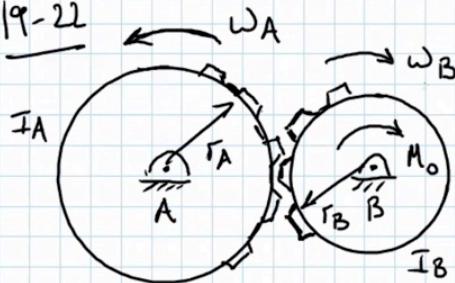
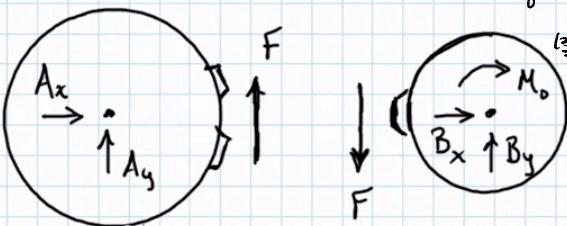


Dynamics Lecture #21

19-22



Free-body diagrams



Homework 9 : 19-31, 19-33, 19-46
Read Ch. 22 Due W 4/22/2020

System initially at rest.

Constant moment M_0 applied at $t = 0$.
Find ω_B and ω_A for time $t > 0$.

$$\begin{aligned} \text{(1)} \quad & \int_0^t M_0 dt - \int_0^t F_{r_A} dt = I_B \omega_B(t) - I_B \omega_B(0) \\ \text{(2)} \quad & \int_0^t F_{r_A} dt = I_A \omega_A(t) - I_A \omega_A(0) \end{aligned}$$

$$\stackrel{(3)}{=} \text{kinematic equation: } \omega_A r_A = \omega_B r_B$$

Rewrite (1) and (2) : take (1) and divide by r_B

$$\frac{1}{r_B} M_0 t - \int_0^t F dt = \frac{1}{r_B} I_B \omega_B(t) \quad \text{(4)}$$

$$\int_0^t F dt = \frac{1}{r_B} I_B \omega_B(t) \quad \text{(5)}$$

Add (4) and (5)

$$\frac{1}{r_B} M_0 t = \frac{1}{r_B} I_B \omega_B + \frac{1}{r_B} I_B \omega_B \quad \stackrel{(6)}{\dots} \text{canceling out } t \text{ since } \omega \text{ is function of time}$$

→ solve using (6) and (3)

$$\frac{1}{r_B} M_0 t = \frac{1}{r_B} I_B \omega_B + \frac{1}{r_B} I_B \left(\frac{r_B}{r_A} \omega_A \right) \quad \dots \text{multiply throughout by } r_B$$

$$M_0 t = I_B \omega_B + \frac{r_B}{r_A} I_B \frac{r_B}{r_A} \omega_B$$

$$\boxed{\omega_B = \frac{M_0 t}{I_B + \left(\frac{r_B}{r_A}\right)^2 I_B}}$$

→ use this to solve $\omega_A = \frac{r_B}{r_A} \frac{M_0 t}{I_B + \left(\frac{r_B}{r_A}\right)^2 I_B}$

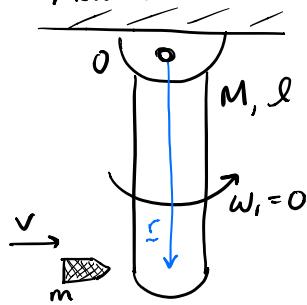
If $I_A = 0$ (gear A not present) then just have motor driving gear B where $\omega_B = M_0 t / I_B$

so when A is present,
the denominator changes

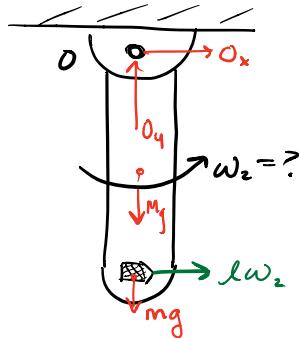
$$I_0 + \left(\frac{r_A}{r_K}\right)^2 I_K = \text{the effective moment of inertia of the two gears, as seen by the motor.}$$

Collision problem w/rigid body
→ collision between pendulum (which we'll model as thin bar) and small projectile/bullet

we'll assume projectile becomes embedded into pendulum



just before collision



just after collision

During collision process, no external moments about O.

$$\therefore (\text{angular momentum})_1 = (\text{angular momentum})_2$$

$$l m v = I_0 \omega_2 + l(m l \omega_2) = (I_0 + m l^2) \omega_2$$

$$\omega_2 = \frac{l m v}{I_0 + m l^2}$$

Is energy conserved in collision?

... analogous to plastic collision, energy is not conserved

$$\begin{array}{l} \text{energy of} \\ \text{system 1} \\ (\text{at particle}) \end{array} : E_1 = \frac{1}{2} m v^2$$



$$E_2 < E_1$$

\therefore energy is not conserved

$$\begin{array}{l} \text{energy} \\ \text{after} \\ \text{collision} \end{array} : E_2 = \frac{1}{2} I_0 \omega_2^2 + \frac{1}{2} m(l \omega_2)^2$$

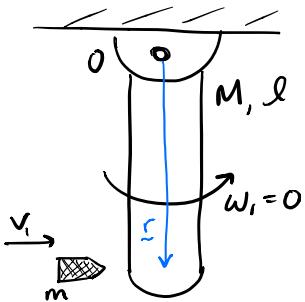
$$= \frac{1}{2} (I_0 + m l^2) \omega_2^2$$

$$= \frac{1}{2} (I_0 + m l^2) \frac{l^2 m^2 v^2}{(I_0 + m l^2)^2}$$

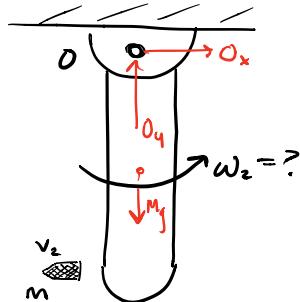
$$= \frac{1}{2} \frac{l^2 m^2 v^2}{(I_0 + m l^2)} = \frac{1}{2} \frac{m l^2}{I_0 + m l^2} m v^2$$

$$E_2 = \frac{1}{2} \left(\frac{m l^2}{I_0 + m l^2} \right) m v^2$$

Now, let's see what happens when we assume energy is conserved ... elastic collision



just before collision



just after collision

$$\text{Angular momentum about } O: I_m v_1 = I_o \omega_2 + I_m v_2 \quad (1)$$

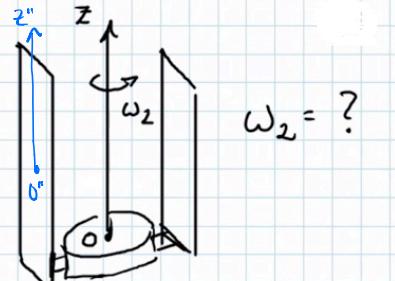
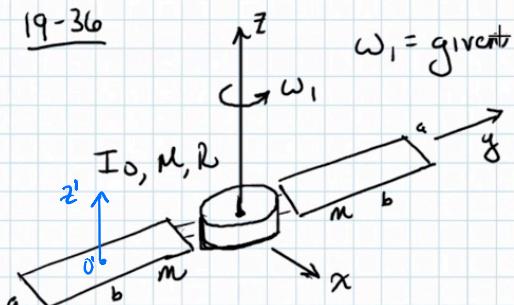
$$\text{energy equation: } \cancel{\frac{1}{2} m v_1^2} = \cancel{\frac{1}{2}} I_o \omega_2^2 + \cancel{\frac{1}{2}} m v_2^2 \quad (2)$$

→ two equations for ω_2 and v_2

19-36 SIMILAR TO HW PROBLEM

satellite w/ a cylinder body, attached to cylinder are two solar panels

19-36



angular velocity is expected to ↑ when panels fold in (like figure skater going faster when fold arms in)

no external moments about O : $(\text{Ang. mom})_1 = (\text{Ang. mom})_2$

$$I_o \omega_1 + 2 \left[I_{o,1} + m \left(R + \frac{b}{2} \right)^2 \right] \omega_1 = I_o \omega_2 + 2 \left[I_{o,2} + m R^2 \right] \omega_2$$

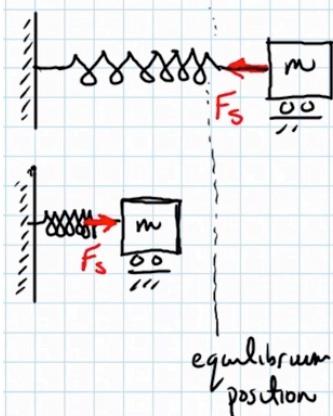
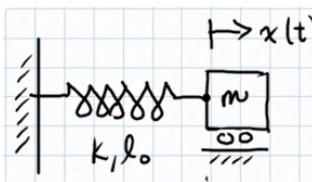
for plates, have to use parallel axis theorem

have to use parallel axis th. for plates

$$\left\{ I_o + 2 \left[\frac{1}{12} m (a^2 + b^2) + m \left(R + \frac{b}{2} \right)^2 \right] \right\} \omega_1 = \left\{ I_o + 2 \left[\frac{1}{12} m a^2 + m R^2 \right] \right\} \omega_2$$

this is smaller than ... so $\omega_2 > \omega_1$

CH. 22 - VIBRATION THEORY

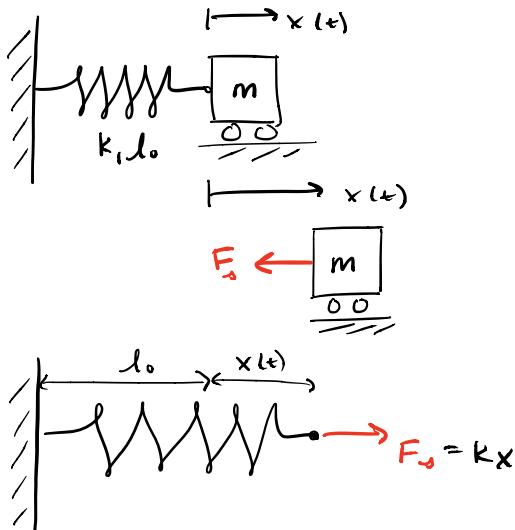


Ch. 22 : Mechanical Vibrations (6)

Vibration: oscillatory (back-and-forth) motion about an equilibrium position.

If mass is disturbed from its equilibrium position, spring exerts a restoring force that pushes mass back toward its equilibrium position.

next, derive the equation of motion for displacement $x(t)$ of mass



For mass:

$$m\ddot{x} = -F_s$$

$$m\ddot{x} = -kx$$

$$m\ddot{x} + kx = 0 \quad (1)$$

\uparrow equation of motion for displacement $x(t)$ of mass

solution of (1) is

$$x(t) = A \sin \omega_n t + B \cos \omega_n t \quad \text{where } \omega_n = \sqrt{\frac{k}{m}}$$

\rightarrow 2nd order, linear DE

ω_n = natural frequency of the spring mass system

can also write: $x(t) = C \sin(\omega_n t + \phi)$

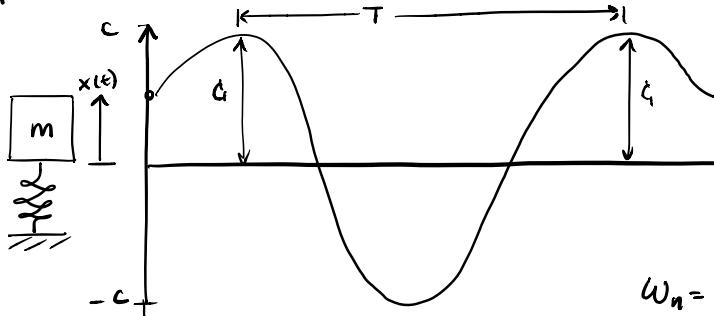
$$\text{where } C = \sqrt{A^2 + B^2}$$

$$\phi = \tan^{-1}(B/A)$$

\uparrow phase angle

$A \pm B$, or C and ϕ are determined by initial conditions

plot $x(t) = C \sin(\omega_n t + \phi)$ vs. t



T = period of oscillation
 C = amplitude of oscillation

$$T = \frac{2\pi}{\omega_n}$$

ω_n = natural frequency (or radian frequency)
... units are radians/second

$$f_n = \frac{1}{T} \quad (1)$$

f_n = cyclic frequency
... units are cycles/second or Hz

$$\omega_n = \frac{2\pi}{T} \quad (2)$$

$$(1) \div (2) \Rightarrow \omega_n = 2\pi f_n$$

↑ rad/sec ↑ Hz

READ VIBRATIONS CHAPTER