



Neutron- and muon- induced background in underground physics experiments

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for the JRA1 and N3 simulation groups

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Background studies in ILIAS

A working group in **ILIAS** is devoted to **background studies** within **N3** and **JRA1** activities



Goals:



- 1) **Simulations** of various types of **background** in underground laboratories; **comparison** of the **results** from MC codes with **each other** and with **experimental data** (MC validation)
- 2) Investigation of methods of **background suppression** and **rejection** (passive shielding, active vetoes etc.); formulation of requirements for **shieldings** and **veto systems** (optimisation). Both items require **detailed Monte Carlo simulations**, including **detector-related effects** (geometry, response)

Active groups: **France** - Saclay, LSM, Lyon, Grenoble; **Germany** - Tuebingen, Munich-Garching, Karlsruhe, Heidelberg; **Spain** - Zaragoza, Canfranc; **Italy** - Milan, LNgS; **United Kingdom** - Imperial College, RAL, Sheffield

Neutron and muon-induced background

Neutrons are a relevant **background source** for most **underground experiments** through different reactions:

- ➡ **Dark Matter:** elastic interactions of fast neutrons have the same signature as **WIMPs**
- ➡ **Double beta decay:** γ -rays emitted in $(n, n'\gamma)$ or (n, γ) interactions may mimic **DBD signature**. **Unstable**

isotopes can be created via $(n, X) \rightarrow$ possible background

Neutrons in **deep underground labs** are **produced** by:

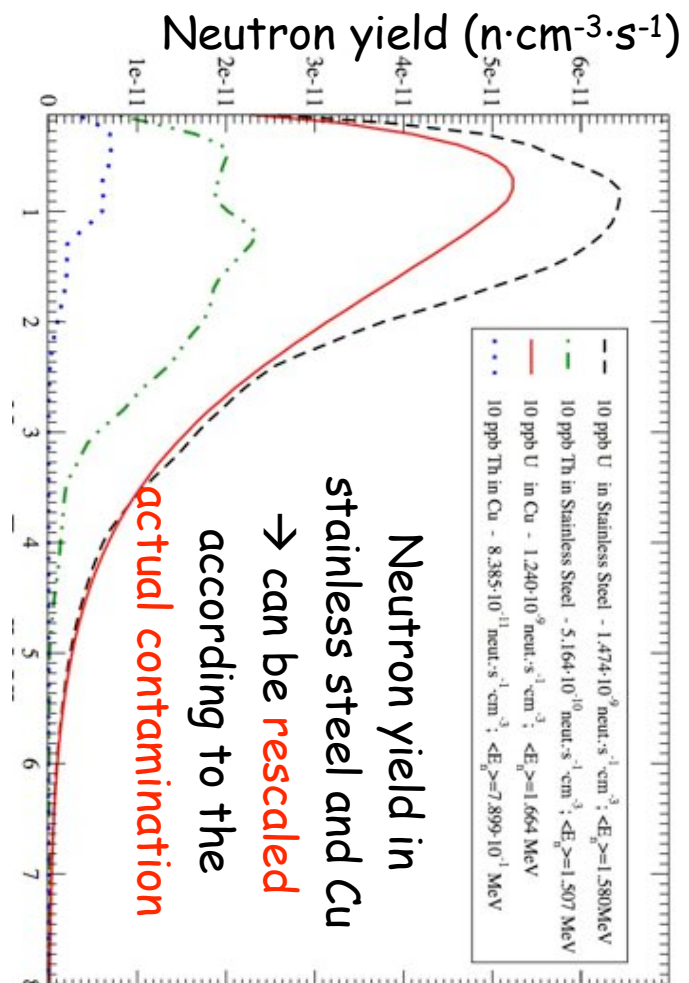
- Natural **radioactivity** (rocks, setup materials), via spontaneous **fission** (^{238}U) and **(α, n) reactions** on light nuclei. **$E_n < 10 \text{ MeV}$**
- Reactions induced by **cosmic ray muons** (in rock or setup materials), as photo-disintegration, spallation.

 **Unstable nuclei** can be also produced (bck for DBD)

Neutron yield from radioactivity

Ingredients for the recipe:

- **chemical composition** of the material (rock, detector component)
- **contamination** in ^{238}U and α emitters (U and Th chains)
- **cross section** of (α, n) interactions and SF half-life
- **propagation** and detection of the produced neutrons in the materials (MORSE and MCNPCE4)



Neutron yield in stainless steel and Cu
→ can be **rescaled**
according to the

actual contamination

Neutron energy (MeV)

Calculation of **neutron yield** per unit volume materials **(MORSE and MCNPCE4)**

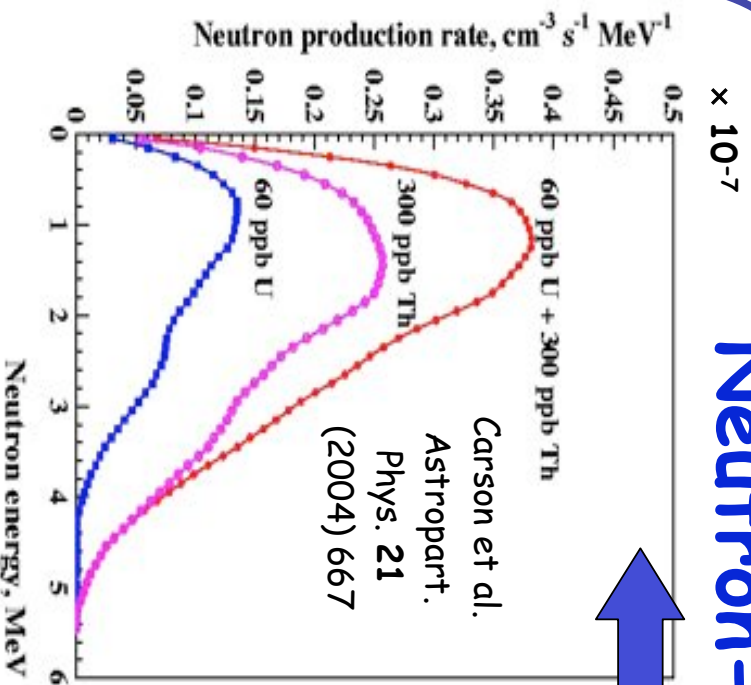
(Wilson et al. SOURCES4A, Technical Report LA-13639-MS, Los Alamos, 1999).

Substantially **improved** within ILIAS in order to increase **energy range**, to extend libraries of **cross-sections** (from codes or data) and of branching ratios to **excited states**.

Carson et al. Astrop. Phys. 21 (2004) 667

Lembrani et al., NIM A 560 (2006) 454

Neutron-induced background



Neutron yield in NaCl \rightarrow neutron production is dominated by (α, n)

Neutron **yield** due to the rock radioactivity and effective **flux** depend on **chemical composition** \rightarrow absorption by H

Typical neutron flux in underground laboratories: **a few $10^{-6} \text{ n}/(\text{cm}^2 \cdot \text{s})$** .

Peak energy between **1 and 2 MeV**

Neutron flux from the **rock** can be **effectively suppressed** by thick **neutron passive shieldings** \rightarrow e.g. $55 \text{ g}/\text{cm}^2$ of CH_2 give < 1

event/(ton \cdot year) for direct search DM experiments

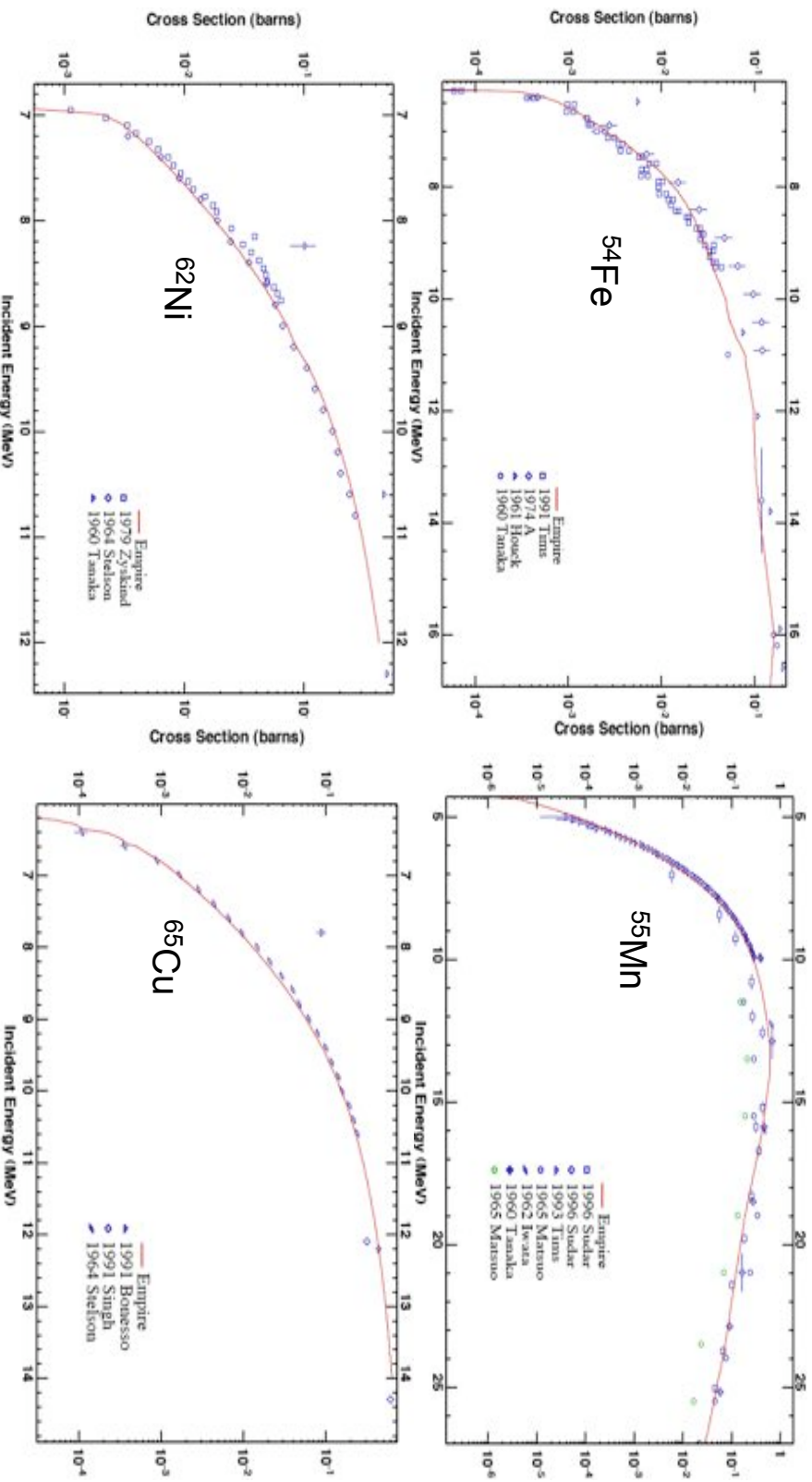
The **dominant source** of neutron background are **detector**

components and **shielding** \rightarrow attention to be paid to **radiopurity** of the internal components and to the **optimization** of passive shielding

Ingredient: (α, n) cross sections

Experimental data are *scarce* or *unavailable* for some target isotopes of interest \rightarrow calculated using the **EMPIRE 2.19 code**.

Good agreement with available experimental data ($\approx 20\%$)

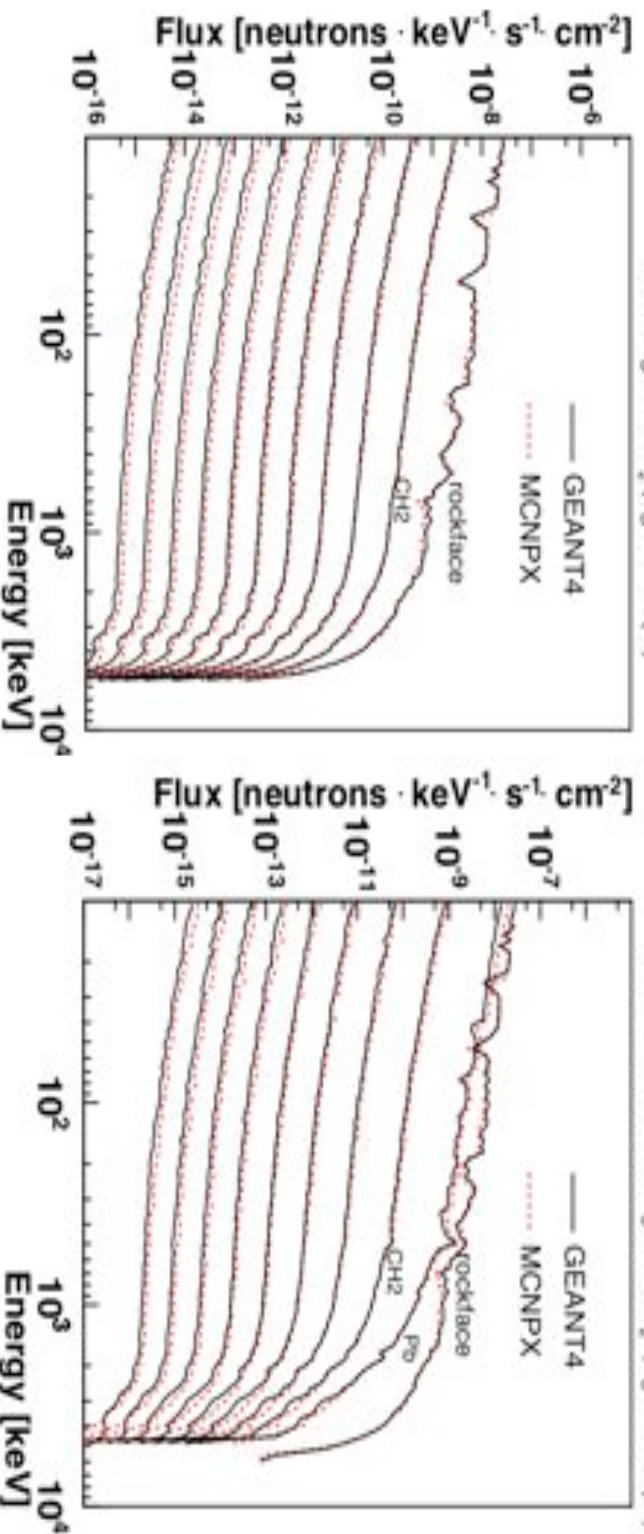


(α, n) interactions in high-Z targets ($> \text{Fe}$) suppressed by Coulomb barrier

Intercomparison and validation

Estimates of neutron background depend on **neutron propagation** in the source material (e.g. rock) and in the passive shielding.

Neutron tracking usually performed with **Geant4** or **MCNPX**



Lemrani
et al,
NIM A,
560
(2006)
454

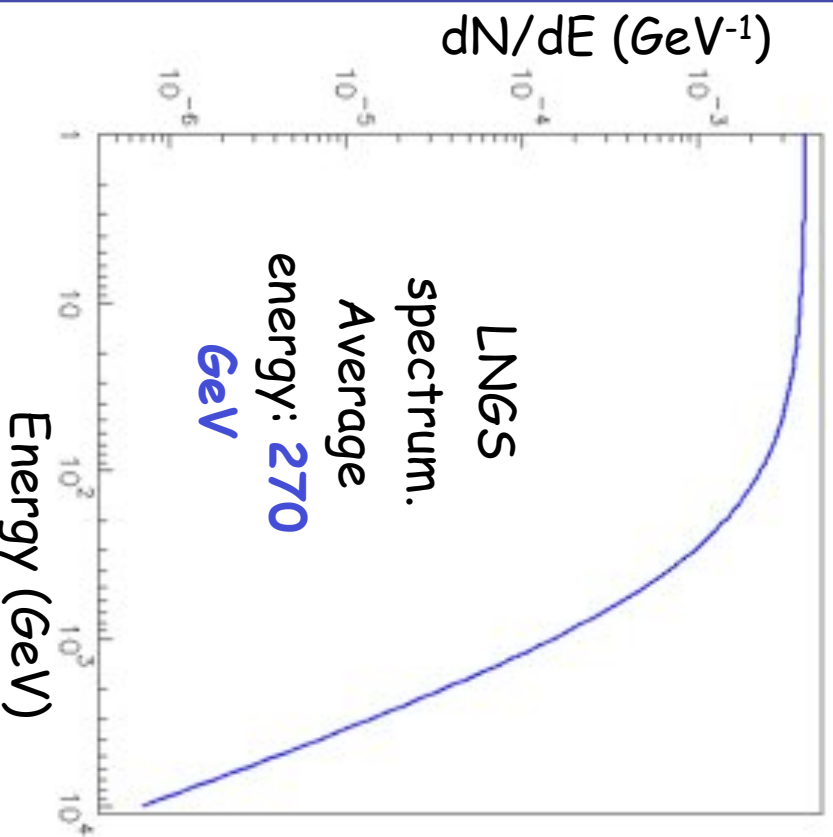
MCNP vs. GEANT4 - **good agreement**: 50% difference after 6

orders of magnitude suppression by shielding

Measurements of environmental **neutron flux** performed or ongoing

within **ILIAS-JRA1** (Tziaferi talk) → require detailed MC

Muon-induced neutrons



Ingredients for the recipe:

- Muon rate (measurements at a the underground site), e.g. 1 $\mu/(m^2 \cdot h)$ at Gran Sasso
- Muon spectrum and angular distribution. Simulations or measurements (if available) - not a problem (for example, **MUSIC** code - Kudryavtsev et al. Phys. Lett. B **471** (1999) 251 or **MUSUN** code - Kudryavtsev et al. NIMA, **505** (2003) 688).

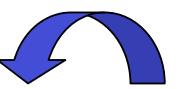
Notice: all particles should be pradd~~sed~~_{ed}, pncp~~ing~~_{ing} (red and 4 det~~ect~~_{ed} field A)

with one code to look for **simultaneous detection** of neutrons and other particles → **background identification** in the detector

Integral neutron yield

Among MC outputs: **integral neutron yield** (normalized per muon and per g/cm² of target material) → can be **compared to experimental data** for cross-check and validation

Fast neutrons (> 0.5-1 MeV) may be measured by **scintillators** or **dark matter massive detectors** → possible only at large distances from the muon track (one should not see the muon). Delayed capture signal may be used for **neutron tagging** (e.g. **Gd-loaded scintillators** or



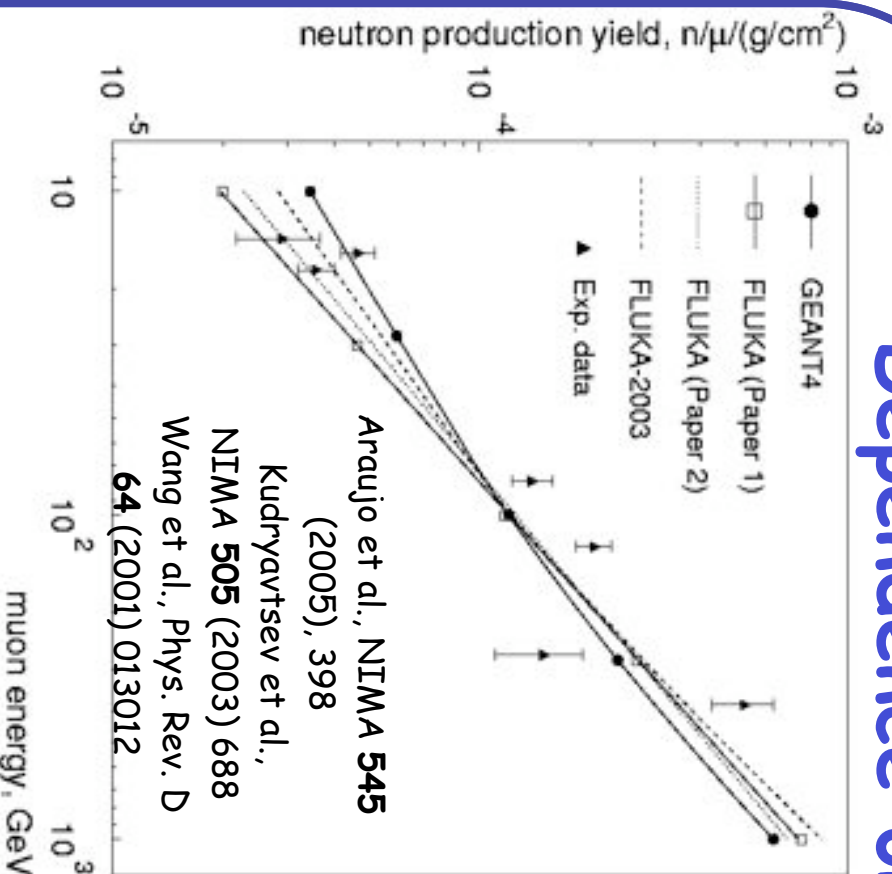
Yield difficult to measure because ³He **neutron counters** **degrades** during transport in matter → data analysis needs a **precise MC simulation**.

Ongoing TA-DUSL project (P2006-13-IUS) to measure **muon-induced neutron flux at Boubly** using ZEPLIN-II veto

Typical flux of fast muon-induced neutrons in deep underground Labs (LNGS, Boubly) $\approx 10^{-9} \text{ cm}^{-2} \text{ s}^{-1}$ (> 1 MeV) → 1000 times smaller than neutrons from rock radioactivity

Depends on **depth** (μ flux, μ energy) and **rock composition**

Dependence on muon energy



→ Neutron production rate in $(CH_n)_n$ (liquid scintillator) Fluka and Geant4 simulations compared with **experimental data**

→ good qualitative agreement, but there is **no precise simulations** of detector geometry and response for **any** of the **experiments at large depths**

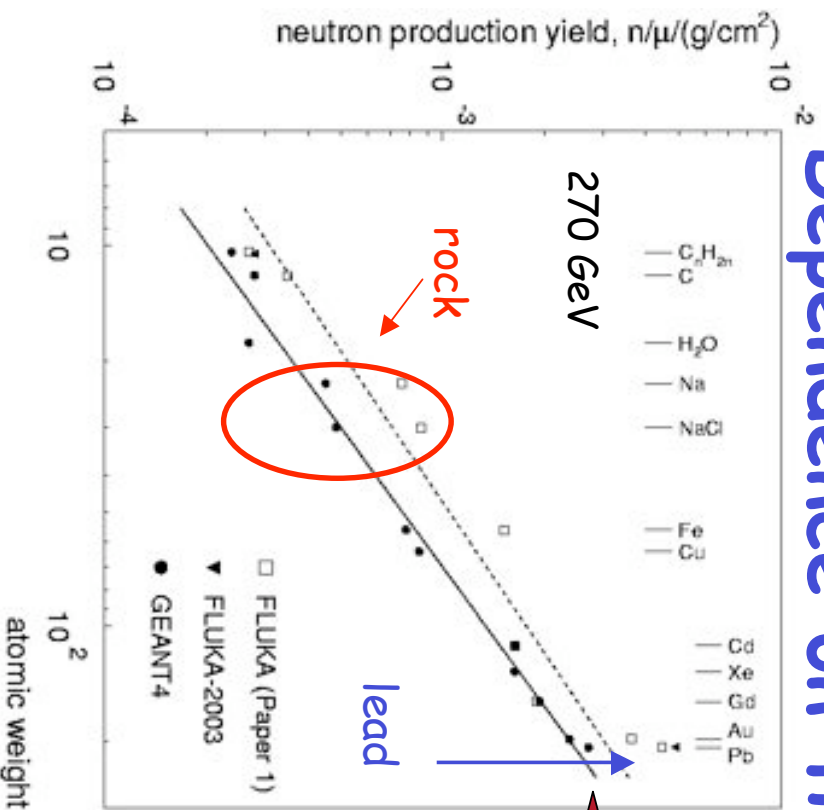
have to be taken into account for a full comparison

Different processes contribute to the neutron production: photonuclear disintegration, nucleon spallation, π^\pm spallation

For Geant4 **nuclear disintegration** by real photons is dominant at

all energies and practically for all materials

Dependence on the A of the target



Comparison of FLUKA vs.
Geant4 for **different targets**

GEANT4: Araujo et al.
NIMA **545** (2005), 398
FLUKA: Kudryavtsev et al.
NIM A **505** (2003) 688

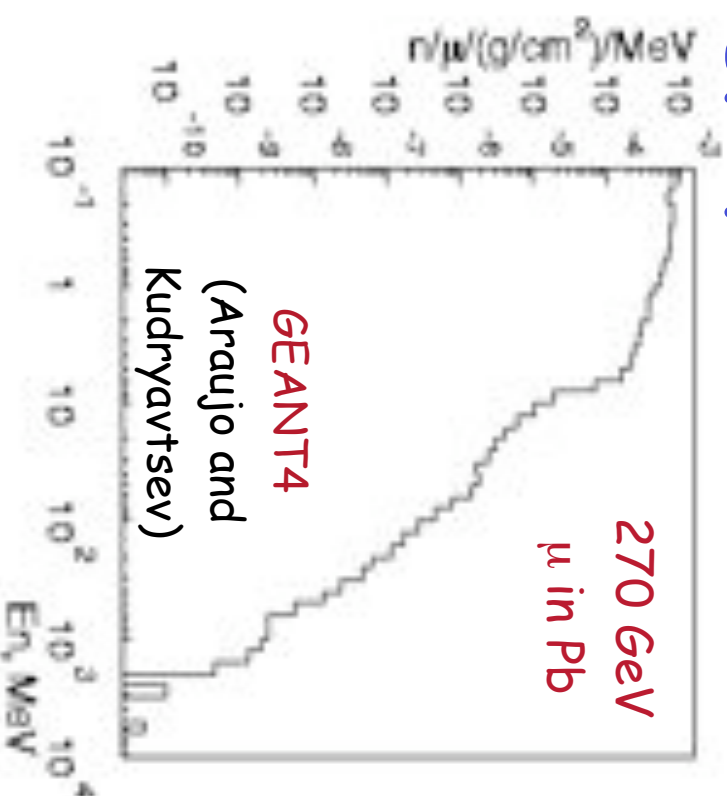
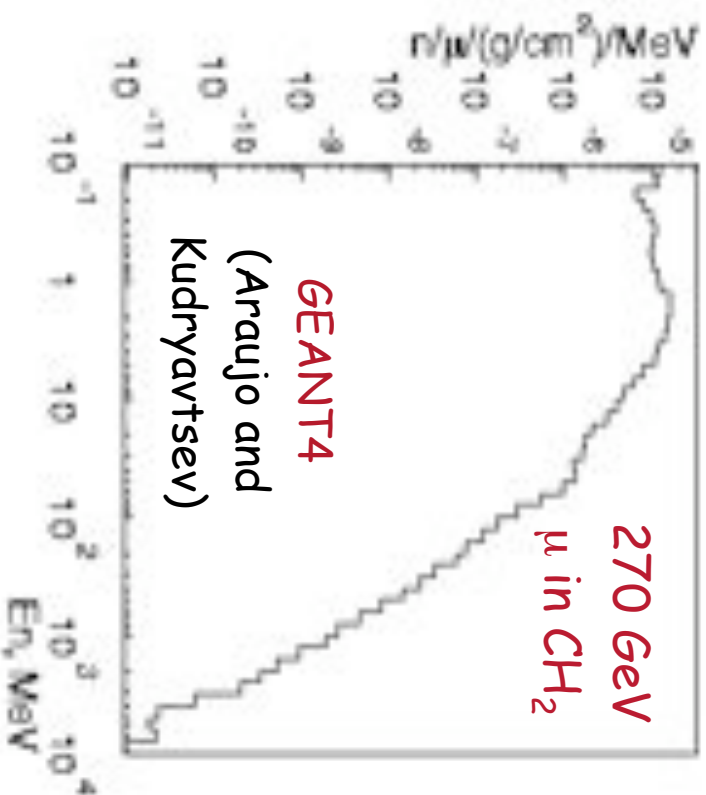
FLUKA → **twice as many neutrons**
as GEANT4 for high- A material.
Better agreement for **lower- A**

$$Y \propto A^{0.8} \cdot \rho$$

High- Z targets (= Pb used for passive γ -ray shielding) have a **higher neutron yield** → **inner** (low- Z) **neutron shielding** (CH_2 , H_2O , LN_2) to reduce muon-induced neutron flux

Notice: neutrons produced in the shielding can be **identified** if the primary μ or other secondaries **interact** in the **detector** and/or in the **veto**

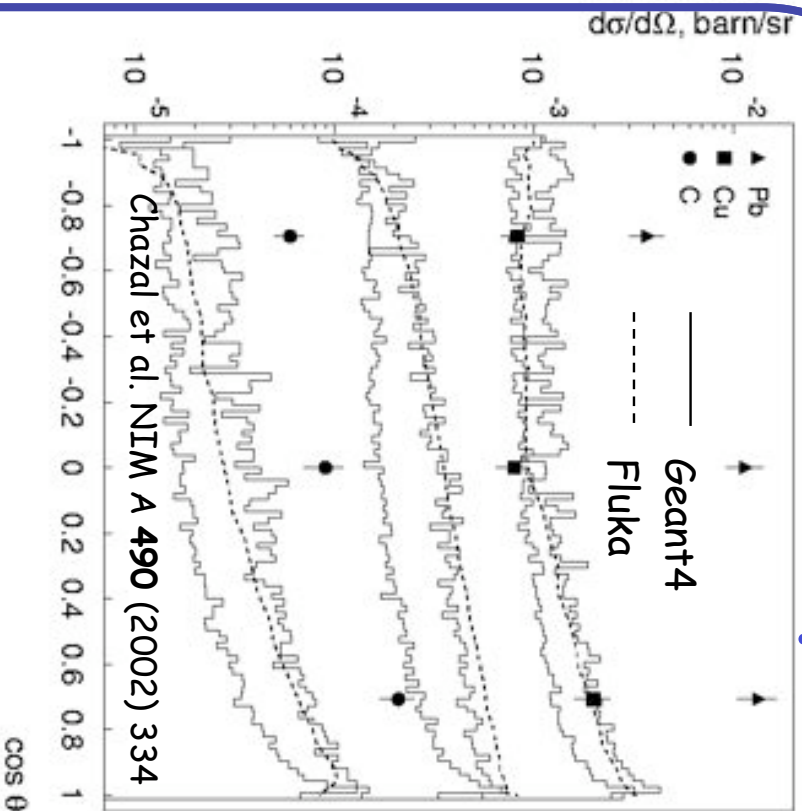
Neutron energy spectra



Energy spectra are **different** for **different** materials:

- main difference is at **low energy**: fast neutron yield (> 20 MeV) is not much different for CH_2 and Pb
- typically **higher** neutron **yield** (= high- A materials) \rightarrow **softer** spectrum

Intercomparison and validation



Differential cross section of neutron production in **thin targets** (C, Cu and Pb) for **190-GeV** muons ($E_n > 10$ MeV)
 NA55 at CERN

While Geant4 and FLUKA **agree** within a factor of two, they both seem to **underestimate**

substantially the **neutron production** Difficult to draw conclusions,

because precise **MC simulations** are **not available** for these experiments
 Other data for Pb are **old** and

controversial but also show significantly

higher neutron production
 Bergamasco et al. Nuovo Cim. A, **13** (1973) 403;
 Gorshkov et al. Sov. J. Nucl. Phys., **18** (1974) 57

Pb is important since it is a common **shielding material**!

New measurements using dark matter detectors or active veto systems, associated with

precise Monte Carlo

simulations → **ILIAS**

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Reduction of neutron background

Neutrons produced by natural radioactivity (rock & materials)

Passive neutron shielding

50 g/cm² of CH₂ reduce neutron flux from the rock by 6 orders of magnitude

Material selection for radiopurity

Neutron flux from the detector components (and shielding) dominant

Neutrons

produced by muon interactions (rock & materials, especially Pb)

Internal neutron shielding

shield μ -induced neutrons from Pb

Simultaneous detection of μ or of other secondaries in the detector (**self-veto**) and/or in active veto
self-veto is effective because neutrons are accompanied by many other secondaries

Detector design (optimization of high-Z passive materials)

Neutron-induced background → below a few events/(year·ton) for DM (10^{-10} pb sensitivity) and 10^{-3} counts/(keV·kg·y) for

DBD at the depth of the **ILIAS laboratories** (2–4 km w.e.)

Site-specific simulations

Example of **site specific** (and **detector specific**) simulation

Sampling μ according to the proper **angular** and **energy distribution** at specific site, propagating muons, generating secondaries, propagating secondaries, everything is detected

Muon-induced background is strongly **dependent** on the **experimental design**

→ **material properties** and **placement** around the detectors are important (affects muon **showering** and

propagation of secondaries)

Shielding design has to be **optimized**:

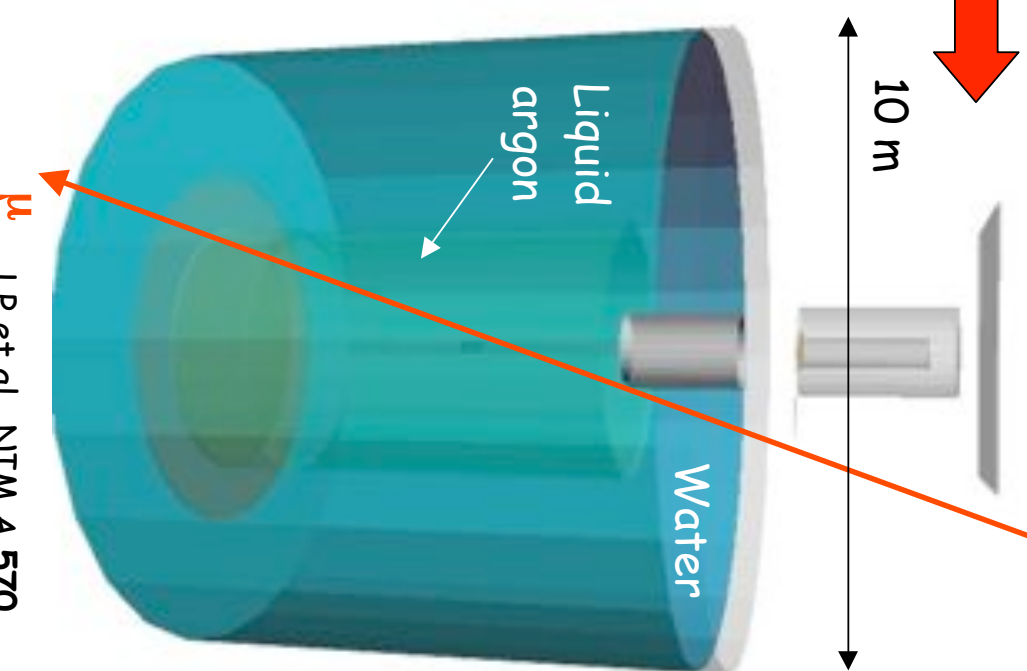
compromise between **suppression** of the **external radiation** and **reduction** of

muon-induced background in the **setup**

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GERDA DBD experiment



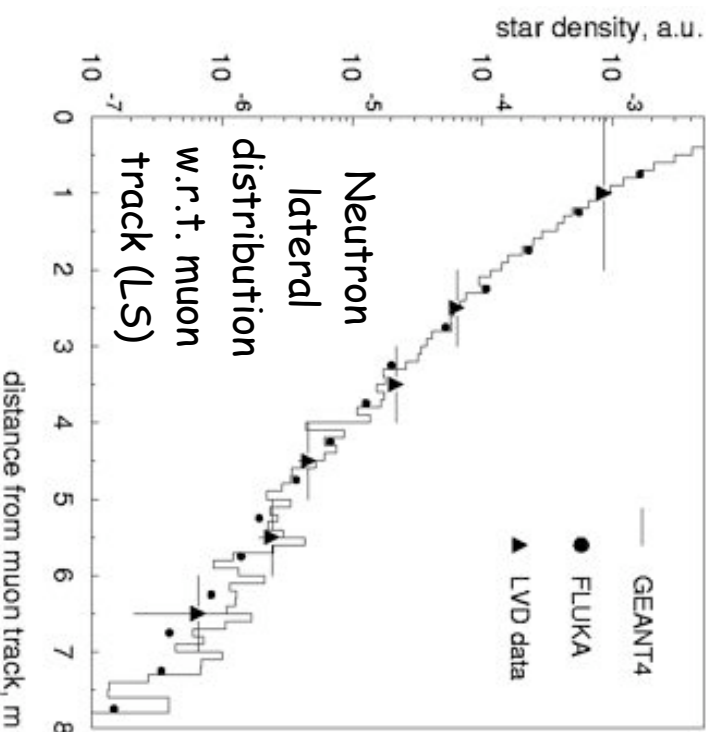
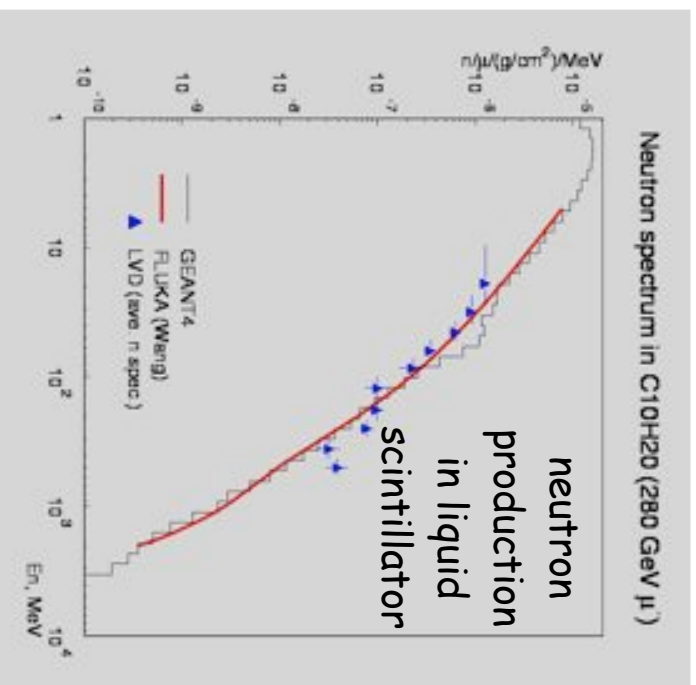
LP et al., NIM A 570
(2007) 149

Conclusions

- Neutrons produced by natural radioactivity and by muons (in **rock** or **detector** setup) can give **background** for underground experiments
- **Monte Carlo** simulations (including detector-specific effects) are critical to **define sensitivity** and **optimize rejection strategy**
- Neutron flux from **radioactivity** → reliably **predicted** (SOURCE54A and Geant4/MCNP). **Detector** components may be the main contribution to neutron background → **radiopurity** issue
- Muon-induced background **strongly dependent** on the **experimental design** (material properties and placement) → neutron shielding inside Pb. **Self-vetoing** effects are relevant
- **Geant4** and **FLUKA** agree within a **factor** of **2** for μ -induced neutrons. Needed **new data**, associated with **precise MC simulations** (→ ILIAS)
- Neutron background can be **suppressed** so that it **does not limit sensitivity** for DM or DBD at the depth of the ILIAS DULs

Backup slides

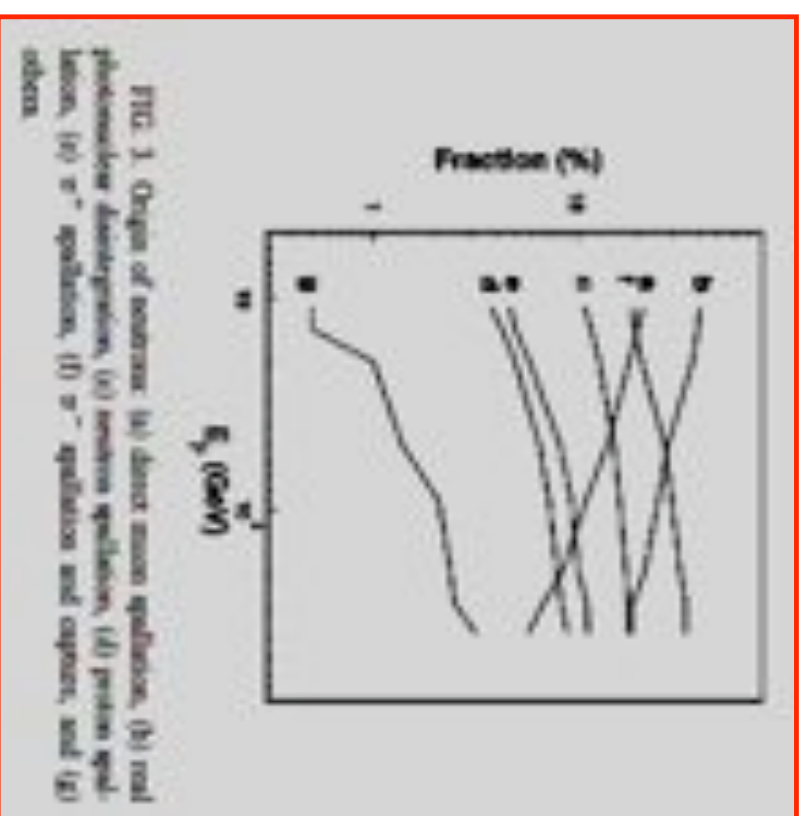
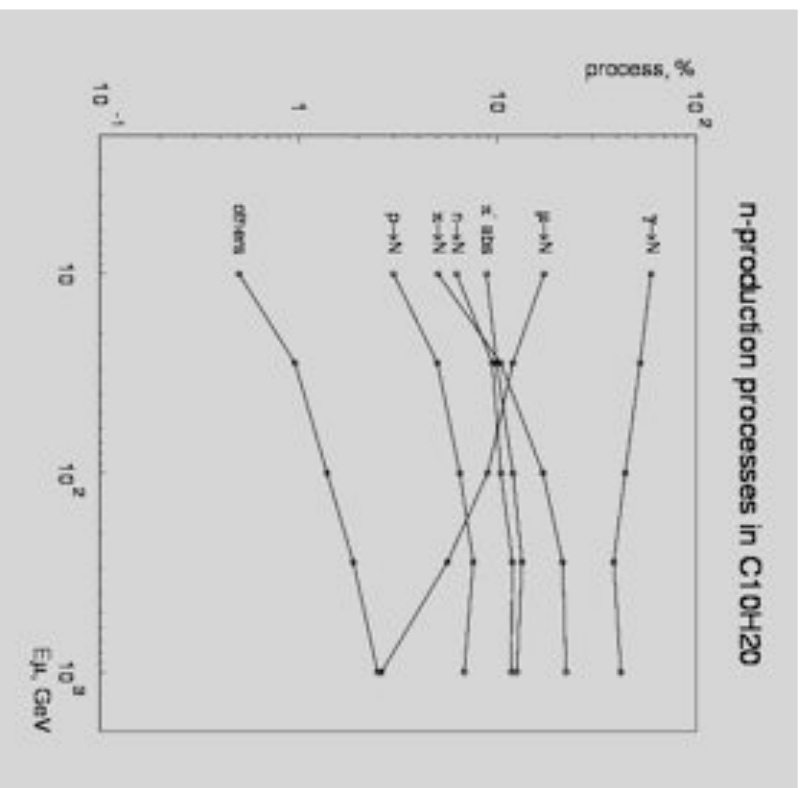
Energy spectrum & lateral distribution



LVD and FLUKA normalised to GEANT4. Spectral shapes are in **good agreement**

Lateral distribution is important for estimating **background rejection capability** → simultaneous detection of the **parent muon** or **other secondary radiation** (in main detector or veto)

Muon-induced neutrons: processes



GEANT4: Araujo et al. NIMA 545 (2005), 398

FLUKA: Wang et al. PRD, 64 (2001) 013012

Real photonuclear disintegration dominates in Geant4 at all energies and for (almost) all materials

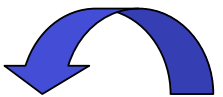
Isotope production & delayed background

n or μ interactions can produce **long-lived** ($T_{1/2} > 1$ s) **unstable isotopes**.

Their decay (**not vetable** because delayed) is a **background for DBD experiments** if $Q > Q_{\beta\beta}$ (not for DM because of γ/n discrimination)

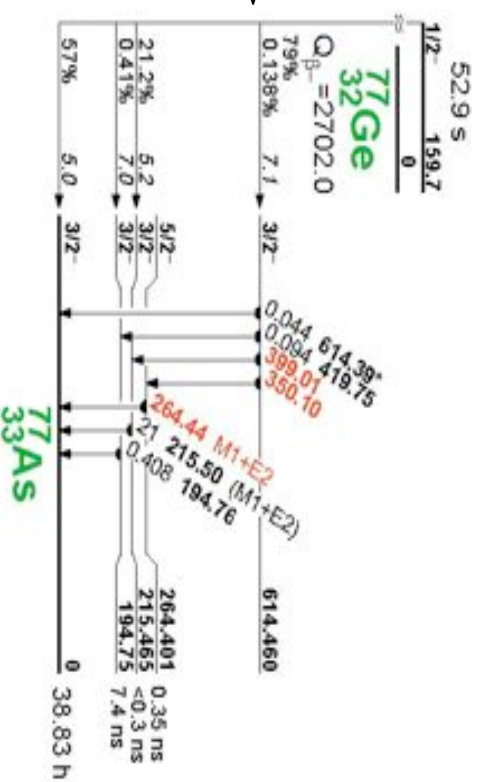
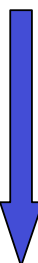
For ^{76}Ge -based DBD experiments

(= GERDA), the most dangerous are ^{77}Ge ($Q=2.7$ MeV, $T_{1/2} = 11.3$ h) and ^{77m}Ge ($Q=2.8$ MeV, $T_{1/2} = 53$ s)



Thermal neutron capture, high cross section (0.14 barn), **scales with enrichment**

Muon veto and self-veto are **uneffective**: **specific cuts** (delayed coincidence using prompt γ -rays) **required** to reduce background



The production rate in GERDA changes by a **factor of 10** using **LAr** or **LN₂**