Using Redundant Internal Coordinates to Optimize Equilibrium Geometries and Transition States

CHUNYANG PENG, PHILIPPE Y. AYALA, and H. BERNHARD SCHLEGEL"

Department of *Chemistry, Wayne State University, Detroit, Michigan 48202*

MICHAEL J. FRISCH

Lorentzian, lnc., North Haven, Connecticut 06473

Received 17 December 1994; accepted 15 May 1995

ABSTRACT

A redundant internal coordinate system for optimizing molecular geometries is constructed from all bonds, all valence angles between bonded atoms, and all dihedral angles between bonded atoms. Redundancies are removed by using the generalized inverse of the *G* matrix; constraints can be added by using an appropriate projector. For minimizations, redundant internal coordinates provide substantial improvements in optimization efficiency over Cartesian and nonredundant internal coordinates, especially for flexible and polycyclic systems. Transition structure searches are also improved when redundant coordinates are used and when the initial steps are guided by the quadratic synchronous transit approach. *0* 1996 by John Wiley & Sons, Inc.

Introduction

eometry optimization is an important aspect of almost all electronic structure calculations. Most of the optimizations are carried out using some variety of quasi-Newton algorithm employing energies and gradients. Geometry optimization has been reviewed in several articles.¹ The efficiency of an optimization depends on a number of factors: (1) the initial geometry, (2) the choice of coordinate system, **(3)** the initial estimate of the Hessian, **(4)** the Hessian updating method, and (5) control of the search direction and step size.

As an estimate for the starting geometry for optimizing equilibrium structures, standard geometries or structures obtained by molecular mechanics minimization are often good. For transition structures, many quasi-Newton methods need to

^{*}Author **to** whom all correspondence should be addressed.

Journal of Computational Chemistry, Vol. 17, No. 1, **49-56 (1996)** *0* **1996 by John Wiley** & **Sons,** Inc. **CCC** 01 92-8651 *196 1* 01 **0049-08**

start fairly close to the quadratic region of the transition state. However, initial estimates for transition structures are more difficult to obtain than for equilibrium structures. Molecular mechanics generally cannot handle transition states involving the making and breaking of bonds. Standard geometries are available for some types of transition states,² but transition state geometries vary much more than equilibrium geometries. Techniques such as synchronous transit, 3 coordinate driving, walking up valleys, and eigenvector following⁴ may be useful in getting close to the transition structure.

Early work in geometry optimization using electronic structure methods used nonredundant internal coordinates (eg, the **Z** matrix internal coordinates used in many molecular orbital programs). Cartesian coordinates were shown to be better for some cyclic molecules, 5 and mixed Cartesian and internal coordinates also had some advantages.6 Cartesian normal modes have also been used successfully.⁷ However, Pulay⁸⁻¹⁰ demonstrated clearly that redundant internal coordinates are the best choice for minimizing polycyclic molecules. Baker¹¹ compared redundant internal to Cartesian coordinates and came to a similar conclusion. In the present work, we adopted a simpler system of redundant internal coordinates (but with more redundancies) than we have used previously to estimate the initial Hessian.¹² We have also extended the use of redundant internal coordinates to transition structure optimization using a synchronous transit guided quasi-Newton method.¹³

For the choice of the initial Hessian, Baker^{5,11} showed that minimizations with Cartesian and redundant internal coordinates were similar if a good initial estimate of the Hessian is used (i.e., a molecular mechanics Hessian instead of a unit Hessian). In this article we use a simple valence force field that is diagonal in the redundant internal coordinate system, similar to the one used in our earlier work.¹² Proper updating of the Hessian is essential for efficient optimizations; most modern updating algorithms give similar results and therefore will not be discussed here.

Choice of the search direction and control of step size are important in determining the efficiency of the optimization.' Trust radius and eigenvector following⁴ methods have been shown to be effective. For transition states, it is highly desirable to monitor the direction of the transition vector (i.e., the eigenvector corresponding to the negative eigenvalue). In previous work¹³ we found that the tangent to the linear or quadratic synchronous transit path 3 was useful in choosing the uphill search direction for optimizing transition states.

Method

GENERATION OF REDUNDANT INTERNAL COORDINATES

Pulay and co-workers $8-10$ defined a natural internal coordinate system that is similar to the coordinates used by vibrational spectroscopists. It minimizes the number of redundancies by using local pseudosymmetry coordinates about each atom and special coordinates for ring deformation, spiro ring fusions, etc. Similar coordinates are used by Baker¹¹ and in TurboMole. 14 To reduce the number of special cases, we use a simpler set of internal coordinates composed of all bond lengths, valence angles, and dihedral angles. This coordinate set, however, has somewhat greater redundancy than the coordinates used by Pulay, but this does not seem to affect the efficiency of the optimization.

Our redundant internal coordinates set is defined in the following manner. First, the interatomic distances are examined to determine which atoms are bonded. Two atoms are considered bonded if their separation is less than 1.3 times the sum of the (single bond) covalent radii of the two atoms. If the molecular system consists of two or more fragments that are not bonded by this criterion, then the shortest distance between the fragments is determined; all interfragment distances that are less than the smaller of 1.3 times this distance, or 2 A, are designated as bonds. **A** hydrogen bond is indicated if the $XH \cdots Y$ distance (X, Y) $= N$, O, F, P, S, Cl) is greater than the sum of the covalent radii, less than 0.9 times the sum of the van der Waals radii, and the $X-H \cdots Y$ angle is greater than 90". A bond stretching coordinate is assigned to each regular, interfragment, and hydrogen bond.

A valence angle bend coordinate is assigned for any two atoms bonded to the same third atom $(\angle A - B - C$, where A is bonded to B and C is bonded to B). Special attention must be given to linear valence angles. If the A -B-C angle is greater than $\sim 175^\circ$, then two orthogonal linear angle bend coordinates are generated (some care is needed so the orientation of the two displacements does not change during the course of the optimization).

A dihedral angle coordinate is assigned for each pair of atoms bonded to opposite ends of a bond $(\angle A - B - C - D$, where A is bonded to B, B is bonded to C, and C is bonded to D) or bonded to a linear array of atoms $(\angle A \rightarrow B \rightarrow C \rightarrow D$ for $A-B-X \cdots Y-C-D$, where $B, X, \cdots Y, C$ are collinear and bonded, e.g., 2-butyne). If one or both of the valence angles involved in a dihedral angle $(\angle A - B - C$ and/or $\angle B - C - D$ in $\angle A - B$ $-C-D$) is linear, then the dihedral angle is omitted. If there are no dihedral angles generated for the molecule by these criteria (e.g., for a tri-coordinate system such as H_2CO or BH_3), the appropriate dihedral angles are added to take care of outof-plane bending. Since the values of the dihedral angles are chosen to be in the range -180 to $+180^\circ$, differences in dihedral angles may need to be incremented by $\pm 360^\circ$ so that the smallest displacement is obtained. In addition to the redundant internal coordinates generated automatically, extra stretch, bend, and dihedral angle coordinates can also be specified in the input.

The coordinate system defined above is based on the identification of bonds. All of the remaining coordinates (valence angles and dihedral angles) are generated from the bonding information. Transition states could pose a problem because they generally contain one or more partially formed or partially broken bonds. For regular transition state optimizations starting from one structure, the bonds being made or broken will need to be specified in the input. For the synchronous transit method for guiding transition state searches,^{3,13} two or three structures are used to start the optimization--one in the reactant valley, the second in the product valley, and optionally a third structure as an initial guess for the transition state geometry. If the third structure is absent, the initial guess for the transition state geometry can be obtained by interpolating between the reactants and products in redundant internal coordinates. The reactant-like structure defines one set of redundant internal coordinates, and the product-like structure defines the other set. The coordinates for the transition state search are taken as the union of the reactantlike and product-like coordinates.

TRANSFORMATION OF THE GRADIENT AND HESSlAN

The transformation of the energy derivatives from Cartesian to redundant internal coordinates is performed in the manner outlined by Pulay and Fogarasi.⁹ If **B** is the Wilson B matrix¹⁵ $(B_{ij} =$

REDUNDANT INTERNAL COORDINATES

 $\partial q_i/\partial x_i$) defining the transformation from Cartesian displacements to redundant internal displacements,

$$
\delta q = B \, \delta x \tag{1}
$$

then the transformed gradient is given by

$$
\mathbf{B}^t \mathbf{g}_q = \mathbf{g}_x \tag{2}
$$

where \mathbf{g}_x is the gradient in Cartesian coordinates and **g,** is the gradient in redundant internal coordinates. Since B is rectangular, the inverse of this transformation is a little more complicated and can be written

$$
\mathbf{g}_q = \mathbf{G}^- \mathbf{B} \mathbf{u} \mathbf{g}_x \tag{3}
$$

where $G = BuB^t$ and **u** is an arbitrary matrix (a unit matrix is used in the present application). The generalized inverse of G is obtained by diagonalizing G and inverting only the nonzero eigenvalues:

$$
\mathbf{V}^t \mathbf{G} \mathbf{V} = \begin{bmatrix} \mathbf{\Lambda} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix}; \quad \mathbf{G}^- = \mathbf{V} \begin{bmatrix} \mathbf{\Lambda}^{-1} & \mathbf{0} \\ \mathbf{0} & \mathbf{0} \end{bmatrix} \mathbf{V}^t \quad (4)
$$

Quasi-Newton optimization methods require an initial estimate of the Hessian. As the default, we use an empirical estimate that is diagonal in the redundant internal coordinate space (see ref. 12 for regular stretch, bend, and torsions, see ref. 16 for hydrogen bonds). Alternatively, the Cartesian Hessian can be calculated analytically at any number of levels of theory and then can be transformed to redundant internal coordinates. By differentiating eq. (2) , it is apparent that the transformation of the Hessian involves the derivative of the B matrix, B' $(B'_{ijk} = \partial^2 q_i / \partial x_i \partial x_k)$:

$$
\mathbf{B}^t \mathbf{H}_q \mathbf{B} + \mathbf{B}'^t \mathbf{g}_q = \mathbf{H}_x \tag{5}
$$

where H_r is the Hessian in Cartesian coordinates and H_a is the Hessian in redundant internal coordinates. The inverse of this transformation is

$$
\mathbf{H}_q = \mathbf{G}^- \mathbf{B} \mathbf{u} \big(\mathbf{H}_x - \mathbf{B}'^t \mathbf{g}_q \big) \mathbf{u}^t \mathbf{B}^t \mathbf{G}^- \qquad (6)
$$

The derivatives of the **B** matrix, **B**', are calculated analytically.

For minimizations, the Hessian in redundant internal coordinates is updated iteratively by applying the Broyden, Fletcher, Goldfarb, Shanno $(BFGS)$ formula¹⁷ using the current point and all previous points, rather than using just the current and the next most recent point. Bofill's update 18 (a mixture of the symmetric Powell^{18b} and the

Murtagh-Sargent^{18c} updates) is used for transition states. Sometimes it may be necessary to calculate a few key rows and columns of the Hessian numerically before proceeding with an optimization. This is accomplished by displacing the appropriate redundant internal coordinates, converting to Cartesian coordinates, calculating the energy and gradient, and then updating the Hessian in redundant internal coordinates using the symmetric Powell method.

OPTIMIZATION STEP IN REDUNDANT INTERNAL COORDINATES

Some care must be taken in generating the Newton step in redundant coordinates so displacements are generated primarily in the nonredundant part of the internal coordinate space. Again following Pulay and Fogarasi? a projector is constructed from the G matrix and its generalized inverse:

$$
\mathbf{P}' = \mathbf{G}\mathbf{G}^{-} = \mathbf{G}^{-}\mathbf{G} \tag{7}
$$

For a constrained optimization, the projector needs to be modified to include the constraints. If C is the projector for the constraints (e.g., a diagonal matrix with 1's on the diagonal for the constraints and 0's elsewhere), the projector with constraints is

$$
\mathbf{P} = \mathbf{P}' - \mathbf{P}'\mathbf{C}(\mathbf{C}\mathbf{P}\mathbf{C})^{-1}\mathbf{C}\mathbf{P}'
$$
 (8)

Both the gradient and the Hessian have to be projected. To prevent displacements in the remainder of the space, the corresponding matrix elements of the Hessian are set to arbitrarily large values:

$$
\tilde{\mathbf{g}}_q = \mathbf{P} \mathbf{g}_q; \n\tilde{\mathbf{H}} = \mathbf{P} \mathbf{H} \mathbf{P} + (\mathbf{1} - \mathbf{P}) \mathbf{A} (\mathbf{1} - \mathbf{P})
$$
\n
$$
= \mathbf{P} \mathbf{H} \mathbf{P} + \alpha (\mathbf{1} - \mathbf{P})
$$
\n(9)

where **A** is the identity matrix times α (a large constant, e.g., 1000 au). The Newton step is then given by

$$
\Delta \mathbf{q} = -\tilde{\mathbf{H}}^{-1} \tilde{\mathbf{g}}_q \tag{10}
$$

From a numerical point of view, this gives essentially the same step as Pulay's approach, which uses the projected generalized inverse of the projected Hessian. The Newton step is not always the best choice. If the predicted step is too large during a minimization, the trust radius method is

used to modify the step size and direction. For transition state searches, the eigenvector following method 4 is used to control the step. In both cases it is desirable to have an invertible Hessian [e.g., eq. (9)].

The eigenvector following method searches uphill along a designated eigenvector of the Hessian and downhill along the remaining eigenvectors. Most often, the eigenvector with the lowest eigenvalue is used. In a previous article¹³ we found that the tangent to a synchronous transit path³ was a useful guide for choosing the correct eigenvector to follow. The same approach is used here. The initial few steps of a transition state optimization are constrained to search for a maximum along the synchronous transit path. In subsequent steps, the tangent to the synchronous transit path is used to choose the best vector for the eigenvector following method (note that the tangent vector must be projected onto the nonredundant part of the internal coordinate space before it is used).

CONVERSION FROM REDUNDANT INTERNAL COORDINATES TO CARTESIAN COORDINATES

Because the displacements are finite and the transformation between redundant internal and Cartesian coordinates is curvilinear, the coordinate conversion must be done in an iterative manner.⁹ The first estimate of the new Cartesian coordinates is given by

$$
\mathbf{x}_1 = \mathbf{x}_0 + \mathbf{u} \mathbf{B}^t \mathbf{G}^- \Delta \mathbf{q} \tag{11}
$$

The values of the internal coordinates are computed from the Cartesian coordinates, and $\mathbf{q}_1 - \mathbf{q}_0$ is compared with Δq (some care must be taken with dihedral angles to avoid extraneous multiples of 360°). The difference, $\Delta \Delta q = \Delta q - (q_1 - q_0)$, is transformed in a similar fashion; the process is repeated until there is no further change in the Cartesian coordinates (root mean square [rmsl change less than 10^{-6}). In the rare cases in which the iteration does not converge, the first estimate of the Cartesian displacements $[eq. (11)]$ is used without subsequent iteration. For a constrained optimization, a small additional displacement, $\Delta q'$, may need to be added to reimpose the constraints. A similar iteration is carried out with $C\Delta q'$ and $C\Delta\Delta q'$, where C is the projector for the constraints used in eq. *(8).*

REDUNDANT INTERNAL COORDINATES

Results and Discussion

The methods outlined earlier have been incorporated into Gaussian **94.19** Table I compares the optimization of equilibrium structures carried out by Baker¹¹ and in the present work. The published starting structures of Baker were used, and the same final energies were obtained (with the exception of hydrazobenzene, for which a lower energy minimum was found). Although the present redundant internal coordinates, the initial estimate of the Hessian, and the convergence criteria are somewhat different, the performance of the two approaches is remarkably similar. Most of the differences can be attributed to the different convergence criteria. In the Gaussian suite of programs, standard convergence requires that the rms gradient in internal coordinates is less than 0.0003 au, the maximum component of the gradient is less than 0.00045 au, the rms of the predicted Cartesian displacement (after rotation and translation to

TABLE I.

Comparison of the Number of Steps Required to Optimize Equilibrium Geometries Using Various Coordinate Systems.

	Degrees of	Cartesians	Redundant Internals		
Molecule ^a	Freedom	Baker ^b	Baker ^b	Present work ^c	
Water	2	5	6	4	
Ammonia	2	6	6	4	
Ethane	3	4	5	4	
Acetylene	2	6	6	4	
Allene	з	5	5	4	
Hydroxysulphane	6	11	8	7	
Benzene	2	4	4	з	
Methylamine	10	5	6	4	
Ethanol	13	6	6	5	
Acetone	8	7	6	5	
Disilyl ether	7	10	8	7	
1,3,5-Trisilacyclohexane	11	8	8	11	
Benzaldehyde	25	6	6	4	
1,3-Difluorobenzene	11	5	5	4	
1,3,5-Trifluorobenzene	4	5	5	4	
Neopentane	3	5	5	4	
Furan	8	8	8	5	
Naphthalene	9	5	5	4	
1,5-Difluoronaphthalene	17	6	6	4	
2-Hydroxybicyclopentane	36	15	15	11	
ACHTAR10	42	11	12	9	
ACANIL01	34	7	8	6	
Benzidine	18	10	9	7	
Pterin	31	9	10	8	
Difluoropyrazine	15	8	9	6	
Mesityl oxide	28	7	$\overline{7}$	5	
Histidine	54	30	19	14	
Dimethylpentane	63	9	12	9	
Caffeine	42	10	12	6	
Methone	81	14	13	11	
ACTHCP	42	28^d	37	27	
Histamine H ⁺	48	26	21	19	
Hydrazobenzene	38	35	19	20 ^e	

^aStarting geometries given in ref. 11; calculated by HF / STO-3G.

 B Ref. 11 using a Hessian from molecular mechanics.

'Present work, using a diagonal Hessian in redundant internal coordinates.

^dConverged to a higher energy final structure than ref. 6.

eConverged to a lower energy final structure than ref. 11: $E(au) = -563.263804$, \angle CNNC = 154.2°, \angle HNNH = 39.8°.

TABLE II. Comparison of the Number of Steps Required to Optimize Equilibrium Geometries Using Various

aStarting geometries given **in** ref. 6; calculated **by** HF / **STO-3G.**

 b Ref. 6.</sup>

'Present **work**

maximum coincidence) is less than 0.0012 au, and the maximum component of the predicted displacement is less than 0.0018 au (i.e., all four must be satisfied). Alternatively, if the maximum gradient and rms gradient are a factor of 100 smaller than their respective thresholds, the optimization is considered converged. Baker requires that the maximum component of the gradient be less than 0.0003 au and either the maximum component of the predicted Cartesian displacement be less than 0.0003 au or the energy change from the previous cycle be less than 10^{-6} au.

Table **I1** summarizes several optimizations of equilibrium structures using **Z** matrix coordinates (i.e., nonredundant internal coordinates), Cartesian coordinates, mixed Cartesian and nonredundant internal coordinates, and redundant internal coordinates. The starting structures and some of the optimization results were published previously.6 For fairly rigid molecules, the results of all four coordinate systems are similar. For more flexible molecules, Cartesian coordinates perform more poorly. For ACTHCP and histamine H^+ , the redundant coordinates show a significant advantage. Additional examples that illustrate the efficiency of redundant coordinate optimizations include a large dye molecule (53 atoms, 153 degrees of freedom, 23 steps from the PM3 geometry to the HF/3-21G optimized geometry) and taxol (113 atoms, 333 degrees of freedom, 58 steps from the PM3 geometry to the HF/STO-3G optimized geometry).

Figure 1 shows a simple example of a constrained optimization using redundant internal coordinates. The potential energy curve for isomerization of cyclohexane from the chair to the twistboat conformation was computed as a relaxed potential surface scan. One of the CCCC dihedral angles was constrained, and all remaining coordinates were optimized; the CCCC dihedral was incremented by 15" and the process was repeated until the scan was completed (six to eight optimization steps per structure). Relaxed potential surface scans of cyclic molecules are considerably more difficult to carry out in Cartesian or nonredundant internal coordinates.

Table **I11** compares the results of a number of transition state optimizations in nonredundant in-

FIGURE 1. Relaxed potential energy surface scan of the conversion of chair cyclohexane to twist boat using redundant internal coordinates.

Comparison of the Number of Steps Required to Optimize Transition State Geometries.										
Reaction ^a	Z Matrix Internals			Redundant Internals						
	Regular ^b	CalcFC ^c	QST3 ^d	CalcFC ^c	QST2 ^e	QST3 ^d				
$CH_4 + F \rightarrow CH_3 + HF$	6									
$CH_3O \rightarrow CH_2OH$	12									
$SiH2 + H2 \rightarrow SiH4$	11									
$CH_5F \rightarrow C_2H_4 + HF$	16	12	15	13						
Diels-Alder reaction	56		23		13	14				
Claisen reaction	38		15		15	15				
Ene reaction	fail	15	28	13	18	18				

TABLE 111.

^a Structures given in ref. 13; all calculations at the HF / 3-21G level.

^bRegular transition state optimization algorithm in Gaussian 92, starting with one structure; empirical estimate of initial Hessian with two to four preliminary steps to calculate key rows and columns of the Hessian by numerical differentiation.

 \textdegree Starting with one structure; initial Hessian calculated analytically at HF / 3-21G.

^d Starting with three structures (reactants, the products, and starting geometry of the transition state); empirical estimate of the Hessian with no preliminary steps; quadratic synchronous transit used to guide the optimization.

^eSame as (d) but starting with two structures (reactant and product); starting geometry transition structure estimated by linear interpolation between reactants and products.

ternal (Z matrix) coordinates and the present set of redundant coordinates. The starting geometries were given previously.^{13, $+$} For the three simplest transition states, all of the approaches give similar results; however, for the larger cyclic transition states, there are significant differences. In general, the redundant internal coordinates perform better than nonredundant internal coordinates. The regular transition state optimization in Z-matrix coordinates takes the most steps. In this case the initial Hessian is estimated empirically, and two to four preliminary gradient calculations (included in the total number of steps) were used to calculate key rows and columns of the Hessian by numerical differentiation. The performance is much improved if the tangent of the quadratic synchronous transit is used to guide the initial steps of the $~$ optimization 13 (the initial Hessian is estimated empirically with no additional numerical differentiation steps). The transition state optimization in redundant internal coordinates takes fewer steps than in **Z** matrix coordinates, and the threestructure quadratic synchronous transit method is somewhat better than the two-structure approach. With either coordinate system, the fewest opti-

'Suitable starting structures can be generated easily with a graphical user interface. The reactants are sketched and minimized in a conformation appropriate for the reaction (for a bimolecular case, the reactants should be constrained so that the bonds being formed are 80-120% longer than equilibrium). The product structures can then be generated by copying the reactants, changing the bonding and constraints so that they are appropriate for the products, and minimizing the resulting structure.

mization steps are taken when the Hessian is calculated analytically. However, the time to compute the Hessian must be added in. If the Hessian calculation takes more than three to five times longer than the gradient, it may be more cost effective to use the three-structure quadratic synchronous transit approach in redundant internal coordinates.

Summary

The use of redundant internal coordinates substantially improves the efficiency of optimizations of equilibrium geometries, particularly for flexible and polycyclic systems. For transition states, the two- and three-structure quadratic synchronous transit guided approach in redundant internal coordinates shows considerable promise.

Acknowledgments

This work was supported by grants from the National Science Foundation (CHE 90-20398) and Gaussian, Inc.

References

1. (a) H. B. Schlegel, In *Modern Electronic Structure Theory,* D. R. Yarkony, Ed., World Scientific Publishing, Singapore, 1995; (b) H. B. Schlegel, *Adv. Chem.* Phys., **67, 249** (1987);

(c) J. D. Head and M. C. Zemer, *Ah. Quantum Chem.,* **20,** 239 (1989); (d) S. Bell and J. S. Crighton, *1. Chem. Phys., 80,* 2464 (1984); (e) H. B. Schlegel, *J. Comp. Chem.,* 3,214 (1982).

- 2. See, for example, K. N. Houk, **Y.** Li, and J. D. Evanseck, *Angew. Chem., Int. Ed. (Engl.),* 31, 682 (1992).
- 3. T. A. Halgren and W. N. Lipscomb, *Chem. Phys. Lett.,* 49, 225 (1977).
- 4. J. Simons and J. Nichols, *Int.* 1. *Quantum Chem. Quantum Chem. Symp.,* **24,** 263 (1990) and references therein; present implementation from J. Baker, *J. Comp. Chem.,* 7,385 (1986).
- 5. J. Baker and W. J. Hehre, *J. Comp. Chem.,* **12,** 606 (1991).
- 6. H. B. Schlegel, *Int. J. Quant. Chem.: Quant. Chem. Symp.,* **26,** 243 (1992).
- 7. H. **L.** Sellers, V. J. Klinkowski, and L. Schafer, *Chem. Phys. Lett., 58,* 541 (1978).
- 8. P. Pulay, G. Fogarasi, F. Pang, and J. E. Boggs, *J. Am. Chem.* Soc., 101, 2550 (1979).
- 9. P. Pulay and G. Fogarasi, *J. Chem. Phys.,* **96,** 2856 (1992).
- 10. G. Fogarasi, **X.** Zhou, P. Taylor, and P. Pulay, *1. Am. Chem. Soc.,* **114,** 8191 (1992).
- 11. J. Baker, *J. Comp. Chem.,* 14, 1085 (1993).
- 12. H. B. Schlegel, *Theor. Chim. Acta,* **66,** 333 (1984).
- 13. C. Peng and H. B. Schlegel, *Israeli J. Chem.,* **33,** 449 (1994).
- 14. R. Ahlrichs, M. Bär, M. Ehrig, M. Häser, H. Horn, and C. Kolmel, *TURBOMOLE,* Biosym Technologies, San Diego, CA, 1992.
- **15.** E. B. Wilson, J. C. Decius, and P. C. Cross, *Molecular Vibrations,* McGraw-Hill, New York, 1955.
- 16. T. H. Fischer and J. Almlof, *J. Phys. Chem.,* **96,** 9768 (1992).
- 17. (a) C. G. Broyden, *J. Inst. Math. Appl.,* **6,** 76 (1970); (b) R. Fletcher, *Comput. J.,* **13,** 317 (1970); (c) D. Goldfarb, *Math. Comput.,* 24, 23 (1970); (d) D. F. Shanno, *Math. Comput.,* **24,** 647 (1970).
- 18. (a) J. M. Bofill, *J. Comp. Chem.,* **15, 1** (1994); (b) 8. A. Murtagh and R. W. H. Sargent, *Comput.* J., 13, 185 (1972); (c) M. J. D. Powell, *Math. Prog.,* **1,** 26 (1971).
- 19. M. J. Frisch, G. W. Trucks, H. B. Schlegel, P. M. W. Gill, B. G. Johnson, M. A. Robb, J, R. Cheeseman, **T.** Keith, G. A. Peterson, J. A. Montgomery, K. Raghavachari, M. A. Al-Laham, V. G. Zakrzewski, J. **V.** Ortiz, J. 8. Foresman, J. Cioslowski, B. 8. Stefanov, A. Nanayakkara, M. Challacombe, C. Y. Peng, P. **Y.** Ayala, W. Chen, M. W. Wong, J. L. Andres, E. S. Replogle, R. Gomperts, R. L. Martin, D. J. Fox, J. S. Binkley, D. J. Defrees, J. Baker, J. P. Stewart, M. Head-Gordon, C. Gonzalez, and J. **A.** Pople, *Gaussian* 94, Gaussian, Inc., Pittsburgh, PA, 1995.