

## Keizer-Levine 1996 Simplified Open-Cell Model

In 1996 Joel Keizer and Leslie Levine investigated the calcium dynamics of the ryanodine receptor (RyR), particularly, receptor adaption and its significance on calcium-induced calcium release-dependent calcium oscillations. To do this Keizer and Levine developed a simplified mechanism that mimics adaption and constructed open-cell and closed-cell models using parameter values which reproduce experimental data for the RyR from cardiac cells. In this report we will reproduce the simplified open-cell model and simulate the model forward in time for a variety of parameters to finally perform a dynamical analysis of model.

### 1 Reproducing the Model

Here we construct a simplified (not definite molecular mechanism) model of the regulation of the RyR by calcium at its cytosolic face ( $[Ca_i^{2+}]$ ) which reproduces semi-quantitatively the observations of adaption.

#### 1.1 Adaption of the Ryanodine Receptor

The mechanism used to mimic adaption of the RyR is shown schematically in Figure 1. States  $C_1$  and  $C_2$  are closed states, with  $C_1$  dominating at low  $[Ca_i^{2+}]$ . States  $O_1$  and  $O_2$  are open state. We have taken these states to be the states of the entire RyR rather than states of the four sub-units for simplicity. Transitions from  $C_1$  to  $O_1$  and from  $O_1$  to  $O_2$  are assumed to be  $Ca^{2+}$  dependent and have been written as depending on the binding of  $n$  and  $m$   $Ca^{2+}$  ions respectively. These steps correspond to calcium induced calcium release (CICR) phenomenon and without the  $O_1$  to  $C_2$  transition the mechanism would not give rise to adaption.

To analyze the mechanism in Figure 1 we translate the schematic diagram into kinetic equations, using the mass action law. Using the notation  $P_{C_1}$  to represent the fraction of RyR that is in state  $C_1$ , we produce the following set of differential equations:

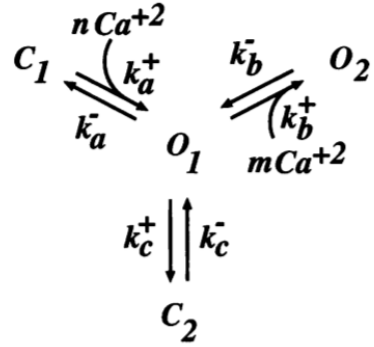


Figure 1: Schematic diagram of transitions among the four states of the RyR used to describe adaptation. States  $C_1$ , and  $C_2$  are closed states, and  $O_1$  and  $O_2$  are open states, assumed to have the same single-channel conductance. The  $k$  are rate constants; only steps  $a$  and  $b$  are  $Ca^{2+}$  dependent. [KL96]

$$\frac{dP_{C_1}}{dt} = -k_a^+ [Ca_i^{2+}]_i^n P_{C_1} + k_a^- P_{O_1} \quad (1)$$

$$\frac{dP_{O_1}}{dt} = k_a^+ [Ca_i^{2+}]_i^n P_{C_1} - k_a^- P_{O_1} - k_b^+ [Ca_i^{2+}]_i^m P_{O_1} + k_b^- P_{O_2} - k_c^+ P_{O_1} + k_c^- P_{C_2} \quad (2)$$

$$\frac{dP_{O_2}}{dt} = k_b^+ [Ca_i^{2+}]_i^m P_{O_1} - k_b^- P_{O_2} \quad (3)$$

$$\frac{dP_{C_2}}{dt} = k_c^+ P_{O_1} - k_c^- P_{C_2} \quad (4)$$

Where subscripts  $a$ ,  $b$  and  $c$  label the three kinetic steps in Figure 1. The  $k_i^\pm$  are rate constant, which along with integers  $n$  and  $m$  are determined by data fitting. The fitting off these parameters relied on the data of Gyorke and Fill [GF93]. Values of the parameters which produce a semi-quantitative fit with experiments are given in Table 1.

Table 1: RyR kinetic constants ( $n = 4$ ,  $m = 3$ )

Rate constant	Value
$k_a^+$	$1500 \mu M^{-4} s^{-1}$
$k_a^-$	$28.8 s^{-1}$
$k_b^+$	$1500 \mu M^{-3} s^{-1}$
$k_b^-$	$385.9 s^{-1}$
$k_c^+$	$1.75 s^{-1}$
$k_c^-$	$0.1 s^{-1}$

On a time scale longer than 10-20 ms "slow time scale" it is possible to approximate the values of  $P_{C_1}$ ,  $P_{O_1}$  and  $P_{O_2}$  by taking advantage of the fact that the kinetic steps  $a$  and  $b$  in Figure 1 rapidly reach equilibrium compared with  $c$ . Hence, for any instantaneous value  $P_{C_2}$  this gives

$$P_{O_1} = \frac{w}{1 + (K_a / [\text{Ca}_i^{2+}])^4 + ([\text{Ca}_i^{2+}] / K_b)^3} \quad (5)$$

Where  $w = 1 - P_{C_2} = P_{C_1} + P_{O_1} + P_{O_2}$ ,  $K_a^4 = k_a^- / k_a^+$ ,  $K_b^3 = k_b^- / k_b^+$  and  $K_c = k_c^- / k_c^+$  (defined for later).

Using Equation 5 combined with a comparable expression for  $P_{O_2}$  gives rise to the expression for  $P_O = P_{O_1} + P_{O_2}$ , the fraction of open channels on the slower time scale:

$$P_O^{\text{slow}} = \frac{w \left( 1 + ([\text{Ca}_i^{2+}] / K_b)^3 \right)}{1 + (K_a / [\text{Ca}_i^{2+}])^4 + ([\text{Ca}_i^{2+}] / K_b)^3} \quad (6)$$

Substituting Equation 5 and the definition of  $w$  into Equation 4 gives

$$\frac{dw}{dt} = -w \frac{k_c^+}{1 + (K_a / [\text{Ca}_i^{2+}])^4 + ([\text{Ca}_i^{2+}] / K_b)^3} - k_c^- (w - 1) \quad (7)$$

With algebra, Equation 7 is rearranged into

$$\frac{dw}{dt} = \frac{-(w - w^\infty ([\text{Ca}_i^{2+}]))}{\tau ([\text{Ca}_i^{2+}])} \quad (8)$$

With definitions

$$w^\infty ([\text{Ca}_i^{2+}]) = \frac{1 + (K_a / [\text{Ca}_i^{2+}])^4 + ([\text{Ca}_i^{2+}] / K_b)^3}{1 + (1/K_c) + (K_a / [\text{Ca}_i^{2+}])^4 + ([\text{Ca}_i^{2+}] / K_b)^3} \quad (9)$$

$$\tau ([\text{Ca}_i^{2+}]) = \frac{w^\infty ([\text{Ca}_i^{2+}])}{k_c^-} \quad (10)$$

Where  $w^\infty$  is the equilibrium value of  $w$  and  $\tau$  is its relaxation time.

## 1.2 Closed-Cell Model

Using the above kinetic model of the RyR to describe CICR from an internal calcium store. This is combined with a passive leak and a sarco/endoplasmic reticulum  $\text{Ca}^{2+}$  ATPase- (SERCA-) type pump [LWB<sup>+</sup>92] that returns calcium to the store. The differential equations for the closed-cell model are the balance equations for  $[\text{Ca}_i^{2+}]$  and the relaxation equation for  $w$ :

$$\begin{aligned} \frac{d[\text{Ca}_i^{2+}]}{dt} &= f_i \left( (\nu_1 + \nu_2) ([\text{Ca}_s^{2+}] - [\text{Ca}_i^{2+}]) - \nu_3 \frac{[\text{Ca}_i^{2+}]^2}{[\text{Ca}_i^{2+}]^2 + (K_3)^2} \right) \\ \frac{dw}{dt} &= \frac{-(w - w^\infty ([\text{Ca}_i^{2+}]))}{\tau ([\text{Ca}_i^{2+}])} \end{aligned} \quad (11)$$

Where  $f_i$  is the fraction of calcium that is unbound (free) in the cytoplasm and  $[\text{Ca}_s^{2+}]$  is the calcium concentration in the store.  $\nu_1$  is the rate constant for RyR,  $\nu_2$  is the rate constant for store leak,  $\nu_3$  maximal rate of SERCA and  $k_3$  is the dissociation constant for SERCA

Since the cell is assumed to be closed to calcium influx and efflux,

$$[\text{Ca}_s^{2+}] = (C_0 - [\text{Ca}_i^{2+}]) \frac{V_i f_s}{V_s f_i} = \frac{C_0 - [\text{Ca}_i^{2+}]}{c_1} \quad (12)$$

Where  $V_i$  and  $V_s$  are the volumes of the cytoplasm and the internal store, respectively,  $f_s$  is the buffering factor for the store, and  $C_0$  is the total amount of free calcium in the cell. The factor  $c_1$  is often referred to as the ratio of effective volume of the store to that of the cytoplasm [Fri95]. The three terms on the right-hand side of Equation 11 are the contributions to the calcium balance in the cytoplasm that come from the RyR, the leak, and the SERCA pump, respectively.

### 1.3 Open-Cell Model

The simplest type of open cell contains voltage-activated  $\text{Ca}^{2+}$  channels for which one can fix the rate of influx by clamping the plasma membrane potential [Mil92]. to mimic this situation we consider an open-cell model with a constant influx of external  $\text{Ca}^{2+}$ ,  $j_{in}$ , that can be controlled parametrically. Although both  $\text{Ca}^{2+}$  pumps and exchange mechanisms are prevalent in the plasma membrane of cells, we treat efflux as occurring via a plasma membrane  $\text{Ca}^{2+}$  - ATPase (PMCA pump) only [Car94]. Adding these terms to the closed-cell model described in the previous section gives us the open-cell  $[\text{Ca}_i^{2+}]$  balance equation:

$$\frac{d[\text{Ca}_i^{2+}]}{dt} = f_i \left( (\nu_1 + \nu_2) ([\text{Ca}_s^{2+}] - [\text{Ca}_i^{2+}]) - \nu_3 \frac{[\text{Ca}_i^{2+}]^2}{[\text{Ca}_i^{2+}]^2 + (K_3)^2} - \nu_{out} \frac{[\text{Ca}_i^{2+}]^2}{[\text{Ca}_i^{2+}]^2 + (K_{out})^2} + j_{in} \right) \quad (13)$$

$$(14)$$

where the final two terms represent the PMCA pump [Car94] and the  $\text{Ca}^{2+}$  influx. Because the cell is open to the external medium, the total free- $\text{Ca}^{2+}$  concentration,  $C_0$ , is no longer fixed but varies according to the equation:

$$\frac{dC_0}{dt} = f_i \left( j_{in} - \nu_{out} \frac{[\text{Ca}_i^{2+}]^2}{[\text{Ca}_i^{2+}]^2 + (K_{out})^2} \right) \quad (15)$$

Where  $j_{in}$  represents influx from a clamped membrane potential in an electrically excitable cell and  $K_{out}$  is the dissociation constant

### 1.4 Simplified Open-Cell Model

Keizer and Levine determined that adaption was not essential for oscillations in the open-cell model, this discovery allowed them to simplify the model further by requiring that:

$$w = w^\infty ([\text{Ca}_i^{2+}]) \quad (16)$$

Thus our simplified model has only two variables,  $[\text{Ca}_i^{2+}]$  and  $C_0$  with the following parameters:

Table 2: **Parameters for simplified open-cell model**

Parameter	Meaning
$f_i$	Buffer factor in cytoplasm
$\nu_1$	Rate constant RyR
$\nu_2$	Rate constant for store leak
$\nu_3$	Maximal rate of SERCA pump
$\nu_{out}$	Maximal rate PMCA pump
$k_3$	Dissociation constant for SERCA
$c_1$	Weighted volume fraction
$K_{out}$	Dissociation constant for PMCA
$j_{in}$	Constant influx rate

## 2 Simulating Simplified Open-Cell Model

To simulate the model, I will be using the popular MATLABtoolbox, MatCont. This toolbox allows you to input a system of equations and define parameters and variables. We define our system as follows.

Name	test
Time	t
Coordinates	Ca_i, C0
Parameters	f_i, v1, v2, v3, K3, c1, v_out, K_out, j_in
Derivatives	NNNN
Userfunctions	(none)
Equations	$ka\_plus = 1500; ka\_minus = 28.8;$ $kb\_plus = 1500; kb\_minus = 385.9;$ $kc\_plus = 1.75; kc\_minus = 0.1;$ $Ka = (ka\_minus/ka\_plus)^(1/4);$ $Kb = (kb\_minus/kb\_plus)^(1/3);$ $Kc = (kc\_minus/kc\_plus);$ $w = (1+(Ka/Ca_i)^4+(Ca_i/Kb)^3)/(1+(1/Kc)+(Ka/Ca_i)^4+(Ca_i/Kb)^3);$ $P\_slow = w*(1 + (Ca_i/Kb)^3) / (1 + (Ka/Ca_i)^4 + (Ca_i/Kb)^3);$ $Ca\_s = (C0 - Ca_i)/c1;$ $Ca\_i' = f_i * (v1 * P\_slow + v2) * (Ca_s - Ca_i) - v3 * (Ca_i^2 / (Ca_i^2 + K3^2)) - v\_out * (Ca_i^2 / (Ca_i^2 + K\_out^2));$ $C0' = f_i * (j\_in - v\_out * (Ca_i^2 / (Ca_i^2 + K\_out^2)));$

Figure 2: A screenshot of the MatCont system used to simulate the simplified open-cell model.

Listing 1: System of Simplified Open-Cell Model

```

ka_plus = 1500; ka_minus = 28.8;
kb_plus = 1500; kb_minus = 385.9;
kc_plus = 1.75; kc_minus = 0.1;
Ka = (ka_minus/ka_plus)^(1/4);
Kb = (kb_minus/kb_plus)^(1/3);
Kc = (kc_minus/kc_plus);
w = (1+(Ka/Ca_i)^4+(Ca_i/Kb)^3)/(1+(1/Kc)+(Ka/Ca_i)^4+(Ca_i/Kb)^3);
P_slow = w*(1 + (Ca_i/Kb)^3) / (1 + (Ka/Ca_i)^4 + (Ca_i/Kb)^3);
Ca_s = (C0 - Ca_i)/c1;
Ca_i' = fi*( (v1*P_slow+v2)*(Ca_s-Ca_i) - v3*(Ca_i^2/(Ca_i^2+K3^2)) - v_out*(Ca_i^2/(Ca_i^2 + K_out^2)) );
C0' = fi*(j_in - v_out*(Ca_i^2/(Ca_i^2 + K_out^2)));

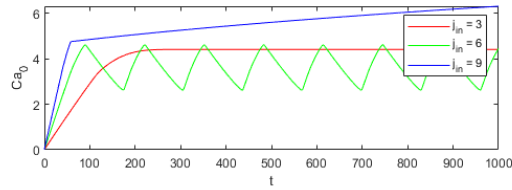
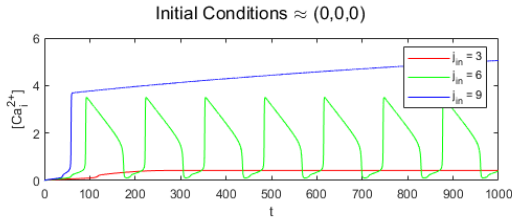
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Using the values from Table 3 we will simulate the model for 3 different values of  $j_{in} = 3, 6, 9$ . From 3 different initial conditions.

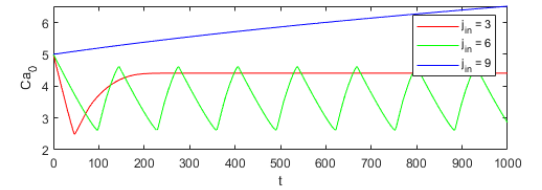
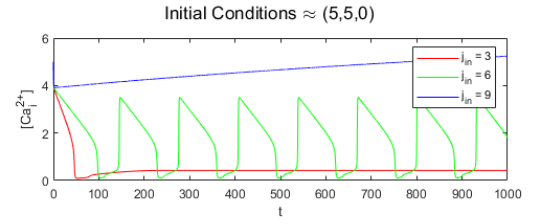
Table 3: **Parameters used for simulations**

Parameter	Value
$f_i$	0.01
$\nu_1$	$40 \text{ s}^{-1}$
$\nu_2$	$0.5 \text{ s}^{-1}$
$\nu_3$	$120 \mu\text{M s}^{-1}$
$\nu_{out}$	$9.0 \mu\text{M s}^{-1}$
$k_3$	$0.30 \mu\text{M}$
$c_1$	0.15
$K_{out}$	$0.60 \mu\text{M}$
$j_{in}$	3.0/6.0/9.0 $\mu\text{M}$

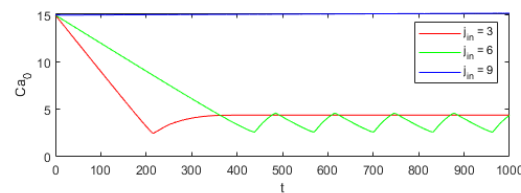
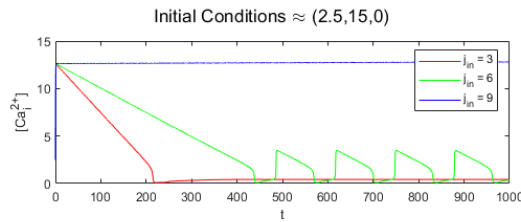
## 2.1 Varied Initial Conditions



(a) Initial Conditions  $([Ca_i^{2+}], C_0, t) = (0, 0, 0)$



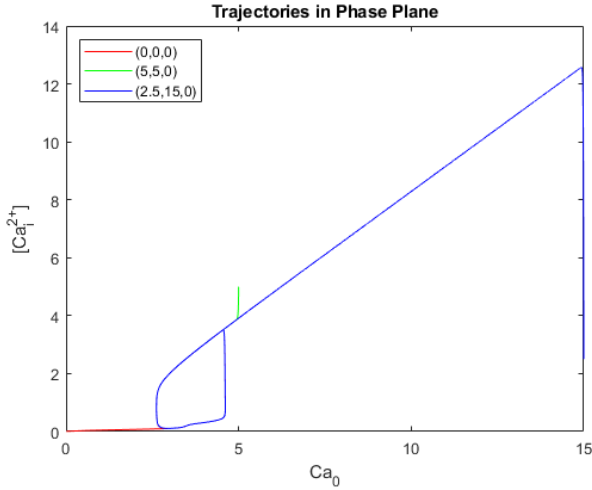
(b) Initial Conditions  $([Ca_i^{2+}], C_0, t) = (5, 5, 0)$



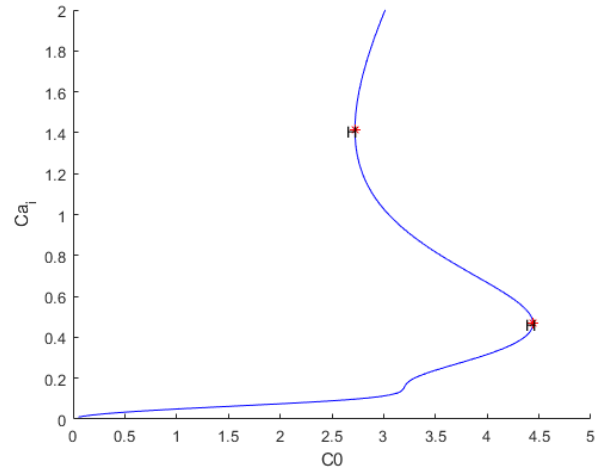
(c) Initial Conditions  $([Ca_i^{2+}], C_0, t) = (2.5, 15, 0)$ . Note the transient behaviour before the oscillations

Figure 3: Simulating the model for 3 different values of  $J_{in}$  showing that oscillations are dependent on  $Ca^{2+}$  influx.

### 3 Bifurcation Analysis



(a) A plot of the phase plane, showing three trajectories with varied initial conditions. We see that they all converged to the same periodic orbit.



(b) This plot shows a curve of equilibrium in the  $[Ca_i^{2+}]$  and  $C_0$  plane as  $J_{in}$  is allowed to varied. The two red crosses represent Hopf bifurcations.

Figure 4: Two figures of the  $([Ca_i^{2+}], Ca_0)$  phase plane

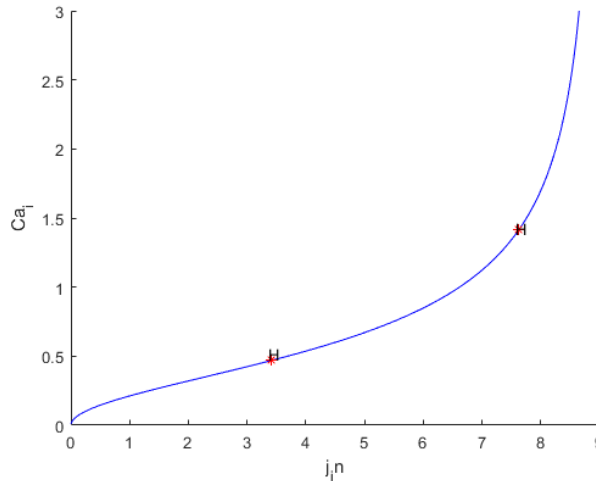


Figure 5: This plot shows a curve of equilibrium in the  $[Ca_i^{2+}]$  and  $j_{in}$  plane with  $C_0$  is allowed to varied. The two red crosses represent Hopf bifurcations. We can see this value range agrees with our earlier plots as oscillations will appear for  $j_{in}$  values between the 2 Hopf bifurcations.

### References

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