Lectures on Random Matrices (Spring 2025) Lecture 6: Steepest descent and local statistics

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1 Recap

1.1 Determinantal structure of the GUE

Last time, we proved the following result:

Theorem 1.1.

2 Double Contour Integral Representation for the GUE Kernel

2.1 One contour integral representation for Hermite polynomials

Recall that the GUE kernel is defined by

$$K_N(x,y) = \sum_{n=0}^{N-1} \psi_n(x)\psi_n(y),$$

with the orthonormal functions

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4},$$

where the (monic, probabilists') Hermite polynomials are given by

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} e^{-x^2/2},$$

and satisfy the generating function

$$\exp\left(xt - \frac{t^2}{2}\right) = \sum_{n \ge 0} p_n(x) \frac{t^n}{n!}.$$

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By Cauchy's integral formula we can write

$$p_n(x) = \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt,$$

where the contour C is a simple closed curve encircling the origin. Therefore,

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4} = \frac{e^{-x^2/4}}{\sqrt{h_n}} \frac{n!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{n+1}} dt.$$

2.2 Another contour integral representation for Hermite polynomials

Note also that

$$\int_{-\infty}^{\infty} e^{-t^2 + \sqrt{2}i t x} dt = \sqrt{\pi} e^{-x^2/2}.$$

Differentiating both sides n times with respect to x (and using the fact that in our convention the Gaussian appears with $x^2/2$) yields

$$\frac{d^{n}}{dx^{n}} \left(e^{-x^{2}/2} \right) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} \left(\sqrt{2}i \, t \right)^{n} e^{-t^{2} + \sqrt{2}i \, t \, x} \, dt.$$

Changing variables via s = it (so that t = -is and dt = -ids) one obtains

$$\frac{d^n}{dx^n} \left(e^{-x^2/2} \right) = \frac{(\sqrt{2})^n}{i\sqrt{\pi}} \int_{-i\infty}^{i\infty} s^n \, e^{s^2 + \sqrt{2} \, s \, x} \, ds.$$

Multiplying by $(-1)^n e^{x^2/2}$ we deduce that

$$p_n(x) = (-1)^n e^{x^2/2} \frac{d^n}{dx^n} \left(e^{-x^2/2} \right) = \frac{i \left(\sqrt{2} \right)^n e^{x^2/2}}{\sqrt{\pi}} \int_{-i\infty}^{i\infty} s^n e^{s^2 - \sqrt{2} s x} ds.$$
 (2.1)

Now, recall that the orthonormal functions are defined as

$$\psi_n(x) = \frac{1}{\sqrt{h_n}} p_n(x) e^{-x^2/4},$$

so that by (2.1)

$$\psi_n(x) = \frac{i e^{x^2/4}}{\sqrt{\pi h_n}} \int_{-i\infty}^{i\infty} (\sqrt{2}s)^n e^{s^2 - \sqrt{2}sx} ds = \frac{i e^{x^2/4}}{\sqrt{2\pi h_n}} \int_{-i\infty}^{i\infty} s^n e^{s^2/2 - sx} ds.$$

2.3 Double contour integral representation for the GUE kernel

We have (Problem ??)

$$h_n = \int_{-\infty}^{\infty} p_n(x)^2 e^{-x^2/2} dx = n! \sqrt{2\pi}.$$

Therefore, we can sum up the kernel (another proof of the Christoffel–Darboux formula):

$$K_n(x,y) = \sum_{k=0}^{n-1} \psi_k(x)\psi_k(y)$$

$$= \sum_{k=0}^{n-1} \frac{e^{-x^2/4}}{\sqrt{h_k}} \frac{k!}{2\pi i} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{k+1}} dt \frac{i e^{y^2/4}}{\sqrt{2\pi h_k}} \int_{-i\infty}^{i\infty} s^k e^{s^2/2 - sy} ds$$

$$= e^{(y^2 - x^2)/4} \sum_{k=0}^{n-1} \frac{1}{4\pi^2} \oint_C \frac{\exp\left(xt - \frac{t^2}{2}\right)}{t^{k+1}} dt \int_{-i\infty}^{i\infty} s^k e^{s^2/2 - sy} ds.$$

We can now extend the sum to $k = -\infty$, and get a formula for the GUE kernel as a double contour integral:

$$K_n(x,y) = \frac{e^{(y^2 - x^2)/4}}{4\pi^2} \oint_C \int_{-i\infty}^{i\infty} \frac{\exp\left\{\frac{s^2}{2} - sy - \frac{t^2}{2} + tx\right\}}{s - t} \left(\frac{s}{t}\right)^n ds dt.$$

Details will be in the next Lecture 6.

Remark 2.1. Many other versions of the GUE / unitary invariant ensembles admit determinantal structure:

- 1. The GUE corners process [JN06]
- 2. The Dyson Brownian motion (e.g., add a GUE to a diagonal matrix) [NF98]
- 3. GUE corners plus a fixed matrix [FF14]
- 4. Corners invariant ensembles with fixed eigenvalues UDU^{\dagger} , where D is a fixed diagonal matrix and U is Haar distributed on the unitary group [Met13]

F Problems (due DATE)

References

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