



# Action Potential Propagation



MNS CP Project WS '19/20  
with Roopa, Andrey, and Luke



# Outline

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- Introduction to the cable equation
- Method of solving the partial differential equation
- Usage with the Hodgkin-Huxley model
- Initiation of an action potential
- Propagation in a myelinated axon

# Motivation

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1. Create a model of an axon with various biological parameters that allow us to simulate action potential conduction.
2. Understand the impact of myelination on the spiking characteristics, as well as propagation velocity.
3. Illustrate what affect multiple action potentials within the same axon simultaneously has.

# The Cable Equation

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# The Cable Equation

$$c_m \frac{\partial V}{\partial t} = \frac{1}{2ar_L} \frac{\partial}{\partial x} \left( a^2 \frac{\partial V}{\partial x} \right) - i_m + i_e$$

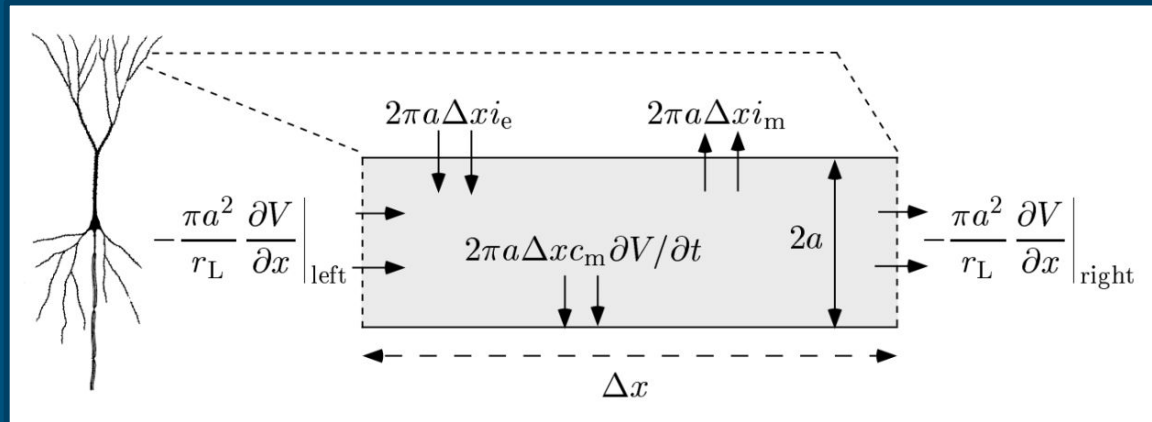


Figure 6.6 Dayan & Abbott, 2001

# The Cable Equation

$$C_m \frac{\partial V}{\partial t} = \frac{1}{2ar_L} \frac{\partial}{\partial x} \left( a^2 \frac{\partial V}{\partial x} \right) - i_m + i_e$$

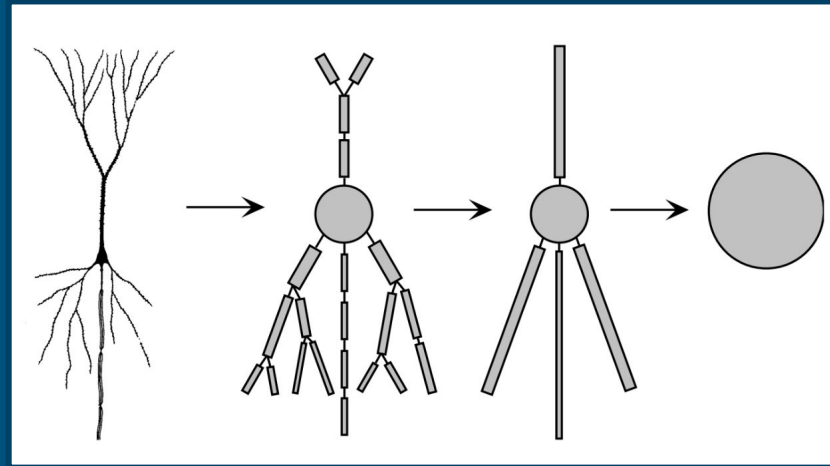


Figure 6.15 Dayan & Abbott, 2001

# The Cable Equation

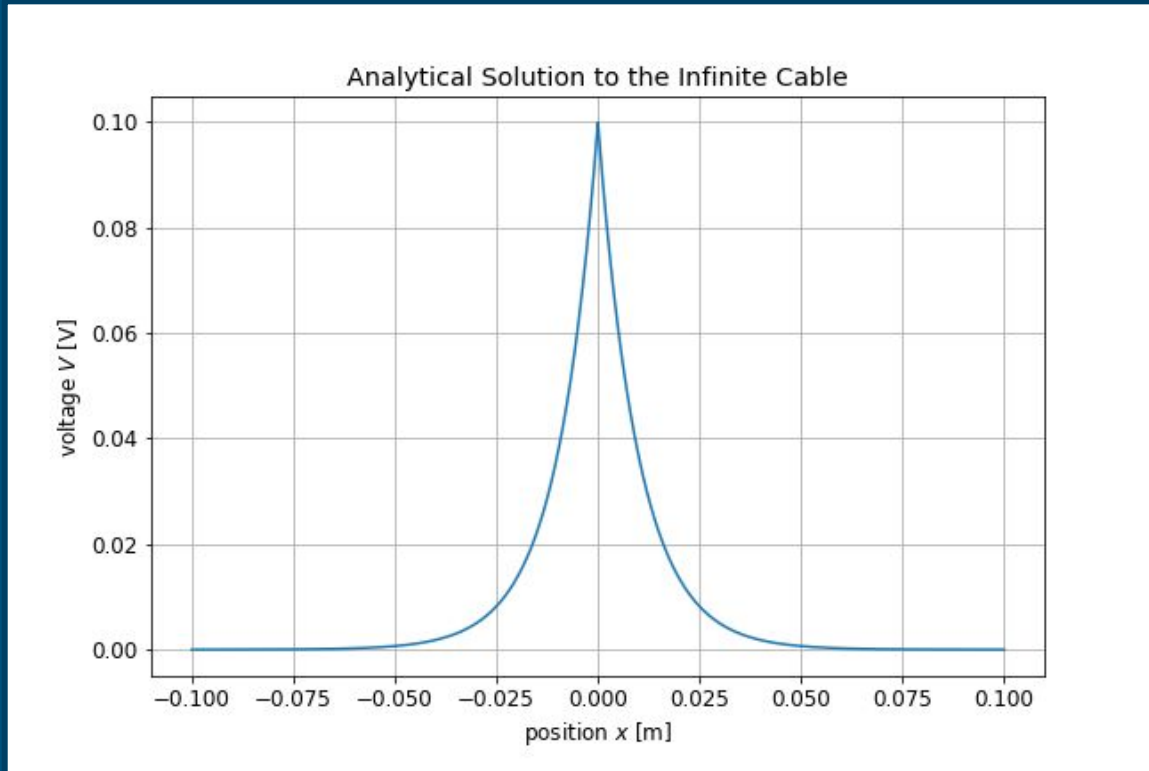
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$$c_m \frac{\partial V}{\partial t} = \frac{1}{2ar_L} \frac{\partial}{\partial x} \left( a^2 \frac{\partial V}{\partial x} \right) - i_m + i_e$$



Assumptions:

- Constant axon radius  $a$
- Identical compartment size
- No branching



The Analytical Solution to the Infinite Cable



# Cable equation reduces to ODE in static case

$$c_m \frac{\partial v}{\partial t} = \frac{1}{2ar_L} \frac{\partial}{\partial x} \left( a^2 \frac{\partial v}{\partial x} \right) - i_m + i_e$$

Current is steady, implying no change over time, therefore:  $\frac{\partial v}{\partial t} = 0$

$$0 = \frac{1}{2ar_L} \frac{\partial}{\partial x} \left( a^2 \frac{\partial v}{\partial x} \right) - i_m + i_e$$

Rearranging terms: 
$$\frac{dv}{dx^2} = \frac{2r_L}{a} (i_m - i_e)$$

Substituting: 
$$i_m = \frac{V - V_{rest}}{r_m} = \frac{v}{r_m}$$

$$\frac{dv}{dx^2} = \frac{2r_L}{ar_m} (v - i_e r_m) \Rightarrow \lambda^2 \frac{dv}{dx^2} = v - i_e r_m$$

[  $i_e = 0$  everywhere except the site of injection ( $x=0$ ) ]

$$\lambda^2 \frac{\partial v}{\partial x^2} = v$$

$$v(x) = B_1 e^{(-x/\lambda)} + B_2 e^{(x/\lambda)}$$

At ends of cable,  $B = 0$ . At  $x=0$ , as we want the solution to be continuous,  $B_1=B_2=B$ , and the solution:  $v(x) = B e^{(-|x|/\lambda)}$

To find  $B$ : current injected = current diffusing away from  $x$

In the region around  $x$ :  $\frac{dv}{dx} = -\frac{B}{\lambda} e^{(-|x|/\lambda)}$ , where  $e^{(-|x|/\lambda)}$  is approx. = 1, giving  $-B/\lambda$ . (Other side is  $B/\lambda$ )

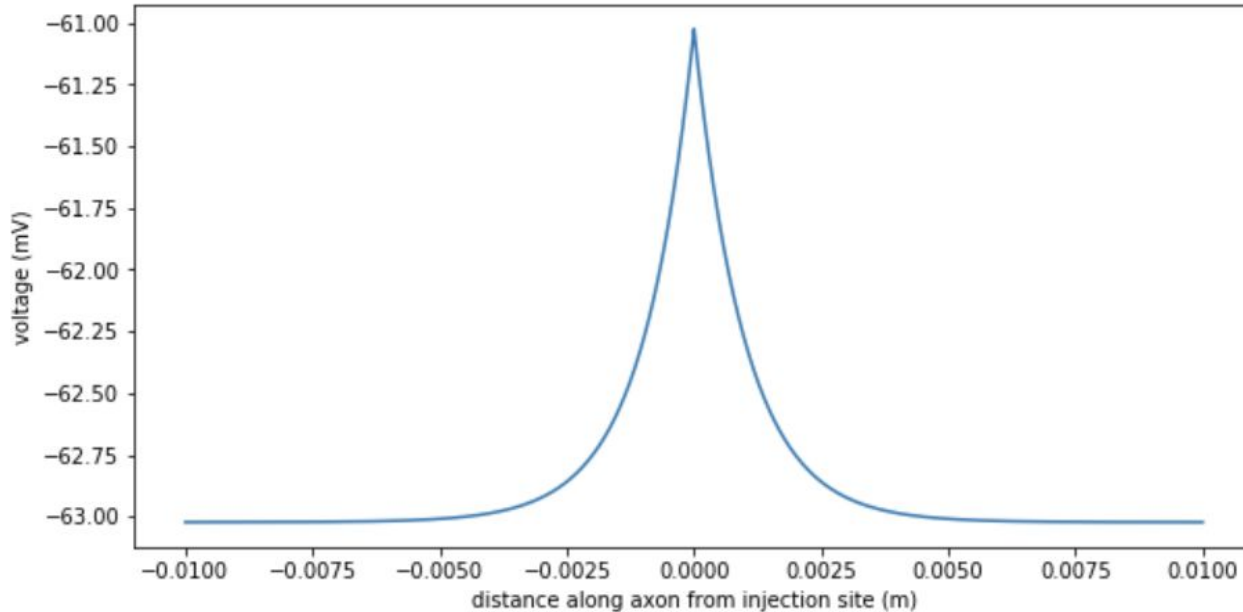
$$\text{Second derivative: } \frac{d^2 v}{dx^2} = -\frac{2B}{\lambda \Delta x}$$

Substituting in cable equation:  $-2\lambda^2 B/\lambda \Delta x = -r_m I_e/2\pi a x$ , which gives:  $B = I_e R_\lambda/2$

$$\text{Therefore: } v(x) = \frac{I_e R_\lambda}{2} e^{(-|x|/\lambda)}$$

# Numerical solution:

Membrane voltage along axon (steady current injection in middle of axon)



# Numerical Solution of PDEs

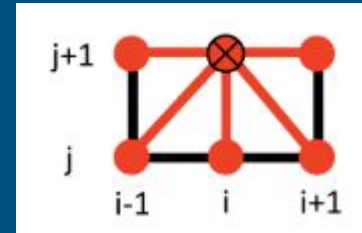
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# Crank-Nicolson Method: Theory

- For problems which involve an update in both time and space, to increase accuracy, "implicit" methods of numerical analysis are preferred, e.g. the Crank-Nicolson scheme.
- time derivative evaluated with the "forward differencing" approach, but spatial derivative evaluated at a time in between the current and next time point
- get a system of linear equations to be solved simultaneously

$$\frac{\partial v}{\partial t} = \frac{\partial v}{\partial x^2}$$

$$\begin{aligned} \frac{v_{i,j+1} - v_{i,j}}{\Delta t} &= \frac{v_{i+1,j+1/2} - 2v_{i,j+1/2} + v_{i-1,j+1/2}}{\Delta x^2} \\ &= \left( \frac{v_{i+1,j+1} - 2v_{i,j+1} + v_{i-1,j+1}}{\Delta x^2} + \frac{v_{i+1,j} - 2v_{i,j} + v_{i-1,j}}{\Delta x^2} \right) \end{aligned}$$



# Crank-Nicolson Method: Application

Coupled equations are solved as follows:

1. starting from compartment #1, we get the voltage,  $\Delta V_\mu$ , in terms of  $\Delta V_{\mu+1}$ .
2. At the last compartment, using the boundary condition, we obtain  $\Delta V_N$ , and substitute it in the previous equation, and continue solving backwards until compartment #1.

$$c_m \frac{dV_\mu}{dt} = -i_m^\mu + \frac{I_e^\mu}{A_\mu} + g_{\mu,\mu+1}(V_{\mu+1} - V_\mu) + g_{\mu,\mu-1}(V_{\mu-1} - V_\mu)$$

$$\frac{dV_\mu}{dt} = A_\mu V_{\mu-1} + B_\mu V_\mu + C_\mu V_{\mu+1} + D_\mu$$

$$A_\mu = \frac{g_{\mu,\mu-1}}{c_m} \quad B_\mu = \frac{\sum_i g_i^\mu + g_{\mu,\mu+1} + g_{\mu,\mu-1}}{c_m}$$

$$C_\mu = \frac{g_{\mu,\mu+1}}{c_m} \quad D_\mu = \frac{\sum_i g_i^\mu E_i + \frac{I_e^\mu}{A_\mu}}{c_m}$$

$$\Delta V_\mu = (A_\mu V_{\mu-1}(t + z\Delta t) + B_\mu V_\mu(t + z\Delta t) + C_\mu V_{\mu+1}(t + z\Delta t) + D_\mu)\Delta t$$





$$\Delta V_{\mu} = (A_{\mu} V_{\mu-1}(t + z\Delta t) + B_{\mu} V_{\mu}(t + z\Delta t) + C_{\mu} V_{\mu+1}(t + z\Delta t) + D_{\mu}) \Delta t$$

$$V_{\mu}(t + z\Delta t) = V_{\mu}(t) + z\Delta V_{\mu}$$

$$a_{\mu} = A_{\mu} z \Delta t$$

$$\Delta V_{\mu} = a_{\mu} \Delta V_{\mu-1} + b_{\mu} \Delta V_{\mu} + c_{\mu} \Delta V_{\mu+1} + d_{\mu}$$

$$b_{\mu} = B_{\mu} z \Delta t$$

$$c_{\mu} = C_{\mu} z \Delta t$$

$$d_{\mu} = (D_{\mu} + A_{\mu} V_{\mu-1}(t) + B_{\mu} V_{\mu}(t) + C_{\mu} V_{\mu+1}(t)) \Delta t$$

Solve  $\Delta V_{\mu}$  equation sequentially for each compartment, in terms of  $\Delta V_{\mu+1}$

$$\Delta V_{\mu-1} = \frac{c_{\mu-1} \Delta V_{\mu} + d'_{\mu-1}}{1 - b'_{\mu-1}}$$

At end of cable use boundary condition to find  $\Delta V_N$ , and substitute backwards to get all  $\Delta V_{\mu}$

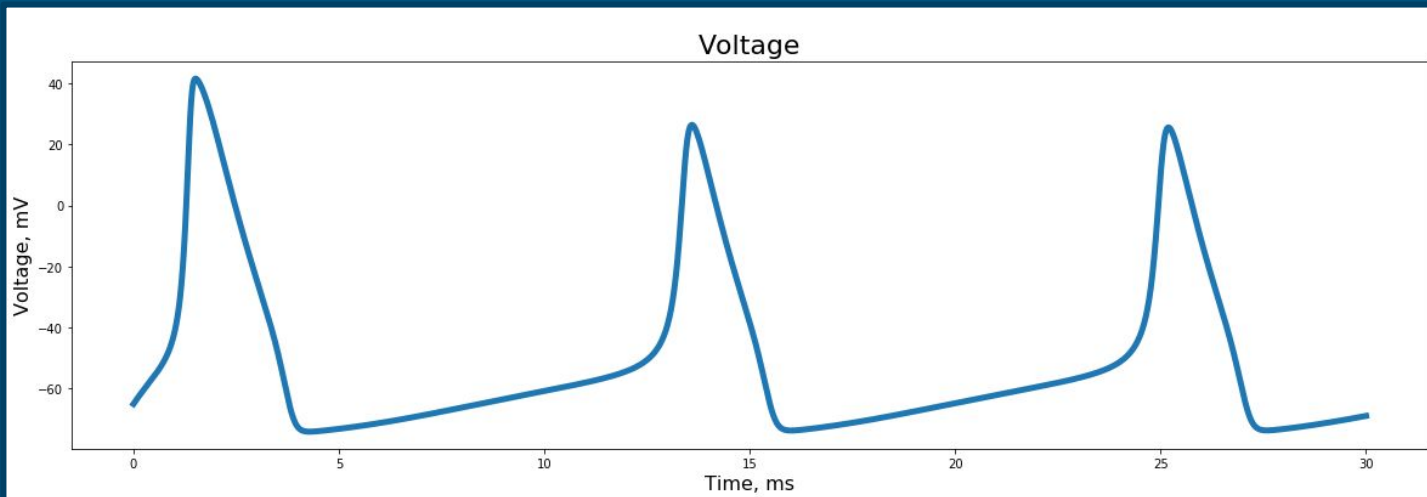
# Hodgkin-Huxley model

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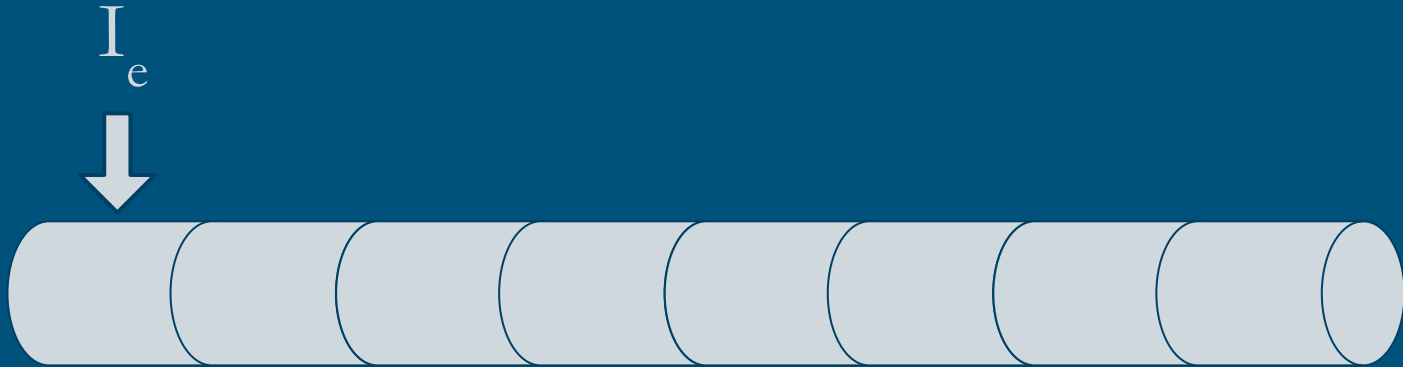
# The Hodgkin-Huxley Model

$$i_m = \bar{g}_L(V - E_L) + \bar{g}_{Na}m^3h(V - E_{Na}) + \bar{g}_Kn^4(V - E_K)$$
$$\frac{dx}{dt} = \alpha_x(V)(1 - x) - \beta_x(V)x$$

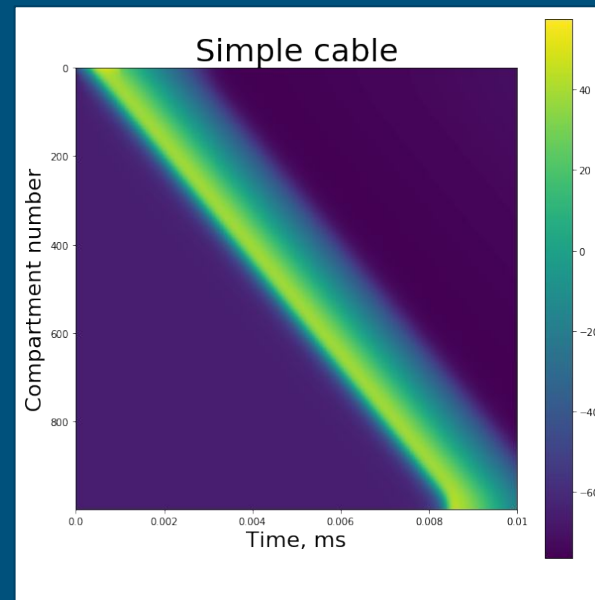
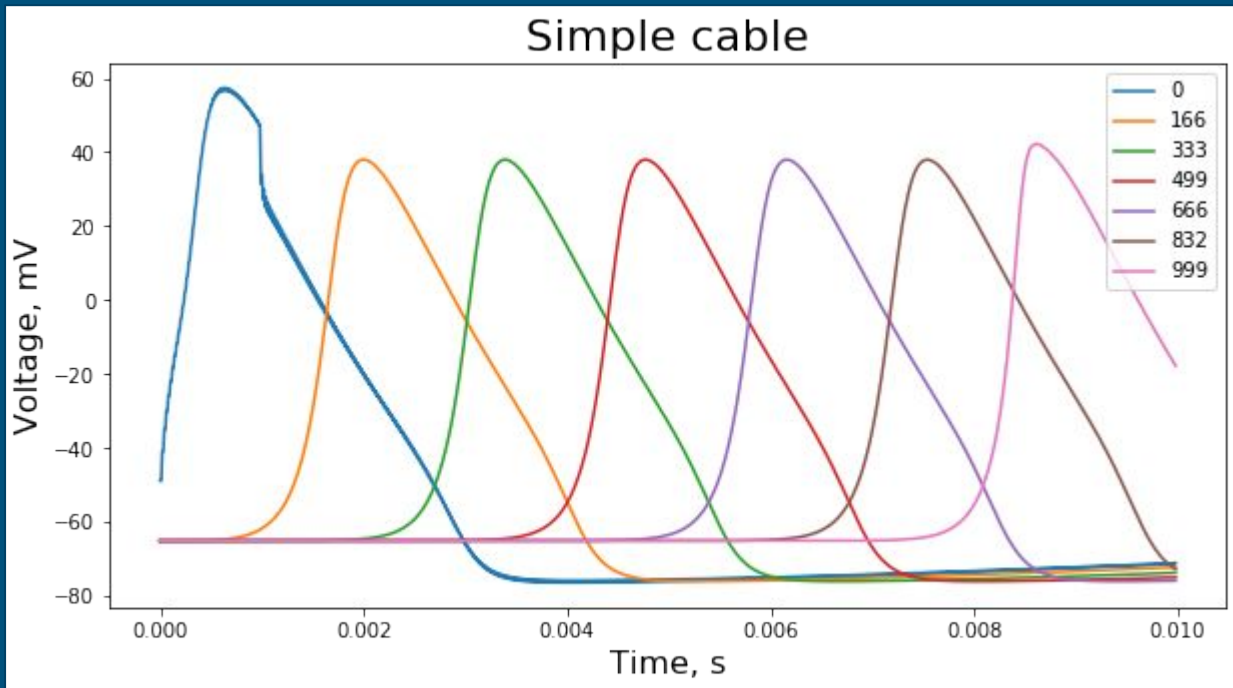


# Initiating an action potential

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# Initiating an action potential



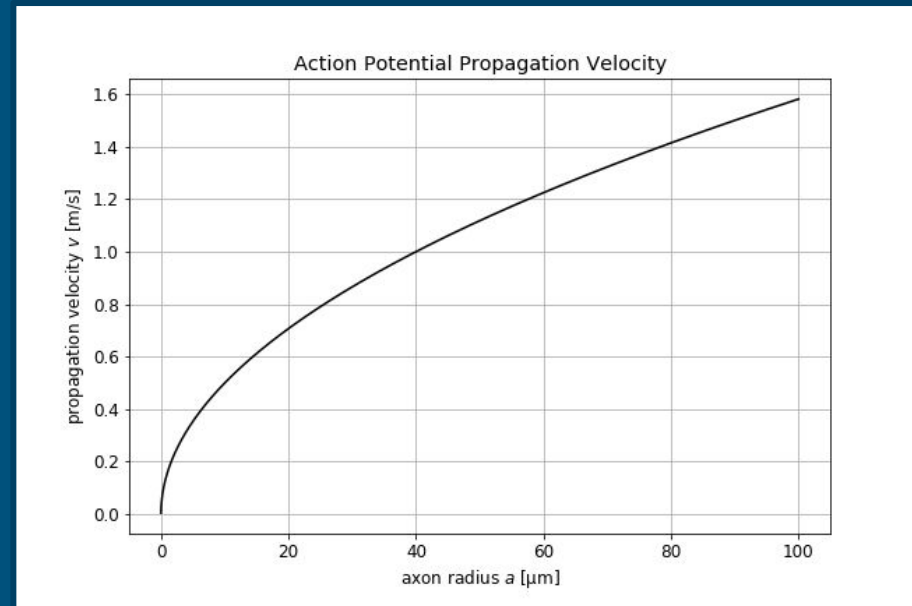
$$v \equiv \frac{x}{t_{max}} = \frac{2\lambda}{\tau_m}$$

$$\tau_m = r_m c_m$$

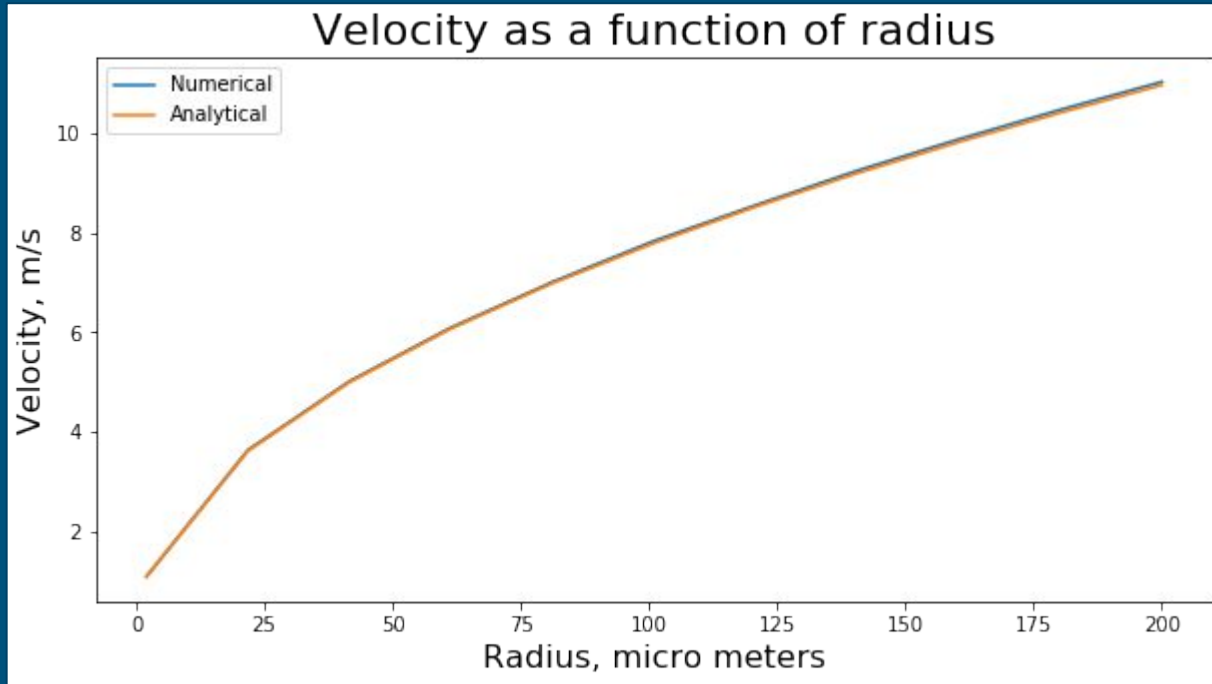
$$\lambda = \sqrt{\frac{r_m a}{2r_L}}$$

$$v = \frac{2}{r_m c_m} \sqrt{\frac{r_m a}{2r_L}} = \frac{1}{c_m} \sqrt{\frac{2a}{r_m r_L}}$$

$$v \propto \sqrt{a}$$



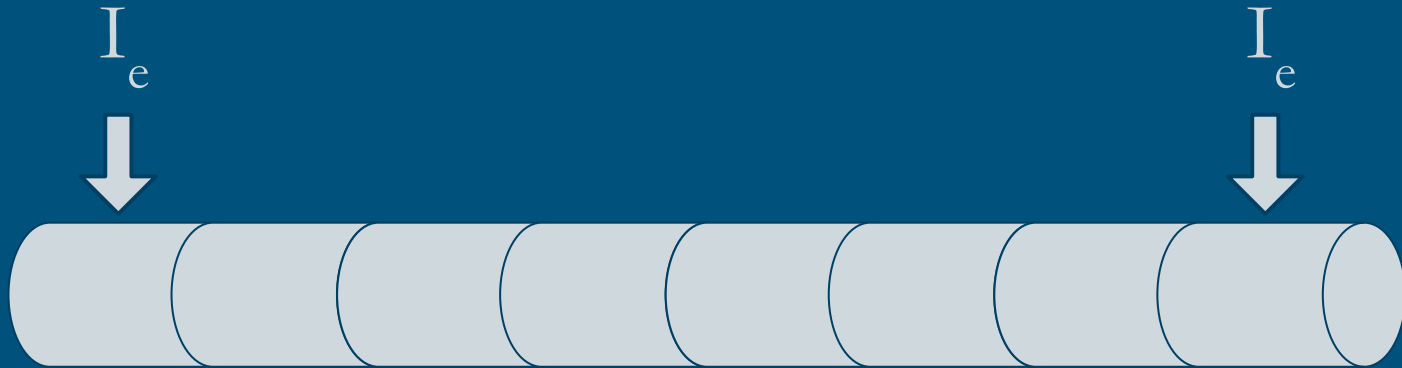
Action Potential Propagation Velocity - Analytically

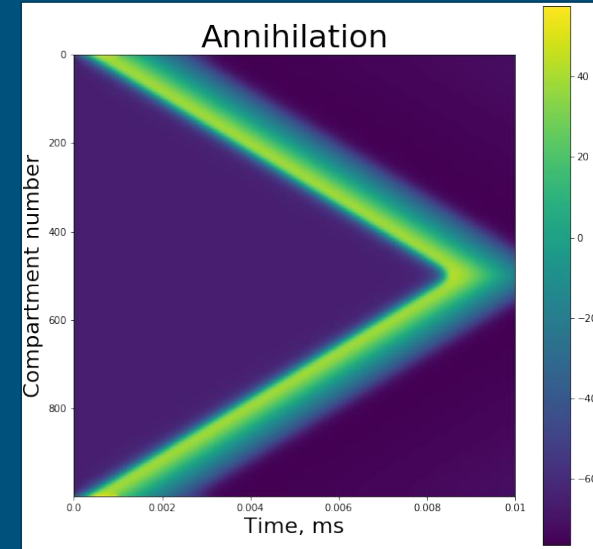
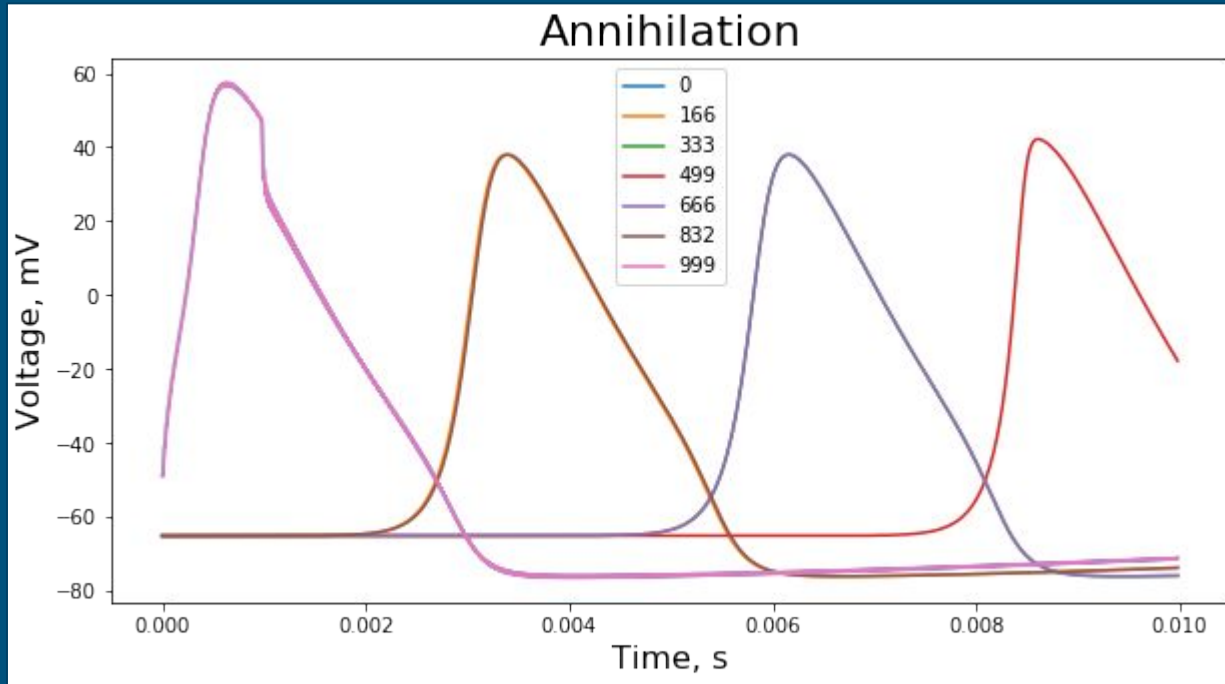


Action Potential Propagation Velocity - Numerically

# Annihilation of Two Action Potentials

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Annihilation of Two Action Potentials

# Myelinated Axons

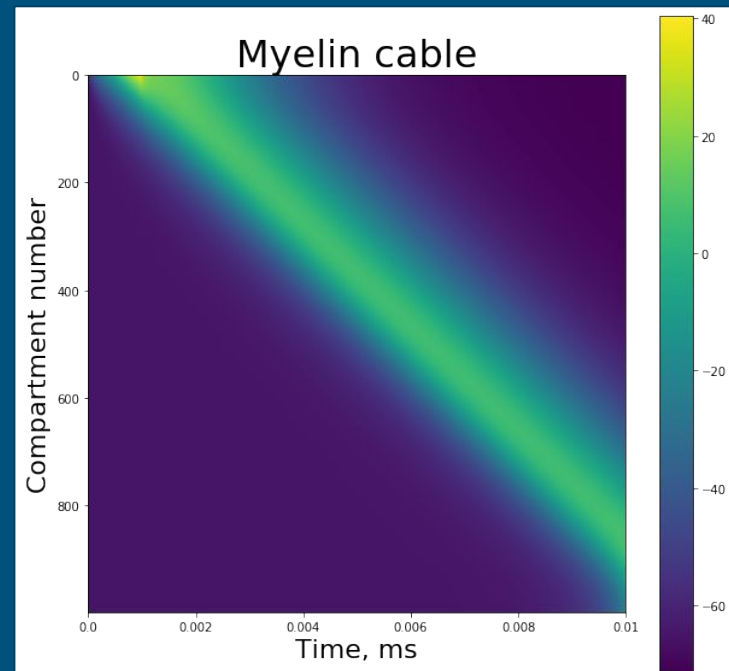
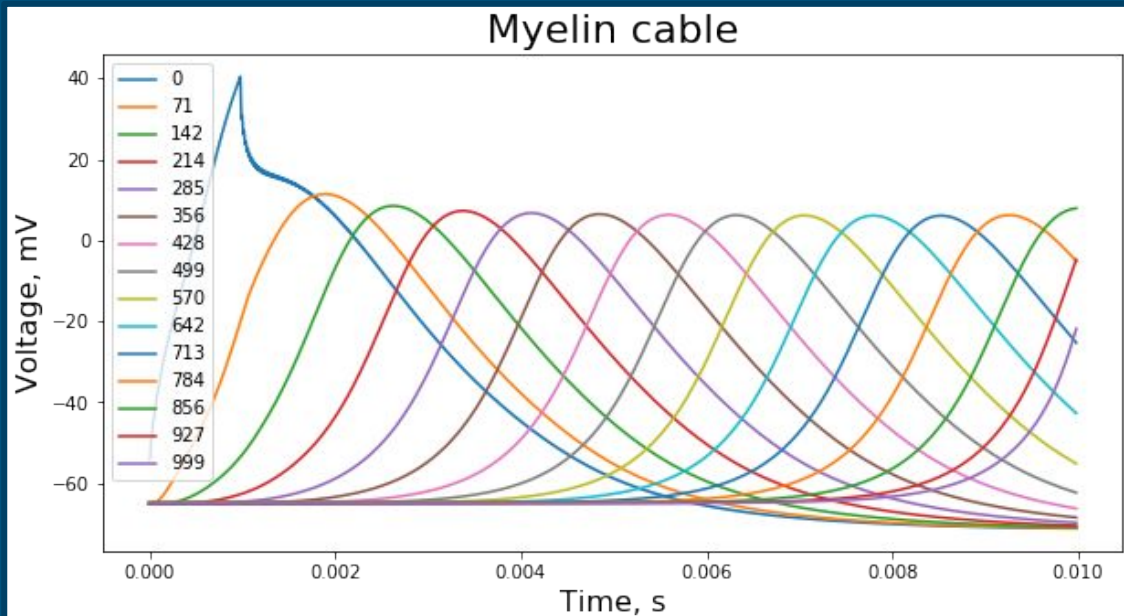


Impact of myelin:

- Increases resistance  $r_m$  by a factor of 5,000
- Decreases capacitance  $c_m$  by 50



# Myelinated Axons



# Branching Cable

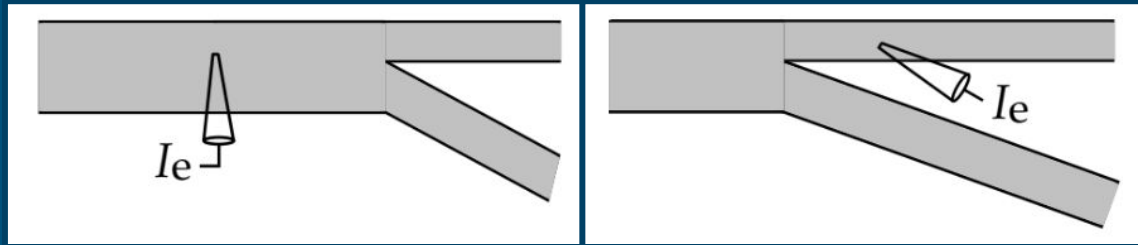
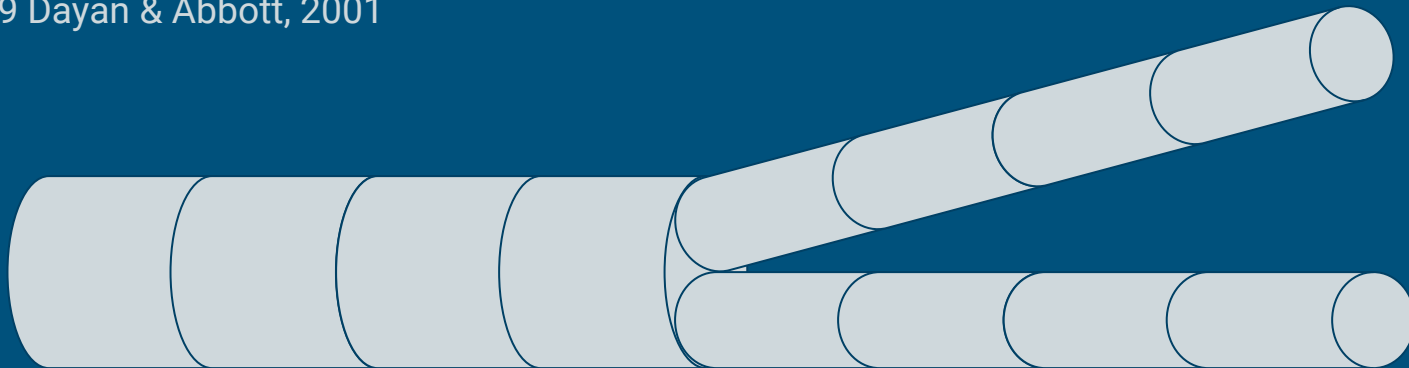


Figure 6.9 Dayan & Abbott, 2001

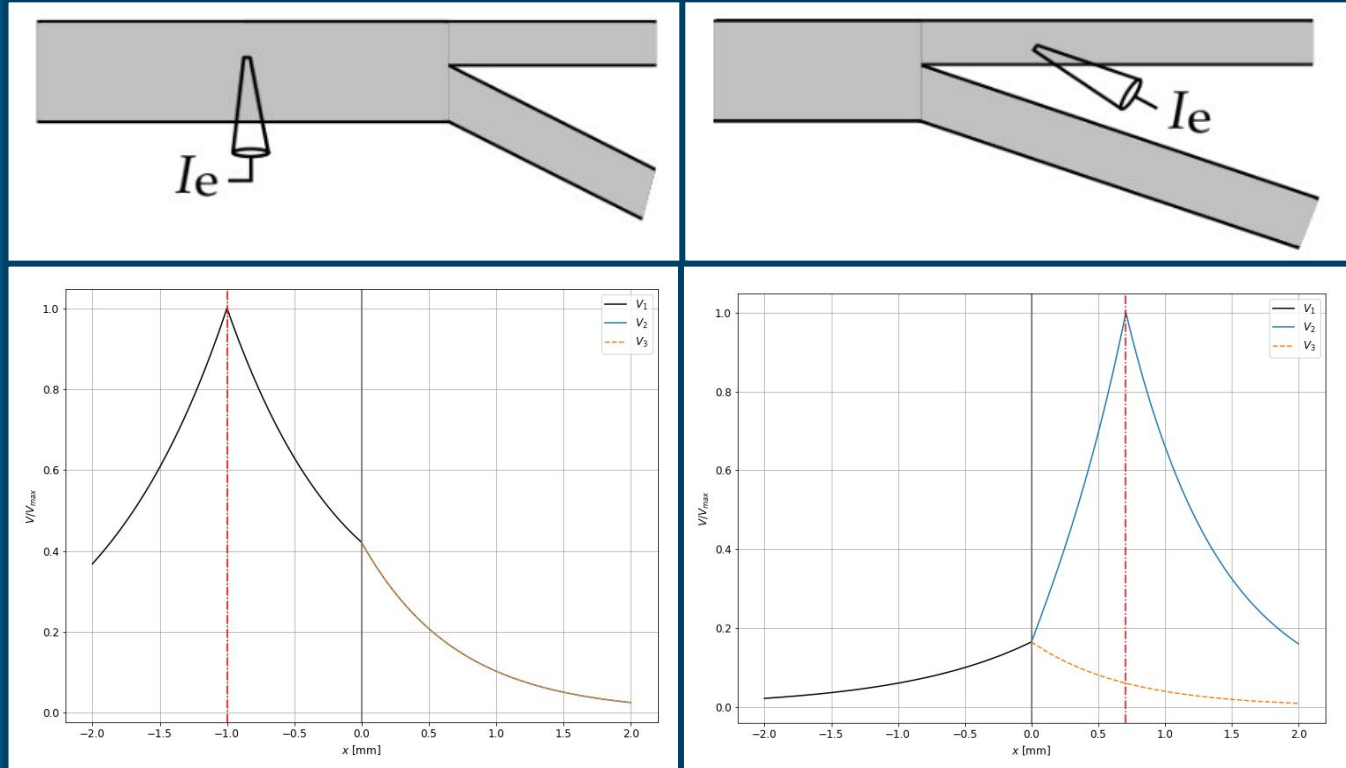


# Branching Cable

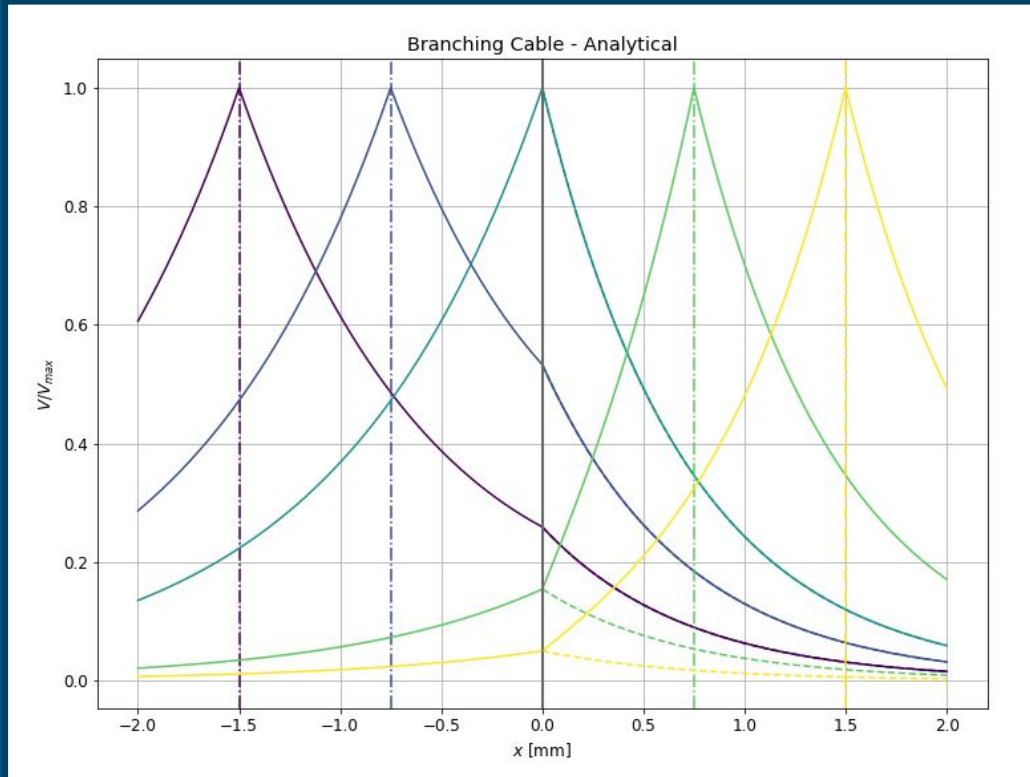
$$\begin{cases} v_1(x) = p_1 I_e R_{\lambda_1} \exp\left(\frac{-x_1}{\lambda_1} - \frac{y}{\lambda_2}\right) \\ v_2(x) = \frac{I_e R_{\lambda_2}}{2} \left[ \exp\left(\frac{-|y-x_2|}{\lambda_2}\right) + (2p_2 - 1) \exp\left(\frac{-(y+x_2)}{\lambda_2}\right) \right] \\ v_3(x) = p_3 I_e R_{\lambda_3} \exp\left(\frac{-x_3}{\lambda_3} - \frac{y}{\lambda_2}\right) \end{cases}$$

$$p_i = \frac{a_i^{3/2}}{a_1^{3/2} + a_2^{3/2} + a_3^{3/2}} \quad \lambda_i = \sqrt{\frac{a_i r_m}{2r_L}} \quad R_{\lambda_i} = \frac{r_L \lambda_i}{\pi a_i^2}$$

Figure 6.9 Dayan & Abbott, 2001



## Branching Cable - Analytical



## Branching Cable - Analytical

# Questions?