

First Mandatory Assignment

MAT-MEK4270

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Problem 1

Consider the POISSON equation

$$\nabla^2 u = f, \quad \mathbf{x} \in \Omega, \quad (1.1)$$

where $\Omega = [0, L] \times [0, L]$. Use a uniform mesh with equal length in each direction, so that the mesh may be described by $\mathbf{x} \times \mathbf{y}$:

$$\mathbf{x} = [i\mathbf{h}]_{i=0}^N, \quad \mathbf{y} = [j\mathbf{h}]_{j=0}^N,$$

where $\mathbf{h} = L/N$. An outline for a script is found in `textttPoisson2D.py`.

a.

Modify the `Poisson2D` class such that `test_convergence_poisson2D` and `test_interpolation` pass.

Implementing the finite difference scheme for the second derivative at the inner points,

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{i+1,j} - 2u_{i,j} + u_{i-1,j}}{\Delta x^2}, \quad \mathbf{x} \in \Omega,$$

in both \mathbf{x} and \mathbf{y} , and a forward and backward EULER scheme at the boundaries, viz.

$$\frac{\partial^2 u}{\partial x^2} \approx \frac{u_{0,j} - 5u_{1,j} + 4u_{2,j} - u_{3,j}}{\Delta x^2}, \quad \mathbf{x} \in \partial\Omega,$$

on the left-most boundary, and similarly for the right-most boundary, only backwards. This can then be written in terms of a matrix, whose inverse acting on $f(\mathbf{x})$ should yield $u(\mathbf{x})$.

$$D2 = \begin{pmatrix} 1 & -5 & 4 & -1 & \cdots & 0 \\ 1 & -2 & 1 & 0 & \cdots & 0 \\ 0 & 1 & -2 & 1 & \cdots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & \cdots & 0 & 1 & -2 & 1 \\ 0 & \cdots & -1 & 4 & -5 & 1 \end{pmatrix}$$

The meshpoints themselves are made up of *linspace*s which are weaved into the ordinals and abscissae of a *meshgrid*.

Problem 2

Consider the wave equation

$$\partial_t^2 \eta = c^2 \nabla^2 \eta, \quad \mathbf{x} \in \Omega, \quad t \in \Theta, \quad (1.2)$$

where Ω is the unit square and $\Theta = [0, T]$. Consider also the COURANT number with the spatial and temporal discretizations $\mathbf{h} = L/N_x$ and $t_m/m = \Delta t = T/N_t$, respectively,

$$Co = \frac{c\Delta t}{\mathbf{h}}.$$

The wave equation (1.2) admits two real stationary waves as solutions. For the DIRICHLET and NEUMANN problems we have the imaginary and real parts respectively of

$$\eta(\mathbf{x}) = e^{i\mathbf{k} \cdot \mathbf{x}} \cos(\omega t),$$

where we restrict $\mathbf{k} = \mathbf{m}\pi$, where $\mathbf{m} = m_x \hat{\mathbf{i}} + m_y \hat{\mathbf{j}}$ for integer m_x, m_y . The numerical scheme for the wave equation is given by

$$\frac{\eta_{n,m}^{m+1} - 2\eta_{n,m}^m + \eta_{n,m}^{m-1}}{\Delta t^2} = c^2 \left(\frac{\eta_{n+1,m}^m - 2\eta_{n,m}^m + \eta_{n-1,m}^m}{\mathbf{h}^2} + \frac{\eta_{n,m+1}^m - 2\eta_{n,m}^m + \eta_{n,m-1}^m}{\mathbf{h}^2} \right) \quad (1.3)$$

a.

Implement the `Wave2D.py` class fully for a homogeneous DIRICHLET boundary condition. Compute also the l^2 error as follows.

$$E = \sqrt{h^2 \sum_{n,m} e_{nm}^m{}^2},$$

where e_{nm}^m is the difference between the numerical and analytical solutions.

b.

Implement the `Wave2D_Neumann` class for NEUMANN boundary conditions.

c.

Show that

$$\eta(\mathbf{x}, t) = e^{i(\mathbf{k} \cdot \mathbf{x} - \omega t)} \quad (1.6)$$

satisfies the wave equation (1.2)

We set $|\mathbf{k}| = k$. We then have that

$$\nabla^2 \eta = -k^2 \eta, \quad \partial_t^2 \eta = -\omega^2 \eta.$$

Using the standard definition of wave celerity, $c = \omega/k$, what we wanted to show is evident.

d.

Assuming isotropic k , we have that the discretization of (1.6) is

$$\eta_{nm}^m = e^{i(kh(m+n) - \omega n \Delta t)}, \quad (1.7)$$

where ω is the numerical approximation of ω . Show that this also solves the wave equation (1.2), and that for COURANT number $\text{Co} = \sqrt{2}/2$, we have that $\omega = \omega$.

Evidently, $\eta_{m+1,n}^m = \eta_{m,n+1}^m$. Cross-multiplying equation (1.3) by Δt , the right-hand side coefficient reduces to unity, given $\text{Co} = \sqrt{2}/2$. Furthermore, we may neglect the common factors $\exp(ikh(m+n))$ and $\exp(-i\omega m \Delta t)$. We must then have that

$$e^{i\omega \Delta t} + e^{-i\omega \Delta t} = e^{ikh} + e^{-ikh}.$$

Clearly,

$$\omega = \frac{kh}{\Delta t} = \frac{\omega h}{c \Delta t \sqrt{2}}.$$

We have used the fact that this system dictates a dispersion relation $\omega = ck\sqrt{2}$. What we wanted to show follows.