

HOFEM-AIRBUS

Transmission/Reflection

Group of Radiofrequency, Electromagnetism, Microwaves
and Antennas (GREMA)

<http://grema.webs.tsc.uc3m.es/>

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Outline

- 1 On the FEM Implementation of TX/RX Conditions in HOFEM
 - Existing FEM Formulation in HOFEM
 - Two-Port Network Parameters

- 2 $[Z]/[Y]$ Approach
 - FEM Formulation
 - Testing
 - HOFEM Implementation

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We start considering different alternatives to implement the TX/RX conditions in the context of the present FEM formulation coded in HOFEM

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FEM Formulation

- Formulation based on double curl vector wave equation (use of **E** or **H**).

$$\nabla \times (f_r^{-1} \nabla \times \mathbf{V}) - k_0^2 g_r \mathbf{V} = -jk_0 H_0 \mathbf{P} + \nabla \times f_r^{-1} \mathbf{L}$$

Table: Formulation magnitudes and parameters

	V	V^d	$\bar{\bar{f}}_r$	$\bar{\bar{g}}_r$	<i>h</i>	P	L
Form. E	E	H	$\bar{\bar{\mu}}_r$	$\bar{\bar{\epsilon}}_r$	η	J	M
Form. H	H	E	$\bar{\bar{\epsilon}}_r$	$\bar{\bar{\mu}}_r$	$-\frac{1}{\eta}$	M	-J

FEM Formulation (cont.)

- Use of $\mathbf{H}(\text{curl})$ spaces:

$$\mathbf{H}(\text{curl})_0 = \{\mathbf{W} \in \mathbf{H}(\text{curl}), \hat{\mathbf{n}} \times \mathbf{W} = 0 \text{ on } \Gamma_D\} \quad (1)$$

$$\mathbf{H}(\text{curl}) = \{\mathbf{W} \in L^2, \nabla \times \mathbf{W} \in L^2\} \quad (2)$$

- and Galerkin method leads to (with respect to the double-curl term only) to:

$$\int_{\Omega} (\nabla \times \mathbf{F}) \cdot \left(\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V} \right) d\Omega + \underbrace{\int_{\Gamma} \mathbf{F} \cdot \left(\hat{\mathbf{n}} \times \underbrace{\left(\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V} \right)}_{-jk_0 h_0 \mathbf{V}^d} \right) d\Gamma}_{\text{"natural" b.c.}}$$

- More precisely

- ▶ Form. \mathbf{E} : $\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V} = -jk_0 \eta_0 \mathbf{H}$
- ▶ Form. \mathbf{H} : $\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V} = +j \frac{k_0}{\eta_0} \mathbf{E}$

FEM Formulation (cont.)

The integral boundary term $\int_{\Gamma} \mathbf{F} \cdot (\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V})) d\Gamma$

- Can be used to weakly impose some boundary conditions of the problem
- For instance, Neumann boundary condition

$$\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V}) = \boldsymbol{\Psi}_N$$
$$\longrightarrow \int_{\Gamma} \mathbf{F} \cdot (\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V})) d\Gamma = \int_{\Gamma} \boldsymbol{\Psi}_N d\Gamma$$

- For instance, ABC boundary condition

$$\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V}) + \gamma \hat{\mathbf{n}} \times \hat{\mathbf{n}} \times \mathbf{V} = \boldsymbol{\Psi}_C$$
$$\longrightarrow \int_{\Gamma} \mathbf{F} \cdot (\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V})) d\Gamma = \int_{\Gamma} \boldsymbol{\Psi}_C d\Gamma + \gamma \int_{\Gamma} (\hat{\mathbf{n}} \times \mathbf{F}) \cdot (\hat{\mathbf{n}} \times \mathbf{V}) d\Gamma$$

FEM Formulation (cont.)

- Note that the implementation of above boundary conditions is straightforward because
 - ▶ Either, the integrand of the boundary term was a known function
 - ▶ or, it can be expressed in terms of the primal unknown \mathbf{V}
- At the end of the day, the contribution of the boundary term is translated into algebra as
 - ▶ extra contributions for the values of the matrix coefficients (related to the existing g_i associated to the corresponding boundary).
 - ▶ I.e., no extra degrees of freedom g_i are needed

What does it happen when that is not the case?

FEM Continuity Conditions

- Assuming $\mathbf{V} = \sum_i g_i \mathbf{N}_i \leftrightarrow \{g\}$

Between Neighbour Elements 1 and 2

- Two continuity conditions:

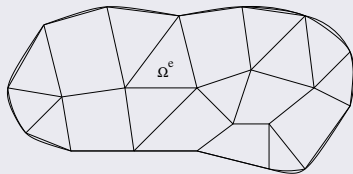
- Strong continuity of $\hat{\mathbf{n}} \times \mathbf{V}$

$$\{g\}_1 = \{g\}_2$$

- Weak continuity of $\hat{\mathbf{n}} \times \mathbf{V}^d$

$$\int_{\Gamma} \mathbf{F} \cdot (\hat{\mathbf{n}}_1 \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V}_1)) d\Gamma \\ + \int_{\Gamma} \mathbf{F} \cdot (\hat{\mathbf{n}}_2 \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V}_2)) d\Gamma = 0$$

We simply omit the term



Note that $\hat{\mathbf{n}} \times \mathbf{F}_1 = \hat{\mathbf{n}} \times \mathbf{F}_2$ and that

$$\mathbf{F} \cdot (\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V})) = (\hat{\mathbf{n}} \times \mathbf{F}) \cdot (\hat{\mathbf{n}} \times (\bar{\bar{f}}_r^{-1} \nabla \times \mathbf{V}))$$

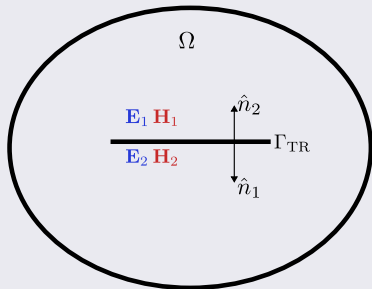
FEM Continuity Conditions (cont.)

Between Elements Simulating a Thin Sheet

- In general, we need to break continuity on
 - ▶ $\hat{\mathbf{n}} \times \mathbf{V} \Rightarrow$ Duplication of degrees of freedom: $\{g\}_1, \{g\}_2$
 - ▶ $\hat{\mathbf{n}} \times \mathbf{V}^d \Rightarrow$ The integral boundary term can not be omitted
- Two linearly independent equations involving four magnitudes $\mathbf{E}_1, \mathbf{E}_2, \mathbf{H}_1, \mathbf{H}_2$:

$$[]_{2 \times 2} \{X\}_{2 \times 1} = \{B\}_{2 \times 1}$$

- Different possibilities depending on the magnitudes chosen as sources (right hand side B) and responses X
- Clear analogy with two-port network parameters of linear circuit analysis:
 $\mathbf{E} \leftrightarrow V \quad \mathbf{H} \leftrightarrow I$



The Simplest (Least Invasive) Approach

The Immittance Approach

- We express $\hat{\mathbf{n}} \times \mathbf{V}^d$ for region 1 and region 2 in terms of $\hat{\mathbf{n}} \times \mathbf{V}$ of region 1 and region 2, i.e.,

$$\begin{Bmatrix} \hat{\mathbf{n}} \times \mathbf{V}_1^d \\ \hat{\mathbf{n}} \times \mathbf{V}_2^d \end{Bmatrix} = \begin{bmatrix} YZ_{11} & YZ_{12} \\ YZ_{21} & YZ_{22} \end{bmatrix} \begin{Bmatrix} \hat{\mathbf{n}} \times \mathbf{V}_1 \\ \hat{\mathbf{n}} \times \mathbf{V}_2 \end{Bmatrix}$$

- ▶ This is equivalent to an immittance characterization of the equivalent two-port network connecting upper and lower regions separated by the sheet.
 - ★ Form. **E**: $\mathbf{V} = \mathbf{E}$, $\mathbf{V}^d = \mathbf{H}$, immittance \equiv admittance
 - ★ Form. **E**: $\mathbf{V} = \mathbf{E}$, $\mathbf{V}^d = \mathbf{H}$, immittance \equiv impedance

The Simplest (Least Invasive) Approach (cont.)

The Immittance Approach

- We substitute. . .

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TITULO

Hola

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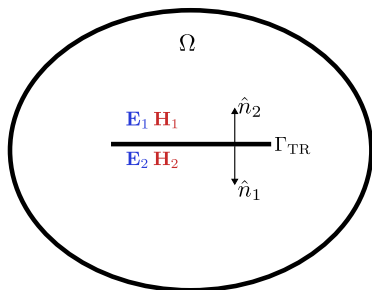
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We describe the implementation of the TX/RX conditions using the characterization of the material sheet in terms of its immittance (impedance/admittance) matrix

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Formulation



$$\hat{n}_1 \times (\mu_r^{-1} \nabla \times \mathbf{E}_1) - \frac{jk_0}{\eta} y_{11} \hat{n}_1 \times (\hat{n}_1 \times \mathbf{E}_1) - \\ - \frac{jk_0}{\eta} y_{12} \hat{n}_2 \times (\hat{n}_2 \times \mathbf{E}_2) = 0,$$

$$\hat{n}_2 \times (\mu_r^{-1} \nabla \times \mathbf{E}_2) - \frac{jk_0}{\eta} y_{21} \hat{n}_1 \times (\hat{n}_1 \times \mathbf{E}_1) - \\ - \frac{jk_0}{\eta} y_{22} \hat{n}_2 \times (\hat{n}_2 \times \mathbf{E}_2) = 0,$$

Note that y_{xx} are relative to the vacuum admittance.

Formulation (cont.)

Find $\mathbf{E} \in \mathbf{H}_0(\text{curl}, \Omega)$ such that

$$\begin{aligned} & \left(\nabla \times \mathbf{w}, \mu_r^{-1} \nabla \times \mathbf{E} \right)_{\Omega} - k_0^2 \left(\mathbf{w}, \varepsilon_r \mathbf{E} \right)_{\Omega} + jk_0 \left\langle \hat{n} \times \mathbf{w}, \hat{n} \times \mathbf{w} \right\rangle_{\Gamma_c} = \\ & \left(\mathbf{w}, \mathbf{F} \right)_{\Omega} - \left\langle \hat{n} \times (\mathbf{w} \times \hat{n}), \boldsymbol{\Psi}_N \right\rangle_{\Gamma_N} - \left\langle \hat{n} \times (\mathbf{w} \times \hat{n}), \boldsymbol{\Psi}_C \right\rangle_{\Gamma_c} \quad \forall \mathbf{w} \in \mathbf{H}_0(\text{curl}, \Omega). \end{aligned}$$

with

$$\begin{aligned} \left(\mathbf{w}, \mathbf{v} \right)_{\Omega} &= \int_{\Omega} \mathbf{w}^* \cdot \mathbf{v} d\Omega, \\ \left\langle \mathbf{w}, \mathbf{v} \right\rangle_{\Gamma} &= \int_{\Gamma} \mathbf{w}^* \cdot \mathbf{v} d\Gamma. \end{aligned}$$

Formulation (cont.)

For *upper* elements on Γ_{TR} (side 1), we have

$$\begin{aligned} \text{LHS}_1 \\ + j \frac{k_0}{\eta} \left\langle \hat{n} \times (\mathbf{w}_1 \times \hat{n}), y_{11} \hat{n} \times (\mathbf{w}_1 \times \hat{n}) \right\rangle_{\Gamma_{\text{TR}}} + j \frac{k_0}{\eta} \left\langle \hat{n} \times (\mathbf{w}_1 \times \hat{n}), y_{12} \hat{n} \times (\mathbf{w}_2 \times \hat{n}) \right\rangle_{\Gamma_{\text{TR}}} = \\ \text{RHS}_1, \end{aligned}$$

whereas for *lower* elements (side 2), we get

$$\begin{aligned} \text{LHS}_2 \\ + j \frac{k_0}{\eta} \left\langle \hat{n} \times (\mathbf{w}_2 \times \hat{n}), y_{21} \hat{n} \times (\mathbf{w}_1 \times \hat{n}) \right\rangle_{\Gamma_{\text{TR}}} + j \frac{k_0}{\eta} \left\langle \hat{n} \times (\mathbf{w}_2 \times \hat{n}), y_{22} \hat{n} \times (\mathbf{w}_2 \times \hat{n}) \right\rangle_{\Gamma_{\text{TR}}} = \\ \text{RHS}_2, \end{aligned}$$

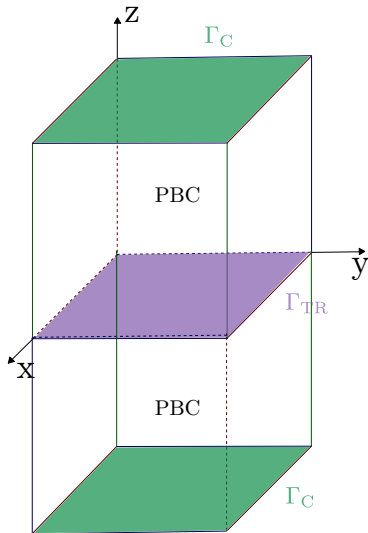
FEM implementation

- The DOFs will be doubled for the faces and the interior edges.
- The exterior edges of Γ_{TR} are not doubled.
 - ▶ Identified by code: the edges associated to two faces are interior.
 - ▶ If the boundaries of the sheet belong to PBC, the edges of Γ_{TR} are also doubled.

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Problem to be solved



Simulation of an infinite medium with transmission/reflection sheet that divides the space into two halves.

- Γ_{TR} : Transmission/reflection sheet defined with

$$\mathbf{Y} = \begin{bmatrix} y_{11} & y_{12} \\ y_{21} & y_{22} \end{bmatrix}.$$

- Γ_C : ABC with excitation with polarization E_y
- The vertical faces are set to PBC

Testbench

- $\mathbf{Y} = \begin{bmatrix} 0 & 0 \\ 0 & 0 \end{bmatrix}$: sanity check, we should get same result as the two halves with a PMC.
- $\mathbf{Y} = \mathbb{I}$: sanity check, we should get same result as the two halves with an ABC.
- Change lower Γ_C by PEC and solve analytic problem with four media: final test.
 - ▶ Obtain parameters for \mathbf{Y} of the equivalent problem.
 - ▶ Get same solutions for the electric field.
 - ▶ Transparent? Puede ser que aproximar con $1e6$. Quizás con ABCD.

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HOFEM implementation

- New boundary condition: TRBC.
 - ▶ We define a normal, \hat{n}_{TRBC} to detect lower and upper side. Upper side is the closer to \hat{n}_{TRBC} .
 - ▶ Definition of y_{11} , y_{12} , y_{21} , and y_{22} as relative values with respect to vacuum admittance.
- Two options for implementation
 - ▶ Integers defined in `tetrahedra_element`.
 - ▶ Allocatable array of $1 \times N_{\text{elem,TR}}$ where the two positions (stored in boundary conditions module, accessible from `mesh_reordering_module` and `elementary_terms_3D`):
 - ① $10 \times$ Neighbor element identifier (to couple \mathbf{w}_2 and \mathbf{w}_1).
 - ② Integer 1,2 (side) (to extract the values of y_{11}, y_{12}, y_{21} , and y_{22}).
- Significant methods involved:
 - ▶ `Postprocessing over reordering_DOF_algorithm_3D`.
 - ▶ `calc_boundary_3D_nxNi_nxNi_term_of_this_element`.
 - ▶ Construction of the MUMPS-related matrix: different number of non-zeros per element, assembly of coupled elements (now single-element assembly).