

Lecture 3: Wasserstein Space

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The material of today's lecture is adapted from Q. Mérigot's lecture notes and [3, 4].

1 Reminders

Let X, Y be compact metric spaces, $c \in \mathcal{C}(X \times Y)$ the cost function and $(\mu, \nu) \in \mathcal{P}(X) \times \mathcal{P}(Y)$ the marginals. In previous lectures, we have seen that the optimal transport problem can be formulated as an optimization over the space of transport plans $\Pi(\mu, \nu)$ — the primal or Kantorovich problem — and as an optimization over potential functions $\{(\varphi, \psi) \in \mathcal{C}(X) \times \mathcal{C}(Y) \mid \varphi \oplus \psi \leq c\}$ — the dual problem. We recall the following results:

- minimizer/maximizers exist for both problems and, for the dual, can be chosen as (φ, φ^c) with φ c -concave.
- at optimality, it holds $\varphi(x) + \psi(y) = c(x, y)$ for γ -almost every (x, y)
- we have the following special cases:
 - for $X = Y \subset \mathbb{R}$ and $c(x, y) = h(y - x)$ with h strictly convex, the optimal transport plan is the (unique) monotone plan, which can be characterized with the quantile functions of μ and ν .
 - for $X = Y$ and $c(x, y) = \text{dist}(x, y)$, we have the Kantorovich-Rubinstein formula

$$\mathcal{T}_c(\mu, \nu) = \sup_{\varphi \text{ 1-Lip}} \int \varphi d(\mu - \nu).$$

- for $X = Y \subset \mathbb{R}^d$ and $c(x, y) = \frac{1}{2}|y - x|^2$, and when μ is absolutely continuous, there exists a unique optimal transport plan. It is of the form $\gamma = (\text{id}, \nabla \tilde{\varphi})_{\#} \mu$ for some $\tilde{\varphi} \in \mathcal{C}(\mathbb{R}^d)$ convex.

2 Wasserstein space

2.1 Definition and elementary properties

Definition 2.1 (Wasserstein space). Let (X, dist) be a compact metric space. For $p \geq 1$, we denote by $\mathcal{P}_p(X)$ the set of probability measures on X endowed with the p -Wasserstein distance, defined as

$$W_p(\mu, \nu) := \left(\min_{\gamma \in \Pi(\mu, \nu)} \int \text{dist}(x, y)^p d\gamma(x, y) \right)^{1/p} = \mathcal{T}_{\text{dist}^p}(\mu, \nu)^{\frac{1}{p}}.$$

This distance is a natural way to build a distance on $\mathcal{P}(X)$ from a distance on X . In particular, the map $\delta : X \rightarrow \mathcal{P}_p(X)$ mapping a point $x \in X$ to the Dirac mass δ_x is an isometry.

Proposition 2.2. W_p satisfies the axioms of a distance on $\mathcal{P}_p(X)$.

Proof. The symmetry of the Wasserstein distance is obvious. Moreover, $W_p(\mu, \nu) = 0$ implies that there exists $\gamma \in \Pi(\mu, \nu)$ such that $\int \text{dist}^p d\gamma = 0$. This implies that γ is concentrated on the diagonal, so that $\gamma = (\text{id}, \text{id})_{\#}\mu$ is induced by the identity map. In other words, $\nu = \text{id}_{\#}\mu = \mu$.

To prove the triangle inequality we will use the gluing lemma below (Lemma 2.3) with $N = 3$. Let $\mu_i \in \mathcal{P}_p(X)$ for $i \in \{1, 2, 3\}$ and let $\gamma_1 \in \Pi(\mu_1, \mu_2)$ and $\gamma_2 \in \Pi(\mu_2, \mu_3)$ be optimal in the definition of W_p . Then, there exists $\sigma \in \mathcal{P}(X^3)$ such that $(\pi_{i,i+1})_{\#}\sigma = \gamma_i$ for $i \in \{1, 2\}$. A fortiori one has $(\pi_1)_{\#}\sigma = \mu_1$ and $(\pi_3)_{\#}\sigma = \mu_3$, so that $(\pi_{13})_{\#}\sigma \in \Pi(\mu_1, \mu_3)$. In particular,

$$\begin{aligned} W_p(\mu_1, \mu_3) &\leq \left(\int_{X^2} \text{dist}(x, y)^p d(\pi_{1,3})_{\#}\sigma(x, y) \right)^{1/p} \\ &= \left(\int_{X^3} \text{dist}(x_1, x_3)^p d\sigma(x_1, x_2, x_3) \right)^{1/p} \\ &\leq \left(\int_{X^3} (\text{dist}(x_1, x_2) + \text{dist}(x_2, x_3))^p d\sigma(x_1, x_2, x_3) \right)^{1/p} \\ &\leq \left(\int_{X^3} \text{dist}(x_1, x_2)^p d\sigma(x_1, x_2, x_3) \right)^{1/p} + \left(\int_{X^3} \text{dist}(x_2, x_3)^p d\sigma(x_1, x_2, x_3) \right)^{1/p} \\ &= W_p(\mu_1, \mu_2) + W_p(\mu_2, \mu_3), \end{aligned}$$

where we used the Minkowski inequality in $L^p(\sigma)$ to get the second inequality, and the property $(\pi_{i,i+1})_{\#}\sigma = \gamma_i$ to get the last equality. \square

Lemma 2.3 (Gluing). *Let X_1, \dots, X_N be complete and separable metric spaces, and for any $1 \leq i \leq N - 1$ consider a transport plan $\gamma_i \in \Pi(\mu_i, \mu_{i+1})$. Then, there exists $\gamma \in \mathcal{P}(X_1, \dots, X_N)$ such that for all $i \in \{1, \dots, N - 1\}$, $(\pi_{i,i+1})_{\#}\gamma = \gamma_i$, where $\pi_{i,i+1} : X_1 \times \dots \times X_N \rightarrow X_i \times X_{i+1}$ is the projection.*

Proof. See Lemma 5.3.2 and Remark 5.3.3 in [1]. \square

Exercise 2.4. Prove the triangle inequality assuming the existence of optimal transport maps between μ_1, μ_2 and μ_2, μ_3 .

Remark 2.5 (Non-compact case). As usual, the compactness assumption is only here for clarity of presentation. In general, when X is a complete and separable metric space, the space $\mathcal{P}_p(X)$ is defined as the set of probability measures such that for some (and thus any) $x_0 \in X$ it holds

$$\int \text{dist}(x_0, y)^p d\mu(y) < \infty.$$

It can be shown that this set endowed with the distance W_p is also a complete and separable metric space. Exercise: show that the Wasserstein distance W_p is finite on this set.

2.2 Comparisons

Comparison between Wasserstein distances Note that, due to Jensen's inequality, since all $\gamma \in \Pi(\mu, \nu)$ are probability measures, for $p \leq q$ we have

$$\left(\int \text{dist}(x, y)^p d\gamma \right)^{\frac{1}{p}} \leq \left(\int \text{dist}(x, y)^q d\gamma \right)^{\frac{1}{q}},$$

which implies $W_p(\mu, \nu) \leq W_q(\mu, \nu)$. In particular, $W_1(\mu, \nu) \leq W_p(\mu, \nu)$ for every $p \geq 1$. On the other hand, for compact (and thus bounded) X , an opposite inequality also holds, since

$$\left(\int \text{dist}(x, y)^p d\gamma \right)^{\frac{1}{p}} \leq \text{diam}(X)^{\frac{p-1}{p}} \left(\int \text{dist}(x, y) d\gamma \right)^{\frac{1}{p}}.$$

This implies that for all $p \geq 1$,

$$W_1(\mu, \nu) \leq W_p(\mu, \nu) \leq \text{diam}(X)^{\frac{p-1}{p}} W_1(\mu, \nu)^{\frac{1}{p}}.$$

Comparison with L^p distances Take $X \subset \mathbb{R}$ with its usual distance. Consider the translation $t_x : y \in \mathbb{R} \mapsto x + y$. Then, since the map t_x is increasing, for any $\mu \in \mathcal{P}_p(X)$ one has $W_p(\mu, t_{x\#}\mu) = |x|$. However,

- if $\rho \in \mathcal{P}(X) \cap L^p(X)$ with $\text{spt}(\rho) \subseteq [0, 1]$, then for all $|x| \geq 1$ one has $\|\rho - t_x \rho\|_{L^p(X)} = 2 \|\rho\|_{L^p(X)}$ where $t_x \rho(y) = \rho(y-x)$. Unlike the $L^p(X)$ norm, the Wasserstein distance is geometry “aware”.
- If $\rho \in \mathcal{P}(X) \cap L^2(X)$ is such that $\|t_x \rho - \rho\|_{L^2(X)} = O(|x|)$, then ρ belongs to the Sobolev space H^1 (see Proposition VIII.3 in [2]). In other words, $\|t_x \rho - \rho\|_{L^2(X)}$ can be much larger than $|x|$ unless ρ is very regular.

Because of these two examples, the Wasserstein distance is a very appealing notion of distance for data analysis (e.g. measuring the distance between signals). The flipside is that the definition of the Wasserstein requires the signal to belong to $\mathcal{P}(X)$ i.e. to be non-negative and with unit mass.

2.3 Topological properties

Theorem 2.6. *Assume that X is compact. For $p \in [1, +\infty[$, we have $\mu_n \rightharpoonup \mu$ if and only if $W_p(\mu_n, \mu) \rightarrow 0$.*

Proof. We only need to prove the result for W_1 thanks to the comparison inequalities between W_1 and W_p in previous section. Let us start from a sequence μ_n such that $W_1(\mu_n, \mu) \rightarrow 0$. Thanks to the duality formula, for every $\varphi \in \text{Lip}_1(X)$, we have $\int \varphi(\mu_n - \mu) \rightarrow 0$. By linearity, the same is true for any Lipschitz function. By density, this holds for any function in $\mathcal{C}(X)$. This shows that convergence in W_1 implies weak convergence.

To prove the opposite implication, let us first fix a subsequence μ_{n_k} that satisfies $\lim_k W_1(\mu_{n_k}, \mu) = \limsup_n W_1(\mu_n, \mu)$. For every k , pick a function $\varphi_{n_k} \in \text{Lip}_1(X)$ such that $\int \varphi_{n_k}(\mu_{n_k} - \mu) = W_1(\mu_{n_k}, \mu)$. Up to adding a constant, which does not affect the integral, we can assume that the φ_{n_k} all vanish at the same point, and they are hence uniformly bounded and equi-continuous. By Ascoli-Arzelà theorem, we can extract a sub-sequence uniformly converging to a certain $\varphi \in \text{Lip}_1(X)$. By replacing the original subsequence with this new one, we have now

$$W_1(\mu_{n_k}, \mu) = \int \varphi_{n_k} d(\mu_{n_k} - \mu) \rightarrow \int \varphi d(\mu - \mu) = 0$$

where the convergence of the integral is justified by the weak convergence $\mu_{n_k} \rightharpoonup \mu$ together with the strong convergence in $\mathcal{C}(X)$ $\varphi_{n_k} \rightarrow \varphi$. This shows that $\limsup_n W_1(\mu_n, \mu) \leq 0$ and concludes the proof. \square

Remark 2.7. In the non-compact case, it can be shown that convergence in $\mathcal{P}_p(X)$ is equivalent to tight convergence (in duality with continuous and bounded functions) and convergence of the p -th order moments i.e. for all $x_0 \in X$,

$$\int \text{dist}(x_0, y)^p d\mu_n(y) \rightarrow \int \text{dist}(x_0, y)^p d\mu(y).$$

3 Geodesics in Wasserstein space

Definition 3.1. Let (X, dist) be a metric space. A constant speed geodesic between two points $x_0, x_1 \in X$ is a continuous curve $x : [0, 1] \rightarrow X$ such that for every $s, t \in [0, 1]$, $\text{dist}(x_s, x_t) = |s - t| \text{dist}(x_0, x_1)$.

Proposition 3.2. Let $\mu_0, \mu_1 \in \mathcal{P}_p(X)$ with $X \subset \mathbb{R}^d$ compact and convex. Let $\gamma \in \Pi(\mu_0, \mu_1)$ be an optimal transport plan. Define

$$\mu_t := (\pi_t)_\# \gamma \text{ where } \pi_t(x, y) = (1 - t)x + ty.$$

Then, the curve μ_t is a constant speed geodesic between μ_0 and μ_1 .

Example 3.3. If there exists an optimal transport map T between μ_0 and μ_1 , then the geodesic defined above is $\mu_t = ((1 - t)\text{id} + tT)_\# \mu_0$.

Remark 3.4. In fact, it can be shown that any geodesic between μ_0 and μ_1 can be constructed as in Proposition 3.2.

Proof. First note that if $0 \leq s \leq t \leq 1$,

$$W_p(\mu_0, \mu_1) \leq W_p(\mu_0, \mu_s) + W_p(\mu_s, \mu_t) + W_p(\mu_t, \mu_1),$$

so that it suffices to prove the inequality $W_p(\mu_s, \mu_t) \leq |t - s| W_p(\mu_0, \mu_1)$ for all $0 \leq s \leq t \leq 1$ to get equality. The inequality is easily checked by building an explicit transport plan using an optimal transport plan γ . Take $\gamma_{st} := (\pi_s, \pi_t)_\# \gamma \in \Pi(\mu_s, \mu_t)$, so that

$$\begin{aligned} W_p(\mu_s, \mu_t)^p &\leq \int \|x - y\|^p d\gamma_{st}(x, y) = \int \|\pi_s(x, y) - \pi_t(x, y)\|^p d\gamma(x, y) \\ &= \int \|(1 - s)x + sy - ((1 - t)x + ty)\|^p d\gamma(x, y) \\ &= \int \|(t - s)(x - y)\|^p d\gamma(x, y) = (t - s)^p W_p(\mu, \nu)^p \quad \square \end{aligned}$$

Corollary 3.5. The space $(\mathcal{P}_p(X), W_p)$ with X compact and convex is a geodesic space, meaning that any $\mu_0, \mu_1 \in \mathcal{P}_p(X)$ can be joined by (at least one) constant speed geodesic.

4 Differentiability of the Wasserstein distance

In this section, we will compute the differential of the Wasserstein distance under additive perturbations.

Theorem 4.1. Let $\sigma, \rho_0, \rho_1 \in \mathcal{P}(X)$. Assume that there exists unique Kantorovich potentials (φ_0, ψ_0) between σ and ρ_0 which are c -conjugate to each other and satisfy $\varphi_0(x_0) = 0$ for some $x_0 \in X$. Then,

$$\frac{d}{dt} \mathcal{T}_c(\sigma, \rho_0 + t(\rho_1 - \rho_0))|_{t=0} = \int \psi_0 d(\rho_1 - \rho_0).$$

Proof. Denote $\rho_t = (1 - t)\rho_0 + t\rho_1 = \rho_0 + t(\rho_1 - \rho_0)$. By Kantorovich duality, we have

$$\mathcal{T}_c(\sigma, \rho_t) \geq \int \varphi_0 d\sigma + \int \psi_0 d\rho_t.$$

This immediately gives

$$\frac{1}{t}(\mathcal{T}_c(\sigma, \rho_t) - \mathcal{T}_c(\sigma, \rho_0)) \geq \int \psi_0 d(\rho_1 - \rho_0).$$

To show the converse inequality, we let (φ_t, ψ_t) be c -conjugate Kantorovich potentials between σ and ρ_t satisfying $\psi_t(x_0) = 0$, giving

$$\frac{1}{t}(\mathcal{T}_c(\sigma, \rho_0) - \mathcal{T}_c(\sigma, \rho_t)) \geq \int \psi_t d(\rho_1 - \rho_0).$$

Moreover, by uniqueness of (φ_0, ψ_0) , we get that φ_t, ψ_t converges uniformly to (φ_0, ψ_0) as $t \rightarrow 0$, thus concluding the proof. \square

The assumption on the uniqueness of the potentials can be guaranteed a priori in the following setting, which corresponds to the distance W_2 (one could prove it for W_p , with $p > 1$ similarly).

Proposition 4.2 (Uniqueness of potentials). *If $X \subseteq \mathbb{R}^d$ is the closure of a bounded and connected open set, $x_0 \in X$, $(\sigma, \rho) \in \mathcal{P}(X)$ satisfies*

$$\text{spt}(\rho) = X \text{ or } \text{spt}(\sigma) = X,$$

then, there exists a unique pair of Kantorovich potentials (φ, ψ) optimal for $c(x, y) = \frac{1}{2} \|x - y\|^2$, c -conjugate to each other, and satisfying $\varphi(x_0) = 0$.

Proof. Assume that $\text{spt}(\sigma) = X$. Since c is Lipschitz on the bounded set X , φ, ψ are Lipschitz and therefore differentiable almost everywhere. Take $(x_0, y_0) \in \text{spt}(\gamma)$ where $\gamma \in \Pi(\sigma, \rho)$ is the optimal transport plan, such that φ is differentiable at $x_0 \in X$. As we have already shown, for any optimal pair (φ, ψ) we necessarily have

$$y_0 = x_0 - \nabla \varphi(x_0),$$

so that if (φ', ψ') is another optimal pair, we should have $\nabla \varphi = \nabla \varphi'$ σ -a.e. Since $\text{spt}(\sigma) = X$ and since X is the closure of a connected open set, this implies $\varphi = \varphi' + C$ for a constant C as desired, and $C = 0$ since $\varphi(x_0) = \varphi'(x_0)$. Moreover, $\psi' = \varphi'^c = \varphi^c = \psi$, allowing to deal with the case where $\text{spt}(\rho) = X$ by symmetry. \square

5 Dynamic formulation of optimal transport

We conclude this lecture with a discussion around a fluid dynamic interpretation of optimal transport. The material in this section is only treated at an informal level and we refer to [3] for a rigorous treatment.

When $X \subset \mathbb{R}^d$, we can interpret the marginals $\mu, \nu \in \mathcal{P}(X)$ as distributions of particles at times $t = 0$ and $t = 1$ respectively. Assume that for each time t , there is a velocity field $v_t : \mathbb{R}^d \rightarrow \mathbb{R}^d$ which moves particles around. The relation between the velocity field and the distribution is given by the continuity equation (satisfied in the sense of distributions)

$$\partial_t \rho_t + \nabla \cdot (\rho_t v_t) = 0.$$

When v_t is regular enough (e.g. Lipschitz continuous in x , uniformly in t), then we can define its flow $T : [0, 1] \times X \rightarrow \mathbb{R}^d$ which is such that $T_t(x)$ gives the position at time t of a particle which is at x at time 0. It solves $T_0(x) = x$ and

$$\frac{d}{dt}T_t(x) = v_t(T_t(x)).$$

Let us denote $\text{CE}(\mu, \nu)$ the set of solutions (ρ, v) to the continuity equation such that $t \mapsto \rho_t$ is weakly continuous and satisfies $\rho_0 = \mu$ and $\rho_1 = \nu$. Consider also the integrated (generalized) “kinetic energy” functional

$$A_p(\rho, v) := \int_0^1 \int_{\mathbb{R}^d} \|v_t(x)\|^p d\mu_t(x) dt.$$

Among all interpolations between μ and ν , it turns out that optimal transport with cost $\|y - x\|^p$ is the one that minimizes A_p . This is called the Benamou-Brenier formulation.

Theorem 5.1 (Dynamic formulation). *Let $\mu, \nu \in \mathcal{P}(\mathbb{R}^d)$ be compactly supported. For $p \geq 1$ it holds*

$$W_p^p(\mu, \nu) = \inf \left\{ A_p(\rho, v) \mid (\rho, v) \in \text{CE}(\mu, \nu) \right\}.$$

Let us give some informal arguments to understand this result.

- Let us first argue that for $(\rho, v) \in \text{CE}(\mu, \nu)$ it holds $A_p(\rho, v) \geq W_p^p(\mu, \nu)$. Assume (ρ, v) is regular enough and consider the flow $T_t(x)$, that satisfies $\rho_t = (T_t)_\# \rho_0$. It holds

$$\begin{aligned} A(\rho, v) &= \int_0^1 \int_{\mathbb{R}^d} \|v_t(T_t(x))\|^p d\rho_0(x) dt \\ &= \int_{\mathbb{R}^d} \left(\int_0^1 \left\| \frac{d}{dt} T_t(x) \right\|^p dt \right) d\rho_0(x) \\ &\geq \int_{\mathbb{R}^d} \|T_1(x) - T_0(x)\|^p d\rho_0(x) \end{aligned}$$

by Jensen’s inequality. Since $(T_1)_\# \rho_0 = \rho_1 = \nu$ and $\rho_0 = \mu$, the last quantity is larger than $W_p^p(\mu, \nu)$.

- Let us build an admissible $(\rho, v) \in \text{CE}(\mu, \nu)$ such that $A(\rho, v) = W_p^p(\mu, \nu)$ using the geodesic between μ and ν . Assume that there exists an optimal transport map T between μ and ν , and set $\rho_t = (T_t)_\# \mu$ with $T_t(x) = (1 - t)x + tT(x)$. Now define the velocity field

$$v_t = \left(\frac{d}{dt} T_t \right) \circ T_t^{-1} = (T - \text{id}) \circ T_t^{-1},$$

which, by construction, is such that (ρ_t, v_t) satisfies the continuity equation in the weak sense. We have the desired equality:

$$A(\rho, v) = \int \|v_t(x)\|^p d\rho_t(x) = \int |T(x) - x|^p d\rho_0(x) = W_p^p(\mu, \nu).$$

Riemannian interpretation. In the case $p = 2$, we can understand (at least at the formal level) the Benamou-Brenier formula as a Riemannian formulation for W_2 (this point of view is due to Otto). In this interpretation, the tangent space at $\rho \in \mathcal{P}_2(X)$ are measures of the form $\delta\rho = -\nabla \cdot (v\rho)$ with a velocity field $v \in L^2(\rho, \mathbb{R}^d)$ and the metric is given by

$$\|\delta\rho\|_\rho^2 = \inf_{v \in L^2(\rho, \mathbb{R}^d)} \left\{ \int \|v(x)\|_2^2 d\rho(x) \mid \delta\rho = -\nabla \cdot (v\rho) \right\}.$$

References

- [1] Luigi Ambrosio, Nicola Gigli, and Giuseppe Savaré, *Gradient flows: in metric spaces and in the space of probability measures*, Springer Science & Business Media, 2008.
- [2] Haïm Brezis, *Analyse fonctionnelle*, Masson, Halsted Press, 1983.
- [3] Filippo Santambrogio, *Optimal transport for applied mathematicians*, Springer, 2015.
- [4] Cédric Villani, *Optimal transport: old and new*, vol. 338, Springer Science & Business Media, 2008.