

Lectures on Random Matrices (Spring 2025)

Lecture 6: Double contour integral kernel. Steepest descent and local statistics

Leonid Petrov

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Notes for the lecturer

1. GUE det structure
2. Formulate Cauchy–Binet and Andreief
3. Recall that $\rho_n = P_n$ and it is $(\det[\psi_i(x_j)]_{n \times n})^2$, then reproduce the proofs here.

1 Recap: Determinantal structure of the GUE

Last time, we proved the following result:

Theorem 1.1. *The GUE correlation functions are given by*

$$\rho_k(x_1, \dots, x_k) = \det \left[K_n(x_i, x_j) \right]_{i,j=1}^k,$$

with the correlation kernel

$$K_n(x, y) = \sum_{j=0}^{n-1} \psi_j(x) \psi_j(y).$$

Here

$$\psi_j(x) = \frac{1}{\sqrt{h_j}} p_j(x) e^{-x^2/4},$$

where $p_j(x)$ are the monic Hermite polynomials, and h_j are the normalization constants so that $\psi_j(x)$ are orthonormal in $L^2(\mathbb{R})$.

For this theorem, we need Cauchy–Binet summation formula and Andreief identity (which is essentially the same as Cauchy–Binet, but when summation is replaced by integration). Having these, we can write

$$\rho_k(x_1, \dots, x_k) = \frac{n!}{(n-k)!} \int_{\mathbb{R}^{n-k}} p(x_1, \dots, x_n) dx_{k+1} \cdots dx_n$$

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$$\begin{aligned}
&= \frac{1}{(n-k)! \widehat{Z}_{n,2}} \sum_{\substack{\sigma, \tau \in S_n \\ \sigma(k+1)=\tau(k+1), \dots, \sigma(n)=\tau(n)}} \text{sgn}(\sigma) \text{sgn}(\tau) \prod_{i=1}^k \psi_{\sigma(i)-1}(x_i) \psi_{\tau(i)-1}(x_i) \\
&= \text{const}_n \sum_{I \subseteq [n], |I|=k} \sum_{\sigma', \tau' \in S(I)} \text{sgn}(\sigma') \text{sgn}(\tau') \prod_{i=1}^k \psi_{\sigma'(i)-1}(x_i) \psi_{\tau'(i)-1}(x_i) \\
&= \text{const}_n \sum_{I \subseteq [n], |I|=k} \det [\psi_{i_\alpha}(x_j)]_{\alpha, j=1}^k \det [\psi_{i_\alpha}(x_j)]_{\alpha, j=1}^k,
\end{aligned}$$

where $I = \{i_1, \dots, i_k\}$ is a subset of $[n]$ of size k , and $S(I)$ is the set of permutations of I . The last sum of products of two determinants is written by the Cauchy–Binet formula as

$$\text{const}_n \cdot \det \left[\sum_{j=0}^{n-1} \psi_j(x_\alpha) \psi_j(x_\beta) \right]_{\alpha, \beta=1}^k,$$

and finally the constant is equal to 1 by Andreief identity.

2 Double Contour Integral Representation for the GUE Kernel

2.1 One contour integral representation for Hermite polynomials

Recall that the GUE kernel is defined by

$$K_N(x, y) = \sum_{n=0}^{N-1} \psi_n(x) \psi_n(y),$$

with the orthonormal functions

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4},$$

where the (monic, probabilists') Hermite polynomials are given by

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} e^{-x^2/2},$$

and satisfy the generating function

$$\exp\left(xt - \frac{t^2}{2}\right) = \sum_{n \geq 0} p_n(x) \frac{t^n}{n!}.$$

By Cauchy's integral formula we can write

$$p_n(x) = \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt,$$

where the contour C is a simple closed curve encircling the origin. Therefore,

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4} = \frac{e^{-x^2/4}}{\sqrt{h_n}} \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt.$$

2.2 Another contour integral representation for Hermite polynomials

Note also that

$$\int_{-\infty}^{\infty} e^{-t^2 + \sqrt{2}i t x} dt = \sqrt{\pi} e^{-x^2/2}.$$

Differentiating both sides n times with respect to x (and using the fact that in our convention the Gaussian appears with $x^2/2$) yields

$$\frac{d^n}{dx^n} \left(e^{-x^2/2} \right) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \left(\sqrt{2}i t \right)^n e^{-t^2 + \sqrt{2}i t x} dt.$$

Changing variables via $s = i t$ (so that $t = -i s$ and $dt = -i ds$) one obtains

$$\frac{d^n}{dx^n} \left(e^{-x^2/2} \right) = \frac{(\sqrt{2})^n}{i\sqrt{\pi}} \int_{-i\infty}^{i\infty} s^n e^{s^2 + \sqrt{2} s x} ds.$$

Multiplying by $(-1)^n e^{x^2/2}$ we deduce that

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} \left(e^{-x^2/2} \right) = \frac{i(\sqrt{2})^n e^{x^2/2}}{\sqrt{\pi}} \int_{-i\infty}^{i\infty} s^n e^{s^2 - \sqrt{2} s x} ds. \quad (2.1)$$

Now, recall that the orthonormal functions are defined as

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4},$$

so that by (2.1)

$$\psi_n(x) = \frac{i e^{x^2/4}}{\sqrt{\pi} h_n} \int_{-i\infty}^{i\infty} (\sqrt{2} s)^n e^{s^2 - \sqrt{2} s x} ds = \frac{i e^{x^2/4}}{\sqrt{2\pi} h_n} \int_{-i\infty}^{i\infty} s^n e^{s^2/2 - s x} ds.$$

2.3 Double contour integral representation for the GUE kernel

We have (Problem ??)

$$h_n = \int_{-\infty}^{\infty} p_n(x)^2 e^{-x^2/2} dx = n! \sqrt{2\pi}.$$

Therefore, we can sum up the kernel (another proof of the Christoffel–Darboux formula):

$$\begin{aligned} K_n(x, y) &= \sum_{k=0}^{n-1} \psi_k(x) \psi_k(y) \\ &= \sum_{k=0}^{n-1} \frac{e^{-x^2/4}}{\sqrt{h_k}} \frac{k!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{k+1}} dt \frac{i e^{y^2/4}}{\sqrt{2\pi} h_k} \int_{-i\infty}^{i\infty} s^k e^{s^2/2 - s y} ds \\ &= e^{(y^2 - x^2)/4} \sum_{k=0}^{n-1} \frac{1}{4\pi^2} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{k+1}} dt \int_{-i\infty}^{i\infty} s^k e^{s^2/2 - s y} ds. \end{aligned}$$

We can now extend the sum to $k = -\infty$, and get a formula for the GUE kernel as a double contour integral:

$$K_n(x, y) = \frac{e^{(y^2-x^2)/4}}{4\pi^2} \oint_C \int_{-\mathrm{i}\infty}^{\mathrm{i}\infty} \frac{\exp\left\{\frac{s^2}{2} - sy - \frac{t^2}{2} + tx\right\}}{s-t} \left(\frac{s}{t}\right)^n ds dt.$$

Details will be in the next [Lecture 6](#).

Remark 2.1. Many other versions of the GUE / unitary invariant ensembles admit determinantal structure:

1. The GUE corners process [[JN06](#)]
2. The Dyson Brownian motion (e.g., add a GUE to a diagonal matrix) [[NF98](#)]
3. GUE corners plus a fixed matrix [[FF14](#)]
4. Corners invariant ensembles with fixed eigenvalues UDU^\dagger , where D is a fixed diagonal matrix and U is Haar distributed on the unitary group [[Met13](#)]

F Problems (due DATE)

References

- [FF14] P. Ferrari and R. Frings, *Perturbed GUE minor process and Warren's process with drifts*, J. Stat. Phys **154** (2014), no. 1-2, 356–377. [arXiv:1212.5534 \[math-ph\]](#). [↑4](#)
- [JN06] K. Johansson and E. Nordenstam, *Eigenvalues of GUE minors*, Electron. J. Probab. **11** (2006), no. 50, 1342–1371. [arXiv:math/0606760 \[math.PR\]](#). [↑4](#)
- [Met13] A. Metcalfe, *Universality properties of Gelfand-Tsetlin patterns*, Probab. Theory Relat. Fields **155** (2013), no. 1-2, 303–346. [arXiv:1105.1272 \[math.PR\]](#). [↑4](#)
- [NF98] T. Nagao and P.J. Forrester, *Multilevel dynamical correlation functions for Dyson's Brownian motion model of random matrices*, Physics Letters A **247** (1998), no. 1-2, 42–46. [↑4](#)

L. PETROV, UNIVERSITY OF VIRGINIA, DEPARTMENT OF MATHEMATICS, 141 CABELL DRIVE, KERCHOF HALL, P.O. BOX 400137, CHARLOTTESVILLE, VA 22904, USA
E-mail: lenia.petrov@gmail.com