
$\mu\mu$	15	18 and 8
$e\mu$	10 and 6	12 and 8

Table 5.2: Di-lepton triggers used for different flavor combinations. The two thresholds listed refer to leading and sub-leading leptons, respectively. The di-muon trigger only requires a single lepton at level-1.

1644 The combination of all triggers shown gives good efficiency for signal events. This efficiency is sum-
 1645 marized in table 5.3. The relative improvement in efficiency by adding the dilepton triggers is also shown
 1646 in the same table. The largest gain comes in the $\mu\mu$ channel. Overall the trigger selection shows a good
 1647 efficiency for $H \rightarrow WW^*$ signal events.

Channel	Trigger efficiency	Gain from 2ℓ trigger
ee	97%	9.1%
$\mu\mu$	89%	18.5%
$e\mu$	95%	8.3%
μe	81%	8.2%

Table 5.3: Trigger efficiency for signal events and relative gain of adding a dilepton trigger on top of the single lepton trigger selection. The first lepton is the leading, while the second is the sub-leading. Efficiencies shown here are for the ggF signal in the $n_j = 0$ category but are comparable for the VBF signal.

1648 5.2.2 MONTE CARLO SAMPLES

1649 Modeling of signal and background processes in the signal region, in particular for the m_T distribution,
 1650 is an important consideration for the final interpretation of the analysis. Therefore, careful consideration
 1651 must be paid to which Monte Carlo (MC) generators are used for specific processes. With the exception
 1652 of the $W + \text{jet}$ and multijet backgrounds, the m_T shape used as the final discriminant is taken from simu-
 1653 lation. (Many backgrounds are normalized from data, as described in section 5.5).

1654 Table 5.4 shows the MC generators used for the signal and background processes, as well as their cross
 1655 sections. In order to include corrections up to next-to-leading order (NLO) in the QCD coupling con-
 1656 stant α_s , the `POWHEG` [74] generator is often used. In some cases, only leading order generators like
 1657 `ACERMC` [75] and `GG2VV` [76] are available for the process in question. If the process requires good
 1658 modeling for very high parton multiplicities, the `SHERPA` [77] and `ALPGEN` [78] generators are used
 1659 to provide merged calculations for five or fewer additional partons. These matrix element level calcula-
 1660 tions must then be additionally matched to models of the underlying event, hadronization, and parton
 1661 shower. There are four possible generators for this: `SHERPA`, `PYTHIA 6`[79], `PYTHIA 8`[80], or `HERWIG`
 1662 [81] + `JIMMY` [82]. The simulation additionally requires an input parton distribution function (PDF).
 1663 The `CT10`[83] PDFs are used for `SHERPA` and `POWHEG` simulated samples, while `CTEQ6Li`[84] is used
 1664 for `ALPGEN + HERWIG` and `ACERMC` simulations. The Drell-Yan samples are reweighted to the `MRST`
 1665 [85] PDFs, as these are found to give the best agreement between data and simulation.

1666 Once the basic hard scattering process is simulated, it must be passed through a detector simulation
 1667 and additional pile-up events must be overlaid. The pile-up events are modeled with `PYTHIA 8`, and the

Process	MC generator	$\sigma \cdot \mathcal{B}$ (pb)
Signal		
ggF	$H \rightarrow WW^*$ POWHEG +PYTHIA 8	0.435
VBF	$H \rightarrow WW^*$ POWHEG +PYTHIA 8	0.0356
VH	$H \rightarrow WW^*$ PYTHIA 8	0.0253
WW		
$q\bar{q} \rightarrow WW$ and $qg \rightarrow WW$	POWHEG +PYTHIA 6	5.68
$gg \rightarrow WW$	GG2VV +HERWIG	0.196
$(q\bar{q} \rightarrow W) + (q\bar{q} \rightarrow W)$	PYTHIA 8	0.480
$q\bar{q} \rightarrow WW$	SHERPA	5.68
VBS $WW + 2$ jets	SHERPA	0.0397
Top quarks		
$t\bar{t}$	POWHEG +PYTHIA 6	26.6
Wt	POWHEG +PYTHIA 6	2.35
$tq\bar{b}$	ACERMC +PYTHIA 6	28.4
$t\bar{b}$	POWHEG +PYTHIA 6	1.82
Other dibosons (VV)		
$W\gamma$ ($p_T^\gamma > 8$ GeV)	ALPGEN +HERWIG	369
$W\gamma^*$ ($m_{\ell\ell} \leq 7$ GeV)	SHERPA	12.2
WZ ($m_{\ell\ell} > 7$ GeV)	POWHEG +PYTHIA 8	12.7
VBS $WZ + 2$ jets	SHERPA	0.0126
($m_{\ell\ell} > 7$ GeV)		
$Z\gamma$ ($p_T^\gamma > 8$ GeV)	SHERPA	163
$Z\gamma^*$ (min. $m_{\ell\ell} \leq 4$ GeV)	SHERPA	7.31
ZZ ($m_{\ell\ell} > 4$ GeV)	POWHEG +PYTHIA 8	0.733
$ZZ \rightarrow \ell\ell\nu\nu$ ($m_{\ell\ell} > 4$ GeV)	POWHEG +PYTHIA 8	0.504
Drell-Yan		
Z ($m_{\ell\ell} > 10$ GeV)	ALPGEN +HERWIG	16500
VBF $Z + 2$ jets	SHERPA	5.36
($m_{\ell\ell} > 7$ GeV)		

Table 5.4: Monte Carlo samples used to model the signal and background processes[21].

ATLAS detector is simulated with GEANT4[86]. Because of the unique phase space of the $H \rightarrow WW^*$ analysis, events are sometimes filtered at generator level to allow for more efficient generation of relevant events. The efficiency of the trigger in MC simulation does not always match the measured efficiency in

1671 data, so trigger scale factors are applied to correct the MC efficiency to the data. These are derived by the
1672 combined performance groups[[72](#), [73](#)].

1673 **5.3 OBJECT SELECTION**

1674 In order to define the signal region, the analysis must first select the objects to be considered. The de-
1675 tails of the object reconstruction algorithms are discussed in Chapter 2, while this section gives specific
1676 selection cuts used in the $H \rightarrow WW^*$ analysis.

1677 The first step in this process is to select a primary vertex candidates. The event’s primary vertex is the
1678 vertex with the largest sum of p_T^2 for associated tracks and is required to have at least three tracks with
1679 $p_T > 450$ MeV. Many of the object selection cuts are then made relative to this chosen primary vertex.

1680 **5.3.1 MUONS**

1681 The analysis uses combined muon candidates, where a track in the Inner Detector has been matched
1682 to a standalone track in the Muon Spectrometer. The track parameters are combined statistically in the
1683 muon reconstruction algorithm[[16](#)]. The muons are required to be within $|\eta| < 2.5$ and have a $p_T >$
1684 10 GeV. To reduce backgrounds coming from mis-reconstructed leptons, there are requirements on
1685 the impact parameter of the muon relative to the primary vertex. The transverse impact parameter d_0
1686 is required to be small relative to its estimated uncertainty, the exact cut value being $d_0/\sigma_{d_0} < 3$. The
1687 longitudinal impact parameter z_0 must satisfy $|z_0 \sin \theta| < 1$ mm.

1688 As discussed previously, the muons must also be isolated. There are two types of lepton isolations
1689 that are calculated: track-based and calorimeter-based. For muons, the track-based isolation is defined
1690 using the scalar sum $\sum p_T$ for tracks with $p_T > 1$ GeV (excluding the muon’s track) within a cone of
1691 $\Delta R = 0.3$ (0.4) for muon with $p_T > 15$ GeV ($10 < p_T < 15$ GeV). The final isolation requirement
1692 is made my requiring that this scalar sum be no more than a certain fraction of the muon’s p_T . This re-
1693 quirement varies with muon p_T and the exact cuts are defined in table 5.5.

1694 The calorimeter-based muon isolation is defined using as a $\sum E_T$ calculated from calorimeter cells
1695 using the same cone size as the track-based isolation but excluding cells with $\Delta R < 0.05$ around the

1696 muon. This requirement is also defined as a cut on the ratio of the sum to the muon p_T and varies with
1697 muon p_T . The cut values are also given in table 5.5.

1698 The isolation requirements loosen as a function of p_T to allow for larger signal acceptance. At low p_T ,
1699 the isolation is tightened to reduce the $W + \text{jets}$ background which arises from a misidentified lepton.

p_T range (GeV)	Calorimeter isolation	Track isolation
10 – 15	0.06	0.06
15 – 20	0.12	0.08
20 – 25	0.18	0.12
> 25	0.30	0.12

Table 5.5: p_T dependent isolation requirements for muons. Muons are required to have the amount of calorimeter or track based cone sums be less than this fraction of their p_T .

1700 5.3.2 ELECTRONS

1701 Electrons are identified by matching reconstructed clusters in the electromagnetic calorimeter with tracks
1702 in the inner detector. The electrons are identified using a likelihood based method[14, 64] which takes
1703 into account the shower shapes in the calorimeter, the matching of tracks to clusters, and the amount of
1704 transition radiation in the TRT. The electrons are required to have $|\eta| < 2.47$, and candidates in the
1705 transition region between the barrel and endcap ($1.37 < |\eta| < 1.52$) are excluded. As the muons, the
1706 electrons are required to have transverse impact parameter significance < 3 , while in the longitudinal
1707 direction they must have $|z_0 \sin \theta| < 0.4$ mm. Some electron requirements also vary with electron E_T ,
1708 and these requirements are summarized in table 5.6.

1709 The isolation for electrons are defined similarly to the muons but with unique cuts on the objects in-
1710 cluded. The track-based isolation is defined using tracks with $p_T > 400$ MeV with cone sizes as defined
1711 previously. The calorimeter-based isolation also uses the same cone size as the muon, but here the cells
1712 within a 0.125×0.175 area in $\eta \times \phi$ around the electron cluster's barycenter are excluded. The other
1713 difference with respect to muons is that the denominator of the isolation ratio is the electron's E_T rather
1714 than p_T . The isolation cuts very with electron E_T and are defined in table 5.6.

1715 The electron is also required to not be consistent with a vertex coming from a photon conversion.

p_T range (GeV)	Quality cut	Calorimeter isolation	Track isolation
10 – 15	Very tight LH	0.20	0.06
15 – 20	Very tight LH	0.24	0.08
20 – 25	Very tight LH	0.28	0.10
> 25	Medium	0.28	0.10

Table 5.6: p_T dependent requirements for electrons. Electrons are required to have the amount of calorimeter or track based cone sums be less than this fraction of their E_T .

1716 5.3.3 JETS

1717 Jets are clustered with the anti- k_T reconstruction algorithm using a radius parameter of $R = 0.4$. They
 1718 are required to have a jet vertex fraction (JVF) of at least 50%, meaning that half of the tracks associated
 1719 with the jet originated from the primary vertex. Jets with no tracks associated (i.e. those outside the ac-
 1720 ceptance of the ID) do not have this requirement applied. Jets are required to have $p_T > 25$ GeV if
 1721 they are within the tracking acceptance ($|\eta| < 2.4$). Jets with $2.4 < |\eta| < 4.5$ are required to have
 1722 $p_T > 30$ GeV. This tighter requirement reduces jets from pileup in the region where JVF requirements
 1723 cannot be applied. The two highest p_T jets in the event are referred to as the “VBF” jets and used to com-
 1724 pute various analysis selections later.

1725 Identification of b -jets is done using the MV1 algorithm and is limited to the acceptance of the ID
 1726 ($|\eta| < 2.5$). The operating point of MV1 that is used is the one that is 85% efficient for identifying true
 1727 b -jets. This operating point has a 10.3% of mis-tagging a light quark jet as a b -jet. In order to improve
 1728 the rejection of b -jets, a lower threshold than the nominal p_T threshold described above is used. For the
 1729 purposes of counting the number of b -jets, jets with p_T down to 20 GeV are used.

1730 5.3.4 OVERLAP REMOVAL

1731 There are some cases where certain reconstructed objects will overlap and one will have to be chosen
 1732 (for example, an electron and a jet in the calorimeter). First, the case of lepton overlap is dealt with. If
 1733 an electron candidate extends into the muon spectrometer, it is removed. If a muon or electron have a
 1734 $\Delta R < 0.1$, the electron is removed and the muon is kept. If two electron candidates overlap within
 1735 the same radius, then the higher E_T electron is kept. Next, the overlap between leptons and jets is con-

1736 sidered. If an electron and jet are within $\Delta R < 0.3$ of one another, the electron is kept and the jet is
1737 removed. However, if a muon and jet overlap within $\Delta R < 0.3$, the jet is kept (as it is likely that the
1738 muon is the result of a semileptonic decay inside the jet).

1739 Once the overlap removal is complete, the final set of objects used in the analysis is defined.

1740 5.4 ANALYSIS SELECTION

1741 The VBF analysis uses two distinct selections. The first is a more standard selection, referred to as “cut-
1742 based”, that applies requirements on the VBF variables and uses m_T as the final discriminating variable.
1743 The second is a looser selection that uses a Boosted Decision Tree (BDT) score as the final discriminator
1744 in order to take advantage of the detailed correlations between the VBF variables. While the BDT analy-
1745 sis is ultimately more sensitive, the cut-based serves as an important component of the analysis. First, the
1746 cut-based allows for confirming the modeling and validity of many variables used as input to the BDT.
1747 Second, because this is the first use of such an MVA technique in the $H \rightarrow WW^*$ analysis, the cut-based
1748 selection allows confirmation of the final BDT result with a more traditional analysis. The cut-based
1749 techniques are the focus of this chapter, but connections to the BDT result will be illustrated when ap-
1750 propriate.

1751 One important note is that because this analysis is dedicated to the measurement of the VBF pro-
1752 duction mode of the Higgs, events coming from gluon fusion production with the Higgs decaying via
1753 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ are treated as background events. This will be seen throughout the various
1754 predictions shown.

1755 5.4.1 COMMON PRE-SELECTION

1756 Both the cut-based and BDT analyses have a common pre-selection that is applied before the main signal
1757 region requirements. The requirements on leptons are common to all n_j bins. The analysis requires two
1758 oppositely charged leptons, with the leading lepton required to have $p_T > 22$ GeV while the subleading
1759 lepton must have $p_T > 10$ GeV. Next, to remove low mass Z/γ^* events, a cut on the dilepton mass
1760 $m_{\ell\ell} > 10$ (12) GeV is applied in the different (same) flavor channel. In the same flavor channels, there is

1761 an additional veto placed on the region around the Z peak, requiring that $|m_{\ell\ell} - m_Z| > 15$ GeV.

1762 There are also requirements on the amount of missing transverse momentum in the event. These
1763 are only applied in the same flavor channels, as in the different flavor channels $t\bar{t}$ is the dominant back-
1764 ground in $n_j \geq 2$. The BDT analysis requires $p_T^{\text{miss}} > 40$ GeV and $E_T^{\text{miss}} > 45$ GeV. The cut-
1765 based analysis must select more tightly on these variables to have maximal sensitivity and thus requires
1766 $p_T^{\text{miss}} > 50$ GeV and $E_T^{\text{miss}} > 55$ GeV.

1767 Finally, because this analysis is focused on VBF, a requirement on the jet multiplicity is placed, with
1768 $n_j \geq 2$. Additionally, the analysis requires that there are no jets identified as b-quarks in the event, or
1769 $n_b = 0$.

1770 5.4.2 CUT-BASED SELECTION

1771 The cut-based selection places sequential requirements on variables reconstructed from the VBF jets in
1772 order to increase the signal to background ratio.

1773 GENERAL BACKGROUND REDUCTION

1774 Top pair production is the primary background in the $n_j \geq 2$ bin. Even though $n_b = 0$ is required,
1775 an additional variable is constructed to further suppress the top background. There is often additional
1776 QCD radiation that accompanies the $t\bar{t}$ system when it is produced. Therefore, a variable which tests for
1777 the presence of this additional radiation, p_T^{sum} , is constructed. It is defined in equation 5.1.

$$p_T^{\text{sum}} = p_T^{\ell\ell} + p_T^{\text{miss}} + \sum p_T^j \quad (5.1)$$

1778 The first cut after pre-selection in the cut-based analysis requires $p_T^{\text{sum}} < 15$ GeV to further suppress
1779 $t\bar{t}$ production.

1780 In the different flavor channels, a cut is made to reduce the contamination from $Z \rightarrow \tau\tau$ decays. The
1781 di- τ invariant mass, $m_{\tau\tau}$, is constructed by assuming that the neutrinos from the τ decays were collinear
1782 with the leptons[?]. The analysis requires that this mass not be consistent with a Z by requiring $m_{\tau\tau} <$
1783 $m_Z - 25$ GeV.

1784 VBF TOPOLOGICAL CUTS

1785 The characteristic feature of VBF production of the Higgs is the presence of two additional forward jets
1786 coming from the incoming partons which radiate the vector bosons that make the Higgs. These jets are
1787 forward because the outgoing partons still carry the longitudinal momentum of the incoming partons.
1788 Figure 5.2 shows the distribution of the η for the leading jet in a VBF event compared to a background
1789 top pair production event. As can be seen, the VBF jets tend to be more forward in η , while the $t\bar{t}$ jets are
1790 more central.

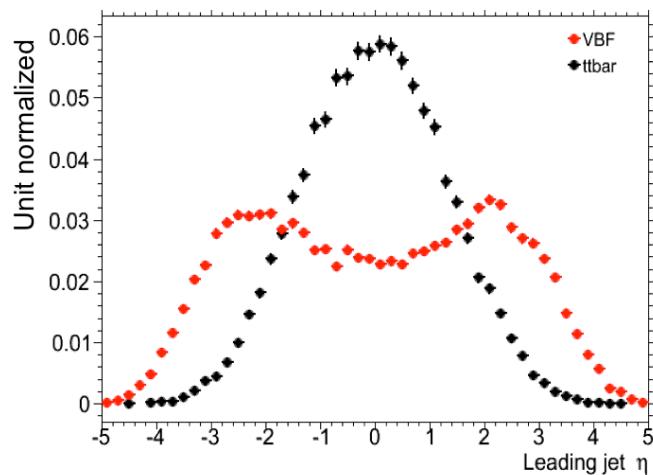


Figure 5.2: Leading jet η in VBF $H \rightarrow WW^*$ (red) and $t\bar{t}$ (black)

1791 Because the cross section for VBF production is an order of magnitude smaller than gluon fusion pro-
1792 duction, these forward jets must be used in order to better reduce background and achieve a good signal
1793 to background ratio. The dedicated VBF search selection requirements are constructed to maximally
1794 exploit the features of the unique VBF topology.

1795 Requirements on the VBF jets are collectively referred to as the “VBF topological cuts”. First, a re-
1796 quirement on the dijet invariant mass of the VBF jets, m_{jj} , is placed, requiring $m_{jj} > 600$ GeV. Next,
1797 the event is required to have a large gap in rapidity between the two VBF jets, or $\Delta y_{jj} > 3.6$. Both
1798 of these cuts put tight requirements on the presence of two forward, high p_T jets moving in opposite

1799 directions in the longitudinal plane.

1800 Beyond requiring the presence of the two forward VBF jets, the analysis also vetoes on the presence
1801 of any additional jets that fall between the two VBF jets. This cut is referred to as the central jet veto, or
1802 CJV. Any events with a third jet with $p_T > 20$ GeV whose rapidity is between the region defined by the
1803 two VBF jets are vetoed. This can be expressed in terms of a variable called the jet centrality, defined in
1804 equation 5.2.

$$C_{j3} = \left| \eta_{j3} - \frac{\eta_{j1} + \eta_{j2}}{2} \right| / \frac{|\eta_{j1} - \eta_{j2}|}{2}, \quad (5.2)$$

1805 Here, η_{j1} and η_{j2} are the pseudorapidities of the leading and subleading jets, respectively, while η_{j3} is
1806 the pseudorapidity of the extra jet in the event (if one exists). Intuitively, C_{j3} is zero when η_{j3} is directly
1807 centered between the two jets and unity when η_{j3} is aligned with either of the VBF jets. Thus, the CJV
1808 can be expressed as a requirement that $C_{j3} > 1$.

1809 The decay products of the Higgs tend to be central as well. Thus, the analysis also requires that both
1810 leptons in the analysis fall within the rapidity gap defined by the jets. This cut is referred to as the outside
1811 lepton veto, or OLV. A quantitative way to define the cut is to require that the centrality of each lepton
1812 (defined analogously to that of the third jet in equation 5.2) correspond to the lepton being within the jet
1813 rapidity gap, or $C_\ell < 1$ for both leptons.

1814 Figure 5.3a-c shows the m_{jj} , Δy_{jj} , and $C_{\ell 1}$ variables at the stage where all previous cuts in the se-
1815 quence have been made. The agreement between data and Monte Carlo is good, and the bottom panels
1816 show their power in discriminating the VBF signal from the background processes.

1817 The final signal region is also split into two bins of m_{jj} , with the first bin corresponding to $600 \text{ GeV} <$
1818 $m_{jj} < 1 \text{ TeV}$ and the second bin corresponding to $m_{jj} > 1 \text{ TeV}$. The first bin has more statistics but
1819 also a larger contribution from background, while the second bin has lower statistics but a 1:1 signal to
1820 background ratio.

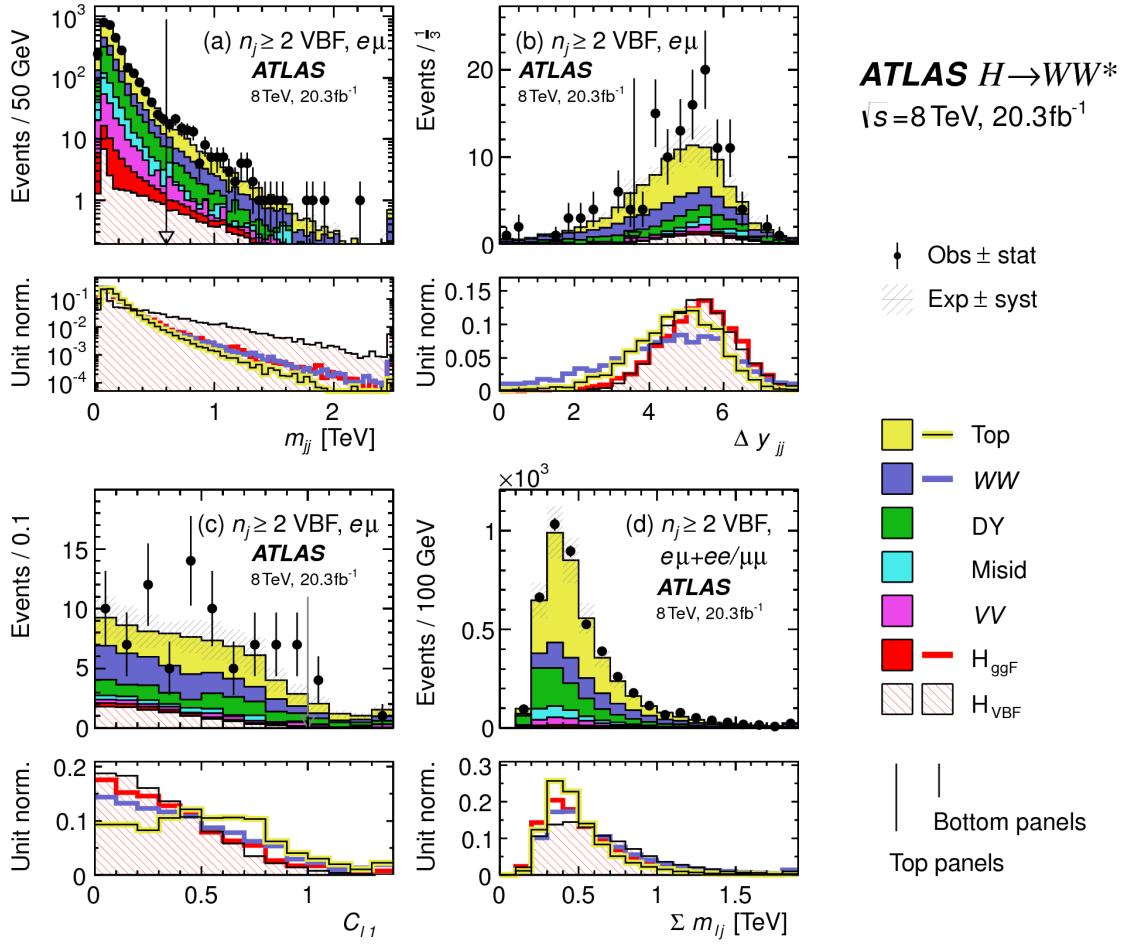


Figure 5.3: Distributions of (a) m_{jj} , (b) Δy_{jj} , (c) $C_{\ell 1}$, and (d) $\sum m_{ij}$, for the VBF analysis. The top panels compare simulation and data, while the bottom panels show normalized distributions for all background processes and signal[21].

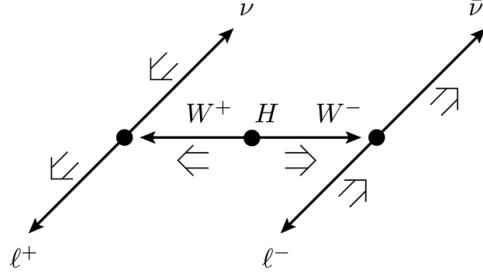


Figure 5.4: A cartoon of the WW final state. Momenta are represented with thin arrows, spins with thick arrows. [21]

1821 **HIGGS TOPOLOGICAL CUTS**

1822 The final state leptons will exhibit unique correlations due to the fact that they are arising from the decay
 1823 of a spin zero resonance. In particular, the spins of the final state leptons and neutrinos must all cancel,
 1824 as shown in figure 5.4. Because the neutrino has a left handed chirality and the anti-neutrino has a right
 1825 handed chirality (in the massless neutrino approximation), the spin and momentum of the particles
 1826 will be anti-aligned and aligned, respectively. In the transverse plane, the momenta of all four final state
 1827 objects must cancel as well. With the constraint of having both the momenta and the spin alignments
 1828 cancel, the final state kinematics strongly prefer having a small angle between the leptons in the trans-
 1829 verse plane (low $\Delta\phi_{\ell\ell}$). This angular correlation will also lead to low values of the di-lepton invariant
 1830 mass $m_{\ell\ell}$. These unique signal final state kinematic correlations will be exploited to define the ultimate
 1831 signal region.

1832 The analysis places additional requirements on the final state leptons. Two requirements on dilepton
 1833 kinematics are made that are common with lower multiplicity jet bins as well. The angle between leptons
 1834 in the transverse plane, $\Delta\phi_{\ell\ell}$, is required to be less than 1.8 radians. Additionally, the dilepton mass $m_{\ell\ell}$
 1835 is required to be less than 50 GeV.

1836 The cut-based analysis uses m_T as the final discriminating variable as in the ggF focused analysis. The
 1837 optimal number of bins in m_T was found to be three bins, with the bin boundaries at 80 and 130 GeV.

1838 Table 5.7 shows the data and estimated signal and background yields from simulation as each cut de-
 1839 scribed above is made. The table shows how each cut reduces specific backgrounds and how the overall

1840 signal to background ratio grows through the cutflow.

1841 Figure 5.5 shows an ATLAS event display of a candidate event in the final signal region.

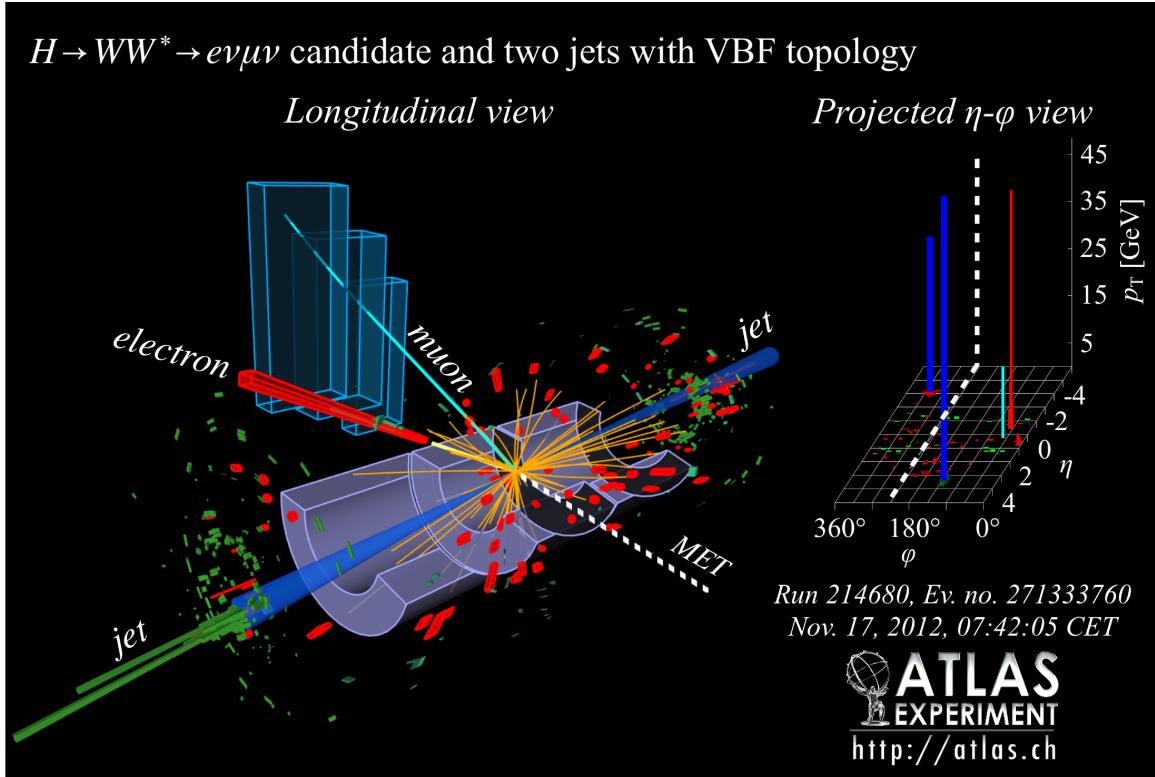


Figure 5.5: Event display of a VBF candidate event[21].

1842 5.4.3 BDT-BASED SELECTION

1843 The boosted decision tree based analysis takes a different philosophy compared to the cut-based. Rather
1844 than cutting sequentially on many variables, the BDT analysis uses many of these variables as inputs
1845 to the BDT and the output BDT score (O_{BDT}) as the final discriminant. The BDT is trained with the
1846 VBF $H \rightarrow WW^*$ simulation as the signal samples and all other processes as background, including ggF
1847 $H \rightarrow WW^*$ production. While the BDT based analysis is treated as a separate result, it has significant
1848 overlap with the cut-based selection.

Table 5.7: Event selection for the $n_j \geq 2$ VBF analysis in the 8 TeV cut-based analysis[21].

Selection	Summary										Composition of N_{bkg}							
	$N_{\text{obs}}/N_{\text{bkg}}$	N_{obs}	N_{bkg}	N_{signal}	N_{ggF}	N_{VBF}	N_{VH}	N_{WW}^{QCD}	N_{WW}^{EW}	$N_{t\bar{t}}$	N_t	N_{top}	N_{Wj}	N_{jj}	N_{VV}	N_{misid}	$N_{\text{Drell-Yan}}$	$N_{ee/\mu\mu} N_{\tau\tau}^{\text{QCD}}$
$e\mu$ sample	1.00 ± 0.00	61434	61180	85	32	26	1350	68	51810	2970	847	308	380	51	3260	46		
$n_b = 0$	1.02 ± 0.01	7818	7700	63	26	16	993	43	3000	367	313	193	273	35	2400	29		
$p_T^{\text{sum}} < 15$	1.03 ± 0.01	5787	5630	46	23	13	781	38	1910	270	216	107	201	27	2010	23		
$m_{\tau\tau} < m_Z - 25$	1.05 ± 0.02	3129	2970	40	20	9.9	484	22	1270	177	141	66	132	7.6	627	5.8		
$m_{jj} > 600$	1.31 ± 0.12	131	100	2.3	8.2	—	18	8.9	40	5.3	1.8	2.4	5.1	0.1	15	1.0		
$\Delta y_{jj} > 3.6$	1.33 ± 0.13	107	80	2.1	7.9	—	11.7	6.9	35	5.0	1.6	2.3	3.3	—	11.6	0.8		
$C_{j3} > 1$	1.36 ± 0.18	58	43	1.3	6.6	—	6.9	5.6	14	3.0	1.3	1.3	2.0	—	6.8	0.6		
$C_{\ell 1} < 1, C_{\ell 2} < 1$	1.42 ± 0.20	51	36	1.2	6.4	—	5.9	5.2	10.8	2.5	1.3	1.3	1.6	—	5.7	0.6		
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_\tau$	2.53 ± 0.71	14	5.5	0.8	4.7	—	1.0	0.5	1.1	0.3	0.3	0.3	0.6	—	0.5	0.2		
<hr/>																		
$ee/\mu\mu$ sample	0.99 ± 0.01	26949	27190	31	14	10.1	594	37	23440	1320	230	8.6	137	690	679	16		
$n_b, p_T^{\text{sum}}, m_{\tau\tau}$	1.03 ± 0.03	1344	1310	13	8.0	4.0	229	12.0	633	86	26	0.9	45	187	76	1.5		
$m_{jj}, \Delta y_{jj}, C_{j3}, C_\ell$	1.39 ± 0.28	26	19	0.4	2.9	0.0	3.1	3.1	5.5	1.0	0.2	0.0	0.7	3.8	0.7	0.1		
$m_{\ell\ell}, \Delta\phi_{\ell\ell}, m_\tau$	1.63 ± 0.69	6	3.7	0.3	2.2	0.0	0.4	0.2	0.6	0.2	0.2	0.0	0.1	1.5	0.3	0.1		

1849 PRE-TRAINING SELECTION AND BDT INPUTS

1850 Before training, the common preselection cuts described in section 5.4.1 are applied. Additionally, the
1851 central jet veto and outside lepton veto described in section 5.4.2 are applied. The BDT has eight input
1852 variables, six of which are also variables that are used in the cut-based analysis. The six shared variables
1853 are p_T^{sum} , m_{jj} , Δy_{jj} , $m_{\ell\ell}$, $\Delta\phi_{\ell\ell}$, and m_T . The seventh variable input in the BDT is a combination of
1854 the variables used to do the OLV in the cut-based analysis. The BDT uses as input the sum of lepton
1855 centralities, or $\sum C_\ell = C_{\ell 1} + C_{\ell 2}$. The final BDT input variable, $\Sigma m_{\ell j}$, is constructed to account for
1856 the correlations between the jets and leptons in the event. It is the sum of the invariant masses of all four
1857 possible lepton-jet combinations.

1858 Figure 5.3d shows the agreement between data and simulation for the $\Sigma m_{\ell j}$ variable, as well as show-
1859 ing its discriminating power. Figure 5.6 shows the distributions of the Higgs topological variables that
1860 are shared between the cut-based and BDT analyses. Figure 5.7 shows the distributions of the VBF topo-
1861 logical variables shared between the cut-based and BDT analyses. In both cases, the VBF yield has been
1862 scaled by a factor of 50 to better show the shape difference compared to the backgrounds.

1863 Table ?? summarizes the cuts applied for the cut-based and analyses, as well as which variables are used
1864 as input to the BDT.

1865 5.5 BACKGROUND ESTIMATION

1866 This section describes the procedures used to estimate backgrounds for the VBF analysis in both the
1867 cut-based and BDT analyses.

1868 5.5.1 GENERAL STRATEGY

1869 Most of the backgrounds in the VBF analysis have shapes estimated from Monte Carlo simulation but
1870 normalizations derived from control regions in data. In essence, a normalization factor (denoted with
1871 β or abbreviated as NF) is derived by scaling the MC yield in the control region to the corresponding
1872 yield in data. Once this factor is derived, it can be used to scale the MC estimate of the background in the

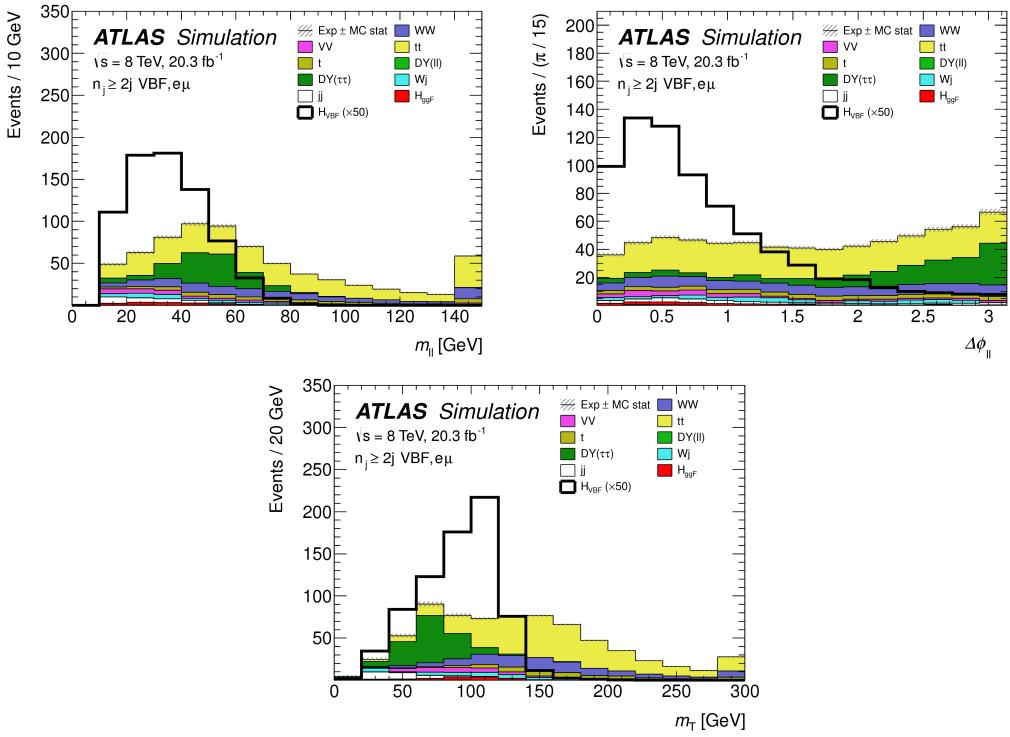


Figure 5.6: Distributions of $m_{\ell\ell}$ (top left), $\Delta\phi_{\ell\ell}$ (top right), and m_T (bottom), Higgs topology variables used in the selection requirements of the cut-based signal region and as inputs to the BDT result. These are plotted after all of the BDT pre-training selection cuts[21].

1873 signal region. This is illustrated in equation 5.3.

$$B_{\text{SR}}^{\text{est}} = B_{\text{SR}} \times \frac{N_{\text{CR}}}{B_{\text{CR}}} \equiv B_{\text{SR}} \times \beta \quad (5.3)$$

1874 Here, B denotes the MC yield prediction in the denoted region, while N denotes the observed num-
1875 ber of events in data in the denoted region.

1876 Another way of writing the same equation, in terms of an extrapolation factor α rather than a nor-
1877 malization factor β . The overall calculation is exactly the same. However, when phrased in this way, it
1878 shows how the uncertainty on the background estimation can be reduced. This is shown in equation 5.4.

$$B_{\text{SR}}^{\text{est}} = N_{\text{CR}} \times \frac{B_{\text{SR}}}{B_{\text{CR}}} \equiv N_{\text{CR}} \times \alpha \quad (5.4)$$

1879 Phrased this way, the equation shows that with enough statistics in the control region, a large theoret-

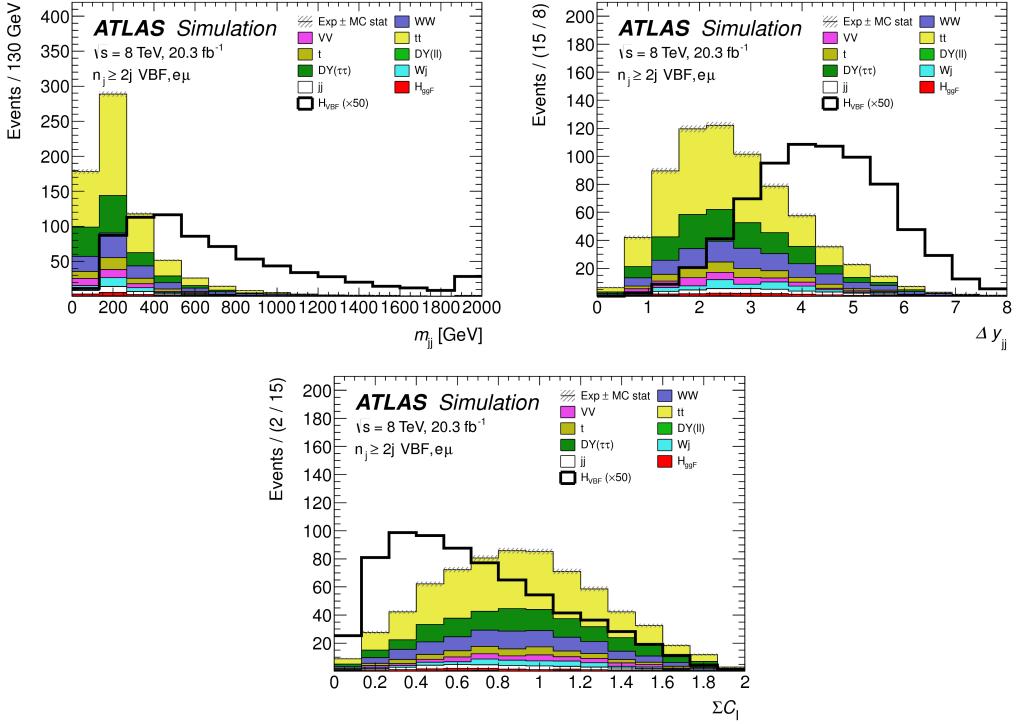


Figure 5.7: Distributions of m_{jj} (top left), Δy_{jj} (top right), $\sum C_l$ (bottom), VBF topology variables used in the selection requirements of the cut-based signal region and as inputs to the BDT result. These are plotted after all of the BDT pre-training selection cuts[21].

ical uncertainty on the overall background yield in the signal region can be replaced by a small statistical uncertainty coming from the number of data events in the CR and a smaller theoretical uncertainty on the extrapolation from the control region to the signal region.

5.5.2 TOP BACKGROUND

The normalization factor β_t for the top background in the VBF analysis is derived in a region required to have one b-tagged jet, or $n_b = 1$. In the cut-based analysis, normalization factors are computed at every stage of the cutflow by applying the appropriate cuts in the CR. These NF are then applied to the $t\bar{t}$ and single top event yields in the SR. In the BDT analysis, a single normalization factor is computed for each bin of O_{BDT} after applying the BDT pre-training cuts described previously. The computed normalization factors are derived with all flavor combinations combined in order to decrease statistical uncertainty. Additionally, in the BDT analysis, BDT bins 2 and 3 are merged for the same reason.

1891 Table 5.8 shows the evolution of the β_t through the cut-based selection. Table 5.9 shows the value
 1892 of the β_t in each bin of O_{BDT} . In all cases, the computed factors are relatively consistent with unity,
 1893 with the largest discrepancy coming in bin 1 of O_{BDT} . The normalization factors in the bins of O_{BDT}
 1894 are also consistent with those derived in the cut-based signal region, increasing confidence in the BDT
 1895 estimation.

Cut	β_t
$p_T^{\text{sum}} < 15 \text{ GeV}$	1.03 ± 0.01
$m_{\tau\tau} < m_Z - 25$	1.05 ± 0.01
$m_{jj} > 600 \text{ GeV}$	0.96 ± 0.06
$\Delta y_{jj} > 3.6$	1.02 ± 0.08
CJV	1.13 ± 0.16
OLV	1.01 ± 0.19
$m_{jj} < 1 \text{ TeV}$	0.94 ± 0.19
$m_{jj} > 1 \text{ TeV}$	1.48 ± 0.66

Table 5.8: Top normalization factors computed at each stage of the cut-based selection. Uncertainties are statistical only.

O_{BDT}	β_t
Bin0	1.09 ± 0.02
Bin1	1.58 ± 0.15
Bin2	0.95 ± 0.31
Bin3	0.95 ± 0.31

Table 5.9: Top normalization factors computed for each bin of O_{BDT} . Uncertainties are statistical only.

1896 Figure 5.8 shows the m_{jj} and O_{BDT} distributions in the top control region. Overall the modeling
 1897 looks consistent with the data.

1898 While these normalization factors can be computed and applied to the expected background yields
 1899 listed in tables like table 5.7, in the end the normalization of the top background is profiled (meaning
 1900 there is a dedicated Poisson constraint) and allowed to float in the final statistical fit.

1901 5.5.3 $Z/\gamma^* \rightarrow \tau\tau$ BACKGROUND

1902 In the different flavor channels, the $Z/\gamma^* \rightarrow \tau\tau$ background is an important one. Di-tau production
 1903 can produce an $e\mu$ final state if each τ lepton decays to a different flavor lepton.

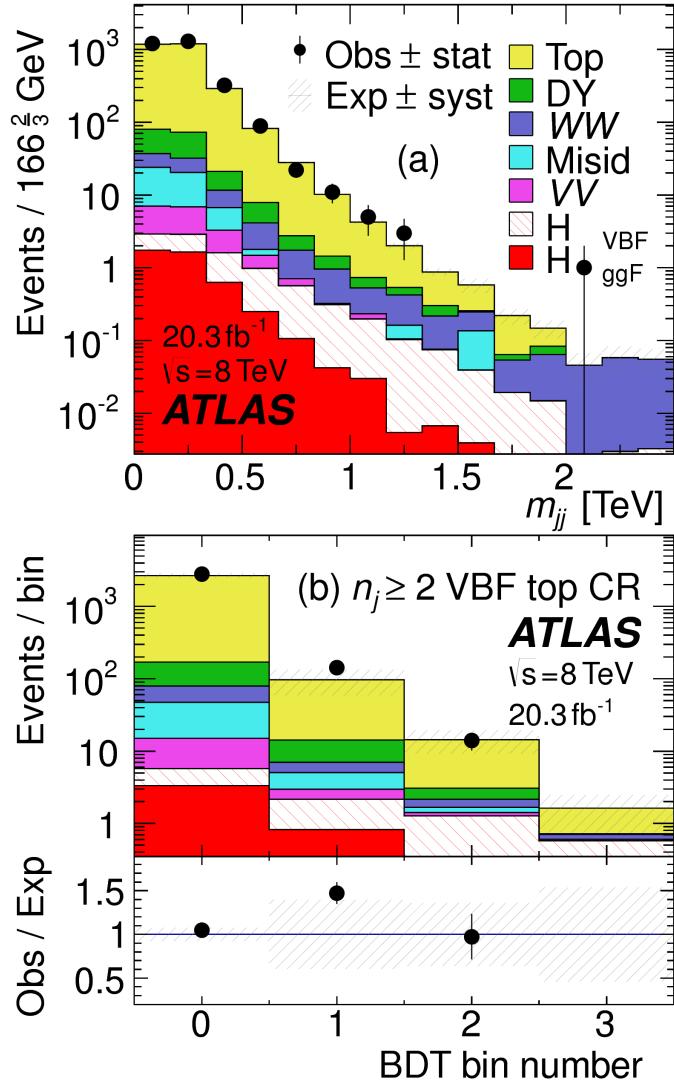


Figure 5.8: Distributions of m_{jj} (a) and O_{BDT} (b) in the VBF $n_b = 1$ top CR [21].

1904 In the BDT analysis, a single normalization factor for the background is derived. A control region is
 1905 defined using the pre-training selection cuts, except requiring that $|m_{\tau\tau} - m_Z| < 25 \text{ GeV}$ so that
 1906 the region is enriched in $Z/\gamma^* \rightarrow \tau\tau$ background. Additional requirements of $m_{\ell\ell} < 80(75) \text{ GeV}$
 1907 in the different (same) flavor channel, as well as $O_{\text{BDT}} > -0.48$ are applied to increase the purity of
 1908 the region. The final $\beta_{Z/\gamma^* \rightarrow \tau\tau}$ is calculated to be 0.9 ± 0.3 (statistical uncertainty only). Because of
 1909 the small contribution of this background in the BDT analysis and the large statistical uncertainty, no
 1910 additional systematics are calculated. The final SR estimate is scaled by this β and not allowed to float in

1911 the fit.

1912 The cut-based corrections are a bit more involved because they need to be applied cut by cut through
1913 the cutflow, as well as in the final signal region for the fit. The region is defined including all SR cuts
1914 up to the $Z/\gamma^* \rightarrow \tau\tau$ veto, which is instead made into a Z mass peak requirement as for the BDT
1915 region. The $m_{\ell\ell}$ cut from the BDT region is included as well. The cut-based approach aims to correct
1916 the normalization of the $Z/\gamma^* \rightarrow \tau\tau$ background in two ways. First, an overall normalization factor is
1917 computed from the control region. However, the VBF topological cuts are not included in this region,
1918 and applying them as is done in the top CR is not feasible due to limited statistics. So, instead, correction
1919 factors (CF) to the cut efficiencies of the VBF cuts are derived in a same flavor $Z \rightarrow \ell\ell$ control region,
1920 which has significantly more statistics. The CF is simply the ratio of the cut efficiencies in data and MC
1921 derived in this region. In the end, the overall background estimate is given by equation 5.5.

$$N_{Z/\gamma^* \rightarrow \tau\tau}^{\text{est}} = B_{Z/\gamma^* \rightarrow \tau\tau}^{\text{SR}} \times \beta_{\tau\tau} \times \frac{\epsilon_{\text{VBF cuts}}^{\text{data}}}{\epsilon_{\text{VBF cuts}}^{\text{MC}}} \quad (5.5)$$

1922 The hypothesis is that while the normalization correction must be derived in a dedicated region, the
1923 efficiency of the VBF cuts should not be sensitive to the type of Z/γ^* process and thus the larger control
1924 region can be exploited to derive the CF. Figure 5.9 shows a shape comparison for the m_{jj} variable in
1925 $Z \rightarrow \tau\tau$ events in the signal region and $Z \rightarrow \ell\ell$ events in the control region. The figure shows that the
1926 shapes are indeed comparable and thus any CF derived in the same flavor control region can reliably be
1927 applied in the signal region.

1928 Table 5.10 shows the overall normalization factor $\beta_{\tau\tau}$ and the efficiency correction factors for the var-
1929 ious VBF topological cuts. In general, the statistical uncertainties on the cut efficiency corrections are
1930 quite good, and the MC tends to underestimate the efficiency of the VBF cuts for the $Z/\gamma^* \rightarrow \tau\tau$ back-
1931 ground. The overall normalization factor is also consistent with that calculated for the BDT analysis.

1932 5.5.4 $Z/\gamma^* \rightarrow \ell\ell$ BACKGROUND

1933 In the same flavor channels, the $Z/\gamma^* \rightarrow \ell\ell$ background is dominant and thus must be estimated
1934 correctly. In both the BDT and cut-based analyses, the background is estimated using the so-called

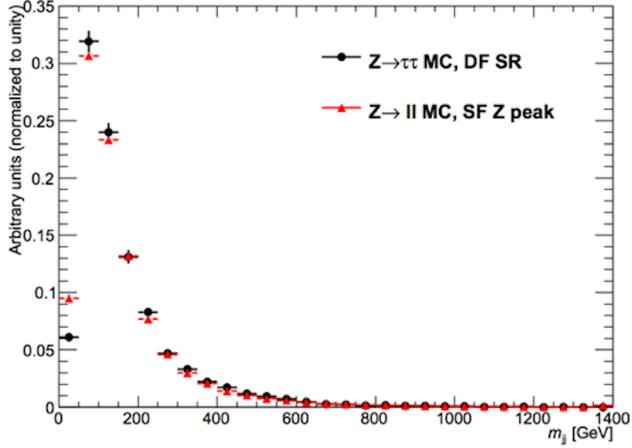


Figure 5.9: Comparison of m_{jj} shape in a same flavor $Z \rightarrow \ell\ell$ control region and the VBF cut-based signal region.

$\beta_{\tau\tau}$	0.97 ± 0.04
Cut	Correction factors
$m_{jj} > 600 \text{ GeV}$	1.09 ± 0.01
$\Delta y_{jj} > 3.6$	1.14 ± 0.02
CJV	1.20 ± 0.02
OLV	1.17 ± 0.03
$m_{jj} < 1 \text{ TeV}$	1.17 ± 0.06
$m_{jj} > 1 \text{ TeV}$	1.18 ± 0.13

Table 5.10: $Z/\gamma^* \rightarrow \tau\tau$ correction factors for the VBF cut-based analysis. Uncertainties are statistical only.

1935 “ABCD” method. The ABCD method creates four different regions by defining cuts on two variables.
 1936 One of the regions (A) is the signal region, while the other regions are defined by inverting one of both
 1937 of the cuts. in this case, the two variables used are $m_{\ell\ell}$ and E_T^{miss} , because inverting either of the SR cuts
 1938 on these variables will give regions rich in the $Z/\gamma^* \rightarrow \ell\ell$ background. Figure 5.10 illustrates the general
 1939 strategy for each region.

1940 In both of the cut-based and BDT analyses, the Z peak region is defined with $|m_{\ell\ell} - m_Z| < 15 \text{ GeV}$.
 1941 In the cut-based analysis, low $m_{\ell\ell}$ corresponds to $m_{\ell\ell} < 50 \text{ GeV}$ (this defines the cut-based SR) while
 1942 in the BDT it is $m_{\ell\ell} < 75 \text{ GeV}$. In the cut-based, high and low E_T^{miss} are defined as opposite ends of
 1943 the 55 GeV cut applied for the signal region definition. The BDT low E_T^{miss} region is between 25 and
 1944 45 GeV, while the high E_T^{miss} region is $E_T^{\text{miss}} > 45 \text{ GeV}$.

1945 Once the regions are defined, the final signal region background estimate is done by taking the esti-

Region A (SR)	Region C
High E_T^{miss}	High E_T^{miss}
Low $m_{\ell\ell}$	Z peak
Region B	Region D
Low E_T^{miss}	Low E_T^{miss}
Low $m_{\ell\ell}$	Z peak

Figure 5.10: General illustration of the ABCD region definitions for $Z/\gamma^* \rightarrow \ell\ell$ background estimation.

mate in region B and extrapolating it to the signal region (A) by multiplying it by the ratio of regions C and D. Effectively, the Z peak region is used to estimate the efficiency of the E_T^{miss} cut in data, and then this efficiency is applied in the low $m_{\ell\ell}$ region. An additional correction is also applied for the non-closure of the method in MC. This is summarized in equations 5.6 and 5.7.

$$N_{Z/\gamma^* \rightarrow \ell\ell}^{\text{SR}} = N_{Z/\gamma^* \rightarrow \ell\ell}^{\text{B}} \times \frac{N_{Z/\gamma^* \rightarrow \ell\ell}^{\text{C}}}{N_{Z/\gamma^* \rightarrow \ell\ell}^{\text{D}}} \times f_{\text{corr}} \quad (5.6)$$

$$f_{\text{corr}} = \frac{B_{\text{MC}}^{\text{A}}/B_{\text{MC}}^{\text{B}}}{B_{\text{MC}}^{\text{C}}/B_{\text{MC}}^{\text{D}}} \quad (5.7)$$

Here, the N refer to data yields in each region with the non Z/γ^* backgrounds subtracted, while B refer to the Z/γ^* yields in MC in each region.

A normalization factor $\beta_{\ell\ell}$ is computed for each analysis as the ratio of the predicted data yield to the MC yield in the SR. The shape of the BDT distribution is taken from data region B, while the shape of the m_T distribution in the cut-based analysis is taken from Z/γ^* MC in the SR. The values of the $\beta_{\ell\ell}$ in the cut-based and BDT analyses from this method are summarized in table 5.11. They are quite consistent with one another within the statistical uncertainties. In the cut-based analysis, the same cut efficiency correction factors shown in table 5.10 are also applied (in product with the $\beta_{\ell\ell}$) in the same flavor channels to this background, as they were derived in the Z peak region.

	β_t
BDT Bin 1	1.01 ± 0.15
BDT Bin 2	0.89 ± 0.28
Cut-based	0.81 ± 0.21

Table 5.11: $Z/\gamma^* \rightarrow \ell\ell$ normalization factors for cut-based and BDT analyses. Uncertainties are statistical only.

1959 5.5.5 WW AND OTHER DIBOSON BACKGROUNDS

1960 The WW and other diboson backgrounds have both their shape and normalization taken from MC
 1961 simulation. They are validated in dedicated control regions and found to agree with data well.
 1962 As WW is the largest of these backgrounds and is irreducible, validating the estimate is of particular
 1963 importance. The validation region is constructed by requiring the pre-selection cuts on leptons and $m_{\ell\ell}$,
 1964 $n_b = 0$, and $m_T > 100$ GeV. The m_{T2} variable[?] is an additional discriminant that will isolate
 1965 the WW background, and a requirement of $m_{T2} > 160$ GeV is placed to define the WW validation
 1966 region. This cut gives a 60% purity for the validation region. The derived normalization factor in the
 1967 region is 1.15 ± 0.19 and is thus consistent with unity. Figure 5.11 shows the m_{T2} distribution and how
 1968 it distinguishes the WW background.

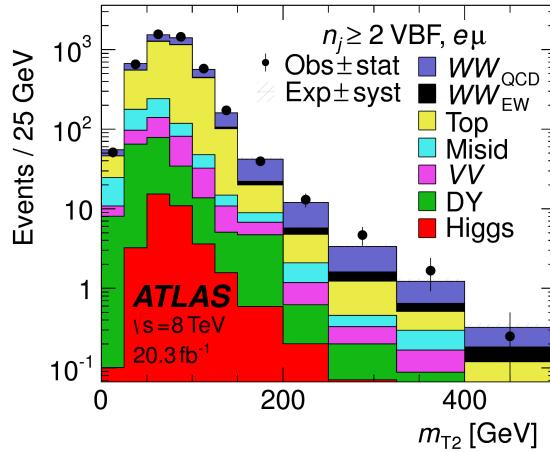


Figure 5.11: Distribution of m_{T2} in the WW validation region of the VBF analysis[21].

1969 5.5.6 HIGGS PRODUCTION VIA GLUON-GLUON FUSION

1970 Because this analysis is dedicated to measuring the VBF contribution to Higgs production, the com-
1971 ponent of Higgs production from gluon-gluon fusion is treated as a background. The shape is taken
1972 directly from simulation, using the generators described in table 5.4. In the final combined fit of all dif-
1973 ferent signal regions, the normalization is controlled by either a combined signal strength parameter μ ,
1974 which controls the normalization of both ggF and VBF production, or a separate parameter μ_{ggF} de-
1975 pending on the interpretation being presented in the final results.

1976 5.5.7 BACKGROUNDS WITH MISIDENTIFIED LEPTONS

1977 As discussed previously, the $W + \text{jets}$ and QCD multijet backgrounds are derived with fully data-driven
1978 methods. These backgrounds do not make a large contribution to the final VBF signal region but their
1979 estimation methods are discussed briefly here.

1980 $W + \text{jets}$ BACKGROUND

1981 The $W + \text{jets}$ background enters the signal region by having one of the jets mis-reconstructed as a lep-
1982 ton. The background is estimated by constructing a control sample with two leptons, where one lepton
1983 passes the usual lepton quality cuts but the second lepton fails one of those cuts (also known as the “anti-
1984 identified” lepton). This control region is rich in the $W + \text{jets}$ contribution because if a second lepton is
1985 reconstructed in a $W + \text{jets}$ event it is likely to be poor quality. The purity of this $W + \text{jets}$ control sample
1986 is 85% to 90% depending on the exact configuration of leptons in the final state.

1987 The signal region estimate of $W + \text{jets}$ is estimated by extrapolation from the control sample to the sig-
1988 nal region using extrapolation factors derived in a $Z + \text{jets}$ control sample in data. The extrapolation fac-
1989 tor is the ratio of the number of lepton candidates satisfying all quality criteria to the number of lepton
1990 candidates anti-identified. This ratio is measured in bins of p_T and η . Thus, the final signal region esti-
1991 mate (binned as the extrapolation factor is binned) is simply the number of events in the anti-identified
1992 lepton control sample multiplied by the extrapolation factor derived from the $Z + \text{jets}$ control sample.
1993 Figure 5.12 shows the extrapolation factors derived for electrons and muons.

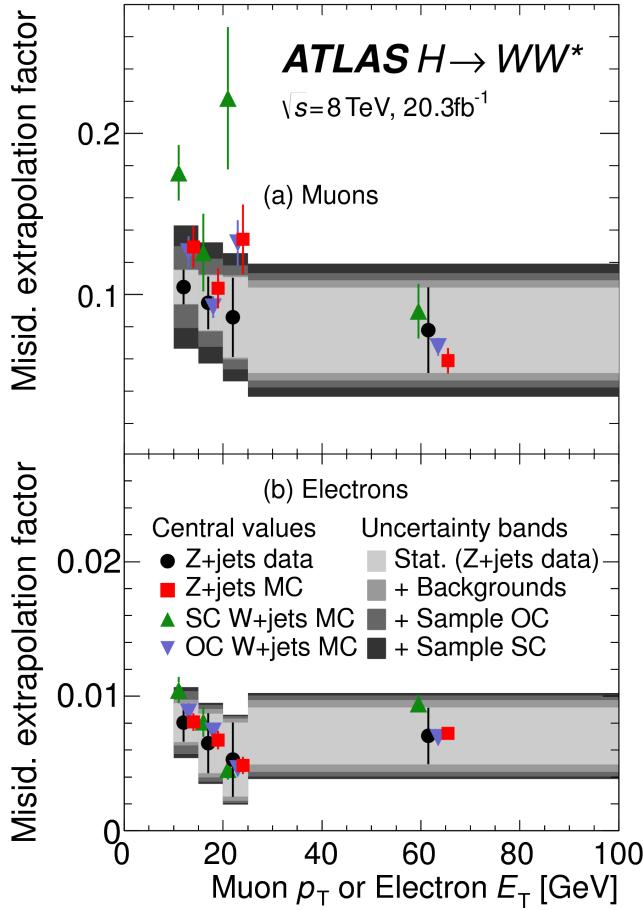


Figure 5.12: Extrapolation factors for the $W + \text{jets}$ estimate derived for muons (a) and electrons (b) as a function of lepton p_T [21].

1994 QCD MULTIJET BACKGROUND

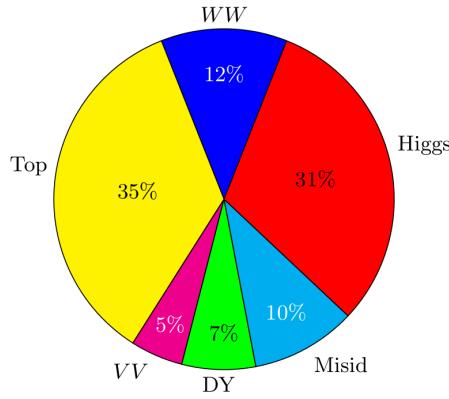
1995 The method for estimating the multijet background is very similar to the $W + \text{jets}$ estimation method.
1996 The control sample in this case has two anti-identified leptons but otherwise satisfies all signal region
1997 requirements. The extrapolation factor is estimated from a multijet sample and applied twice to the
1998 control sample.

1999 5.5.8 BACKGROUND COMPOSITION IN FINAL SIGNAL REGION

2000 After all of these estimation procedures, the final signal region composition can be calculated. The esti-
2001 mated yields are all shown in table 5.7. Figure 5.13 shows the relative percentages of the different back-

ground for the different flavor and same flavor final states. In $e\mu$, the leading backgrounds are top backgrounds, ggF Higgs, and SM WW production. In $ee/\mu\mu$, the leading background is Drell-Yan, followed by top and ggF Higgs.

(a) $n_j \geq 2$ VBF, $e\mu$



(b) $n_j \geq 2$ VBF, $ee/\mu\mu$

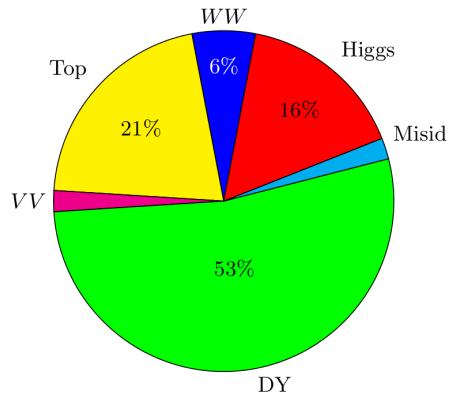


Figure 5.13: Background composition in final VBF signal region[21].

5.6 SYSTEMATIC UNCERTAINTIES

There are two main types of systematic uncertainties that are assessed for the analysis. First, theoretical uncertainties associated with the various signal and background yield estimates are discussed. Then, experimental uncertainties due to detector effects are shown. Normalization uncertainties refer to uncertainties that affect the cross section of the process in question in the signal region being probed. Shape uncertainties refer to systematic uncertainties that affect the shape of the final discriminating variable (either m_T or O_{BDT}).

5.6.1 THEORETICAL UNCERTAINTIES

There are four main components to theoretical uncertainties assigned to signal and background processes taken from Monte Carlo. Each one is a different source of variation in the overall acceptance for that process. The first involves variation of the QCD renormalization and factorization scales used in the calculation. In this case, the two scales are varied independently and simultaneously by factors of two high or low and quantifying the resulting variation in normalization and shape for the process. This

2018 approximates the correction to the cross section that would come from including the next order of the
 2019 QCD calculation (referred to as scale uncertainty). Next, there is an uncertainty associated with the PDF
 2020 set used in generating the events. The uncertainty eigenvectors for the given PDF set are studied, and the
 2021 envelope of maximal variation is taken as an uncertainty. Finally, there are two uncertainties associated
 2022 with the choice of MC software (referred to as PDF uncertainty). An uncertainty associated with the
 2023 generator chosen for the hard scattering process is evaluated by keeping the parton showering software
 2024 constant but varying the matrix element generator and taking the maximal variation as an uncertainty
 2025 (referred to as the generator uncertainty). The converse variation can also be done, where the matrix ele-
 2026 ment generator remains constant and the generator used for the underlying event/parton shower mod-
 2027 eling is varied (referred to as the UE/PS uncertainty). In cases where the background is normalized in a
 2028 control region, the systematic uncertainty arises from variations of the extrapolation factor α between
 2029 the CR and the SR, which can affect the normalization of the background in the SR.

2030 There are two additional uncertainties that are applied to the Higgs processes as well. First, there are
 2031 uncertainties assigned to the Higgs total production cross section. Then, there are uncertainties assigned
 2032 based on the fact that the analysis is done in exclusive jet bins and it is possible for signal events to mi-
 2033 grate from one bin to the next depending on the presence or absence of jets. These are assigned using the
 2034 Jet Veto Efficiency (JVE) procedure[?] for ggF events and the Stewart-Tackmann (ST) method[87] for
 2035 VBF production.

2036 Table 5.12 shows the total theory uncertainties on the backgrounds in the cut-based analysis. These are
 2037 the sum in quadrature of the uncertainties from each of the variations described above.

Process	Theory syst. (%)
ggF H	48
Top	26
QCD WW	37
$Z/\gamma^* \rightarrow \tau\tau$	6.1

Table 5.12: Systematic uncertainties for various processes in the cut-based VBF analysis, given in units of % change in yield.
 Values are given for the low m_{jj} signal region.

2038 Figures 5.14 and 5.15 show the variations in the extrapolation factor from the PDF and QCD uncer-
 2039 tainties on the top background estimate, binned in m_T , for the cut-based analysis. In both cases, there

2040 was no significant shape uncertainty but normalization uncertainties were assigned according to the
 2041 maximal variation. These uncertainties enter into the 26% total uncertainty on top quoted in table 5.12.
 2042 While the estimate for the same-flavor $Z/\gamma^* \rightarrow \ell\ell$ background is data-driven, there is still a sys-
 2043 tematic uncertainty taken for the non-closure of the method in Monte Carlo. This is taken as the max-
 2044 imum of the deviation of the non-closure factor f_{corr} from unity and its uncertainty, or $\max(|1 -$
 2045 $f_{\text{corr}}|, \delta f_{\text{corr}})$. For the cut-based analysis this non-closure uncertainty 23%, while for the BDT analy-
 2046 sis it is 17%.

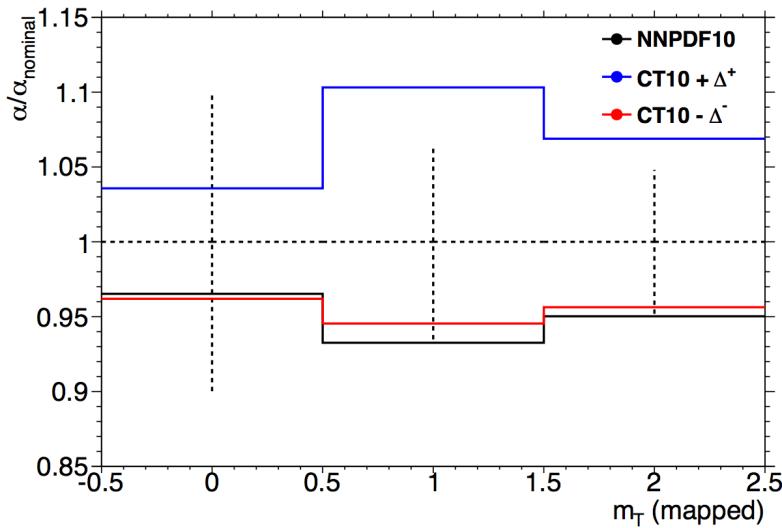


Figure 5.14: Variations in the top background extrapolation factor in the cut-based analysis due to PDF uncertainties, binned in m_T .

2047 5.6.2 EXPERIMENTAL UNCERTAINTIES

2048 In this analysis, the theoretical uncertainties end up being the most dominant, but there are some ex-
 2049 perimental uncertainties that make a contribution as well. The first is the uncertainty on the measured
 2050 integrated luminosity, which affects backgrounds whose normalization is taken from MC and is mea-
 2051 sured to be 2.8% in the 8 TeV dataset [88]. The dominant sources of uncertainty overall are uncertainties
 2052 on the jet energy scale and resolution and the b -tagging efficiency. Additional sources include lepton
 2053 uncertainties on identification, resolution, and trigger efficiency, as well as uncertainties on the missing
 2054 transverse momentum .

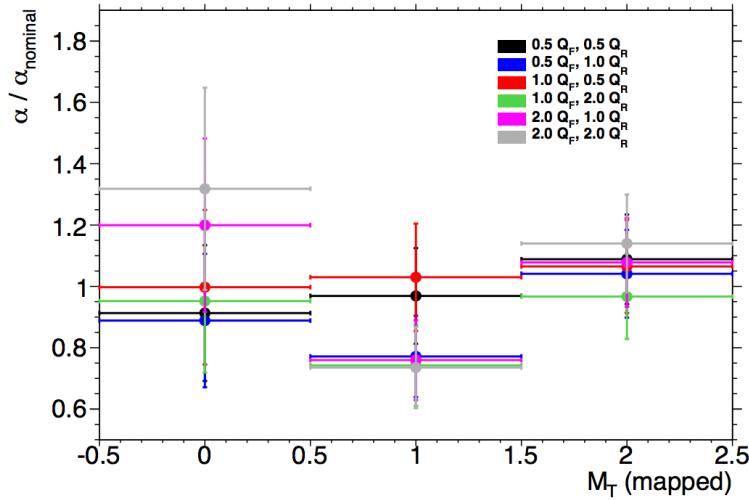


Figure 5.15: Variations in the top background extrapolation factor in the cut-based analysis due to QCD scale uncertainties, binned in m_T .

2055 The jet energy scale uncertainty is split into several independent components, including jet-flavor
 2056 dependent calorimeter response uncertainties, uncertainties on modeling of pile-up interactions, uncer-
 2057 tainties on extrapolation from the central to forward detector regions, and MC non-closure [89]. The
 2058 uncertainty on energy scale for jets used in this analysis ranges from 1% to 7% depending on the jet p_T
 2059 and η . The jet energy resolution varies from 5% to 20%, with uncertainties ranging from 2% to 40%
 2060 (the largest uncertainties occurring at the selection threshold).

2061 The b-tagging efficiency is independently measured in data samples enriched in dileptonic decays of
 2062 $t\bar{t}$ events or in events where a muon is reconstructed in the vicinity of a jet[90, 91]. The efficiencies and
 2063 their uncertainties are binned in p_T and decomposed into uncorrelated components using an eigenvec-
 2064 tor method[?]. Uncertainties on the efficiency range from 1% to 7.8%. The uncertainty on the rate of
 2065 misidentification of c -jets as b -jets ranges from 6-14%, while the uncertainty on the rate of light jet mis-
 2066 tagging ranges from 9-19% depending on p_T and η .

2067 The total experimental uncertainties on different signal and background components are summa-
 2068 rized in table 5.13. They are compared to the level of other statistical and systematic uncertainties as well.
 2069 Overall, the experimental uncertainties are sub-dominant compared to the statistical and theoretical
 2070 uncertainties.

Sample	Total error	Stat. error	Expt. syst. err.	Theo. syst. err.
$n_j \geq 2$ VBF-enriched				
N_{sig}	13	-	6.8	12
N_{bkg}	9.2	4.7	6.4	4.5
N_{WW}	32	-	14	28
N_{top}	15	9.6	7.6	8.5
N_{misid}	22	-	12	19
N_{VV}	20	-	12	15
$N_{\tau\tau} (\text{DY})$	40	25	31	2.9
$N_{ee/\mu\mu} (\text{DY})$	19	11	15	-

Table 5.13: Composition of the post-fit uncertainties (in %) on the total signal (N_{sig}), total background (N_{bkg}), and individual background yields in the VBF analysis[21].

2071 5.7 RESULTS

2072 While the combined results of all the $H \rightarrow WW^*$ sub-analyses will be discussed in the next chapter, this
 2073 section presents the results of the VBF specific analysis and interpretations.

2074 As table 5.7 shows, the final cut-based signal region contains 20 events in data with $m_T < 150$ GeV,
 2075 14 coming from the $e\mu$ channel and 6 coming from the $ee + \mu\mu$ channel. The BDT analysis has many
 2076 more candidates due to its looser selection, and the yields in each bin of O_{BDT} are shown in table 5.14.

2077 Figure 5.16(a) shows the final distribution of data candidates compared to the expected m_T distri-
 2078 bution for signal and background. The data are very consistent with a VBF Higgs hypothesis. Fig-
 2079 ure 5.16(b) shows where the data candidates fall in the two-dimensional binning of m_T and m_{jj} used
 2080 in the fit for the cut-based analysis.

2081 Figure 5.17 shows the distributions of O_{BDT} and m_T in the VBF BDT analysis. Again the data are
 2082 quite consistent with a VBF Higgs hypothesis.

2083 Because the cut-based result is used as a validation for the BDT analysis and the two signal regions are
 2084 not fully orthogonal, it is interesting to explore which events overlap between the two analyses. Of the
 2085 twenty events in the cut-based signal region, only seven were not selected by the BDT analysis, while the
 2086 other thirteen also enter the BDT signal region. Figure ?? shows where the different analysis candidates
 2087 lie in the m_{jj} - m_T plane. This shows clearly that the advantage of the BDT analysis is that it can extract

(a) Before the BDT classification

Selection	Summary						Composition of N_{bkg}											
	$N_{\text{obs}}/N_{\text{bkg}}$	N_{obs}	N_{bkg}	N_{signal}	N_{ggF}	N_{VBF}	N_{VH}	N_{WW}	$N_{\text{WW}}^{\text{NEW}}$	$N_{\text{WW}}^{\text{QCD}}$	N_t	N_{missid}	N_{Wj}	N_{ij}	N_{VV}	$N_{e\mu/\mu\mu}$	$N_{\tau\tau}^{\text{QCD}}$	$N_{\tau\tau}^{\text{NEW}}$
$e\mu$ sample	1.04 ± 0.04	718	689	13	15	2.0	90	II	327	42	29	23	31	—	2.2	130	2	—
$ee/\mu\mu$ sample	1.18 ± 0.08	469	397	6.0	7.7	0.9	37	3	132	17	5.2	1.2	10.1	168	23	1	—	—

(b) Bins in O_{BDT}

$e\mu$ sample	Bins in O_{BDT}																	
	Bin 0 (not used)	Bin 1	Bin 2	Bin 3	Bin 4	Bin 5	Bin 6	Bin 7	Bin 8	Bin 9	Bin 10	Bin 11	Bin 12	Bin 13	Bin 14	Bin 15	Bin 16	Bin 17
$ee/\mu\mu$ sample	1.91 ± 0.08	396	345	3.8	1.3	0.8	33	2	123	16	4.1	1.1	8.8	137	20.5	0.5	—	—
Bin 0 (not used)	0.82 ± 0.14	53	45	1.5	2.2	0.1	3.0	0.5	10.4	1.8	0.8	0.2	0.9	26	1.7	0.1	—	—
Bin 1	1.77 ± 0.49	14	7.9	0.6	2.5	—	0.8	0.3	1.1	0.2	0.2	—	0.3	4.4	0.3	0.1	—	—
Bin 2	6.52 ± 2.87	6	0.9	0.2	1.7	—	0.1	0.2	0.2	—	—	—	—	0.7	—	—	—	—
Bin 3	5.41 ± 2.32	6	1.1	0.4	3.1	—	0.3	0.2	0.3	0.1	—	—	0.1	—	0.1	0.1	0.1	0.1

Table 5.14: Event selection for the VBF BDT analysis. The event yields in (a) are shown after the preselection and the additional requirements applied before the BDT classification (see text). The event yields in (b) are given in bins in O_{BDT} after the classification[21].

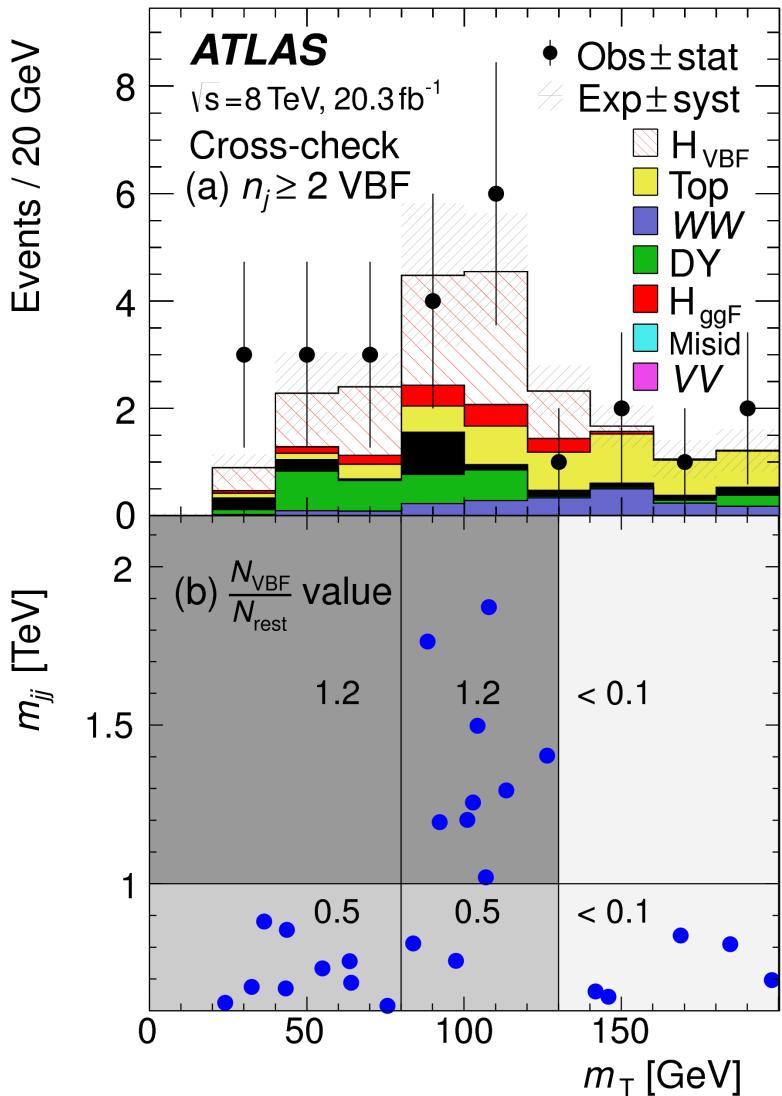


Figure 5.16: Postfit distributions in the cut-based VBF analysis. Panel (a) shows the one-dimensional m_T distribution, while (b) shows the data candidates split into the bins of m_T and m_{jj} used in the final fit[21].

2088 signal candidates lower m_{jj} region due to its ability to recognize correlations with other variables.
 While the context of these results in the broader $H \rightarrow WW^*$ statistical analysis will be presented
 2089 in the next chapter, the significance of the VBF observation can be shown here. In the BDT analysis,
 2090 the expected signal significance was 2.7σ , while the observed significance was 3.1σ . In the cut-based
 2091 analysis, the expected significance was 2.1σ and the observed significance was 3.0σ . The compatibility
 2092 between these two results can be evaluated by computing the probability of observing a larger difference
 2093 between these two results can be evaluated by computing the probability of observing a larger difference

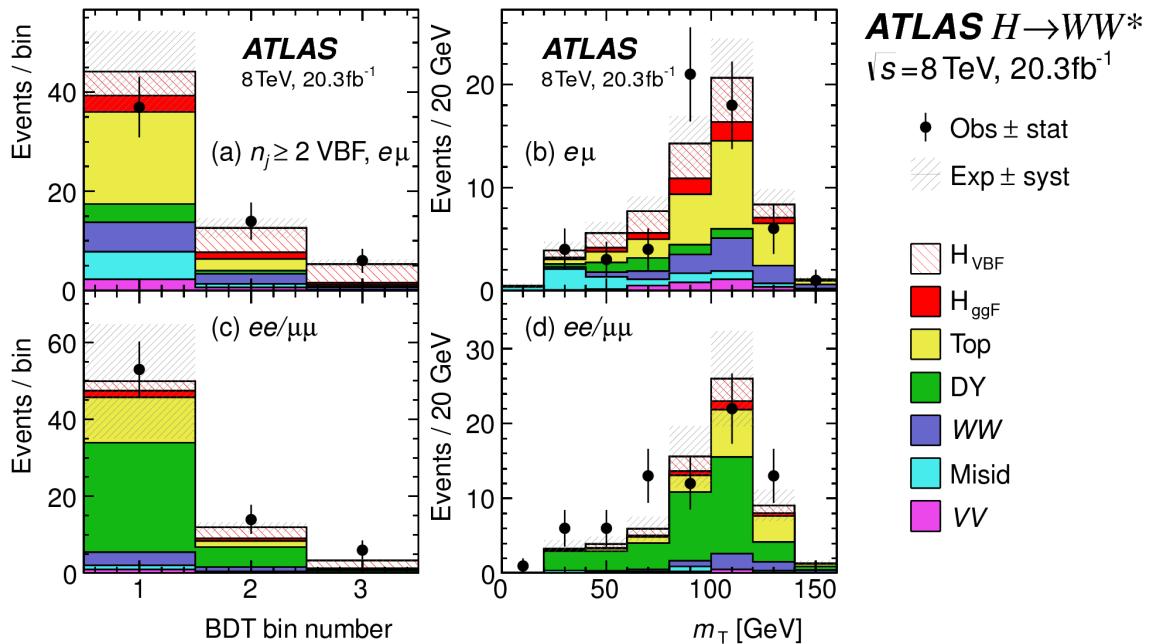


Figure 5.17: Postfit distributions in the BDT VBF analysis[21].

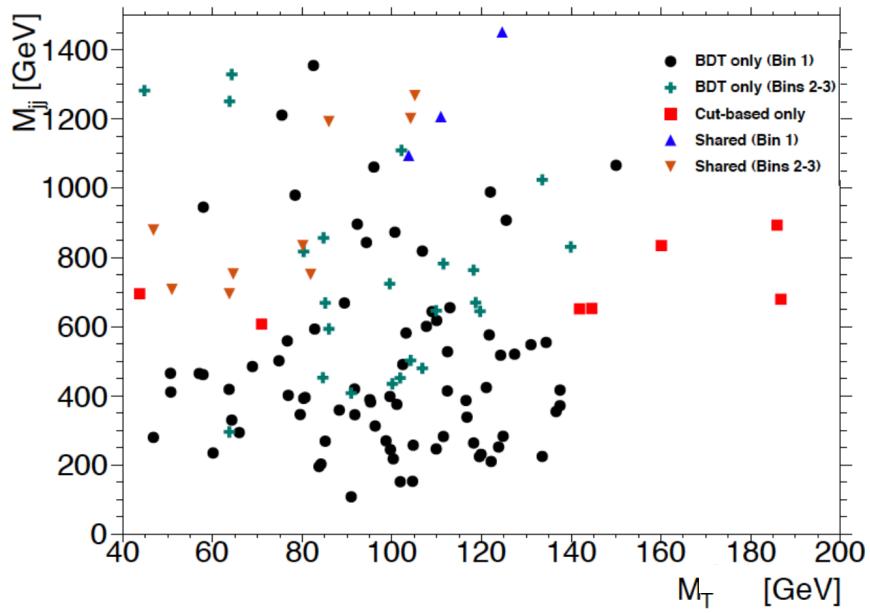


Figure 5.18: Overlap between cut-based and BDT VBF signal region candidates in the m_{jj} - m_T plane.

2094 in Z_0 values than the one measured. Using toy Monte Carlo with the ggF signal strength fixed to unity
2095 and considering only statistical uncertainties, this probability is computed to be 79%, indicating good
2096 agreement between the analyses.

2097 This result represents the first observation of the vector boson fusion production of a Higgs boson.

*The feeling is less like an ending than just another
starting point.*

Chuck Palahniuk

6

2098

2099

Combined Run I $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$

2100

results

2101 6.1 INTRODUCTION

2102 In the final statistical analysis of $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$, the dedicated gluon-gluon fusion and vector
2103 boson fusion sensitive signal regions are all combined into a single fit to determine the main parameters
2104 of interest, the Higgs signal strength μ and mass m_H . Therefore, while the specific requirements applied
2105 for the VBF sensitive analysis are discussed in chapter 5, the final measurement of these parameters can
2106 only be discussed in combination with the results of the ggF dedicated analysis. For example, because
2107 ggF Higgs production is considered a background in the VBF analysis, the ggF dedicated signal regions
2108 can actually constrain the normalization of this background in the VBF dedicated region.

2109 This chapter presents the combined interpretation of results in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis

SR category i			Fit var.
n_j , flavor	$\otimes m_{\ell\ell}$	$\otimes p_T^{\ell 2}$	
$n_j = 0$			
$e\mu$	$\otimes [10, 30, 55]$	$\otimes [10, 15, 20, \infty]$	$\otimes [e, \mu]$
$ee/\mu\mu$	$\otimes [12, 55]$	$\otimes [10, \infty]$	m_T
$n_j = 1$			
$e\mu$	$\otimes [10, 30, 55]$	$\otimes [10, 15, 20, \infty]$	$\otimes [e, \mu]$
$ee/\mu\mu$	$\otimes [12, 55]$	$\otimes [10, \infty]$	m_T
$n_j \geq 2$ ggF			
$e\mu$	$\otimes [10, 55]$	$\otimes [10, \infty]$	m_T
$n_j \geq 2$ VBF			
$e\mu$	$\otimes [10, 50]$	$\otimes [10, \infty]$	O_{BDT}
$ee/\mu\mu$	$\otimes [12, 50]$	$\otimes [10, \infty]$	O_{BDT}

Table 6.1: All signal regions definitions input into final statistical fit[21].

for gluon fusion and vector boson fusion Higgs production. First, the results of the dedicated gluon fusion search are presented. Then, a comparison of the individual production mode signal strengths (μ_{ggF} and μ_{VBF} and a measurement of the combined signal strength (μ) are shown. Subsequently, the measured values of the Higgs couplings to fermions and vector bosons is presented. Finally, the cross section measurement for ggF and VBF production are shown.

6.2 RESULTS OF DEDICATION GLUON FUSION $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ SEARCH

The details of the dedicated gluon fusion $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ search are not discussed in this thesis and instead left to more comprehensive sources[21]. However, a brief summary of the results are essential for describing the results of the full analysis and interpreting the results of the dedicated VBF search in this broader context.

Table 6.1 shows the individual signal regions that were input into the final statistical fit. The ggF dedicated bins use m_T as their discriminating variable and are separated into bins of p_T of the subleading lepton as well. The VBF dedicated bin uses the O_{BDT} distribution as its final discriminant.

Table 6.2 shows the yields in the various signal regions in both data and expected signal and back-

2124 grounds. The yields for signal and background are all scaled according to the final normalizations calcu-
2125 lated in the fit.

	N_{obs}	N_{bkg}	N_{ggF}	N_{VBF}
$n_j = 0$	3750	3430 ± 90	300 ± 50	8 ± 4
$n_j = 1$	1596	1470 ± 40	102 ± 26	17 ± 5
$n_j \geq 2, \text{ggF } e\mu$	1017	960 ± 40	37 ± 11	13 ± 1.4
$n_j \geq 2, \text{VBF}$	130	99 ± 9	7.7 ± 2.6	21 ± 3

Table 6.2: Post-fit yields in the different ggF and VBF dedicated signal regions[21].

2126 Figure 6.1 shows the final post-fit m_T distribution in the $n_j \leq 1$ regions. The data are very consistent
2127 with the hypothesis of ggF Higgs production.

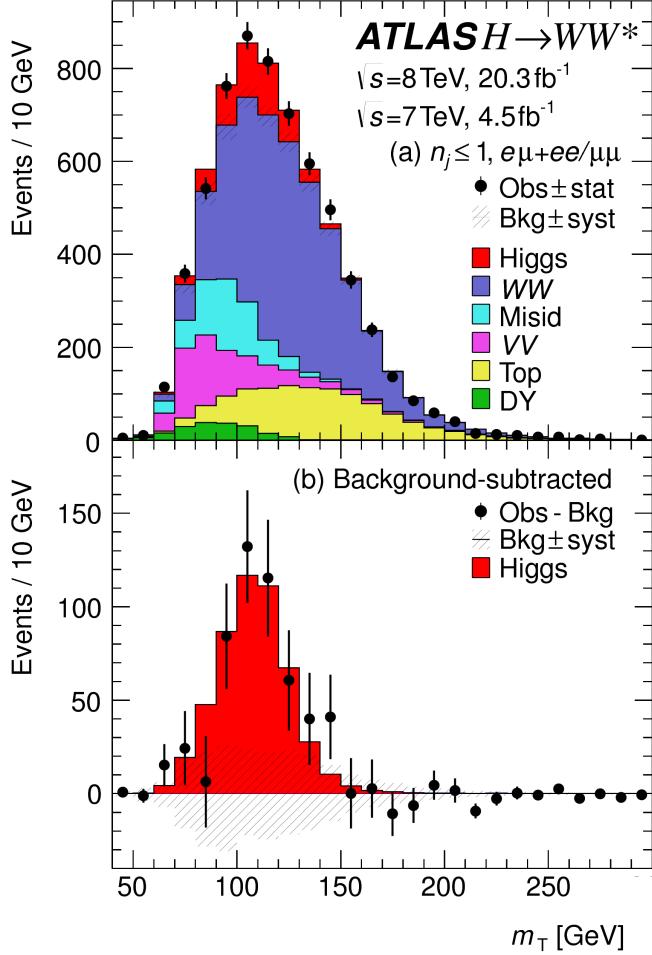


Figure 6.1: Post-fit m_T distribution in the $n_j \leq 1$ regions[21].

2128 These yields are used as input, along with the VBF results in chapter 5, for the physical interpretation
2129 of results presented in subsequent sections.

2130 **6.3 SIGNAL STRENGTH MEASUREMENTS IN ggF AND VBF PRODUCTION**

2131 When all of the signal regions are combined in the fit, there can be a combined measurement of the sig-
2132 nal strength as well as the individual ggF and VBF signal strengths. The combined signal strength is the
2133 ratio of the sum of the gluon fusion and VBF cross sections to the theory prediction, or a singal strength
2134 for the total Higgs production cross section that this analysis is sensitive to. The final measured com-
2135 bined signal strength μ is measured shown in equation 6.1.

$$\begin{aligned} \mu &= 1.09 \quad {}^{+0.16}_{-0.15} \text{ (stat.)} \quad {}^{+0.08}_{-0.07} \left(\begin{array}{l} \text{expt} \\ \text{syst} \end{array} \right) \quad {}^{+0.15}_{-0.12} \left(\begin{array}{l} \text{theo} \\ \text{syst} \end{array} \right) \quad \pm 0.03 \left(\begin{array}{l} \text{lumi} \\ \text{syst} \end{array} \right) \\ &= 1.09 \quad {}^{+0.16}_{-0.15} \text{ (stat)} \quad {}^{+0.17}_{-0.14} \text{ (syst)} \\ &= 1.09 \quad {}^{+0.23}_{-0.21}. \end{aligned} \tag{6.1}$$

2136 Figure 6.2 gives the best fit signal strength $\hat{\mu}$ as a function of hte hypothesized Higgs mass. The value
2137 at 125.36 GeVcorresponds to the μ quoted in equation 6.1. This value of the Higgs mass is used because
2138 it is the most precise mass measurement from ATLAS, a result of the combined $\gamma\gamma$ and ZZ mass mea-
2139 surements[?].

2140 As explained in chapter 3, a probability p_0 can be computed using the test statistic q_0 to quantify the
2141 probability that the background could fluctuate to produce an excess at least as large as the one observed
2142 in the data. The local p_0 value is shown in figure 6.3 as a function of m_H . The minimum p_0 value is
2143 at $m_H = 130$ GeV and coresponds to a significance of 6.1σ . The curve is relatively flat and the sig-
2144 nificance is the same at 125.36 GeVwithin the quoted precision. The expected significance for a signal
2145 with strength $\mu = 1.0$ is 5.8σ . This represents the first discovery level significance measurement in the
2146 $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ analysis.

2147 All the results presented so far in this section have been for the combined gluon fusion and VBF pro-
2148 duction modes. However, each signal strength can be calculated separately in the likelihood as well.
2149 There are two ways to do this. First, the likelihood can be parameterized in terms of a single parameter,

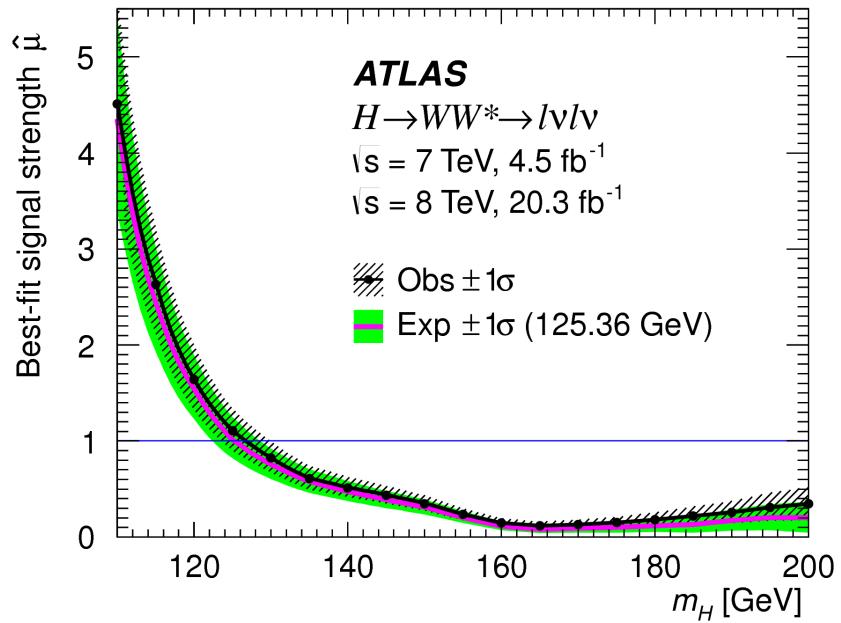


Figure 6.2: Best fit signal strength $\hat{\mu}$ as a function of hypothesized m_H [21].

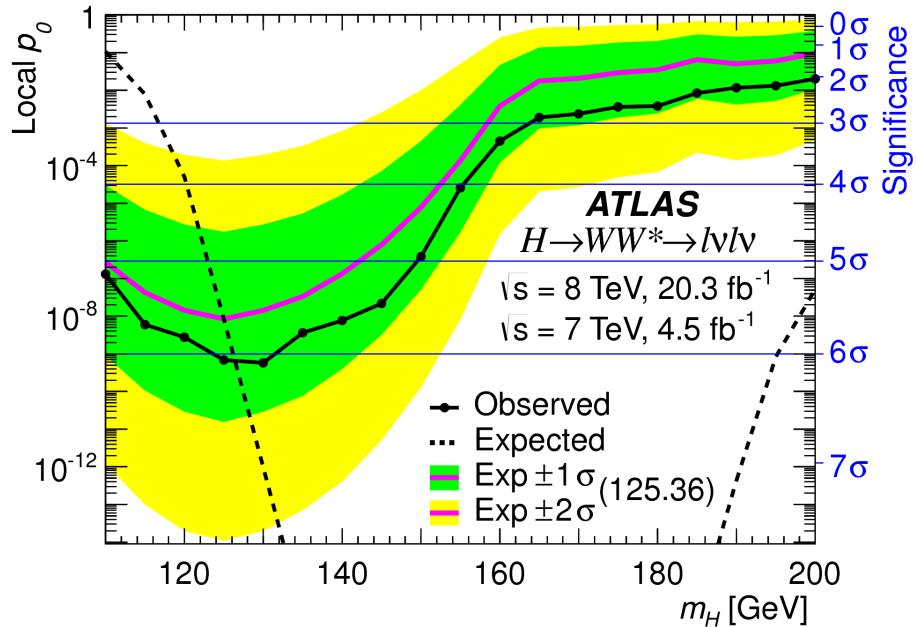


Figure 6.3: Local p_0 as a function of m_H [21].

the ratio of the VBF and gluon fusion signal strengths. With this method, the significance of the VBF observation can be evaluated. Figure 6.4 shows the likelihood as a function of the ratio $\mu_{\text{VBF}}/\mu_{\text{ggF}}$.

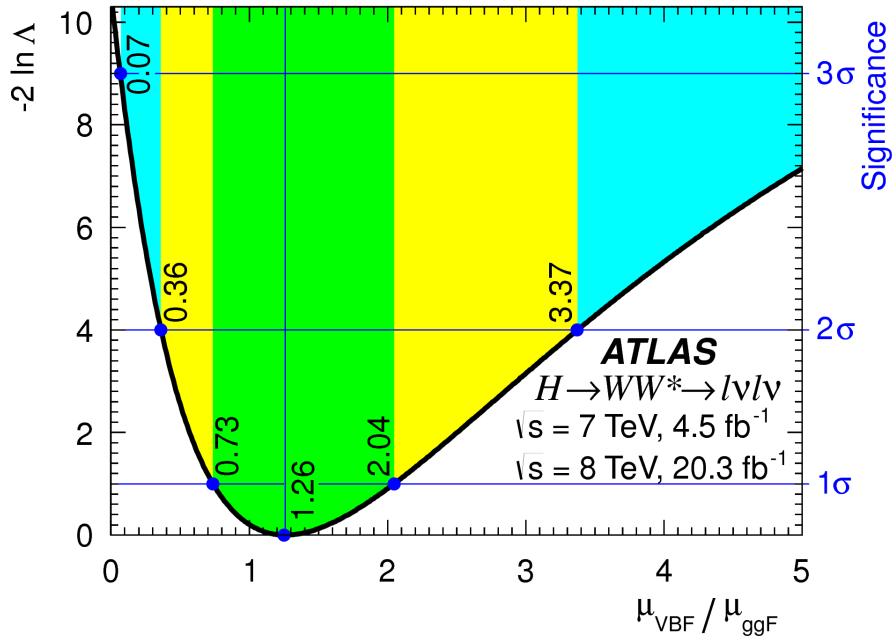


Figure 6.4: Likelihood as a function of $\mu_{\text{VBF}} / \mu_{\text{ggF}}$ [21].

2152 The best fit value of the ratio of signal strengths is shown in equation 6.2. Within the quoted uncer-
 2153 tainties, it is consistent with a ratio of unity.

$$\frac{\mu_{\text{VBF}}}{\mu_{\text{ggF}}} = 1.26^{+0.61}_{-0.45} \text{ (stat.)}^{+0.50}_{-0.26} \text{ (syst.)} = 1.26^{+0.79}_{-0.53} \quad (6.2)$$

2154 The null hypothesis for VBF production corresponds to a ratio of $\mu_{\text{VBF}} / \mu_{\text{ggF}} = 0$. The likelihood
 2155 in figure 6.4 gives a significance of 3.2σ at $\mu_{\text{VBF}} / \mu_{\text{ggF}} = 0$, as quoted in chapter 5.

2156 In addition to the ratio of signal strengths, each signal strength can be varied independently in the
 2157 likelihood as well. Figure 6.5 shows the two dimensional likelihood scan in the $\mu_{\text{ggF}}-\mu_{\text{VBF}}$ plane. The
 2158 best fit values of the two signal strengths are shown in equation 6.3. Both are consistent with unity
 2159 within their uncertainties.

$$\begin{aligned} \mu_{\text{ggF}} &= 1.02 \pm 0.19^{+0.22}_{-0.18} = 1.02^{+0.29}_{-0.26} \\ \mu_{\text{VBF}} &= 1.27 \pm 0.40^{+0.44}_{-0.40} \pm 0.21^{+0.29}_{-0.21} = 1.27^{+0.53}_{-0.45} \end{aligned} \quad (6.3)$$

(stat.) (syst.)

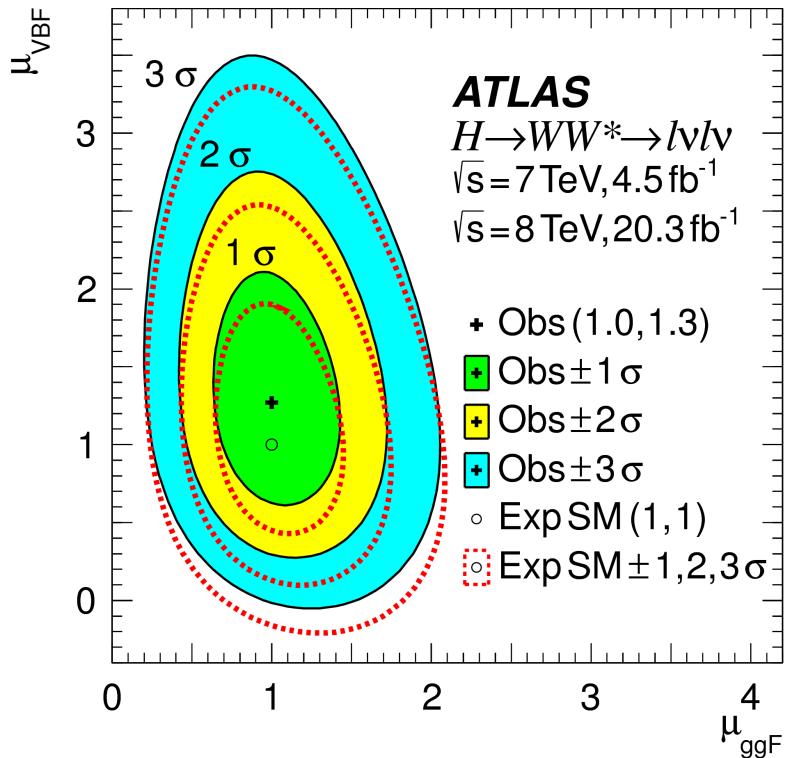


Figure 6.5: Likelihood scan as a function of μ_{VBF} and μ_{ggF} [21].

2160 6.4 MEASUREMENT OF HIGGS COUPLINGS TO VECTOR BOSONS AND FERMIONS

2161 Similar to the parameterization of signal strength, the couplings of the Higgs to fermions and bosons can
 2162 also be parameterized. The parameter of interest in this case is κ , or the ratio of the measured coupling
 2163 to the standard model expectation. Both the fermion and boson couplings have these so-called scale fac-
 2164 tors, κ_F for fermions and κ_V for bosons. Gluon fusion production is sensitive to the fermion couplings
 2165 through the top quark loops in its production, while VBF production is sensitive to the vector boson
 2166 couplings in its production. Both modes are sensitive to the vector boson couplings in their decays. The
 2167 signal strengths will have dependence on the coupling scale factors as described in equation 6.4[2].

$$\begin{aligned} \mu_{\text{ggF}} &\propto \frac{\kappa_F^2 \cdot \kappa_V^2}{(\mathcal{B}_{H \rightarrow f\bar{f}} + \mathcal{B}_{H \rightarrow gg}) \kappa_F^2 + (\mathcal{B}_{H \rightarrow VV}) \kappa_V^2} \\ \mu_{\text{VBF}} &\propto \frac{\kappa_V^4}{(\mathcal{B}_{H \rightarrow f\bar{f}} + \mathcal{B}_{H \rightarrow gg}) \kappa_F^2 + (\mathcal{B}_{H \rightarrow VV}) \kappa_V^2}. \end{aligned} \quad (6.4)$$

2168 Figure 6.6 shows the two-dimensional likelihood scan of κ_F and κ_V . The best-fit values are given in
2169 equation 6.5. The best-fit values are consistent with unity within their uncertainties.

$$\begin{aligned} \kappa_F &= 0.93 & +0.24 & +0.21 & = 0.93 & +0.32 \\ && -0.18 & -0.14 && -0.23 \\ \kappa_V &= 1.04 & +0.07 & +0.07 & = 1.04 & \pm 0.11. \end{aligned} \quad (6.5)$$

(stat.) (syst.)

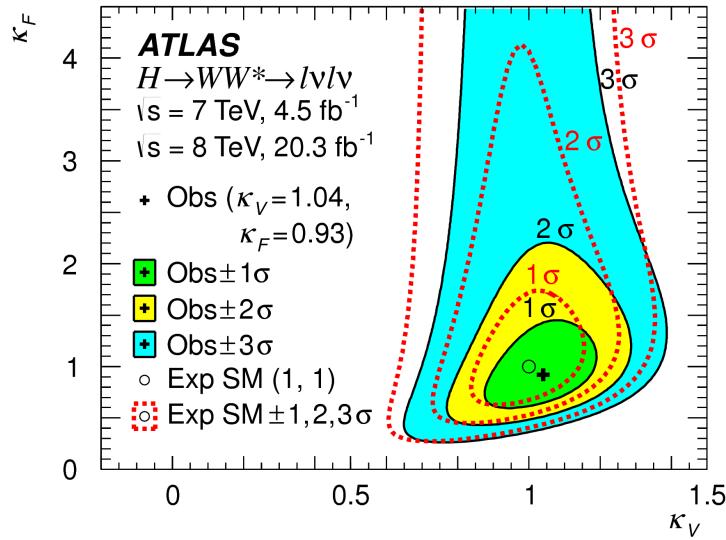


Figure 6.6: Likelihood scan as a function of κ_F and κ_V [21].

2170 **6.5 HIGGS PRODUCTION CROSS SECTION MEASUREMENT**

2171 Another measurement that comes naturally from the signal strength numbers quoted earlier is the pro-
2172 duction cross section and 7 and 8 TeV for both gluon fusion and VBF production. The general equa-
2173 tion for calculating the cross section is given in equation 6.6.

$$\begin{aligned} (\sigma \cdot \mathcal{B}_{H \rightarrow WW^*})_{\text{obs}} &= \frac{(N_{\text{sig}})_{\text{obs}}}{\mathcal{A} \cdot \mathcal{C} \cdot \mathcal{B}_{WW \rightarrow \ell\nu\ell\nu}} \cdot \frac{1}{\int L dt} \\ &= \hat{\mu} \cdot (\sigma \cdot \mathcal{B}_{H \rightarrow WW^*})_{\text{exp}} \end{aligned} \quad (6.6)$$

2174 $(N_{\text{sig}})_{\text{obs}}$ is the number of events observed in data. \mathcal{A} is the geometric and kinematic acceptance of
 2175 the detector, while \mathcal{C} is the efficiency of the signal region selection for events that are reconstructed in the
 2176 detector. The branching ratio of a WW system to leptons must also be divided out. The production
 2177 cross section depends on the center of mass energy and the production mode desired (gluon fusion or
 2178 VBF), and so three separate cross section measurements are quoted in equation 6.7.

$$\begin{aligned}
 \sigma_{\text{ggf}}^{\text{7TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 2.0 \pm 1.7 {}^{+1.2}_{-1.1} = 2.0 {}^{+2.1}_{-2.0} \text{ pb} \\
 \sigma_{\text{ggf}}^{\text{8TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 4.6 \pm 0.9 {}^{+0.8}_{-0.7} = 4.6 {}^{+1.2}_{-1.1} \text{ pb} \\
 \sigma_{\text{VBF}}^{\text{8TeV}} \cdot \mathcal{B}_{H \rightarrow WW^*} &= 0.51 {}^{+0.17}_{-0.15} {}^{+0.13}_{-0.08} = 0.51 {}^{+0.22}_{-0.17} \text{ pb.}
 \end{aligned} \tag{6.7}$$

(stat.) (syst.)

2179 The predicted cross section values for gluon fusion are 3.3 ± 0.4 pb at 7 TeV and 4.2 ± 0.5 pb
 2180 at 8 TeV, consistent with the measured values within their uncertainties. For vector boson fusion, the
 2181 predicted cross section is 0.35 ± 0.02 pb, again consistent with the measured value.

2182 6.6 CONCLUSION

2183 The combined analysis of the gluon fusion and vector boson fusion processes in $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
 2184 in the 7 and 8 TeV datasets has yielded the first discovery level significance for Higgs production in this
 2185 decay channel. Additionally, precise measurements of the couplings to vector bosons and fermions are
 2186 given. Finally, signal strengths and cross sections for each production mode are measured. Figure 6.7
 2187 shows the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ measurements in comparison with other Higgs decay channels in
 2188 ATLAS. The measurement of signal strength from this channel remains the most sensitive in both the
 2189 gluon fusion and VBF production modes for the Run 1 dataset.

ATLAS

Individual analysis

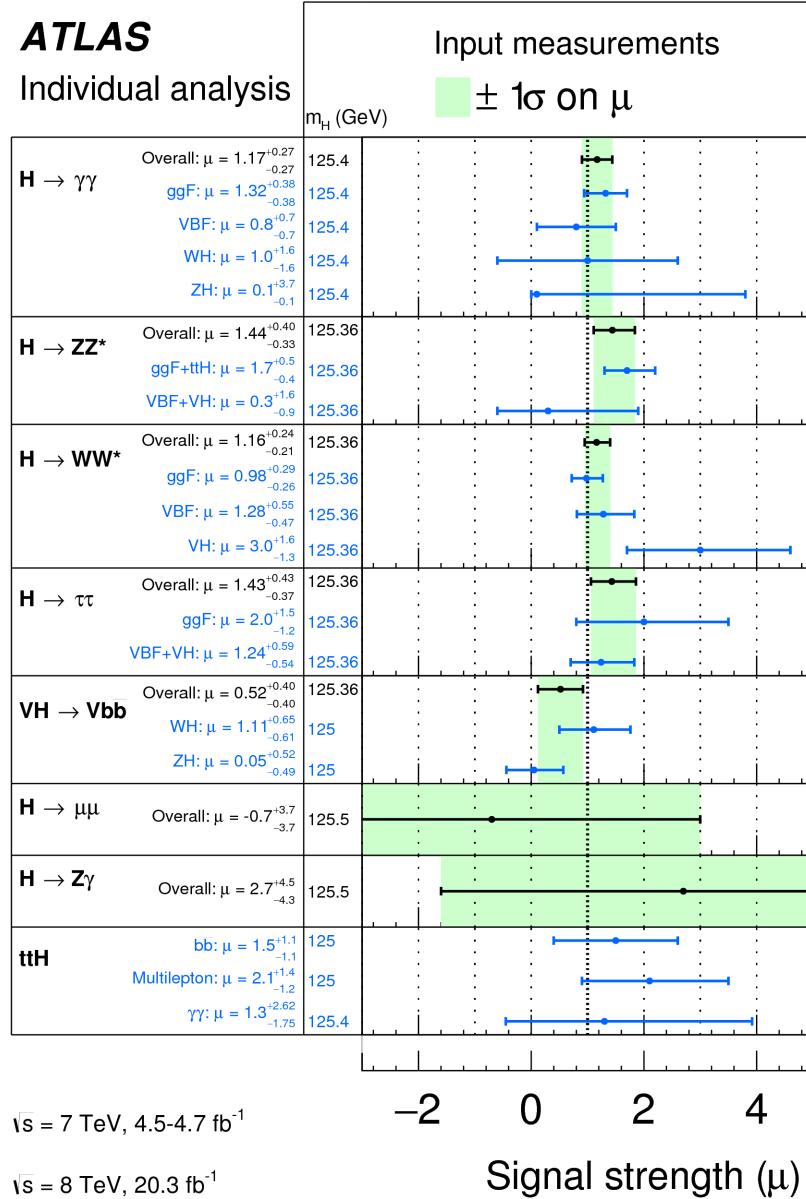


Figure 6.7: Comparison of signal strength measurements in different Higgs decay channels on ATLAS[23].

Part III

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Search for Higgs pair production in the

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$HH \rightarrow b\bar{b}b\bar{b}$ channel in LHC Run 2 at $\sqrt{s} =$

2193

13 TeV

*Passion is in all great searches and is necessary to all
creative endeavors.*

W. Eugene Smith

7

2194

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Search for Higgs pair production in boosted $b\bar{b}b\bar{b}$ final states

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7.1 INTRODUCTION

2198 After the discovery of the Higgs boson in the ATLAS Run 1 dataset and the subsequent measurements
2199 of its properties, the Higgs transformed into a potential tool in searches for physics beyond the Stan-
2200 dard Model. The pair production cross section of the Higgs can be enhanced through BSM physics.
2201 Studying di-Higgs production also probes the Higgs self-coupling, shedding light on the structure
2202 of the Higgs potential. This chapter presents a search for resonant production of a Higgs pair in the
2203 $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ final state in 3.2 fb^{-1} of data collected at $\sqrt{s} = 13 \text{ TeV}$. In particular, this
2204 chapter focuses on a search for this final state in the regime where m_X is large ($\gtrsim 1 \text{ TeV}$) and the Higgs
2205 bosons in the decay are significantly boosted. A tailored selection for this boosted selection, using novel

2206 techniques in jet substructure and b -tagging, is discussed. Then, the data-driven background estimate is
 2207 presented. Finally, the results of the search are shown. The signal models used as benchmarks are a spin-
 2208 2 Randall Sundrum graviton (RSG) and a narrow width spin-0 resonance. These models are described
 2209 in more detail in Chapter 1. Limits on signal models are reserved for the next chapter where the results of
 2210 this chapter are combined with the results of a separate selection dedicated to the lower m_X regime.

2211 7.2 MOTIVATION

2212 With the center of mass energy increase from $\sqrt{s} = 8$ TeV to $\sqrt{s} = 13$ TeV, the LHC and ATLAS
 2213 are able to probe new resonances at higher mass scales than previously accessible in Run 1. This is a
 2214 powerful motivator for searching for a new resonance in the early 13 TeV data. Figure 7.1 shows the
 2215 ratios of parton luminosities between 8 and 13 TeV for different resonance masses. For a resonance of
 2216 $M_X = 2$ TeV, the cross section at $\sqrt{s} = 13$ TeV is roughly a factor of 10 larger than at $\sqrt{s} = 8$ TeV.

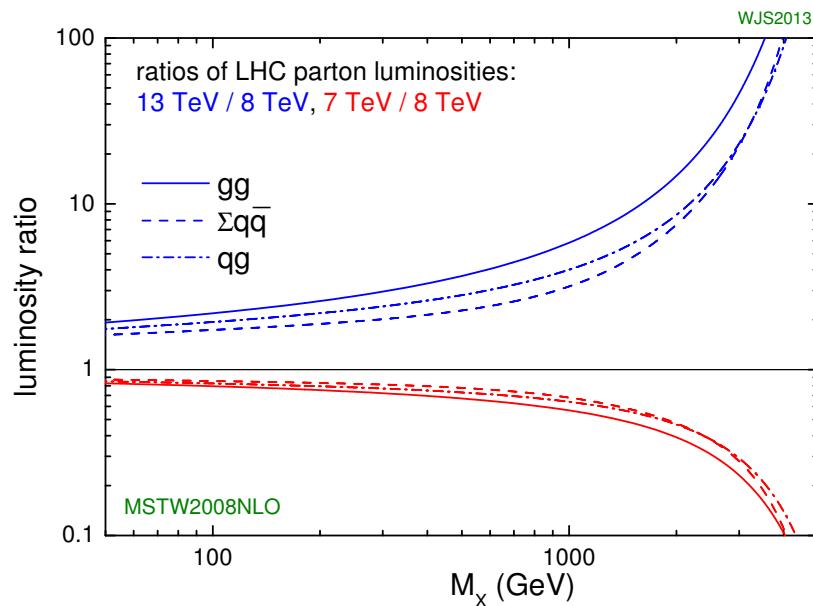


Figure 7.1: Parton luminosity ratios as a function of resonance mass M_X for 13/8 TeV and 7/8 TeV [24].

2217 Higgs pair production offers a vast array of unprobed regions of phase space where searches for BSM
 2218 physics can be made. Chapter 1 discusses some possibilities for both resonant and non-resonant enhance-

ment of the di-Higgs production cross section. Given the increased mass reach of the LHC in Run 2, it is particularly important to focus on resonant searches at high m_X . One consideration when conducting a search in the HH final state is which decay modes of the Higgs to consider. Figure 7.2 shows the branching ratio of the HH final state for different combinations of decays of each individual Higgs. As the largest branching ratio for the 125 GeV Higgs is $H \rightarrow b\bar{b}$, the $HH \rightarrow b\bar{b}b\bar{b}$ branching ratio is also the largest at 33%.

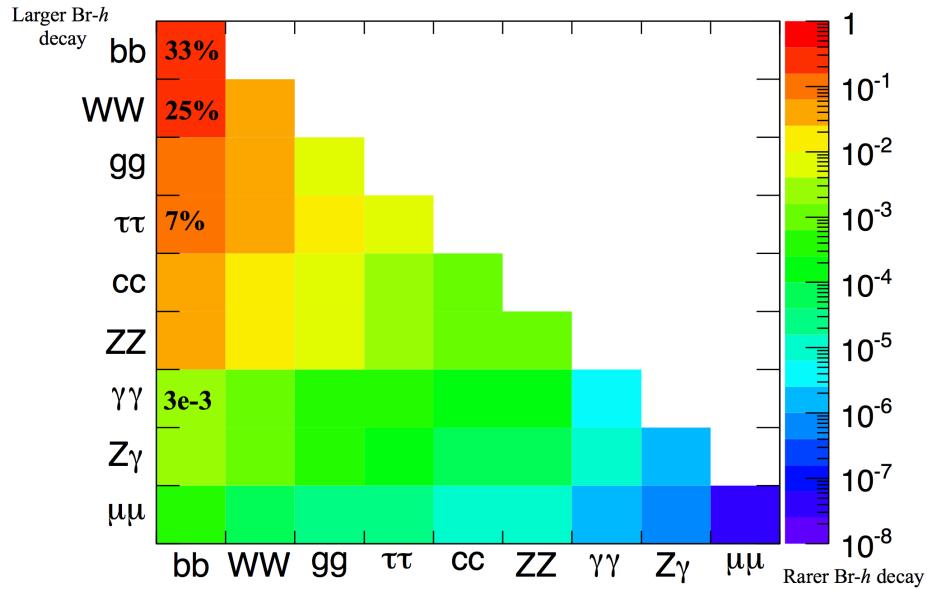


Figure 7.2: Summary of HH branching ratios [25].

At high m_X , the Higgs bosons resulting from the decay of a heavy resonance will have large p_T^* . The ΔR between the decay products of the Higgs is inversely proportional to the Higgs p_T , as shown in equation ??.

$$\Delta R \approx \frac{2m}{p_T} \quad (7.1)$$

Figure 7.3 shows the minimum ΔR between truth level B decay vertices in simulation samples for Randall-Sundrum gravitons of different masses. The figure shows that as the mass of the graviton increases, the ΔR distribution between the b quarks in the Higgs decay tends to shift to lower values. Be-

*In the limit that $m_H \ell \ell m_X$, the Higgs p_T is roughly $m_X/2$.

cause of this effect, it is necessary to tailor a selection to target these merged b -jets.

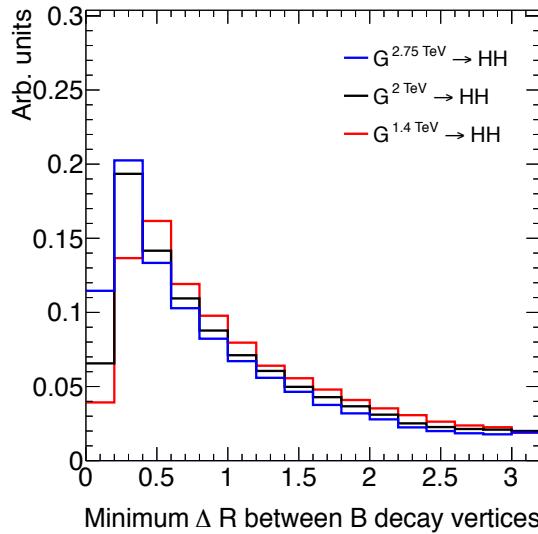


Figure 7.3: Minimum ΔR between B decay vertices for different RSG masses in a $G_{\text{KK}}^* \rightarrow HH \rightarrow 4b$ sample with $c = 1$

7.3 DATA AND SIMULATION SAMPLES

7.3.1 SIGNAL MODELS

While the resonance search is by its nature generic (as it is a simple search for a peak in the $4b$ invariant mass spectrum), there are two signal models that the selection requirements have been optimized for. The first is Randall-Sundrum (RSG) model, where a tower of massive spin-2 Kaluza-Klein gravitons is predicted. The second is a heavy narrow spin-0 resonance, the so-called “heavy Higgs”. This type of resonance arises, for example, in the two Higgs doublet model (2HDM). More details about the physics of these models and their motivation is given in chapter 1.

Signal graviton (G_{KK}^*) events are generated at leading order (LO) with **MADGRAPH5 v2.2.2** [92].

The PDF set used is the **NNPDF2.3 LO** set [93]. For modeling parton shower and hadronization in jets, **PYTHIA 8.186** is used with the A14 tune [80, 94]. The free parameters in the RSG model are the graviton mass and the coupling constant $c \equiv k/\bar{M}_{\text{Pl}}$ [†]. Both the production cross section and width of the

[†] k is the curvature constant for the warped extra dimension and \bar{M}_{Pl} is the Planck mass divided by 8π

graviton are proportional to c^2 . Samples are generated at both $c = 1$ and $c = 2$ for a variety of mass points between 300 GeV and 3 TeV.

The second signal sample is a heavy spin-0 resonance H with a fixed width of $\Gamma_H = 1$ GeV. This is generated with **MADGRAPH5** and uses the **CT10** PDF set [83]. The parton shower and hadronization are handled by **HERWIG ++** with the **CTEQ6L1** PDF set and the **UEEE5** event tune [84, 95, 96]. Because the width and branching ratios depend on 2HDM parameters, each mass point generated with this fixed width corresponds to a different point in the 2HDM parameter phase space. Mass points are generated between 300 GeV and 1 TeV as with the RSG signal samples.

7.3.2 BACKGROUND SAMPLES

While the dominant **QCD** multijet background is estimated with a fully data-driven method, the sub-dominant backgrounds $t\bar{t}$ and $Z + \text{jets}$ are modeled with some input from simulation.

$t\bar{t}$ events are simulated at next-to-leading order (NLO) with the **POWHEG-BOX** version 1 generator using the **CT10** PDF set [97]. The parton shower, hadronization, and underlying event are simulated with **PYTHIA 6.428** with the **CTEQ6L1** PDF set [79]. The Perugia 2012 tune is used [98]. NNLO **QCD** corrections to the cross sections are computed in **Top++ 2.0** [99]. The top quark mass is set to 172.5 GeV. The shapes of distributions in $t\bar{t}$ are taken from MC while the normalization is taken from data.

Finally, the $Z + \text{jets}$ background is simulated with **PYTHIA 8.186** and the **NNPDF2.3** LO PDF set. This background is negligible compared to the others and is taken fully from MC.

7.3.3 DATA SAMPLE AND TRIGGER

This analysis is done on 3.2 fb^{-1} of data taken in 2015 at $\sqrt{s} = 13$ TeV. The details of the machine conditions during this time can be found in Chapter 2. Only data which was taken during stable beam conditions with all detectors functioning is used. Events must pass a trigger which requires a single 360 GeV large radius ($R = 1.0$) jet to be reconstructed in the HLT. Figure 7.4 shows the trigger efficiency for various trigger options as a function of graviton mass. Above $m_G > 1$ TeV, the single large radius jet trigger is 99% efficient for events passing the signal selection.

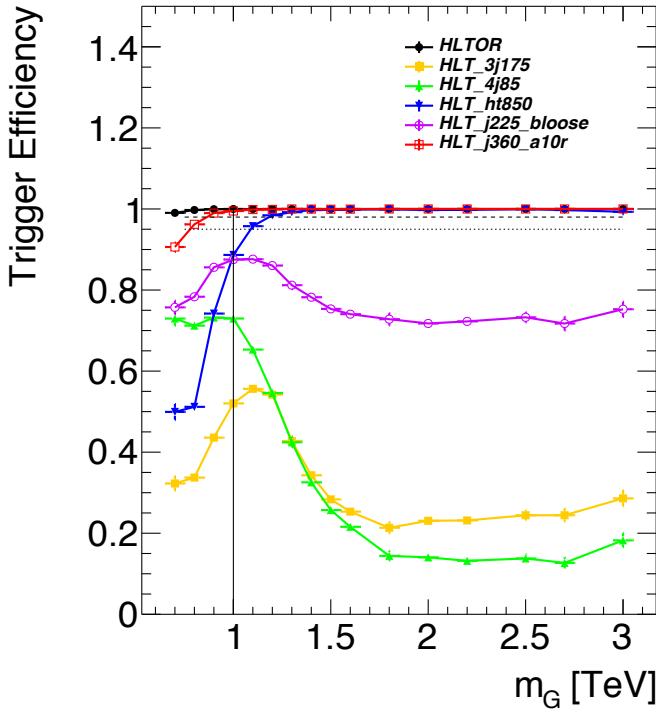


Figure 7.4: Trigger efficiency for events passing all signal region selections as a function of mass in $G_{\text{KK}}^* \rightarrow HH \rightarrow 4b$ samples with $c = 1$ [26]. In the trigger names, “j” refers to a jet or jets. “ht” refers to H_T , the scalar sum of transverse momenta in the event. “bloose” refers to a loose b -tagging requirement applied to the jet. “a10r” refers to anti- k_T jets with $R = 1.0$. The numbers at the end are the thresholds on the given quantity in GeV.

2269 7.4 EVENT RECONSTRUCTION AND OBJECT SELECTION

2270 The boosted selection first begins by defining a unique set of objects that can be exploited to increase
 2271 signal efficiency in the kinematic regime where the final state b -jets are very merged.

2272 7.4.1 LARGE RADIUS ($R = 1.0$) JETS

2273 The first step towards reconstructing the final state is to define objects that can be used to measure the
 2274 kinematics of the Higgs bosons. In the boosted selection anti- k_T jets with a radius parameter of 1.0
 2275 are used. These jets are much larger in angular size than the typical $R = 0.4$ jets and are intended to
 2276 encompass both jets resulting from the Higgs decay[‡]. The jets are built from clusters in the calorimeter

[‡]This is in contrast to the resolved selection, which uses two $R = 0.4$ anti- k_T jets for each Higgs

2277 calibrated with local calibration weighting [17].

2278 Because of the large extent of these jets, great care must be taken to remove potential contributions of
2279 calorimeter clusters from pile-up. This is done using a technique called jet trimming [100]. With trim-
2280 ming, the constituents of the large radius jet are re-clustered with a smaller radius with the k_T algorithm.
2281 Then, these so-called subjets are removed from the larger jet if $p_T^{\text{subjet}}/p_T^{\text{jet}} < f_{\text{cut}}$. In this analysis, the
2282 subjet radius is $R = 0.2$ and $f_{\text{cut}} = 0.05$. Trimming has been shown to improve the mass resolution
2283 of large radius jets. Figure 7.5 shows the effect of trimming on the large radius jet mass (M_J). Because
2284 the large radius jet fully contains the higgs decay products, its invariant mass should correspond to the
2285 125 GeV mass of the Higgs. The trimming algorithm brings the jet mass much closer to the expected
2286 Higgs mass and improves the mass resolution.

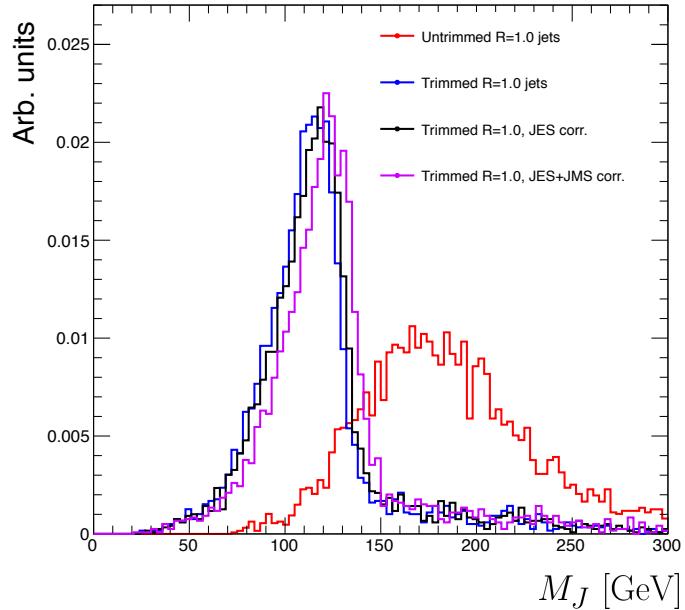


Figure 7.5: Comparison of untrimmed and trimmed jet masses for large radius jets in a RSG sample with $m_{G_{\text{KK}}^*} = 1 \text{ TeV}$. JES (JMS) refers to the standard jet energy (mass) scale calibration for ATLAS [17].

2287 The large radius jets are required to satisfy $250 < p_T < 1500 \text{ GeV}$. They must also be within
2288 $|\eta| < 2.0$ in order to ensure that the full jet is within the inner detector tracking volume. Finally, they
2289 are required to have $M_J > 50 \text{ GeV}$. The upper p_T cut and lower threshold on mass are applied to
2290 correspond to the kinematic range where uncertainties are available in ATLAS calibrations [101, 102].

2291 7.4.2 TRACK JETS AND b -TAGGING

2292 Because the b -jets from boosted Higgs decays are so close together (as illustrated in figure 7.3), narrow ra-
2293 dius jets are required to fully resolve both b -jets. The minimum radius feasible for jets based on calorime-
2294 ter deposits is determined by the calorimeter granularity. However, because b -tagging relies on informa-
2295 tion from the inner detector, it is possible to define another type of jet that can have a smaller radius and
2296 better b -tagging resolution. These jets are called “track jets” [27, 102].

2297 Track jets are formed by applying the usual anti- k_T clustering algorithm to tracks that are required
2298 to be consistent with the primary vertex. After the jet axis has been determined using these tracks, a sec-
2299 ond step of track association is also performed to add tracks that can be useful for b -tagging [27]. In this
2300 analysis, the tracks are clustered with a radius parameter of $R = 0.2$. This radius has been shown to
2301 give good performance in boosted Higgs tagging [27, 102]. Figure 7.6 shows a comparison among dif-
2302 ferent track jet radii of the efficiency for reconstructing two b -jets from each Higgs in a RSG sample as a
2303 function of mass. Track jets with radius of 0.2 give the best performance, especially at high mass.

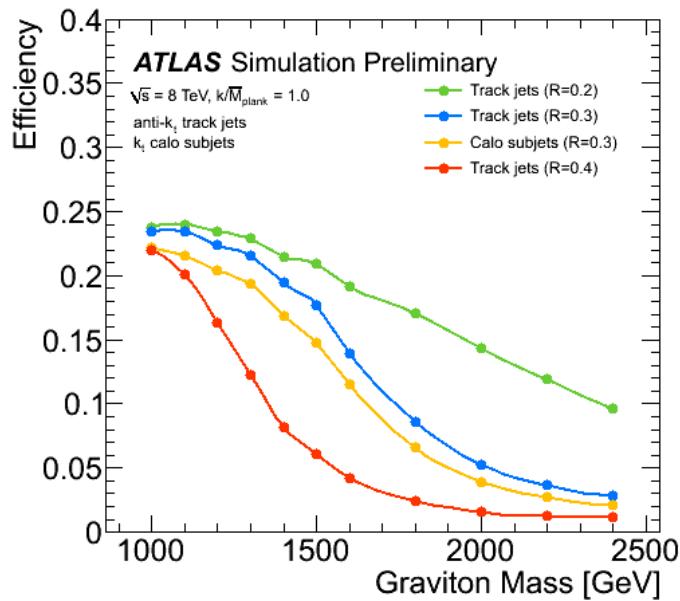


Figure 7.6: Efficiency of finding two b -jets from each Higgs in an RSG event using calorimeter jets with $R = 0.3$ or different track jet radii [27]

2304 In this analysis, track jets are required to have $p_T > 10$ GeV and $|\eta| < 2.5$. They must also have at

2305 least two tracks.

2306 **7.4.3 MUONS**

2307 Muons are used in this study to correct the four-momenta of calorimeter jets by accounting for semi-
2308 leptonic b decays. The muons used are combined ID and MS muons which must satisfy tight identifica-
2309 tion requirements [16]. The muons must have $p_T > 4 \text{ GeV}$ and $|\eta| < 2.5$. Table 7.1 summarizes the
2310 object requirements described in this section.

	R	p_T	$ \eta $	M
Calorimeter jets	1.0	$250 < p_T < 1500 \text{ GeV}$	< 2.0	$> 50 \text{ GeV}$
Track jets	0.2	$> 10 \text{ GeV}$	< 2.5	-
Muons	-	4 GeV	< 2.5	-

Table 7.1: Summary of requirements on objects used in the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search

2311 **7.5 EVENT SELECTION**

2312 The first requirement in the boosted selection is for ≥ 2 large radius jets satisfying the selections outlined
2313 above. The two highest momentum large-R jets in the event are referred to as “Higgs candidates”. The
2314 leading jet is required to have $p_T > 350 \text{ GeV}$.

2315 Track jets satisfying the object selections are matched to Higgs candidate jets via ghost association [103].

2316 Each Higgs candidate must have at least 2 track jets associated with it. These basic requirements are illus-
2317 trated in figure 7.7

2318 The QCD multijet background produces less central jets than high mass resonances, so there is an ad-
2319 ditional requirement that the two Higgs candidates be close together in η . The large-R jets are required
2320 to satisfy $|\Delta\eta(JJ)| < 1.7$.

2321 The final set of requirements ensures that the Higgs candidates are consistent with expected proper-
2322 ties of the 125.0 GeV Higgs. First, a variable (X_{hh}) is defined to measure the consistency of both of the

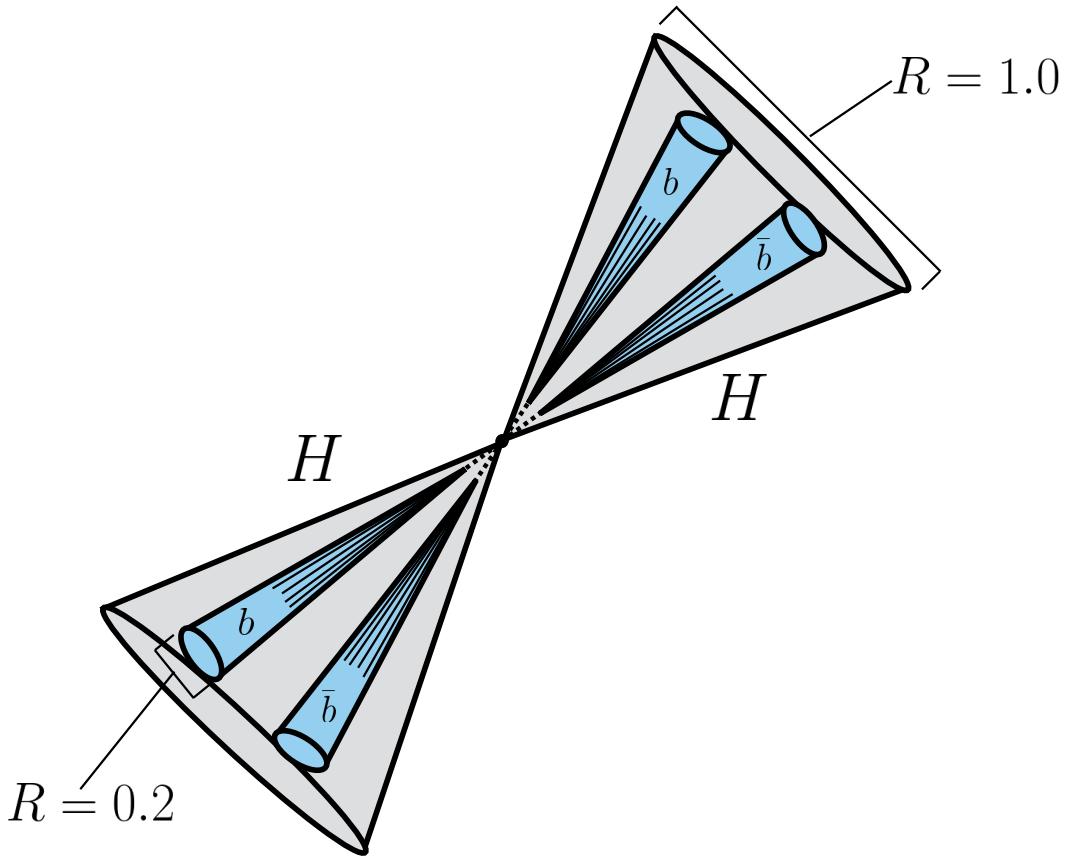


Figure 7.7: Illustration of the boosted selection requirements on Higgs candidates. Each large-radius calorimeter jet (Higgs candidate) must contain two track jets

2323 Higgs candidate jets with the SM Higgs mass. This is shown in equation 7.2.

$$X_{hh} = \sqrt{\left(\frac{M_J^{\text{lead}} - 124 \text{ GeV}}{0.1 M_J^{\text{lead}}}\right)^2 + \left(\frac{M_J^{\text{sublead}} - 115 \text{ GeV}}{0.1 M_J^{\text{sublead}}}\right)^2} \quad (7.2)$$

2324 The mass values in the X_{hh} formula are optimized to maximize signal efficiency. The sub-leading jet
2325 typically has a lower mass due to semi-leptonic b decays and final state radiation. X_{hh} effectively acts as
2326 a χ^2 measurement of the consistency of the two Higgs candidate masses with the signal hypothesis. The
2327 denominators of each term ($0.1 M$) give the uncertainty on the mass measurement for the large radius
2328 jets. Events are required to satisfy $X_{hh} < 1.6$.

2329 The last requirement applied is on the number of b -tagged track jets. There are two signal regions de-

fined. The first requires exactly four b -tagged track jets, two in each Higgs candidate (known as the 4 signal region). At high resonance masses, this requirement is inefficient, so an additional signal region requiring only three b -tagged track jets is also defined (known as the 3 b signal region). While this has a larger background it is also more efficient for high resonance masses. For both signal regions, threshold on MV2 score is chosen such that the algorithm is 77% efficiency in finding true b -jets[§]. Different working points were tested and this was found to be optimal. Appendix A has more details on this optimization.

Before making the requirement on X_{hh} , the masses of the Higgs candidates are corrected for semi-leptonic b decays using muons with the criteria outlined in the previous section. Any muons within a $\Delta R < 0.2$ of a b -tagged track jet have their four-momenta added to the four-momentum of the Higgs candidate. This correction does not affect the pre-selection requirements but does affect the X_{hh} requirement and the final invariant mass distribution used.

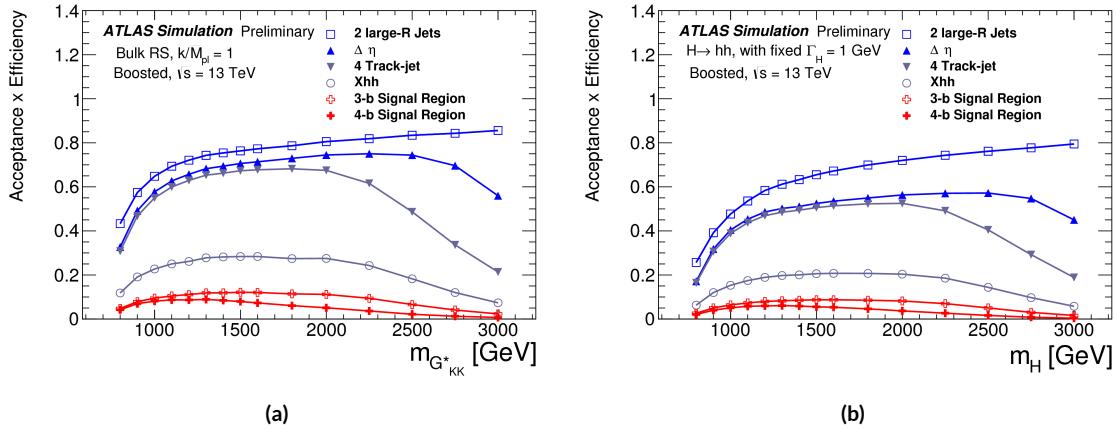


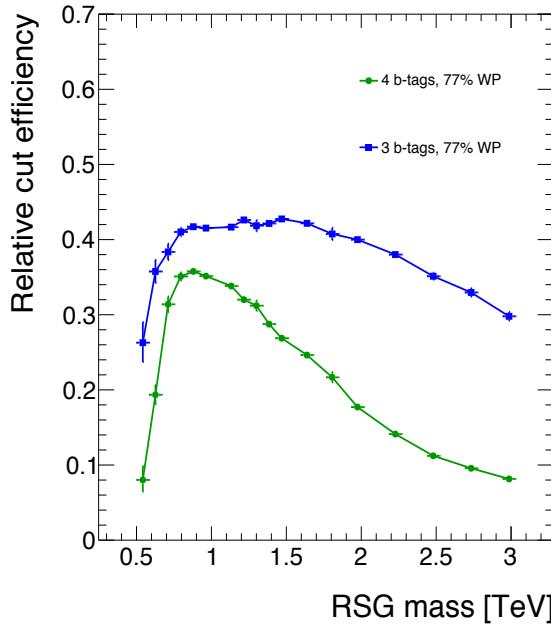
Figure 7.8: Acceptance \times efficiency as a function of mass for (a) RSG and (b) narrow heavy scalar signal models [28].

Figure 7.8 shows the product of acceptance and efficiency as a function of mass for both the RSG and narrow heavy scalar resonance signal models. After $m_X > 1$ TeV, the efficiency of the 4 b requirement begins to decline. After $m_X > 2$ TeV, the efficiency of requiring two track jets in each Higgs candidate begins to decline as well. Both of these behaviors illustrate the difficulty of resolving the merged decay

[§]The specific MV2 algorithm chosen is MV2c20, where the fraction of charm events used in the training is 20%

2346 products at high mass. More details on the degradation of the b -tagging efficiency at high masses are
2347 shown in appendix B.

2348 Figure 7.9 shows a more detailed comparison of the signal efficiency in the $3b$ vs $4b$ signal regions for
2349 the RSG model. The efficiencies shown here are relative to all prior selection requirements.



2350 **Figure 7.9:** Efficiency of requiring 3 or 4 b -tagged track jets vs. RSG mass. The efficiency quoted is relative to the previous
2351 selection requirements (rather than an absolute efficiency).

2352 The final discriminating variable used in the boosted analysis is M_{2J} , the invariant mass of the two
2353 Higgs candidates. In order to improve the mass resolution, the four-momenta of each Higgs candidate
2354 are scaled by m_h/M_J . The effect of this correction is small in the boosted analysis but is done for consistency
2355 with the resolved selection.

2356 Table 7.2 shows the effect of the selection requirements on signal and background simulations as well
2357 as data.

2358 7.6 DATA-DRIVEN BACKGROUND ESTIMATION

2359 The largest background to this final state is QCD multijet production, constituting 80-90% of the total
2360 background. Because of the difficulties in modeling higher order QCD processes, this background is

Selection	Data	$m_{G_{KK}^*} = 1\text{TeV}$	$m_{G_{KK}^*} = 2\text{TeV}$	$t\bar{t}$	$Z+\text{jets}$
N(fiducial large-R jets) ≥ 2	2202396	23.3	0.48	32345.2	4255.7
leading large-R jet $p_T > 350\text{ GeV}$	1873741	22.9	0.48	26511.7	3649.9
Both large-R jet $m > 50\text{ GeV}$	1854625	21.2	0.47	24369.8	3575.8
Both large-R jet $p_T < 1500\text{ GeV}$	1853601	21.2	0.46	24346.5	3572.9
$ \Delta\eta(JJ) < 1.7$	1435273	20.8	0.44	20751.0	3265.8
≥ 2 track-jets per large-R jet	1224727	19.8	0.40	18234.5	2692.6
$3 b\text{-tags}, X_{hh} < 1.6$	316	3.4	0.067	46.7	2.0
$4 b\text{-tags}, X_{hh} < 1.6$	20	2.9	0.030	1.4	0.0

Table 7.2: Effect of boosted selection on data, RSG signal models, $t\bar{t}$, and $Z+\text{jets}$. The numbers from simulation are normalized with the MC generator cross section and do not take into account the data driven estimates described in section 7.6 [33].

estimated with a fully data-driven method. The only other non-negligible background is $t\bar{t}$, constituting the other 10-20%. Due to the presence of $t\bar{t}$ in the sideband region where the QCD background will be estimated, the normalization of the QCD and $t\bar{t}$ backgrounds are simultaneously estimated.

7.6.1 MASS REGION DEFINITIONS

The first step in the data-driven background estimate is to define a sideband mass region where the background normalization can be derived. Additionally, a control region is defined where the background estimate can be validated. The control (CR) and sideband (SB) regions are defined using a radial distance in the two-dimensional large-R jet mass plane, R_{hh} , which is defined in equation 7.3.

$$R_{hh} = \sqrt{(M_J^{\text{lead}} - 124\text{ GeV})^2 + (M_J^{\text{sublead}} - 115\text{ GeV})^2} \quad (7.3)$$

Events in the sideband region are required to fail the signal region $X_{hh} < 1.6$ requirement and have $R_{hh} > 35.8\text{ GeV}$. The control region consists of those events which are not in the signal or sideband regions. Figure 7.10 shows the definition of the signal, control, and sideband mass regions.

Table 7.3 summarizes the mass region selections for the three different regions used in the analysis.

*The $Z+\text{jets}$ background is a sub-percent level contribution

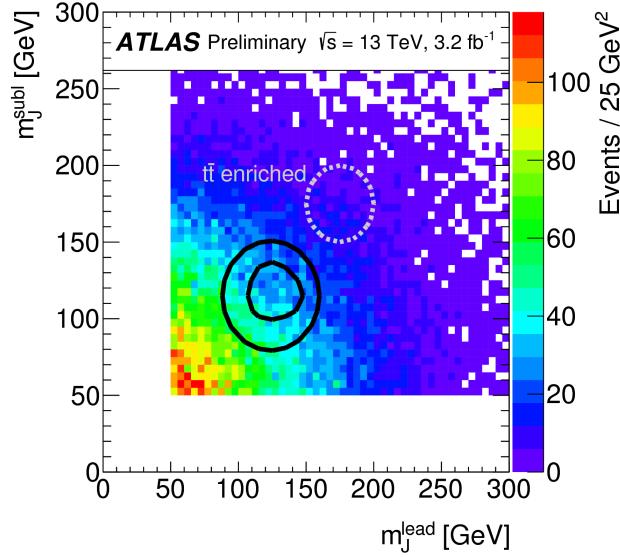


Figure 7.10: M_J^{sublead} vs. M_J^{lead} in a 2 b -tag data sample. The signal region is defined by the inner black contour ($X_{hh} < 1.6$) and the sideband region is defined by the outer contour ($R_{hh} > 35.8 \text{ GeV}$). The region between the black contours is the control region. The mass region which is enriched in $t\bar{t}$ background is also shown for illustration. [28]

Region	Requirement	Notes
Signal Region (SR)	$X_{hh} < 1.6$	-
Control Region (CR)	$R_{hh} < 35.8 \text{ GeV}$ and $X_{hh} > 1.6$	Used for validation of background estimates
Sideband Region (SB)	$R_{hh} > 35.8 \text{ GeV}$	Used to derive background normalization

Table 7.3: Mass region definitions used for background estimation

7.6.2 BACKGROUND ESTIMATION

The method for estimating the background in this analysis is similar to the ABCD method presented in Chapter 5. In this case, the two handles used to define different regions for the estimate are the number of b -tagged track jets and the mass regions. A region requiring exactly two b -tagged track jets in one large-R jet (referred to as the 2-tag or 2 b region) is defined for use in the background estimate. The number of expected background events in the 3 b and 4 b signal regions is then given by equation 7.4.

$$N_{\text{bkg}}^{3(4)-\text{tag},\text{SR}} = \mu_{\text{Multijet}} N_{\text{Multijet}}^{2-\text{tag},\text{SR}} + \beta_{t\bar{t}} N_{t\bar{t}}^{3(4)-\text{tag},\text{SR}} + N_{Z+\text{jets}}^{3(4)-\text{tag},\text{SR}} \quad (7.4)$$

2377 In this equation, $N_{\text{bkg}}^{3(4)\text{-tag}}$ is the expected number of background events in the $3b$ or $4b$ signal re-
 2378 gions. $N_{\text{Multijet}}^{2\text{-tag}}$ is the number of multijet events in the 2-tag region. $N_{t\bar{t}}^{3(4)\text{-tag}}$ is the number of $t\bar{t}$
 2379 events predicted in the MC for the $3b$ or $4b$ signal region, and the variable is similarly defined for the
 2380 $Z+\text{jets}$ background. The $\beta_{t\bar{t}}$ parameter is a scale factor used to correct the normalization of the $t\bar{t}$ esti-
 2381 mate in the signal region. μ_{Multijet} is an extrapolation factor that is derived in the sideband region and
 2382 used to estimate the ratio of 2-tag events to 3(4)-tag events in the signal region. It is defined in equa-
 2383 tion 7.5.

$$\mu_{\text{Multijet}} = \frac{N_{\text{Multijet}}^{3(4)\text{-tag,SB}}}{N_{\text{Multijet}}^{2\text{-tag,SB}}} = \frac{N_{\text{data}}^{3(4)\text{-tag,SB}} - \beta_{t\bar{t}} N_{t\bar{t}}^{3(4)\text{-tag,SB}} - N_{Z+\text{jets}}^{3(4)\text{-tag,SB}}}{N_{\text{data}}^{2\text{-tag,SB}} - \beta_{t\bar{t}} N_{t\bar{t}}^{2\text{-tag,SB}} - N_{Z+\text{jets}}^{2\text{-tag,SB}}} \quad (7.5)$$

2384 The $t\bar{t}$ scale factor ($\beta_{t\bar{t}}$) and the QCD multijet extrapolation factor (μ_{Multijet}) are estimated together
 2385 in a simultaneous fit in the sideband region. Then, the number of events in the 2-tag signal region is
 2386 used, along with the $t\bar{t}$ estimate in the $3b$ and $4b$ signal regions and μ_{Multijet} , to estimate the total num-
 2387 ber of background events in the two final signal regions. The shape of the final discriminant M_{2J} is also
 2388 taken from the 2-tag signal region where there are more statistics. This method is illustrated graphically
 2389 in figure 7.11.

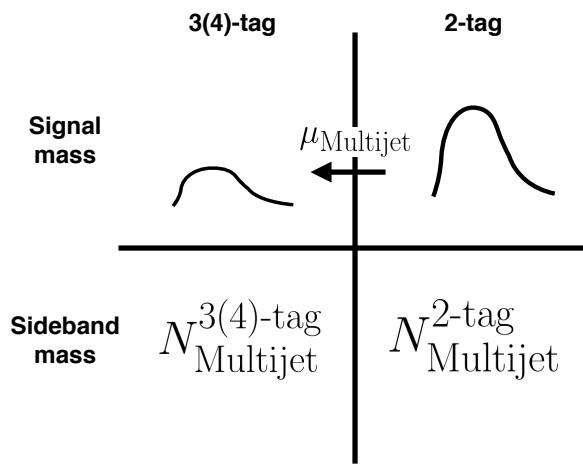
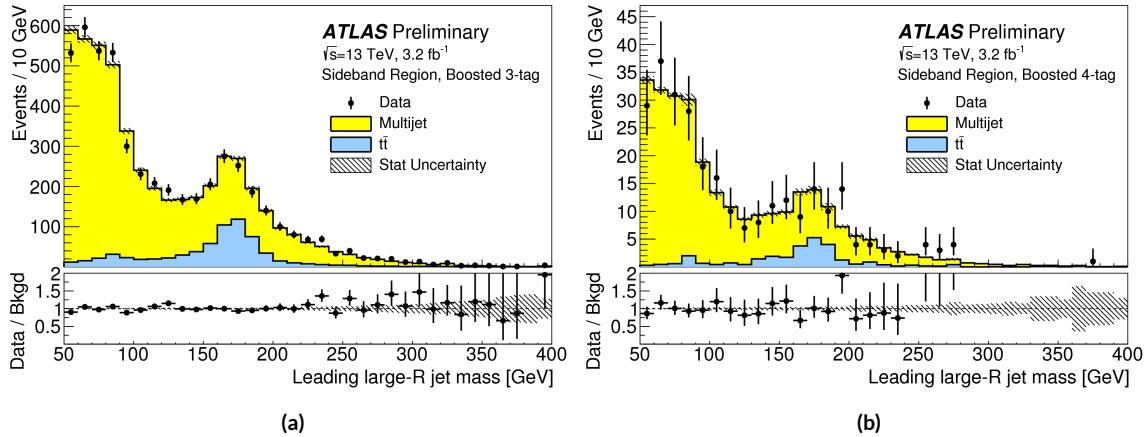


Figure 7.11: An illustration of the data-driven background estimation technique for the boosted analysis

2390 In the $3b$ region, the fit yields values of $\mu_{\text{Multijet}} = 0.160 \pm 0.03$ and $\beta_{t\bar{t}} = 1.02 \pm 0.09$. In the $4b$

2391 region, the fit gives $\mu_{\text{Multijet}} = 0.0091 \pm 0.0007$ and $\beta_{t\bar{t}} = 0.82 \pm 0.39$. The uncertainties quoted
 2392 are statistical only. The larger uncertainties in the $4b$ values indicate the lower statistics available in that
 2393 region.

2394 Figure 7.12 shows the distributions of data and background estimates in the $3b$ and $4b$ sideband re-
 2395 gions after the background fit has been done. The normalizations are constrained from the fit to match
 2396 that of the data, but good modeling of the shape of the mass of the leading large-R jet is seen as well. The
 2397 shapes of the kinematic distributions in the $4b$ region are taken from the $3b$ region due to the better MC
 2398 statistics in that region.



2399 **Figure 7.12:** Leading large-R jet mass in the $3b$ (a) and $4b$ (b) sideband regions. The multijet and $t\bar{t}$ backgrounds are
 2400 estimated using the data-driven methods described above. Because their normalizations are derived in the sideband region,
 2401 the total background normalization is constrained by default to match the normalization of the data [28].

2402 7.6.3 BACKGROUND SHAPE FIT

2403 As mentioned in the previous section, the background shape in the 3-tag and 4-tag signal regions is taken
 2404 from the 2-tag signal mass region. Due to the limited statistics available, the background shapes are ad-
 2405 ditionally smoothed after being extrapolated to the 3-tag and 4-tag signal regions. Only the data in the
 2406 range $900 < M_{2J} < 2000$ GeV is included in the fit due to the limited statistics available above 2 TeV.
 2407 Both the $t\bar{t}$ and QCD multijet background are independently fit with an exponential shape, $y = e^{ax+b}$.
 2408 Other shapes are considered and used for the systematic uncertainties. Table 7.4 shows the fit values for

the parameters. Because both the $3b$ and $4b$ QCD shapes come from the 2-tag region, the slopes derived are very similar.

	a	b
QCD ($4b$)	0.00545 ± 0.00021	5.44 ± 0.24
$t\bar{t}$ ($4b$)	0.00746 ± 0.00021	4.88 ± 0.36
QCD ($3b$)	0.00545 ± 0.00021	8.30 ± 0.24
$t\bar{t}$ ($3b$)	0.00746 ± 0.00021	8.58 ± 0.36

Table 7.4: Parameters derived for exponential fit to background M_{2J} shape in the $3b$ and $4b$ signal regions [33]

7.6.4 VALIDATION OF BACKGROUND ESTIMATE

The background estimate can be validated by using the method to estimate the number of events in the control mass region rather than the signal mass region. Figure 7.13 shows the M_{2J} distribution in the $3b$ and $4b$ control regions, comparing data and background estimates. In both cases, both the background shape and normalization are consistent with the data, indicating good agreement. The ratio of data to the background estimates is also fit to a line in the figure to test for any shape difference. The slope of the line is within 1σ (from the fit uncertainties) of flat, further indicating that the data is consistent with the background estimate in the control region.

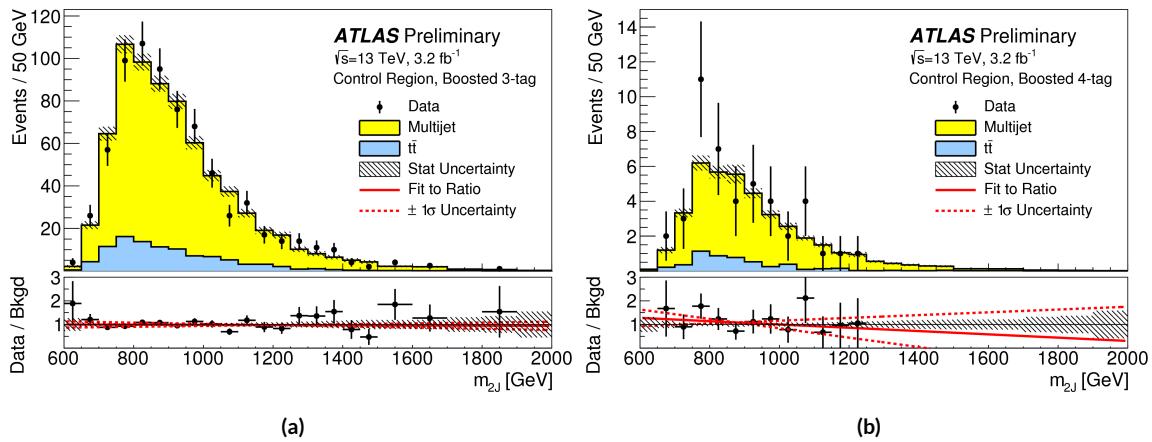


Figure 7.13: Di-jet invariant mass (M_{2J}) in the $3b$ (a) and $4b$ (b) control regions. The multijet and $t\bar{t}$ backgrounds are estimated using the data-driven methods described above [28].

Table 7.5 shows the yields in data and background estimates in the 3-tag and 4-tag sideband and con-

²⁴¹⁷ trol regions. Again, here, it can be seen that the total number of predicted background events from the
²⁴¹⁸ data driven method is consistent with the number of data events in the region.

Sample (3-tag)	Sideband Region	Control Region
Multijet	4328 ± 27	607 ± 10
$t\bar{t}$	683.5 ± 8.1	99.6 ± 3.1
$Z+jets$	31.8 ± 3.7	7.7 ± 1.8
Total	5043 ± 28	715 ± 11
Data	5043	724
Sample (4-tag)	Sideband Region	Control Region
Multijet	247.4 ± 1.5	34.7 ± 0.6
$t\bar{t}$	28.4 ± 1.5	5.1 ± 0.7
$Z+jets$	3.4 ± 1.2	0.6 ± 0.5
Total	279.2 ± 2.5	40.3 ± 1.0
Data	279	45

Table 7.5: The number of events in data and predicted background events in the boosted 3-tag and 4-tag sideband and control regions. The uncertainties shown are statistical only.^[28]

²⁴¹⁹ 7.7 SYSTEMATIC UNCERTAINTIES

²⁴²⁰ The systematic uncertainties in this analysis can be divided into two broad categories. The first type is
²⁴²¹ uncertainties associated with the modeling of the signal processes. The second type of uncertainty is
²⁴²² associated with both the shape and normalization of the background prediction.

²⁴²³ 7.7.1 SIGNAL MODELING UNCERTAINTIES

²⁴²⁴ The signal modeling uncertainty has three main components: theoretical uncertainty on the acceptance,
²⁴²⁵ experimental uncertainties on the large-R jets, and experimental uncertainties on the track jets related to
²⁴²⁶ b -tagging. In this analysis the experimental uncertainties are the most significant.

²⁴²⁷ The first uncertainty on signal modeling is the theoretical uncertainty on the acceptance. As explained
²⁴²⁸ in section 5.6.1, there are four components to this uncertainty. The first is related to missing higher order
²⁴²⁹ terms from the matrix element calculations which is estimated by varying the QCD renormalization and

2430 factorization scales. The second is uncertainty due to the PDF set used. The third is a generator uncer-
2431 tainty which is estimated by modifying the generator used to model the underlying event and hadroniza-
2432 tion. Finally, there is an uncertainty associated with the modeling of the initial state and final state radia-
2433 tion (ISR/FSR). The total theoretical uncertainty on the signal yield is 3%, and this is dominated by the
2434 ISR/FSR modeling.

2435 There are uncertainties on the large-R jets in both the jet energy scale (JES) and jet energy resolution
2436 (JER) as well as the jet mass scale (JMS) and jet mass resolution (JMR). These are evaluated using $\sqrt{s} =$
2437 8 TeV data from Run 1 of ATLAS and extrapolated to the Run 2 beam and detector conditions using
2438 MC[¶]. The details of these uncertainties can be found in reference [104].

2439 Uncertainties on the track jets are related to the b -tagging efficiency. The total uncertainty on the sig-
2440 nal yield due to b -tagging is evaluated by propagating variations of the b -tagging efficiency through the
2441 boosted selection requirements. The uncertainties are calculated jet-by-jet and parameterized as a func-
2442 tion of b -jet p_T and η [105]. For high p_T b -jets (with $p_T > 300$ GeV), the uncertainties are extrapolated
2443 using MC simulation from the lower p_T b -jets [106].

2444 Table 7.6 shows the systematic uncertainties on the signal normalization for models with $m_{G_{KK}^*} =$
2445 1.5 TeV and both $c = 1$ and $c = 2$ as well as a narrow width heavy scalar. The dominant uncertainty
2446 comes from b -tagging and this uncertainty is larger in the 4-tag region than the 3-tag region.

2447 7.7.2 BACKGROUND UNCERTAINTIES

2448 Uncertainties on the QCD multijet background normalization and shape are estimated using the control
2449 mass region. As shown previously, the background predictions in the control region match with the
2450 data yields within the statistical uncertainty in both the 3-tag and 4-tag control regions. As an additional
2451 protection, the statistical uncertainty on the background prediction in the control region is assigned as a
2452 systematic uncertainty on the normalization of the QCD background.

2453 Additional robustness tests are done by varying the definition of the control mass region and the b -
2454 tagging requirements used to define the 2-tag sample. In all cases, the effect of the variations is found to

[¶]The uncertainties are correspondingly larger due to the uncertainty of this extrapolation.

Source	Background		G_{KK}^*	H
	$c = 1$	$c = 2$		
Luminosity	-	5.0	5.0	5.0
3-tag				
JER	< 1	< 1	< 1	< 1
JES	2	< 1	< 1	< 1
JMR	1	12	12	11
JMS	5	14	13	17
b -tagging	1	23	22	23
Theoretical	-	3	3	3
Multijet Normalization	3	-	-	-
Statistical	2	1	1	1
Total	7	31	30	33
4-tag				
JER	< 1	< 1	< 1	< 1
JES	< 1	< 1	< 1	< 1
JMR	4	12	13	13
JMS	5	13	13	14
b -tagging	2	36	36	36
Theoretical	-	3	3	3
Multijet Normalization	14	-	-	-
Statistical	3	1	1	1
Total	15	42	42	43

Table 7.6: Summary of systematic uncertainties in the total background and signal event yields (expressed in %) in the boosted 3-tag and 4-tag signal regions. Systematic uncertainties on the signal normalization are shown for models with $m_{G_{KK}^*} = 1.5$ TeV and both $c = 1$ and $c = 2$ as well as a narrow width heavy scalar.

2455 be within the statistical uncertainties on the background normalization in the control region.

2456 Shape uncertainties on the background are evaluated using two techniques. First, as shown in figure 7.13, the ratio between the data and background prediction is fit with a linear function. The uncertainties on the slope of this fit are assigned as shape uncertainties. An additional uncertainty is assigned 2457 by using alternate power law fit functions for the smoothing of the background shape. Table 7.7 shows 2458 the alternate shapes used. The largest difference between the nominal fit function and the alternates, 2459 taking into account the 1σ uncertainty band on each fit as well, is taken as a shape uncertainty.

2460

Functional Form
$f_1(x) = p_0(1-x)^{p_1}x^{p_2}$
$f_2(x) = p_0(1-x)^{p_1}e^{p_2 x^2}$
$f_3(x) = p_0(1-x)^{p_1}x^{p_2} x$
$f_4(x) = p_0(1-x)^{p_1}x^{p_2} \ln x$
$f_5(x) = p_0(1-x)^{p_1}(1+x)^{p_2} x$
$f_6(x) = p_0(1-x)^{p_1}(1+x)^{p_2} \ln x$
$f_7(x) = \frac{p_0}{x}(1-x)^{p_1-p_2} \ln x$
$f_8(x) = \frac{p_0}{x^2}(1-x)^{p_1-p_2} \ln x$

Table 7.7: Alternate fit functions used to model the M_{JJ} distribution in the QCD multijet background. In the equations, $x = M_{JJ}/\sqrt{s}$.

2462 The uncertainties on the $t\bar{t}$ background are obtained by propagating the various experimental varia-
 2463 tions (JES, JER, JMS, JMR, b -tagging) through the analysis selection requirements. Table 7.6 summarizes
 2464 the background uncertainties in the 3-tag and 4-tag regions.

2465 7.8 RESULTS

2466 Table 7.8 shows the observed yields in the 3-tag and 4-tag signal regions for the boosted analysis com-
 2467 pared to the predicted number of background events. In the 3-tag region, 316 events are observed with
 2468 a predicted background of 285 ± 19 . In the 4-tag region, 20 events are observed with a predicted back-
 2469 ground of 14.6 ± 2.4 . Figure 7.14 shows the M_{JJ} distribution in the 3-tag and 4-tag regions. There are
 2470 some small excesses in the data, in particular in the 3-tag region around $M_{JJ} \approx 900$ GeV and in the
 2471 region of $1.6 < M_{JJ} < 2.0$ TeV. The significance of these excesses will be evaluated in the next chapter
 2472 in the statistical combination with the resolved results.

Sample	Signal Region (3-tag)	Signal Region (4-tag)
Multijet	235 ± 14	13.5 ± 2.4
$t\bar{t}$	48 ± 22	1.2 ± 1.0
$Z+jets$	2.0 ± 2.2	-
Total	285 ± 19	14.6 ± 2.4
Data	316	20
$G_{KK}^*(1000 \text{ GeV}), c = 1$	3.4 ± 0.9	2.9 ± 1.1

Table 7.8: Observed yields in the 3-tag and 4-tag signal regions for the boosted analysis compared to the predicted number of background events Errors correspond to the total uncertainties in the predicted event yields. The yields for a graviton with $m_{G_{KK}^*} = 1 \text{ TeV}$ and $c = 1$ are also shown. [28]

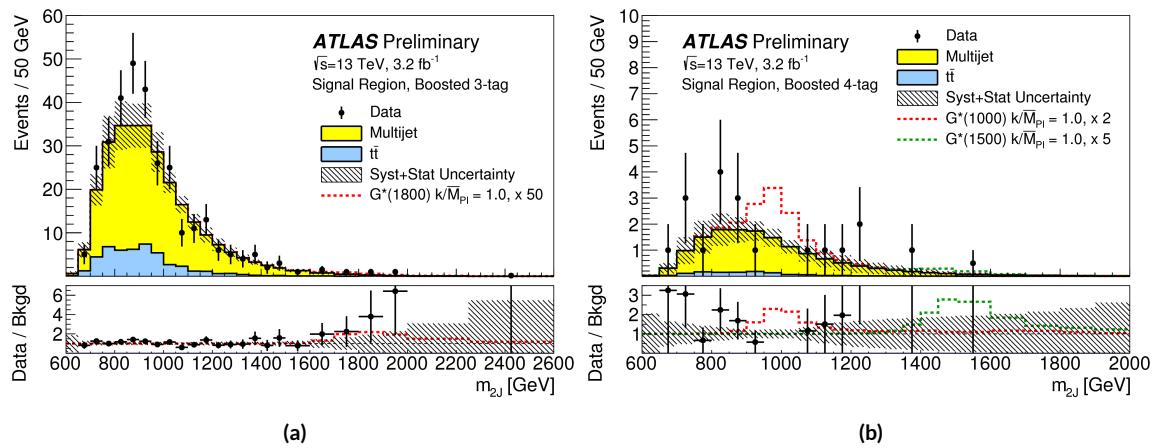


Figure 7.14: Di-jet invariant mass (M_{2J}) in the $3b$ (a) and $4b$ (b) signal regions. The multijet and $t\bar{t}$ backgrounds are estimated using the data-driven methods described above. In the $3b$ region, a graviton signal with $m_{G_{KK}^*} = 1.8 \text{ TeV}$ and $c = 1$ is overlaid, with the cross section multiplied by a factor of 50 so that the signal is visible. In the $4b$ region, signals with $m_{G_{KK}^*} = 1.0 \text{ TeV}$ and $m_{G_{KK}^*} = 1.5 \text{ TeV}$ are overlaid, both with $c = 1$ and the yields multiplied by factors of 2 and 5 respectively [28].

This is a really enlightening quote.

Tomo Lazovich

8

2473

2474

Combined limits from boosted and resolved searches

2475

2476 8.1 INTRODUCTION

2477 In order to cover the full mass range of possible resonances decaying to di-Higgs final states, two distinct
2478 tailored selections were produced. The resolved selection is more sensitive in the mass range of $400 < m_X < 1100$ GeV while the boosted selection is more sensitive to masses in the range $1100 < m_X <$
2479 3000 GeV. Chapter 7 presents the details of the boosted selection and results. In setting limits on spin-2
2480 Randall-Sundrum graviton (RSG) and narrow width heavy scalar (H) models, the results of the boosted
2481 selection are combined with the results of the resolved selection to cover the full mass range.
2482

2483 This chapter presents limits on signal models resulting from the $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search in both
2484 the resolved and boosted selections. It first presents a brief overview of the resolved results that go into

2485 the limit setting. Then, an overview of the statistical methods used for the search and limit setting is
 2486 given. Finally, limits on the RSG and heavy scalar models are presented.

2487 **8.2 RESOLVED RESULTS**

2488 The details of the resolved selection will not be presented here and can be found in reference [28]. In
 2489 basic terms, the selection searches for four $R = 0.4$ b-tagged calorimeter jets (where each pair of jets is
 2490 one Higgs candidate). This is distinct from the boosted methodology which searches for merged decay
 2491 products. The backgrounds to the resolved selection are the same as those presented in Chapter 7 for the
 2492 boosted analysis.

2493 Table 8.1 shows the results for data yields and expected background in the resolved signal region. Fig-
 2494 ure 8.1 shows the M_{2J} distribution in the resolved signal region. The total number of events is consis-
 2495 tent with the prediction and no significant excess is seen. One event in the boosted 4-tag signal is shared
 2496 with the resolved signal region and has a mass of 852 GeV.

Sample	Signal Region Yield
Multijet	43.3 ± 2.3
$t\bar{t}$	4.3 ± 3.0
$Z + \text{jets}$	-
Total	47.6 ± 3.8
Data	46
SM hh	0.25 ± 0.07
$G_{\text{KK}}^*(800 \text{ GeV}), c = 1$	5.7 ± 1.5

Table 8.1: Observed yields in the resolve selection 4-tag signal region compared to the predicted number of background events Errors correspond to the total uncertainties in the predicted event yields. The yields for a graviton with $m_{G_{\text{KK}}^*} = 800 \text{ GeV}$ and $c = 1$ are also shown. [28]

2497 **8.3 SEARCH TECHNIQUE AND RESULTS**

2498 The statistical technique used for the search in this analysis is the same as that used in the $H \rightarrow WW^*$
 2499 analysis presented in section 3.6.2. The test statistic q_0 is used to define the p -values which measure the

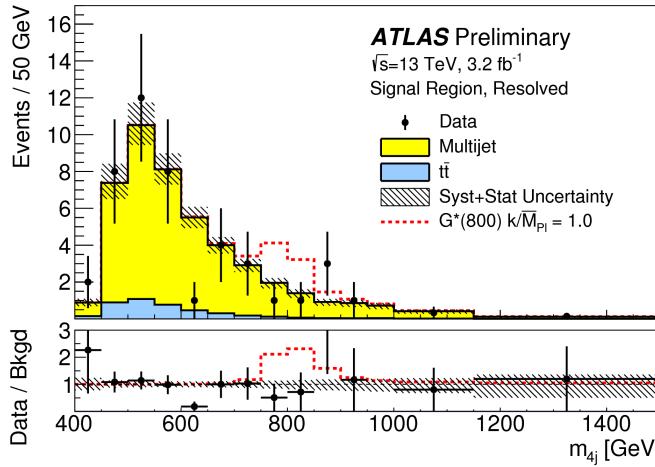


Figure 8.1: Di-jet invariant mass (M_{2J}) in the resolved signal region. Agraviton signal with $m_{G^*_\text{KK}} = 800 \text{ GeV}$ and $c = 1$ is overlaid. [28].

2500 compatibility of the data with the background-only hypothesis corresponding to a signal strength $\mu =$
 2501 0.

2502 Local p_0 values are computed to quantify the probability that the background could produce a fluc-
 2503 tuation greater than or equal to the one observed in the data. In the resolved analysis, no significant ex-
 2504 cesses are observed. The largest discrepancy with respect to the background only hypothesis occurs near
 2505 a resonance mass of 900 GeV and is found to be less than 2σ in significance.

2506 In the boosted selection, the largest local excess is a broad excess in the $3b$ signal region that begins
 2507 near $M_{2J} \approx 1.7 \text{ GeV}$. Assuming a G^*_KK with this mass and $c = 1.0$, the local significance of this excess
 2508 is 2.0σ .

2509 8.4 LIMIT SETTING

2510 In the absence of any significant excess observed in the data, limits on different signal models can be set.
 2511 This section describes the limit setting procedure and presents combined results of the resolved and
 2512 boosted analyses.

2513 8.4.1 LIMIT SETTING PROCEDURE

2514 The procedure used for setting exclusion limits in this analysis is the CL_s method [107]. The first step in
 2515 setting the limits is to define a test statistic which will be used. For limit setting, the test statistic is shown
 2516 in equation 8.1.

$$\tilde{q}_\mu = \begin{cases} -2 \ln \frac{L(\mu, \hat{\theta}(\mu))}{L(0, \hat{\theta}(0))} & \hat{\mu} < 0 \\ -2 \ln \frac{L(\mu, \hat{\theta}(\mu))}{L(\hat{\mu}, \hat{\theta})} & 0 \leq \hat{\mu} < \mu \\ 0 & \hat{\mu} > \mu \end{cases} \quad (8.1)$$

2517 In the above equation, μ is the value of the signal strength under test, $\hat{\mu}$ is the best fit μ , $\hat{\theta}$ is the
 2518 best fit value of the nuisance parameters, $\hat{\theta}$ is the best fit value of the nuisance parameters under the fixed
 2519 μ value, and L is the Poisson likelihood of the data (as described in section 3.6.2).

2520 The test statistic \tilde{q}_μ is constructed to protect against two interesting corner cases when setting the
 2521 upper limit on the cross section. First, it protects against negative signal strengths μ which are unphys-
 2522 ical. Second, it does not count excesses in the data larger than those expected by a signal strength μ as
 2523 evidence against the μ hypothesis.

2524 The CL_s statistic is constructed by taking a ratio of two probabilities. CL_{s+b} is the probability that
 2525 the signal+background hypothesis would produce a value of the test statistic that is less than or equal
 2526 to the observed value*. CL_b is the probability that the background only hypothesis will pro-
 2527 duce a value of the test statistics less than or equal to the observed. The CL_s statistic is then the ratio
 2528 $\text{CL}_{s+b}/\text{CL}_b$. A 95% upper limit on the cross section is set at the value of μ that makes the CL_s statistic
 2529 less than 5%.

2530 In practice, the limits are computed numerically within an asymptotic approximation for the distri-
 2531 bution of the test statistic \tilde{q}_μ . The details of this approximation can be found in reference [70].

2532 The resolved and boosted analyses are combined using a very simple procedure rather than a full sta-
 2533 tistical combination. For each mass point tested, the limit which gives the most stringent constraint is

*Lower values of \tilde{q}_μ mean better compatibility

used. This means that for mass points below 1.1 TeV the resolved signal region is used, while at and above this point the combination of the orthogonal $3b$ and $4b$ boosted signal regions is used.

8.4.2 LIMIT SETTING RESULTS

Figure 8.2 shows the combined 95% upper bounds as a function of mass for three different models: G_{KK}^* with $c = 1$, G_{KK}^* with $c = 2$, and a narrow heavy scalar H .

The cross section of $\sigma(pp \rightarrow G_{\text{KK}}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $c = 1$ is constrained to be less than 70 fb for masses in the range $600 < m_{G_{\text{KK}}^*} < 3000$ GeV. For the RSG model with $c = 2$, cross sections limits between 40 fb and 200 fb are set for the mass range of $500 < m_{G_{\text{KK}}^*} < 3000$ GeV. Masses in the range of $475 < m_{G_{\text{KK}}^*} < 785$ GeV are excluded with $c = 1$ (with an exclusion of the range 465 to 745 GeV expected). Masses less than 980 GeV are excluded with $c = 2$ (with an exclusion for masses less than 1 TeV expected).

In the heavy Higgs model, the cross section upper limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges from 30 to 300 fb in the mass range of $500 < m_H < 3000$ GeV.

The resolved analysis can also set an upper limit on the Standard Model di-Higgs production cross section discussed in chapter 3. The upper limit on $\sigma(pp \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ in the Standard Model is constrained to be less than 1.22 pb.

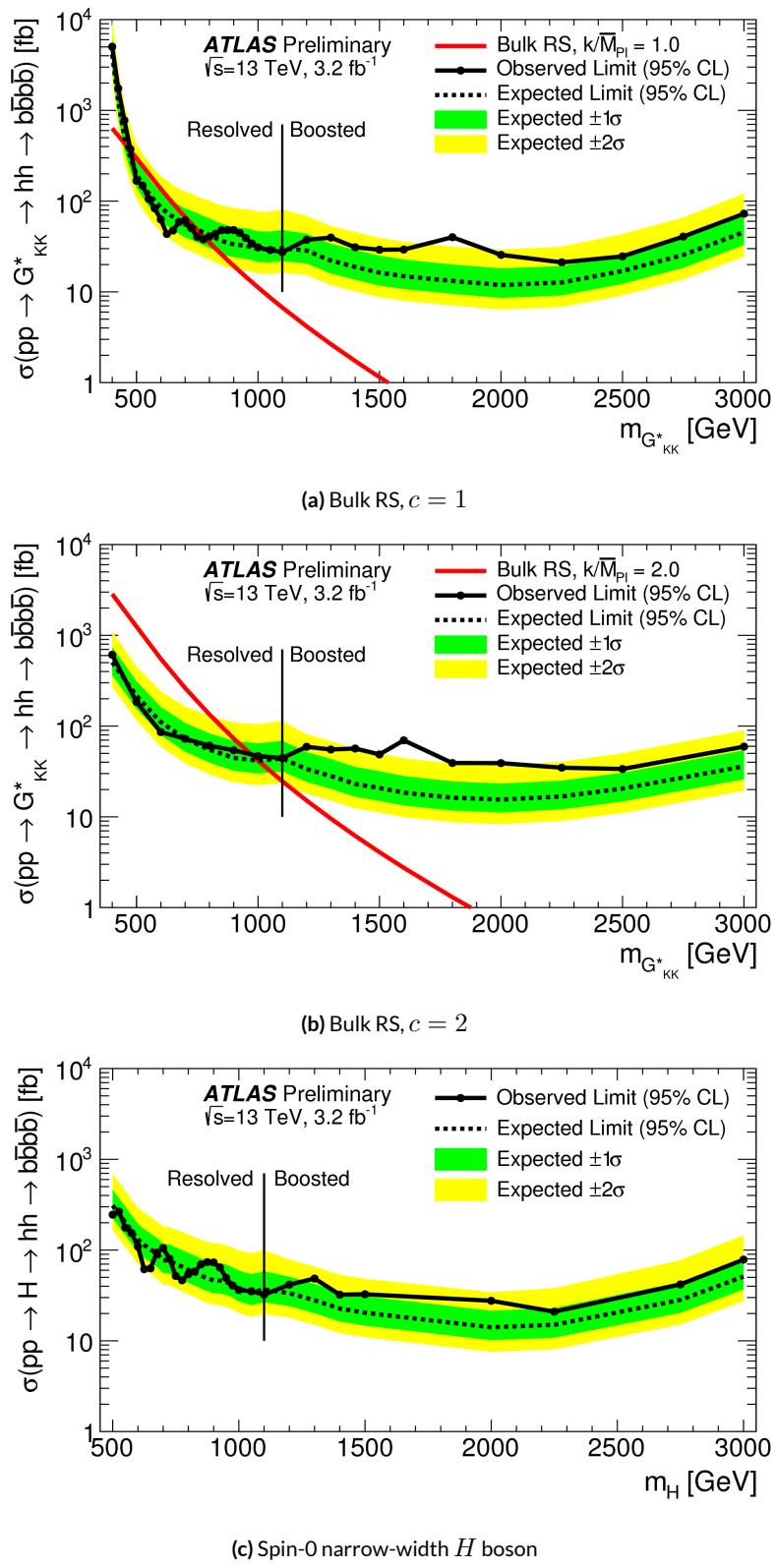


Figure 8.2: Expected and observed upper limit as a function of mass for G^*_{KK} in the RSG model with (a) $c = 1$ and (b) $c = 2$, as well as (c) H with fixed $\Gamma_H = 1$ GeV, at the 95% confidence level in the CL_s method. [28]

2550

Part IV

2551

Looking ahead

9

2552

2553

Conclusion

2554 This dissertation presented two distinct studies: the observation and measurement of the Higgs boson
2555 in the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$ channel at $\sqrt{s} = 7$ TeV and $\sqrt{s} = 8$ TeV and a search for Higgs pair
2556 production in the $HH \rightarrow b\bar{b}b\bar{b}$ channel at $\sqrt{s} = 13$ TeV with the ATLAS detector in pp collisions at
2557 the Large Hadron Collider.

2558 In the $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$, results from both the discovery of the Higgs boson and the full
2559 ATLAS Run 1 dataset were presented. With the full 20.3 fb^{-1} at $\sqrt{s} = 8$ TeV and 4.2 fb^{-1} at $\sqrt{s} =$
2560 7 TeV, ATLAS achieved discovery level significance in the $H \rightarrow WW^*$ channel alone and obtained
2561 the first observation of vector boson fusion production in that channel. The combined signal strength
2562 is measured to be $\mu = 1.09^{+0.23}_{-0.21}$. The total observed significance of the $H \rightarrow WW^*$ process is ob-
2563 served to be 6.1σ (with 5.8σ expected). Advanced methods for background reduction and estimation,
2564 particularly in same-flavor lepton final states, are shown. The VBF signal strength is measured to be
2565 $\mu_{\text{VBF}} = 1.27^{+0.53}_{-0.45}$ with an observed significance of 3.2σ (with 2.7σ expected).

2566 These results required many novel innovations. The increase of pileup interactions in the higher in-
2567 stantaneous luminosity LHC conditions of 2012 led to a degradation of missing transverse momentum
2568 resolution. As a result, the prominent Z/γ^* +jets background of the same flavor $H \rightarrow WW^* \rightarrow \ell\nu\ell\nu$
2569 final states increased greatly. New variables, including a track-based missing transverse momentum and
2570 a measurement of the balance between the dilepton system and recoiling jets, allowed for significant re-
2571 duction of this background. In the VBF channel, selections were optimized to exploit the unique VBF
2572 final state topology. Incorporating these variables into a boosted decision tree technique allowed the
2573 analysis to exceed the 3σ observation threshold.

2574 At $\sqrt{s} = 13$ TeV, a search for Higgs pair production in the $b\bar{b}b\bar{b}$ final state with 3.2 fb^{-1} was con-
2575 ducted. A signal region optimized for the boosted final states arising from high mass resonances was
2576 constructed. This signal region utilized large-radius calorimeter jets and b -tagging with small radius track
2577 jets to maximize the signal acceptance. No significant excesses were observed, and upper limits on cross
2578 sections are placed for spin-2 Randall Sundrum gravitons (RSG) and narrow spin-0 resonances. The
2579 cross section of $\sigma(pp \rightarrow G_{KK}^* \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ with $k/\bar{M}_{\text{Pl}} = 1$ is constrained to be less than 70 fb for
2580 masses in the range $600 < m_{G_{KK}^*} < 3000 \text{ GeV}$. For the RSG model with $k/\bar{M}_{\text{Pl}} = 2$, cross sections
2581 limits between 40 fb and 200 fb are set for the mass range of $500 < m_{G_{KK}^*} < 3000 \text{ GeV}$. The cross
2582 section upper limits for $\sigma(pp \rightarrow H \rightarrow hh \rightarrow b\bar{b}b\bar{b})$ ranges from 30 to 300 fb in the mass range of
2583 $500 < m_H < 3000 \text{ GeV}$.

A

2584

2585

Optimization of b -tagging working point in

2586

$$X \rightarrow HH \rightarrow b\bar{b}b\bar{b} \text{ search}$$

2587 To the $3b$ and $4b$ signal regions in the boosted $X \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ search, the MV₂ algorithm with a
2588 20% fraction of charm events in training is used (MV_{2c20}). Once the algorith is selected, an efficiency
2589 working point must also be chosen. This working point defines the efficiency with which true b -jets
2590 are tagged and also fixes the overall background rejection of the algorithm. Higher efficiency working
2591 points accept more true b -jets but also allow for more background. Five different working points (70%,
2592 77%, 80%, 85%, 90%) are tested. With each working point, the full data driven background estimation
2593 method is run to quantify the amount of background that will be present in the final signal region. The
2594 significance is quantified using the median discovery significance for signal and background with Poisson

2595 errors, given in equation A.1 [108].

$$Z = \sqrt{2 \left((s + b) \ln \left(1 + \frac{s}{b} \right) - s \right)} \quad (\text{A.1})$$

2596 Note that in the limit where s is much smaller than b , this equation reduces to the more well known
 2597 s/\sqrt{b} .

2598 Figure A.1 shows the estimated significance as a function of signal mass in RSG $c = 1$ models for
 2599 the $3b$ and $4b$ signal regions. The 77% working point gives the best performance over a wide range of
 2600 masses in the $4b$ signal region. As this is the region which contributes the most to the total discovery
 2601 significance, the 77% efficiency working point is chosen for the analysis.

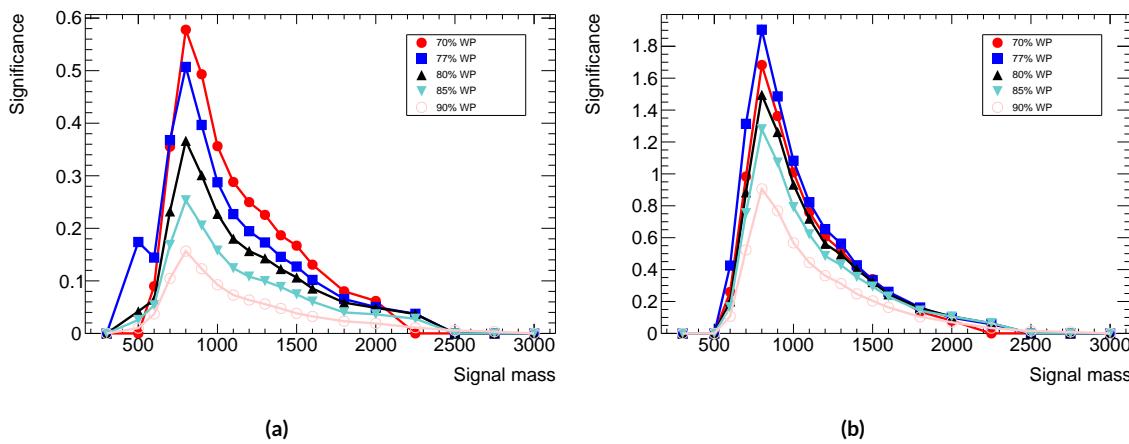


Figure A.1: Estimated significance as a function of signal mass for RSG $c = 1$ models in the $3b$ (a) and $4b$ (b) regions for different b -tagging efficiency working points

B

2602

2603

b-tagging performance at high p_T

2604 One of the limiting factors of the signal acceptance in the *FourB full* search at high resonance masses is
2605 the degradation of the *b*-tagging efficiency for high p_T jets. This appendix presents a study of the under-
2606 lying causes of this degradation.

2607 B.I MV2 ALGORITHM OVERVIEW

2608 The MV2 algorithm is a boosted decision tree incorporating twenty-four input variables constructed
2609 from three lower level input algorithms: IPxD, SV1, and JetFitter. IPxD uses the two and three dimen-
2610 sional impact parameter information of tracks in the jet to construct templates for light, charm, and bot-
2611 tom quarks and compute likelihood ratios. SV1 is a secondary vertex reconstruction algorithm. JetFitter
2612 attempts to fit to the full decay chain of the B hadron, looking for multiple decay vertices aligned along a
2613 single axis. Figure B.I summarizes the inputs to MV2.

IP2D and IP3D (6 inputs)	SV1 (8 inputs)	JetFitter (8 inputs)
$\log(p_b/p_u)$	Mass	Mass
$\log(p_b/p_c)$	Energy fraction	Energy fraction
$\log(p_c/p_u)$	# tracks at vertex	# vertices
	# 2 track vertices	# tracks at vertex
	Lxy	# 1 track vertices
	L3d	# 2 track vertices
	3D significance	3D significance
	ΔR	ΔR
Kinematics (2 inputs)		
	p_T	
	η	

Figure B.1: Summary of the inputs to the MV2 b -tagging algorithm

2614 **B.2 CHANGES IN MV2 SCORE AT HIGH p_T**

2615 The degradation of b -tagging at high p_T was studied in particular in the context of RSG models at high
 2616 mass. Figure B.2 shows the p_T of the leading track jet inside of the leading calorimeter jet in RSG events.
 2617 At high $m_{G_{KK}^*}$, the p_T spectrum of track jets is much harder than at lower masses due to the increased
 2618 Higgs p_T .

2619 Figure B.3 shows the MV2c20 algorithm score for the leading and subleading track jets inside of the
 2620 leading calorimeter jet. In both cases, it can be seen that at higher RSG masses the MV2 score shifts to-
 2621 wards more background like (negative) values. Additionally, this effect is more pronounced in the lead-
 2622 ing track jet than the subleading.

2623 To understand what is causing this change in the MV2c20 score, the same comparisons can be made
 2624 for the input variables of MV2c20. The focus in these comparisons will be on the leading track jet as
 2625 this is the one seen to have the largest difference in MV2 score. Figure B.4 shows the log likelihood ratio
 2626 $\log(p_b/p_u)$ from the IP3D (three dimensional impact parameter) algorithm. At higher masses, the IP3D
 2627 likelihood ratio distribution does become more background-like.

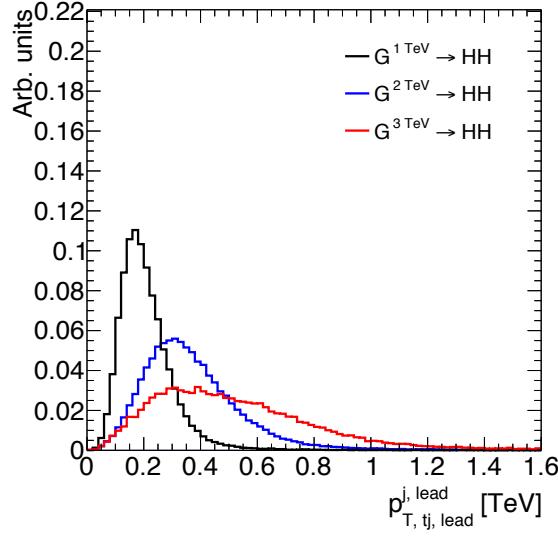


Figure B.2: p_T of the leading track jet in the leading calorimeter jet for different signal masses in RSG $c = 1$ models

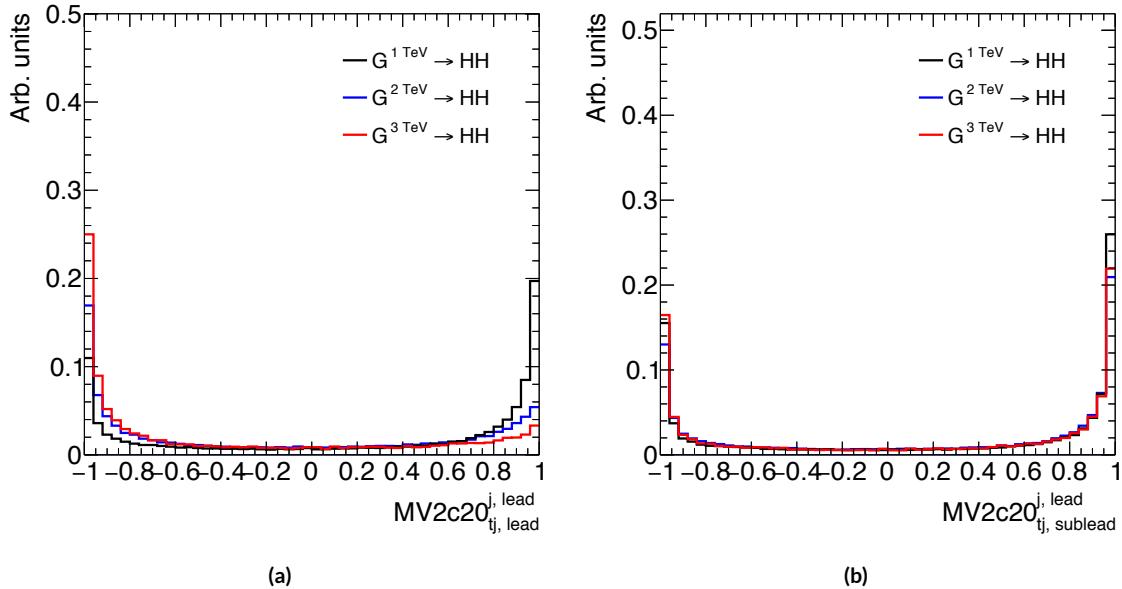


Figure B.3: MV2c20 score for the leading track jet (a) and subleading track jet (b) of the leading calorimeter jet for different signal masses in RSG $c = 1$ models

2628 Figure B.5 shows the mass and number of tracks at the secondary vertex computed by the SVI algo-
 2629 rithm. When there is no secondary vertex found, the algorithm assigns a default negative value for these
 2630 quantities. Both of these distributions show that there is a significantly larger fraction of jets where no
 2631 secondary vertex is found in the high mass samples compared to the $m_{G_{KK}^*} = 1$ TeV sample. The SVI

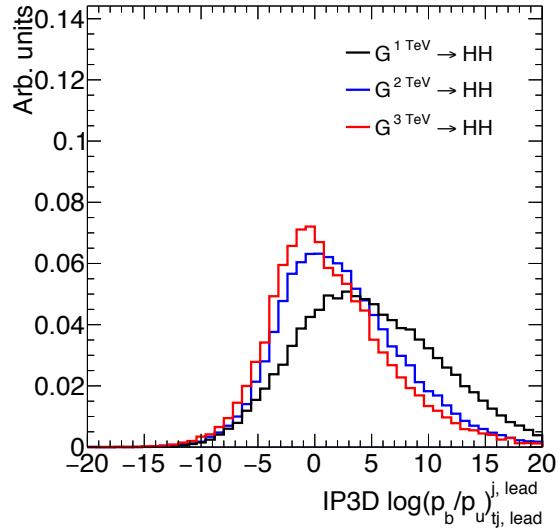


Figure B.4: IP3D log-likelihood ratio ($\log(p_b/p_u)$) of the leading track jet in the leading calorimeter jet for different signal masses in RSG $c = 1$ models

algorithm's inability to find a secondary vertex could be an important factor in the overall MV₂ score shift, as this eliminates eight of the input variables that would normally contribute information to the algorithm.

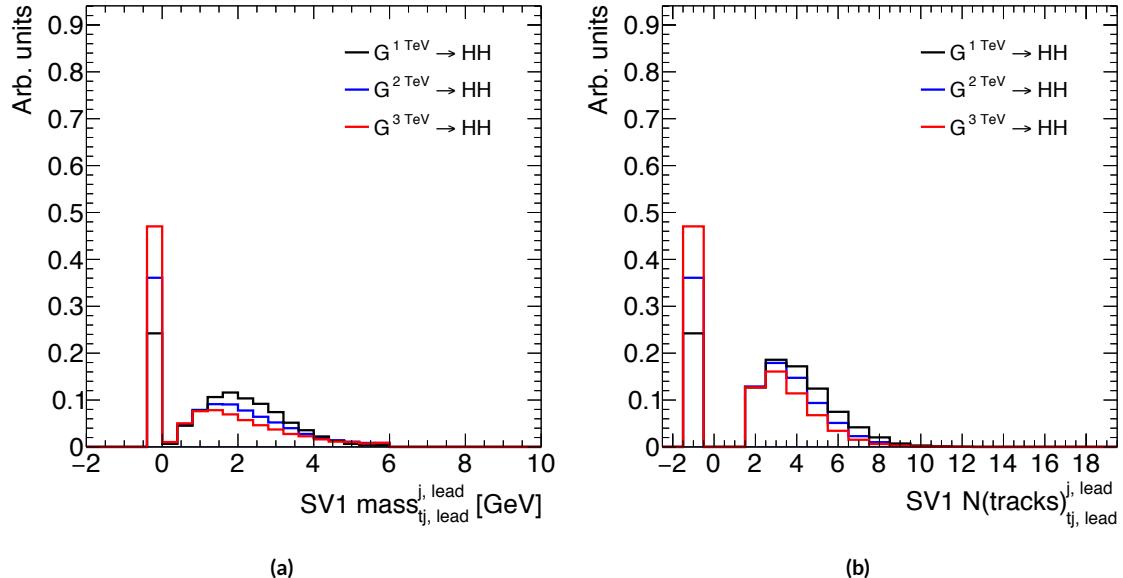


Figure B.5: Mass (a) and number of tracks (b) for the secondary vertices computed with the SV1 algorithm. When no secondary vertex is found, the quantites are assigned to default negative values.

2635 Figure B.6 shows the same quantities for the JetFitter algorithm. In this case, there is also a change in
 2636 the fraction of jets which have their secondary vertices successfully reconstructed, but this change is not
 2637 as drastic as that seen in SV1. There is also an increase in the number of jets which have high values of
 2638 mass.

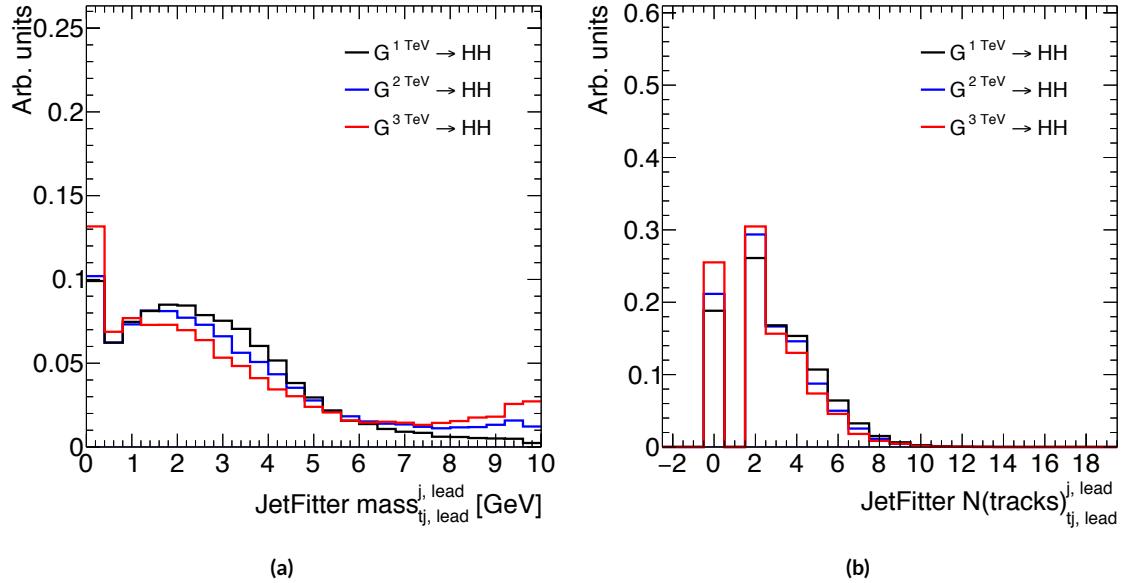


Figure B.6: Mass (a) and number of tracks (b) for vertices computed with the JetFitter algorithm. When no vertices are found, the quantities are assigned to default negative values.

2639 B.3 TAGGING EFFICIENCY BY INDIVIDUAL JET

2640 In the last section, the largest changes in MV2 score were seen for the leading track jet inside the leading
 2641 calorimeter jet. To confirm that the overall 4*b* tagging efficiency is indeed degrading because of degra-
 2642 dation of the leading track jet efficiency, the tagging efficiency for each individual jet as a function of mass
 2643 can be compared. This is shown in figure B.7. The figure shows that the leading jet tagging efficiency
 2644 in both calorimeter jets degrades heavily, while the sub-lead jet tagging efficiency remains relatively con-
 2645 stant.

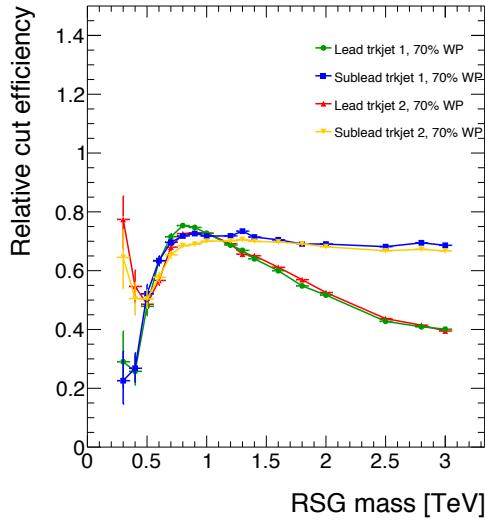


Figure B.7: MV2c20 b -tagging efficiency for each of the four track jets in the boosted $4b$ selection as a function of RSG mass for $c = 1$ models.

B.4 EFFECT OF MULTIPLE b -QUARKS INSIDE ONE JET

One hypothesis for why the efficiency of b -tagging the leading track jet degrades is that at high masses, the b quarks get close enough together that both of them are inside of the leading track jet. Because MV2 is not tuned for tagging multiple b quarks inside one jet, the tagging efficiency could degrade. Figure B.8 shows MV2 scores and SV1 mass for cases where there are two b quarks at truth level within the radius of the leading track jet compared to cases where there is only one true b^* .

B.5 CHANGES IN TRACK QUALITY AT HIGH p_T

To investigate the change in MV2 score further, individual track quantities within the jet can be compared.

*When two truth b quarks are required in the leading jet, the subleading jet is required to have zero. When one is required for the leading, one is also required for the subleading.

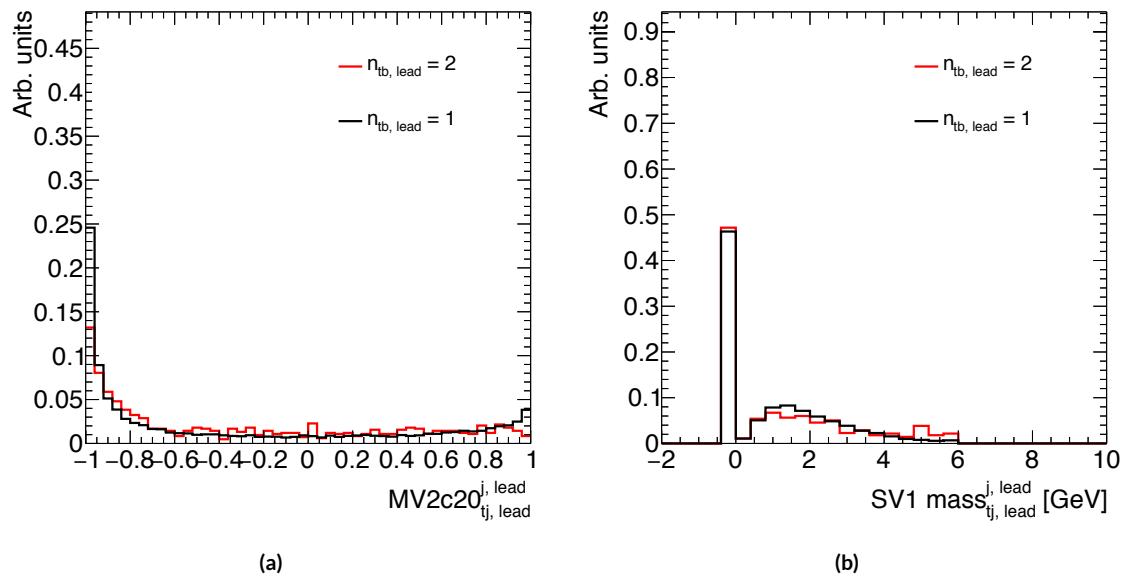


Figure B.8: MV2c20 score (a) and SV1 mass (b) for leading track jets with two truth b quarks ($n_{tb, \text{lead}} = 2$) compared to those with only one truth b ($n_{tb, \text{lead}} = 1$).

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