## Lectures on Random Matrices (Spring 2025)

# Lecture 6: Double contour integral kernel. Steepest descent and local statistics

Leonid Petrov

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#### Notes for the lecturer

- GUE det structure
- Formulate Cauchy–Binet and Andreief
- Recall that  $\rho_n = P_n$  and it is  $(\det[\psi_i(x_j)]_{n \times n})^2$ , then reproduce the proofs here.
- Recall the Christoffel–Darboux formula:

$$K_n(x,y) = \frac{e^{-\frac{x^2+y^2}{4}}}{\sqrt{2\pi}h_{n-1}} \frac{p_n(x)p_{n-1}(y) - p_{n-1}(x)p_n(y)}{x - y},$$

here  $h_{n-1} = \sqrt{2\pi}(n-1)!$ .

## 1 Recap: Determinantal structure of the GUE

Last time, we proved the following result:

**Theorem 1.1.** The GUE correlation functions are given by

$$\rho_k(x_1,\ldots,x_k) = \det\left[K_n(x_i,x_j)\right]_{i,j=1}^k,$$

with the correlation kernel

$$K_n(x,y) = \sum_{j=0}^{n-1} \psi_j(x)\psi_j(y).$$

Here

$$\psi_j(x) = \frac{1}{\sqrt{h_j}} p_j(x) e^{-x^2/4},$$

where  $p_j(x)$  are the monic Hermite polynomials, and  $h_j$  are the normalization constants so that  $\psi_j(x)$  are orthonormal in  $L^2(\mathbb{R})$ .

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For this theorem, we need Cauchy–Binet summation formula and Andreief identity (which is essentially the same as Cauchy–Binet, but when summation is replaced by integration). Having these, we can write

$$\rho_k(x_1, \dots, x_k) = \frac{n!}{(n-k)!} \int_{\mathbb{R}^{n-k}} p(x_1, \dots, x_n) \, dx_{k+1} \cdots dx_n$$

$$= \frac{1}{(n-k)!} \sum_{\substack{\sigma, \tau \in S_n \\ \sigma(k+1) = \tau(k+1), \dots, \sigma(n) = \tau(n)}} \operatorname{sgn}(\sigma) \operatorname{sgn}(\tau) \prod_{i=1}^k \psi_{\sigma(i)-1}(x_i) \psi_{\tau(i)-1}(x_i)$$

$$= \operatorname{const}_n \sum_{I \subseteq [n], |I| = k} \sum_{\sigma', \tau' \in S(I)} \operatorname{sgn}(\sigma') \operatorname{sgn}(\tau') \prod_{i=1}^k \psi_{\sigma'(i)-1}(x_i) \psi_{\tau'(i)-1}(x_i)$$

$$= \operatorname{const}_n \sum_{I \subseteq [n], |I| = k} \det \left[ \psi_{i_\alpha}(x_j) \right]_{\alpha, j=1}^k \det \left[ \psi_{i_\alpha}(x_j) \right]_{\alpha, j=1}^k,$$

where  $I = \{i_1, \dots, i_k\}$  is a subset of [n] of size k, and S(I) is the set of permutations of I. The last sum of products of two determinants is written by the Cauchy-Binet formula as

$$\operatorname{const}_n \cdot \det \left[ \sum_{j=0}^{n-1} \psi_j(x_\alpha) \psi_j(x_\beta) \right]_{\alpha, \beta = 1}^k,$$

and finally the constant is equal to 1 by Andreief identity.

## 2 Double Contour Integral Representation for the GUE Kernel

#### 2.1 One contour integral representation for Hermite polynomials

Recall that the GUE kernel is defined by

$$K_N(x,y) = \sum_{n=0}^{N-1} \psi_n(x)\psi_n(y),$$

with the orthonormal functions

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4},$$

where the (monic, probabilists') Hermite polynomials are given by

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} e^{-x^2/2}.$$
 (2.1)

Note that the monic Hermite polynomials are uniquely defined by the orthogonality property. We are not proving (2.1) here, it is an exercise.

**Lemma 2.1** (Generator function for Hermite polynomials). We have

$$\exp\left(xt - \frac{t^2}{2}\right) = \sum_{n>0} p_n(x) \frac{t^n}{n!}.$$

The series converges for all t since the left-hand side is an entire function of t.

*Proof.* Write the generating function as

$$\sum_{n>0} p_n(x) \frac{t^n}{n!} = \sum_{n>0} \frac{(-1)^n t^n}{n!} e^{x^2/2} \frac{d^n}{dx^n} e^{-x^2/2}.$$

Since the factor  $e^{x^2/2}$  does not depend on n, we can factor it out:

$$\sum_{n\geq 0} p_n(x) \frac{t^n}{n!} = e^{x^2/2} \sum_{n\geq 0} \frac{(-t)^n}{n!} \frac{d^n}{dx^n} e^{-x^2/2}.$$

Now, recall Taylor's theorem: for any analytic function f we have

$$f(x-t) = \sum_{n>0} \frac{(-t)^n}{n!} f^{(n)}(x).$$

Applying this with  $f(x) = e^{-x^2/2}$ , we deduce that

$$\sum_{n>0} \frac{(-t)^n}{n!} \frac{d^n}{dx^n} e^{-x^2/2} = e^{-(x-t)^2/2}.$$

Thus, our generating function becomes

$$\sum_{n>0} p_n(x) \frac{t^n}{n!} = e^{x^2/2} e^{-(x-t)^2/2},$$

as desired.  $\Box$ 

By Cauchy's integral formula we can write using Lemma 2.1:

$$p_n(x) = \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt,$$
 (2.2)

where the contour C is a simple closed curve encircling the origin. Indeed, here we use the complex analysis property

$$\frac{1}{2\pi i} \oint_C \frac{1}{z^{k+1}} dz = \begin{cases} 1, & \text{if } k = 0, \\ 0, & \text{if } k \neq 0, \end{cases}$$

so (2.2) is simply a complex analysis version of the operation of extracting the coefficient of  $t^n$  in the Taylor expansion.

Therefore,

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4} = \frac{e^{-x^2/4}}{\sqrt{h_n}} \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt.$$

#### 2.2 Another contour integral representation for Hermite polynomials

We start with the Fourier transform identity

$$\int_{-\infty}^{\infty} \exp\left(-\frac{t^2}{2} + i t x\right) dt = \sqrt{2\pi} e^{-x^2/2}.$$

Differentiating both sides n times with respect to x yields

$$\frac{d^n}{dx^n} \left( e^{-x^2/2} \right) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (i\,t)^n \, e^{-t^2/2 + i\,t\,x} \, dt.$$

Recalling the definition

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} \left(e^{-x^2/2}\right),$$

we obtain

$$p_n(x) = \frac{(-1)^n e^{x^2/2}}{\sqrt{2\pi}} \int_{-\infty}^{\infty} (i t)^n e^{-t^2/2 + i t x} dt.$$

Next, perform the change of variable

$$s = it$$
, so that  $t = -is$ ,  $dt = -ids$ .

Under this substitution the factors transform as follows:

$$(it)^n = s^n,$$

and the exponent becomes

$$-\frac{t^2}{2} + itx = -\frac{(-is)^2}{2} + i(-is)x = \frac{s^2}{2} + sx.$$

Thus, the integral transforms into

$$\int_{-\infty}^{\infty} (it)^n e^{-t^2/2 + itx} dt = -i \int_{-i\infty}^{i\infty} s^n e^{s^2/2 + sx} ds.$$

Substituting back we have

$$p_n(x) = \frac{(-1)^n e^{x^2/2}}{\sqrt{2\pi}} (-i) \int_{-i\infty}^{i\infty} s^n e^{s^2/2 + sx} ds.$$

That is,

$$p_n(x) = \frac{i(-1)^{n+1} e^{x^2/2}}{\sqrt{2\pi}} \int_{-i\infty}^{i\infty} s^n e^{s^2/2 + s x} ds.$$

Finally, change the sign of s, and we get:

$$p_n(x) = \frac{i e^{x^2/2}}{\sqrt{2\pi}} \int_{-i\infty}^{i\infty} s^n e^{s^2/2 - s x} ds.$$

Therefore,

$$\psi_n(x) = \frac{i e^{x^2/4}}{\sqrt{2\pi h_n}} \int_{-i\infty}^{i\infty} s^n e^{s^2/2 - s x} ds.$$

#### 2.3 Normalization of Hermite polynomials

Lemma 2.2. We have

$$h_n = \int_{-\infty}^{\infty} p_n(x)^2 e^{-x^2/2} dx = n! \sqrt{2\pi}.$$

*Proof.* Multiply the generating function

$$\exp\left(xt - \frac{t^2}{2}\right) = \sum_{n>0} p_n(x) \frac{t^n}{n!}$$

with a second copy (with parameter s):

$$\exp\left(xs - \frac{s^2}{2}\right) = \sum_{m>0} p_m(x) \frac{s^m}{m!}.$$

Then,

$$\exp\left(xt - \frac{t^2}{2}\right) \exp\left(xs - \frac{s^2}{2}\right) = \sum_{n,m>0} p_n(x)p_m(x)\frac{t^n s^m}{n!m!}.$$

Integrate both sides against  $e^{-x^2/2} dx$ . Using the orthogonality

$$\int_{-\infty}^{\infty} p_n(x)p_m(x)e^{-x^2/2}dx = h_n\delta_{nm},$$

the right-hand side becomes

$$\sum_{n>0} \frac{h_n}{(n!)^2} (ts)^n.$$

On the left-hand side, we have

$$\int_{-\infty}^{\infty} e^{-x^2/2} \exp\left(x(t+s) - \frac{t^2 + s^2}{2}\right) dx.$$

Completing the square in x or recalling the standard Gaussian integral yields

$$\sqrt{2\pi} \exp\left(\frac{(t+s)^2 - (t^2 + s^2)}{2}\right) = \sqrt{2\pi} \exp(ts).$$

Thus, we obtain

$$\sqrt{2\pi} \exp(ts) = \sum_{n\geq 0} \frac{h_n}{(n!)^2} (ts)^n.$$

Expanding the left side as

$$\sqrt{2\pi} \sum_{n \ge 0} \frac{(ts)^n}{n!},$$

and comparing coefficients, we conclude that

$$\frac{h_n}{(n!)^2} = \frac{\sqrt{2\pi}}{n!} \implies h_n = n!\sqrt{2\pi}.$$

This completes the proof.

#### 2.4 Double contour integral representation for the GUE kernel

We can sum up the kernel (essentially, this is another proof of the Christoffel–Darboux formula):

$$K_n(x,y) = \sum_{k=0}^{n-1} \psi_k(x)\psi_k(y)$$

$$= \frac{e^{\frac{x^2 - y^2}{4}}}{(2\pi)^2} \oint_C dt \int_{-i\infty}^{i\infty} ds \exp\left\{-\frac{t^2}{2} + xt + \frac{s^2}{2} - ys\right\} \underbrace{\sum_{k=0}^{n-1} s^k t^{-k-1}}_{\underbrace{1 - (s/t)^n \atop t - s}}.$$
(2.3)

Here we used the two contour integral representations for Hermite polynomials, and the explicit norm (Lemma 2.2). At this point, the t contour is a small circle around 0, and the s contour is a vertical line in the complex plane. Their mutual position can be arbitrary at this point — the s contour goes along the imaginary line. Indeed, the fraction  $\frac{1-(s/t)^n}{t-s}$  does not have a singularity at s=t due to the cancellation.

Let us now move the s contour to be to the left of the t contour, as in Figure 1. On the new contours, we have |s| > |t|. Now we can add the summands  $s^k t^{-k-1}$  for all  $k \le -1$  into the sum in (2.3). Indeed, for |s| > |t|, the series in k converges, while the summand  $s^k t^{-k-1}$  has zero residue at 0 and thus adding the summands does not change the value of the integral.

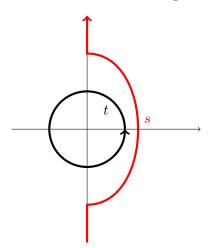


Figure 1: Integration contours for the GUE kernel (2.4).

With this extension of the sum, formula (2.3) becomes

$$K_n(x,y) = \frac{e^{(y^2 - x^2)/4}}{(2\pi)^2} \oint_C dt \int_{-i\infty}^{i\infty} ds \frac{\exp\left\{\frac{s^2}{2} - sy - \frac{t^2}{2} + tx\right\}}{s - t} \left(\frac{s}{t}\right)^n. \tag{2.4}$$

**Remark 2.3.** The s contour passes to the right of the t contour, but it might as well pass to the left of it. Indeed, one can deform the s contour to the left while picking the residue at s = t:

$$\operatorname{Res}_{s=t} \frac{\exp\left\{\frac{s^2}{2} - sy - \frac{t^2}{2} + tx\right\}}{s - t} \left(\frac{s}{t}\right)^n = -e^{t(x - y)}.$$

This function is entire in t, and its integral over the t contour is zero. Therefore, there is no difference where the s contour passes with respect to the t contour.

#### 2.5 Extensions

Many other versions of the GUE / unitary invariant ensembles admit determinantal structure:

- 1. The GUE corners process [JN06]
- 2. The Dyson Brownian motion (e.g., add a GUE to a diagonal matrix) [NF98]
- 3. GUE corners plus a fixed matrix [FF14]
- 4. Corners invariant ensembles with fixed eigenvalues  $UDU^{\dagger}$ , where D is a fixed diagonal matrix and U is Haar distributed on the unitary group [Met13]

We will discuss the corners process structure in the next Lecture 7.

### F Problems (due 2025-03-12)

## References

- [FF14] P. Ferrari and R. Frings, Perturbed GUE minor process and Warren's process with drifts, J. Stat. Phys 154 (2014), no. 1-2, 356–377. arXiv:1212.5534 [math-ph]. ↑7
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- L. Petrov, University of Virginia, Department of Mathematics, 141 Cabell Drive, Kerchof Hall, P.O. Box 400137, Charlottesville, VA 22904, USA E-mail: lenia.petrov@gmail.com