

NOMENCLATURE

LATIN SYMBOLS

a	leaf area density	$\text{m}^2 \text{m}^{-3}$
A	area	m^2
A_f	frontal area	m^2
A_L	net leaf area	m^2
A_n	assimilation rate	$\text{mol m}^{-2} \text{s}^{-1}$
b	Vogel exponent	-
c_d	leaf drag coefficient	-
C_d	drag coefficient	-
c_p	specific heat (at constant pressure)	$\text{J kg}^{-1} \text{K}^{-1}$
c_a^*	reference CO_2 concentration	mol mol^{-1}
C	aerodynamic resistance coefficient	$\text{s}^{0.5} \text{m}^{-1}$
$C_{o,a}$	oxygen concentration	mol mol^{-1}
D	diameter	m
D	vapor pressure deficit	kPa
D	vapor pressure deficit	$\text{Pa Pa}^{-1}, \text{mol mol}^{-1}$
e_m	maximum quantum efficiency of the leaf	-
E	leaf transpiration flux	kg s^{-1}
f_c	CO_2 assimilation (photosynthesis) rate	$\text{mol m}^{-2} \text{s}^{-1}$
f_v	transpiration rate	$\text{mol m}^{-2} \text{s}^{-1}$
F	force	N
F_d	drag force	N
g	gravitational acceleration	$\text{m}^2 \text{s}^{-1}$
g_l	liquid water flux	$\text{kg m}^{-2} \text{s}^{-1}$
g_v	water vapor flux	$\text{kg m}^{-2} \text{s}^{-1}$
$g_{v,leaf}$	water vapor mass flux	$\text{kg m}^{-2} \text{s}^{-1}$
$g_{v,root}$	root water uptake	$\text{kg m}^{-2} \text{s}^{-1}$
$G_{v,root}$	net root water uptake	kg s^{-1}
$G_{v,xylem}$	net xylem water flux	kg s^{-1}
h	enthalpy	J kg^{-1}
$h_{c,h}$	convective heat transfer coefficient	$\text{W m}^{-2} \text{K}^{-1}$

$h_{c,m}$	convective mass transfer coefficient	m s^{-1}
H	height	m
HU	Hounsfield units	-
I	turbulence intensity	%
k	turbulent kinetic energy	$\text{m}^2 \text{s}^{-2}$
k_{st}	stomatal conductance	$\text{mol m}^{-2} \text{s}^{-1}$
k_{st}^*	effective stomatal conductance	$\text{mol m}^{-2} \text{s}^{-1}$
K	permeability	m^2
K	hydraulic conductivity	m s^{-1}
K_{lp}	liquid water permeability	s
K_{vp}	water vapor permeability	s
K_{vT}	water vapor permeability due to temperature	s
K_c	Michaelis constant for CO_2	???
K_o	Michaelis constant for O_2	???
l	length	m
l	characteristic leaf size	m
\mathcal{L}	Lagrangian	-
L_v	latent heat of vaporization for water ($= 2.5 \times 10^6$)	J kg^{-1}
m	mass	kg
M	molar mass	kg mol^{-1}
M_a	molar mass of dry air ($= 0.028\,966$)	kg mol^{-1}
M_v	molar mass of water vapor ($= 0.018\,015\,34$)	kg mol^{-1}
n	Van Genuchten parameter	-
p_c	capillary pressure	Pa
p_l	liquid pressure	Pa
p_l	gas phase pressure	Pa
p_v	partial vapor pressure	Pa
$p_{v,i}$	intercellular vapor pressure	Pa
$p_{v,sat}$	saturation vapor pressure	Pa
P	pressure	Pa
Pr	Prandtl number	-
Pr_t	turbulent Prandtl number	-
q	heat flux	W m^{-2}
$q_{lat,lat}$	latent heat flux from leaf	W m^{-2}
q_r	net radiative heat flux	W m^{-2}
$q_{r,sw}$	short-wave radiative heat flux	W m^{-2}
$q_{r,lw}$	long-wave radiative heat flux	W m^{-2}

$q_{rad,lat}$	radiative heat flux into leaf	$W m^{-2}$
$q_{sen,lat}$	sensible heat flux from leaf	$W m^{-2}$
Q_p	flux of incoming PAR	???
r	root area density	$m^2 m^{-3}$
r_a	aerodynamic resistance	$s m^{-1}$
r_s	stomatal resistance	$s m^{-1}$
R	universal gas constant (= 8.314 459 8)	$J mol^{-1} K^{-1}$
RAI	root area index	$m^2 m^{-2}$
Re	Reynolds number	-
R_v	specific gas constant of dry air (= 287.042)	$J kg^{-1} K^{-1}$
R_v	specific gas constant of water vapor (= 461.524)	$J kg^{-1} K^{-1}$
Sc_t	turbulent Schmidt number	-
s_ϵ	volumetric TDR source	$W m^{-3} s^{-1}$
s_ρ	volumetric mass source	$kg m^{-3} s^{-1}$
s_k	volumetric TKE source	$W m^{-3}$
$s_{q,r}$	volumetric radiative source	$W m^{-3}$
s_r	volumetric root water uptake source	$kg m^{-3} s^{-1}$
s_T	volumetric temperature source	$K m^{-3}$
s_u	volumetric momentum source	$N m^{-3}$
s_w	volumetric humidity source	$kg kg^{-1} s^{-1}$
S_l	liquid saturation	-
t	time	s
T	temperature	K
T_g	ground temperature	K
T_L	leaf temperature	K
T_{sky}	sky temperature	K
TR	net hourly transpiration rate	$g h^{-1}$
u	velocity	$m s^{-1}$
u_*	friction velocity	$m s^{-1}$
U	mean wind speed	$m s^{-1}$
U_{ref}	reference velocity	$m s^{-1}$
$UTCI$	universal thermal climate index	$^{\circ}C$
V_{cmax}	maximum carboxylation capacity	???
w	moisture content	$kg m^{-3}$
w_a	dry air content	$kg m^{-3}$
w_{cap}	capillary moisture content	$kg m^{-3}$
w_l	liquid water content	$kg m^{-3}$

w_s	solid matrix water content	kg m^{-3}
w_v	water vapor content	kg m^{-3}
z	vertical height	m
z_0	aerodynamic roughness height	m

GREEK SYMBOLS

α	aerodynamic porosity	-
α^{2D}	2D aerodynamic porosity	-
α_p	leaf absorptivity of PAR	-
β	optical porosity	-
β	volumetric thermal expansion coefficient	K^{-1}
β_p	fraction of MKE converted to TKE	-
β_d	fraction of TKE shortcut to TDR	-
γ	apparent quantum yield	-
δ_v	water vapor diffusion coefficient	s
ε	turbulent kinetic energy dissipation rate	$\text{m}^2 \text{s}^{-3}$
κ	von Kármán constant (= 0.41)	-
λ	thermal conductivity	$\text{W m}^{-1} \text{K}^{-1}$
λ	Lagrange multiplier	mol mol^{-1}
μ	attenuation coefficient	m^{-1}
μ	chemical potential	J mol^{-1}
μ	dynamic viscosity	$\text{kg m}^{-1} \text{s}^{-1}$
ν	kinematic viscosity	$\text{m}^2 \text{s}^{-1}$
ν_t	turbulent viscosity	$\text{m}^2 \text{s}^{-1}$
ρ	density	kg m^{-3}
ρ_a	density of air	kg m^{-3}
ρ_l	density of liquid water (= 1000)	kg m^{-3}
ρ_s	density of solid matrix	kg m^{-3}
σ_v	Schmidt number	-
σ_{v_t}	turbulent Schmidt number	-
ϕ_0	open porosity	$\text{m}^3 \text{m}^{-3}$
ψ	shelter parameter	-
ψ_g	gravitational potential	Pa
ψ_L	leaf water potential	Pa
ψ_R	root water potential	Pa

ψ_s	soil water potential	Pa
Ω	domain	-
Ω_a	air domain	-
Ω_s	soil domain	-

SUBSCRIPTS

<i>eff</i>	effective
<i>g</i>	gas
<i>g</i>	ground
<i>i</i>	intercellular
<i>l</i>	liquid (water)
<i>lw</i>	long-wave
<i>L</i>	leaf
<i>o</i>	open
<i>pore</i>	pore
<i>ref</i>	reference
<i>root</i>	root
<i>sat</i>	saturated
<i>R</i>	root
<i>s</i>	solid
<i>s</i>	soil
<i>sky</i>	sky
<i>sw</i>	short-wave
<i>t</i>	time
<i>v</i>	vapor
<i>xylem</i>	xylem

ACRONYMS

ABL	atmospheric boundary layer
CFD	computational fluid dynamics
CHTC	convective heat transfer coefficient
CMTC	convective mass transfer coefficient

CT	computational tomography
DEHS	Di-Ethyl-Hexyl-Sebacat
DIG	diagonal-based incomplete Cholesky
DNS	direct numerical simulation
ETHZ	Eidgenössische Technische Hochschule Zürich
FOV	field of view
FFT	fast Fourier transformation
FVM	finite volume method
HPC	high performance computing
LAI	leaf area index
LAI	leaf area density
LDPE	low-density Polyethylene
LES	large eddy simulation
LHS	left hand side
MOE	modulus of elasticity
PAR	photosynthetically active radiation
PCG	preconditioned conjugate gradient
PIV	particle image velocimetry
PPM	parts per million
RAI	root area index
REV	representative elementary volume
RANS	Reynolds-averaged Navier-Stokes
RHS	right hand side
ROI	region of interest
SPIV	stereoscopic particle image velocimetry
TDR	turbulent dissipation rate
TKE	turbulent kinetic energy
UHI	urban heat island

UTCI	universal thermal climate index
WUE	water use efficiency

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NUMERICAL MODEL FOR VEGETATION IN URBAN AREAS

1.1 INTRODUCTION

The simplified vegetation model is described in ?? introduction the governing equation for moist flow vegetation. The vegetation model provides the necessary source/sink terms for heat, mass and momentum exchanges between vegetation and air as described in ?. Furthermore, radiation transfer within vegetation is also described in ?. For reader, a detailed derivation of the thermodynamic of the moist air is given in ? and a detailed derivation of the governing equation of moist flow is given in ?.

In this chapter, the numerical method for modeling vegetation inside an urban area is described. The chapter focuses on coupling of the vegetation model with the of heat (incl. radiation) and mass fluxes of the urban surfaces. The governing equation for the the coupled heat and moisture transport in the porous material is first described.

1.2 GOVERNING EQUATIONS OF COUPLED HEAT AND MOISTURE TRANSPORT

1.2.1 *Composition of the porous material*

The building materials and the soil is considered as a porous material consisting of three phases: solid phase (denoted with s), liquid phase referring to liquid water (l) and the air phase which is split into dry air (a) and water vapor v (Janssen 2002; Carmeliet 2005; Defraeye 2011; Saneinejad 2013). The open porosity ϕ_o ($\text{m}^3 \text{m}^{-3}$) of the porous material is defined as:

$$\phi_o = \frac{V_{pore}}{V} \quad (1.1)$$

where V_{pore} (m^3) is the volume of open pores and V (m^3) is the total volume of the porous material S . The solid material content w_s (kg m^{-3}) is defined as:

$$w_s = (1 - \phi_o) \rho_s \quad (1.2)$$

where ρ_s (kg m^{-3}) is the solid material matrix density. Similarly the dry air w_a , water vapor w_v , liquid water w_l contents are defined as:

$$w_l = \phi_o S_l \rho_l \quad (1.3)$$

$$w_a = \phi_o (1 - S_l) \rho_a \quad (1.4)$$

$$w_v = \phi_o (1 - S_l) \rho_v \quad (1.5)$$

where they are related to the degree of liquid saturation S_l of the porous open pores:

$$S_l = \frac{\phi_{o,l}}{\phi_o} \quad (1.6)$$

with $\phi_{o,l}$ being the amount of liquid water occupied inside the open pores and ρ_a, ρ_l being the air and liquid water densities, respectively. The total moisture content w (kg m^{-3}) inside the porous material is simply the sum of liquid water and water vapor:

$$w = w_l + w_v \quad (1.7)$$

1.2.2 Water potential

Water potential is a universal parameter to determine the *water status* in any medium (Nobel 2009). In the thesis, we use it to define the water status of multiple domains such as solid porous materials (soil and building facades), plant xylem, and air. The water potential ψ (Pa) describes the chemical potential of water μ (J mol^{-1}) with respect to chemical potential of pure water $\mu^{o,l}$ (J mol^{-1}) at the same temperature, standard atmosphere and at zero level:

$$\psi = \frac{\mu - \mu^{o,l}}{V_l^o} \quad (1.8)$$

where $V_l^o = 18.0510^{-3} \text{ m}^3 \text{ mol}^{-1}$ is the molar volume of pure water in liquid phase. The water potential for pure water at standard atmosphere is $\psi = 0$ Pa. The water tends to move towards a region where $\mu - \mu^{o,w}$ is lower, i.e. in the direction of $-\nabla\psi$. The water potential is related to pressure potential, osmotic potential, matrix potential, and gravitational potential. In our study, we assume that the pressure potential gradient is negligible in the solid and we ignore the influence of osmotic potential ψ_o (Pa) as we assume we have a non-saline porous material. Therefore, only the matrix potential and gravitation potential influences the water transport:

$$\psi = \underbrace{p_c}_{\psi_c} + \underbrace{\rho_l g z}_{\psi_g} \quad (1.9)$$

where $\psi_c = p_c$ (Pa) is the capillary potential due to capillary pressure, which represents the contribution of the matrix potential, and $\psi_g = \rho_l g h$ (Pa) is the gravitational potential with $g = |g|$ (m s^{-2}) where g is the gravitational acceleration. The capillary

pressure p_c is defined as the difference between liquid and gas phase pressure, p_l and p_g , respectively:

$$p_c = p_l - p_g \quad (1.10)$$

and is related to relative humidity ϕ by the Kelvin's law:

$$p_c = \rho_l R_v T \ln(\phi) \quad (1.11)$$

The gravitational potential ψ_g (Pa) is defined as:

$$\psi_g = -\rho g \cdot x = \rho g z \quad (1.12)$$

where $g = |g|$ with z oriented upward. Thus, by taking the capillary and gravitational water potentials into account, the transport of water can be described in building materials and, more importantly, soil where the plant roots are present.

1.2.3 Coupled transport of heat and mass

Conservation of mass

The conservation of mass in the solid domain is defined as:

$$\frac{\partial w_s}{\partial t} = 0 \quad (1.13)$$

$$\frac{\partial w_a}{\partial t} + \nabla \cdot (w_a \mathbf{u}_a) = 0 \quad (1.14)$$

$$\frac{\partial w_l + w_v}{\partial t} + \nabla \cdot (w_l \mathbf{u}_l + w_v \mathbf{u}_v) = 0 \quad (1.15)$$

assuming that solid matrix does not move, mass of different phases only change due to evaporation or condensation. Other phenomena such as melting, freezing, sublimation and deposition are neglected. Assuming dry air does not contribute to moisture storage, i.e., $\partial w_a / \partial t = 0$, the conservation of mass simplifies to rate of change of moisture content $w = w_l + w_v$, from Eq. (1.15):

$$\frac{\partial w}{\partial t} = -\nabla \cdot (\mathbf{g}_l + \mathbf{g}_v) \quad (1.16)$$

where $\mathbf{g}_l \equiv w_l \mathbf{u}_l$ ($\text{kg m}^{-2} \text{s}^{-1}$) and $\mathbf{g}_v \equiv w_v \mathbf{u}_v$ ($\text{kg m}^{-2} \text{s}^{-1}$) are defined as the liquid water and water vapor fluxes, respectively. Additionally, the contribution of root water uptake due to plant transpiration is introduced through the source term s_r ($\text{kg m}^{-3} \text{s}^{-1}$):

$$\frac{\partial w}{\partial t} = -\nabla \cdot (\mathbf{g}_l + \mathbf{g}_v) + s_r \quad (1.17)$$

The sink term due to root water uptake is explained in detail later in Section 1.3. In this thesis, the conservation of mass is solved using the p_c -form Richards equation, where Eq. (1.17) becomes:

$$\frac{\partial w}{\partial p_c} \frac{\partial p_c}{\partial t} = -\nabla \cdot (\mathbf{g}_l + \mathbf{g}_v) + s_r \quad (1.18)$$

with $C_{mm} \equiv \partial w / \partial p_c$ ($\text{kg m}^{-3} \text{Pa}^{-1}$) is defined as the moisture capacity. Thus, the change in water content in the porous material is simply due to the liquid and vapor fluxes, and root water uptake. The liquid water flux \mathbf{g}_l in porous media is given by:

$$\mathbf{g}_l = -K_{lp} \nabla (p_c + \rho_l g z) \quad (1.19)$$

and assumes the air pressure effects be negligible with respect to capillary and gravitational effects, where K_{lp} (s) the liquid water permeability. We assume that the liquid water permeability is only due to pressure gradient, and the influence of thermal gradient is neglected (Carmeliet 2005). The water vapor flux in the porous media is given by:

$$\mathbf{g}_v = K_{vp} \nabla p_c + K_{vT} \nabla T \quad (1.20)$$

where

$$K_{vp} = -\delta_v \frac{p_v}{\rho_l R_v T} \quad (1.21)$$

is the water vapor permeability (s) due to pressure,

$$K_{vT} = -\delta_v \frac{p_v}{\rho_l R_v T^2} (\rho_l L_v - p_c) \quad (1.22)$$

is the water vapor permeability (s) due to temperature, and

$$\delta_v = \frac{D_{va,mat}}{R_v T} \quad (1.23)$$

is the water vapor diffusion coefficient (s) where $D_{va,mat}$ ($\text{m}^2 \text{s}^{-1}$) is the binary apparent diffusion coefficient between dry air and water vapor (Carmeliet 2005; Defraeye 2011; Saneinejad 2013; Kubilay 2014). Thus, substituting the fluxes Eqs. (1.19) and (1.20) into Eq. (1.18), the expanded form of conservation of mass is given as:

$$C_{mm} \frac{\partial p_c}{\partial t} = \nabla \cdot \left(K_{lp} \nabla (p_c + \rho_l g z) + K_{vp} \nabla p_c + K_{vT} \nabla T \right) + s_r \quad (1.24)$$

Conservation of energy

The conservation of energy is given as:

$$\frac{\partial h}{\partial t} + \nabla \cdot (h \mathbf{u}) = -\nabla \cdot \mathbf{q} \quad (1.25)$$

where h (J kg^{-1}) is the enthalpy of total solid domain:

$$h = \sum_i w_i h_i = w_s h_s + w_a h_a + w_l h_l + w_v h_v \quad (1.26)$$

and heat conduction \mathbf{q} (W m^{-2}) is given by Fourier's law as:

$$\mathbf{q} = -\lambda \nabla T \quad (1.27)$$

where T (K) is temperature, λ ($\text{W m}^{-1} \text{K}^{-1}$) is thermal conductivity.

Substituting Eq. (1.26) into Eq. (1.25) expands to:

$$\frac{\partial}{\partial t} (w_s h_s + w_a h_a + w_l h_l + w_v h_v) + \nabla \cdot (w_a \mathbf{g}_a + w_l \mathbf{g}_l + w_v \mathbf{g}_v) = -\nabla \cdot \mathbf{q} \quad (1.28)$$

and assuming dry air and water vapor does not contribute to heat storage, i.e., $\partial w_a h_a / \partial t \approx 0$ and $\partial w_v h_v / \partial t \approx 0$, and convection term of dry air is negligible $\nabla \cdot w_a \mathbf{g}_a \approx 0$, Eq. (1.28) simplified to:

$$\frac{\partial (w_s h_s + w_l h_l)}{\partial t} + \nabla \cdot (w_l \mathbf{g}_l + w_v \mathbf{g}_v) = -\nabla \cdot \mathbf{q} \quad (1.29)$$

The enthalpies of solid h_s , liquid water h_l and water vapor h_v are defined as:

$$h_s = c_{ps} (T - T_{ref}) \quad (1.30)$$

$$h_l = c_{pl} (T - T_{ref}) \quad (1.31)$$

$$h_v = c_{pv} (T - T_{ref}) + L_v \quad (1.32)$$

and L_v is the latent heat of vaporization of water. Substituting Eqs. (1.30) to (1.32) into Eq. (1.29), the conservation of energy is given as:

$$\begin{aligned} & \left(c_{ps} w_s + c_{pl} w \right) \frac{\partial T}{\partial t} + \left[c_{pl} (T - T_{ref}) \frac{\partial w}{\partial p_c} \right] \frac{\partial p_c}{\partial t} = \\ & -\nabla \cdot \left\{ \mathbf{q} + \underbrace{c_{pl} (T - T_{ref}) \mathbf{g}_l}_{q_l} + \underbrace{\left[c_{pv} (T - T_{ref}) + L_v \right] \mathbf{g}_v}_{q_v} \right\} \end{aligned} \quad (1.33)$$

where:

$$C_{TT} = c_{ps} w_s + c_{pl} w \quad (1.34)$$

$$C_{Tp} = c_{pl} (T - T_{ref}) \frac{\partial w}{\partial p_c} \quad (1.35)$$

are the thermal capacity terms. Substituting liquid and vapour fluxes, Eqs. (1.19) and (1.20), respectively, thermal capacities Eqs. (1.34) and (1.35) and heat conduction Eq. (1.27) into conservation of energy expands Eq. (1.33) to:

$$\begin{aligned}
 C_{TT} \frac{\partial T}{\partial t} + C_{Tp} \frac{\partial p_c}{\partial t} = \nabla \cdot \bigg(& \lambda \nabla T + K_{lp} c_{pl} (T - T_{ref}) \nabla p_c \\
 & + K_{lp} c_{pl} (T - T_{ref}) \rho_l g z \\
 & - K_{vp} [c_{pv} (T - T_{ref}) + L_v] \nabla p_c \\
 & - K_{vT} [c_{pv} (T - T_{ref}) + L_v] \nabla T \bigg)
 \end{aligned} \tag{1.36}$$

1.2.4 Linearized heat and mass transport equation

In the present study, the conservation of mass and energy is coupled together as a heat and mass model (HAM) according to (Janssen 2002; Carmeliet 2005; Defraeye 2011; Saneinejad 2013; Kubilay et al. 2018). Due to the shape of the water retention curve and the hydraulic conductivity curve, as shown in Fig. 1.1, the Richards equation is highly non-linear. Therefore, the numerical solution of the equation is very sensitive to convergence tolerance and requires linearization techniques to maintain accuracy and computational efficiency. So, methods such as fixed-point Picard iterations is used to solve the non-linear equations. More details on the discretization is provided in Janssen (2002), Liu (2012), and Kubilay et al. (2018).

The conservation of mass (i.e., the moisture transport equation), Eq. (1.24), in the linearized form is given as:

$$\begin{aligned}
 C_{mm}^{n+1,k} \frac{p_c^{n+1,k+1} - p_c^n}{\Delta t} = \nabla \cdot \bigg(& K_{lp}^{n+1,k} \nabla (p_c^{n+1,k+1} + \rho_l g z) \\
 & + K_{vp}^{n+1,k} \nabla p_c^{n+1,k+1} \\
 & + K_{vT}^{n+1,k} \nabla T^{n+1,k} \bigg) \\
 & + s_r^{n+1,k+1}
 \end{aligned} \tag{1.37}$$

where the capacity, permeabilities and temperature are determined from the previous Picard iteration (k). The subscript (n) denotes the global (i.e., outer) time step of t and k denoting the internal Picard iteration step. When the Picard solution approaches convergence, i.e., $p_c^{n+1,k} \rightarrow p_c^{n+1,k+1}$, $\partial w / \partial p_c$ becomes negligible resulting in larger mass

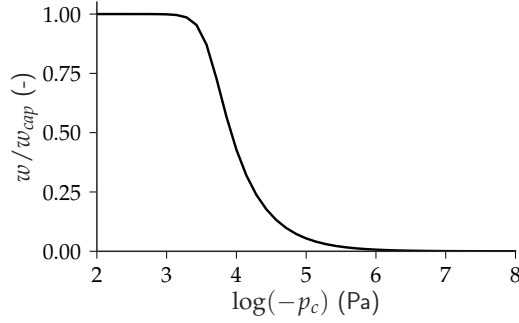


Figure 1.1: Typical non-linear moisture water retention curve for a material which is capillary saturated.

conservation errors. The error is minimized by formulating the Richard's equation in mixed-form (Liu 2012):

$$\begin{aligned}
 C_{mm}^{n+1,k} \frac{p_c^{n+1,k+1} - p_c^n}{\Delta t} = \nabla \cdot \left(K_{lp}^{n+1,k} \nabla \left(p_c^{n+1,k+1} + \rho_l g z \right) \right. \\
 \left. + K_{vp}^{n+1,k} \nabla p_c^{n+1,k+1} \right. \\
 \left. + K_{vT}^{n+1,k} \nabla T^{n+1,k} \right) \\
 + s_r^{n+1,k+1} - \frac{w^{n+1,k+1} - w^n}{\Delta t}
 \end{aligned} \tag{1.38}$$

where an additional moisture change term is added.

The linearized form of the heat equation is given as:

$$\begin{aligned}
 C_{TT}^{n+1,k} \frac{T^{n+1,k+1} - T^n}{\Delta t} = \nabla \cdot \left(\lambda \nabla T^{n+1,k+1} \right. \\
 + K_{lp}^{n+1,k} c_{pl} \left(T^{n+1,k+1} - T_{ref} \right) \nabla p_c^{n+1,k+1} \\
 + K_{lp}^{n+1,k} c_{pl} \left(T^{n+1,k+1} - T_{ref} \right) \rho_l g z \\
 - K_{vp}^{n+1,k} \left[c_{pv} \left(T^{n+1,k+1} - T_{ref} \right) + L_v \right] \nabla p_c^{n+1,k} \\
 \left. - K_{vT}^{n+1,k} \left[c_{pv} \left(T^{n+1,k+1} - T_{ref} \right) + L_v \right] \nabla T^{n+1,k+1} \right)
 \end{aligned} \tag{1.39}$$

where the capillary pressure time derivative term is ignored. The mixed-form the heat equation is given as:

$$\begin{aligned}
C_{TT}^{n+1,k} \frac{T^{n+1,k+1} - T^n}{\Delta t} = & \nabla \cdot \left(\lambda \nabla T^{n+1,k+1} \right. \\
& + K_{lp}^{n+1,k} c_{pl} \left(T^{n+1,k+1} - T_{ref} \right) \nabla p_c^{n+1,k+1} \\
& + K_{lp}^{n+1,k} c_{pl} \left(T^{n+1,k+1} - T_{ref} \right) \rho_l g z \\
& - K_{vp}^{n+1,k} \left[c_{pv} \left(T^{n+1,k+1} - T_{ref} \right) + L_v \right] \nabla p_c^{n+1,k} \\
& \left. - K_{vT}^{n+1,k} \left[c_{pv} \left(T^{n+1,k+1} - T_{ref} \right) + L_v \right] \nabla T^{n+1,k+1} \right) \\
& - \frac{C_{TT}^{n+1} T^{n+1} - C_{TT}^n T^n}{\Delta t}
\end{aligned} \tag{1.40}$$

The system of linear equations is solved by Krylov subspace iteration solver, i.e. preconditioned conjugate gradient (PCG) with diagonal-based incomplete Cholesky (DIG) preconditioning. The convergence criteria for the Picard iteration is user-defined:

$$|p_c^{n+1,k+1} - p_c^{n+1,k}| \leq \delta p_c \tag{1.41}$$

$$|T^{n+1,k+1} - T^{n+1,k}| \leq \delta T \tag{1.42}$$

where $\delta p_c = \delta T = 10^{-2}$.

1.3 SOIL-PLANT-ATMOSPHERE CONTINUUM

In the present study, the components of the water potential inside the plants are not directly determined. The soil-plant-atmosphere continuum model that is integrated into the vegetation model is described in this section, implemented according the state-of-art techniques: (Idso 1977; Farquhar et al. 1980; Manzoni et al. 2011; Volpe et al. 2013; Manoli et al. 2014; Launiainen et al. 2015). The root-system of the plants are represented as a network-like structure assuming cooperative strategy among the individual roots and a bulk plant transpiration through a single xylem. We assume no water storage inside the plant and therefore, the water flux from soil to root $G_{v,root}$, from root to leaf through xylem $G_{v,xylem}$, and from leaf to air $G_{v,leaf}$ is in equilibrium, as depicted in Fig. 1.2:

$$G_{v,root} = G_{v,xylem} = G_{v,leaf} \tag{1.43}$$

and so, the atmospheric evaporative demand (AED) is dependent on the water availability near the roots of the plants.

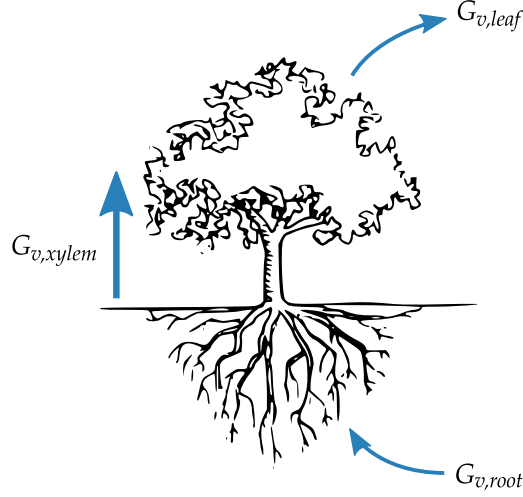


Figure 1.2: Soil-Plant-Atmosphere Continuum: Water balance

1.3.1 Water transport in soil-root system

The sink in soil moisture s_r ($\text{kg m}^{-3} \text{s}^{-1}$) due to root water uptake as given in Eq. (1.24) is simply defined as:

$$s_r = r g_{v,root} \quad (1.44)$$

where r ($\text{m}^2 \text{m}^{-3}$) is the root area density and $g_{v,root}$ ($\text{kg m}^{-2} \text{s}^{-1}$) is the root water uptake, i.e., the flux of water from root to soil. It is defined as:

$$g_{v,root} = k_{sr}^* (\psi_s - \psi_R) \quad (1.45)$$

where k_{sr}^* (s m^{-1}) is effective conductance of the soil-root system (i.e., rhizosphere), ψ_s (Pa) is the soil water potential, and ψ_R (Pa) is the (bulk) root water potential. The effective conductance of the soil-root system (or rhizosphere) k_{sr}^* (s m^{-1}) is given as:

$$k_{sr}^* = \frac{1}{|g|} \frac{k_s k_r}{k_s + k_r} \quad (1.46)$$

where g (m s^{-2}) is the gravitational acceleration, k_s (s^{-1}) is the soil conductance in the root region, and k_r (s^{-1}) is the conductance of the root system. The soil conductance in the root region k_s (s^{-1}) is defined as:

$$k_s = \alpha K r \quad (1.47)$$

where

$$\alpha = \sqrt{\left(\frac{L}{RAI}\right) \frac{1}{d}} \quad (1.48)$$

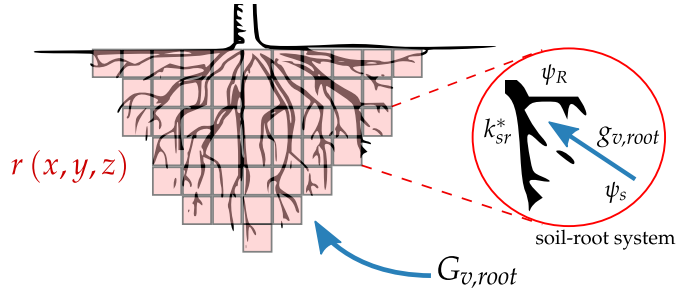


Figure 1.3: Soil-Plant-Atmosphere Continuum: Water transport in soil-root system

d (m) is the root diameter, $RAI = \int r \, dz$ is the root area index ($\text{m}^2 \text{m}^{-2}$). The hydraulic conductivity in soil K (m s^{-1}) is given as:

$$K = K_{lp} |g| \quad (1.49)$$

where K_{lp} (s) being the liquid permeability defined in the Section 1.2.3. Finally, the conductance of the root system k_r (s^{-1}) is given as:

$$k_r = r \frac{\Delta z}{\beta} \quad (1.50)$$

where Δz is the vertical height of the root layer mesh and $\beta = 3 \times 10^8$ s.

Thus, the net root water uptake from the soil domain $G_{v,root}$ (kg s^{-1}) is given as:

$$G_{v,root} = \int_{\Omega_s} s_r \, dV = \int_{\Omega_s} r \, g_{v,root} \, dV \quad (1.51)$$

where Ω_s denotes the soil domain.

1.3.2 Water transport in xylem

The water flux through the plant in xylem $g_{v,xylem}$ ($\text{kg m}^{-2} \text{s}^{-1}$) is defined as:

$$g_{v,xylem}(\psi_L) = k_x^* (\psi_R - \psi_L) \quad (1.52)$$

where k_x^* (s m^{-1}) is the effective xylem conductance, ψ_R (Pa) is the (bulk) root water potential, and ψ_L (Pa) is the (bulk) leaf water potential. The net water flux $G_{v,xylem}$ (kg s^{-1}) is given as:

$$G_{v,xylem} = \int_{\partial\Omega_{x|s}} g_{v,xylem} \, dA = g_{v,xylem} A_x \quad (1.53)$$



Figure 1.4: Soil-Plant-Atmosphere Continuum: Water transport in xylem

where A_x (m^2) is the xylem cross-sectional area. The effective xylem conductance k_x^* (s m^{-1}) of water is:

$$k_x^* = k_x \rho_l \quad (1.54)$$

where plant xylem conductance k_x ($\text{m Pa}^{-1} \text{s}^{-1}$) is modeled using a “*vulnerability curve*” approach. The xylem conductance becomes exponentially smaller with increasing leaf water potential (Volpe et al. 2013). This empirical model is based on plant response to the vulnerability to xylem cavitation and embolism that could occur at high water potential gradients. The xylem conductance k_x ($\text{m Pa}^{-1} \text{s}^{-1}$) is defined by:

$$k_x = k_{x,max} \exp \left\{ - \left(- \frac{\psi_L}{d} \right)^c \right\} \quad (1.55)$$

where $k_{x,max}$ ($\text{m Pa}^{-1} \text{s}^{-1}$) is the maximum xylem conductance, and c and d are fit-coefficients (Volpe et al. 2013).

1.3.3 Water transport from leaf to air

The leaf transpiration rate $g_{v,leaf}$ ($\text{kg m}^{-2} \text{s}^{-1}$) is defined in ??, and is simply defined as:

$$g_{v,leaf} = k_{st,v}^* (p_{v,leaf} - p_v) \quad (1.56)$$

where $k_{st,v}^*$ (s m^{-1}) is the effective stomatal conductance to water vapor, $p_{v,leaf}$ (Pa) is the leaf surface vapor pressure and p_v (Pa) is the atmospheric vapor pressure. The net plant transpiration rate is given as:

$$G_{v,leaf} = \int_{\Omega_a} a g_{v,leaf} dV \quad (1.57)$$

where a ($\text{m}^2 \text{m}^{-3}$) is the leaf area density. The water mass flux to the atmosphere is assumed to be in equilibrium with the water vapor flux through xylem:

$$G_{v,xylem} = G_{v,leaf} \quad (1.58)$$

and so:

$$k_x^* (\psi_R - \psi_L) A_x = \int_{\Omega_a} a g_{v,leaf} dV \quad (1.59)$$

Therefore, the root water potential ψ_R can be determined once leaf water potential ψ_L (Pa), net plant transpiration rate $G_{v,leaf}$ (kg s^{-1}), effective xylem conductance k_x^* (s m^{-1}) and xylem cross-section area A_x (m^2):

$$\psi_R = \psi_L + \frac{G_{v,leaf}}{A_x k_x^*} \quad (1.60)$$

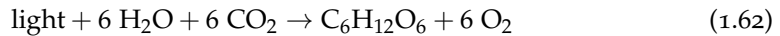
The assumption is that water extract by roots only feed the transpiration and storage inside the plant is neglected. Furthermore, one can also take in account of the gravitation potential change due to tree height:

$$\psi_R = \psi_L + \frac{G_{v,leaf}}{A_x k_x^*} + \underbrace{\rho_l g H}_{\psi_g} \quad (1.61)$$

however, with a tree height of $H = 10$ m, the additional potential is only $\psi_g = 0.1$ MPa.

1.3.4 Improved stomatal model

The photosynthetic reaction takes light, water and CO_2 and creates carbohydrate and oxygen:



and additional moisture is lost by evaporation when stomatal cavity is exposed to atmosphere. Therefore, the photosynthetic process is directly related to atmospheric condition such as CO_2 concentration, availability of light and temperature. Furthermore, the transpiration rate is also dependent on the atmospheric humidity and the availability of water for transpiration (Ball 1987; Leuning et al. 1995). Based on these conditions, a generally accepted theory is the plant regulates the stomatal aperture to optimize the photosynthetic rate for a given transpiration rate. Moreover, the function of vegetation can be simplified as just maximizing the photosynthesis (or CO_2 assimilation) for a given transpiration rate (water use) (Medlyn et al. 2011). The water use efficiency or WUE quantifies the efficiency of the plant of reaching this target. The WUE is defined as:

$$\text{WUE} = \frac{f_c}{f_v} \quad (1.63)$$

where f_c ($\text{mol m}^{-2} \text{s}^{-1}$) is the CO_2 assimilation rate (i.e., denoted also as A_n (in plant-science) or $G_{c,leaf}$ (in building physics), also known as photosynthesis rate) and f_v ($\text{mol m}^{-2} \text{s}^{-1}$) is the transpiration rate (i.e. denoted also as f_e (in plant-science) or $G_{v,leaf}$ (in building physics), also known as water use).

The stomatal optimality model reflects the theory of the stomatal behavior (Cowan 1978). The optimal stomatal control is derived from the minimization problem described by the Lagrangian:

$$\mathcal{L}(k_{st}) = f_c - \lambda f_v \quad (1.64)$$

where λ (mol mol^{-1}) is a Lagrange multiplier and represents the marginal water cost of plant carbon gain (Katul et al. 2010; Medlyn et al. 2011; Manoli et al. 2014) and $f_c = f_c(k_{st})$ and $f_v = f_v(k_{st})$ where both assimilation rate and transpiration rate are both dependent on the stomatal conductance k_{st} ($\text{mol m}^{-2} \text{s}^{-1}$) (i.e. g_s (in plant-science) or $h_{c,m}$ (in building-physics) or $1/r_s$ where r_s is stomatal resistance). Cowan (1978) shows that optimal stomatal behaviour is at the minima of the Lagrangian:

$$\frac{\partial \mathcal{L}}{\partial k_{st}} = 0 \quad (1.65)$$

leading to the following constraint:

$$\lambda = \frac{\partial f_v}{\partial k_{st}} \frac{\partial k_{st}}{\partial f_c} \quad (1.66)$$

or simply:

$$\lambda = \frac{\partial f_v}{\partial f_c} \quad (1.67)$$

Following, these constraints, the stomatal conductance is can be determined with additional closure model for assimilation rate and transpiration rate. The assimilation rate f_c can be described from the perspective of photochemical reaction model and the Fickian diffusion model from the stomatal cavity.

The Farquhar model of photosynthesis describing the biochemical demand function is given as:

$$f_c = \frac{a_1 c_i}{a_2 + s c_a} \quad (1.68)$$

where c_i (mol mol^{-1}) is the intercellular CO_2 concentration, c_a (mol mol^{-1}) is the ambient CO_2 concentration, and a_1 and a_2 are parameters dependent on whether photosynthetic reaction rate is limited by light or Rubisco (Rubilose biphosphate (RuBP) carboxylase-oxygenase) (Farquhar et al. 1980; Katul et al. 2010) and $s = 0.7$ is the constant representing long-term intercellular to ambient CO_2 concentration ratio (Volpe et al. 2013). Note that we use a linearized model assuming $c_p \ll c_i$ where c_p is the CO_2 compensation point. As the photosynthesis can either be light-limited or Rubisco limited, the true assimilation rate f_c is given as:

$$f_c = \min(f_c^l, f_c^r) \quad (1.69)$$

where f_c^l is the light-limited assimilation rate and f_c^r is the Rubisco limited assimilation rate. Note that it is also possible to incorporate the dark (or night) respiration and in that case $f_c = \min(f_c^l, f_c^r) - r_d$, but is simplified in our study.

Light-limited

When the assimilation (or photosynthesis) rate is *light-limited*:

$$a_1(x) = \alpha_p e_m Q_p = \gamma Q_p(x) \quad (1.70)$$

and

$$a_2(x) = 2c_p(x) \quad (1.71)$$

where α_p is the leaf absorptivity of photosynthetically active radiation (PAR), e_m is the maximum quantum efficiency of the leaf, $\gamma = 0.015$ is the apparent quantum yield, Q_p ($\text{mol m}^{-2} \text{s}^{-1}$) is the flux of incoming PAR and c_p (mol mol^{-1}) is the the CO_2 compensation point with:

$$c_p(x) = \frac{K_c(x)}{2K_o(x)} C_{o,a} \frac{k_o}{k_c} \quad (1.72)$$

where $k_c = 2.5 \text{ s}^{-1}$ and $k_o = 0.18k_c$ (Farquhar et al. 1980). Therefore, the light-limited assimilation rate is:

$$f_c^l(x) = \frac{\gamma Q_p(x) c_i(x)}{2c_p(x) + s_{c_a}(x)} \quad (1.73)$$

Rubisco-limited

If the assimilation rate is *Rubisco-limited*:

$$a_1(x) = V_{cmax}(x) \quad (1.74)$$

and

$$a_2(x) = K_c(x) \left(1 + \frac{C_{o,a}}{K_o(x)} \right) \quad (1.75)$$

where V_{cmax} is the maximum carboxylation capacity (referenced at 25 °C), K_c and K_o are Michaelis constant for CO_2 and O_2 inhibition (referenced at 25 °C), and $C_{o,a} = 0.21 \text{ mol mol}^{-1}$ is the oxygen concentration in the atmosphere. The maximum carboxylation capacity is given as:

$$V_{cmax}(x) = V_{cmax,25} \frac{\exp \{0.088 (T_l(x) - 298.15)\}}{1 + \exp \{0.29 (T_l(x) - 314.15)\}} \quad (1.76)$$

and K_c and K_o are

$$K_c(x) = K_{c,25} \exp \{ \gamma_c (T_l(x) - 298.15) \} \quad (1.77)$$

$$K_o(x) = K_{o,25} \exp \{ \gamma_o (T_l(x) - 298.15) \} \quad (1.78)$$

where T_l (K) is the leaf temperature, $V_{cmax,25} = 5.9 \times 10^{-5} \text{ mol m}^{-2} \text{ s}^{-1}$, $\gamma_c = 0.074$, $\gamma_o = 0.015$, $K_{c,25} = 3 \times 10^{-3} \text{ mol mol}^{-1}$ and $K_{o,25} = 0.3 \text{ mol mol}^{-1}$. Therefore, the Rubisco-limited assimilation rate is:

$$f_c^r(\mathbf{x}) = \frac{V_{cmax}(\mathbf{x})c_i}{K_c(\mathbf{x}) \left(1 + \frac{C_{o,a}}{K_o(\mathbf{x})}\right) + sc_a(\mathbf{x})} \quad (1.79)$$

The Fickian diffusion through stomata is given as:

$$f_c = k_{st} (c_a - c_i) \quad (1.80)$$

$$f_v = k_{st,v} \left(\frac{p_{v,i} - p_v}{p} \right) \quad (1.81)$$

where $k_{st,v}$ ($\text{mol m}^{-2} \text{ s}^{-1}$) is the stomatal conductance to water vapor:

$$k_{st,v} = a_c k_{st} \quad (1.82)$$

where $a_c = 1.6$ is the relative diffusion of water vapor to CO_2 . Furthermore, $p_{v,i} = p_{v,sat}(T_l)$ (Pa) is the intercellular vapor pressure inside the stomatal cavity assumed to be at saturation at the leaf temperature T_l . Thus, equating Fickian CO_2 flux to the Farquhar biochemical demand, we have:

$$f_c = \frac{a_1 c_i}{a_2 + sc_a} = k_{st} (c_a - c_i) \quad (1.83)$$

where c_i is the unknown. Rewriting, we get:

$$c_i = c_a \frac{a_2 + sc_a}{a_1/k_{st} + a_2 + sc_a} \quad (1.84)$$

and substituting c_i into the biochemical demand function, the assimilation rate is a closed-problem as:

$$f_c = \frac{k_{st} a_1 c_a}{a_1 + k_{st} (a_2 + sc_a)} \quad (1.85)$$

Thus, the stomatal conductance can be finally obtained from the minimizing the problem:

$$\frac{\partial \mathcal{L}}{\partial k_{st}} = \frac{\partial f_c}{\partial k_{st}} - \lambda \frac{\partial f_v}{\partial k_{st}} = 0 \quad (1.86)$$

which becomes:

$$\frac{\partial}{\partial k_{st}} \left[\left(\frac{k_{st} a_1 c_a}{a_1 + k_{st} (a_2 + sc_a)} \right) - \lambda a_c k_{st} VPD \right] = 0 \quad (1.87)$$

where $VPD = (p_{v,i} - p_v) / p$ (mol mol^{-1}). The derivative becomes:

$$\frac{a_1^2 c_a}{[a_1 + k_{st} (a_2 + s c_a)]^2} - \lambda a_c VPD = 0 \quad (1.88)$$

Therefore, solving for k_{st} we obtain:

$$k_{st}(x) = \frac{a_1(x)}{a_2(x) + s c_a(x)} \left(-1 + \sqrt{\frac{c_a(x)}{a_c \lambda(\psi_l) VPD(x)}} \right) \quad (1.89)$$

where the marginal water use is empirically related to the leaf water potential $\lambda = \lambda(\psi_l)$ (Katul et al. 2010; Manoli et al. 2014). Therefore, the stomatal response change to water availability is reflected through the change in leaf water potential ψ_l . Additionally, in literature it is known that stomata does not completely close during night allowing for respiration. Therefore, taking this into account:

$$k_{st}(x) = \frac{a_1(x)}{a_2(x) + s c_a(x)} \left(-1 + \sqrt{\frac{c_a(x)}{a_c \lambda(\psi_l) VPD(x)}} \right) + k_{st,n} \quad (1.90)$$

where $k_{st,n}$ ($\text{mol m}^{-2} \text{s}^{-1}$) is the nocturnal stomatal conductance ($k_{st,n} = 0.018 \text{ mol m}^{-2} \text{s}^{-1}$ (Manoli et al. 2014)). The intercellular CO_2 concentration simplifies to:

$$c_i = c_a \left(1 - \sqrt{\frac{a \lambda VPD}{c_a}} \right) \quad (1.91)$$

and thus closing the (with the exception of λ) the photosynthetic rate. So far, the stomatal conductance model is derived by neglecting the contribution of boundary layer conductance k_b (i.e. g_b (in plant-science), or inverse of boundary layer resistance r_b , assumed to be equivalent to aerodynamic resistance r_a). Therefore, the effective stomatal conductance k_{st}^* is defined as:

$$k_{st}^* = \frac{k_{st} k_b}{k_{st} + k_b} \quad (1.92)$$

assuming the resistance are in series. Therefore, the plant fluxes into atmosphere become:

$$f_c = k_{st}^* f_c \quad (1.93)$$

$$f_v = k_{st,v}^* f_v \quad (1.94)$$

Furthermore, the fluxes in units $\text{kg m}^{-2} \text{s}^{-1}$ can be simply determined as:

$$g_{c,leaf} = M_c f_c \quad (1.95)$$

$$g_{v,leaf} = M_v f_v \quad (1.96)$$

where $M_c = 4.401 \times 10^{-2} \text{ kg mol}^{-1}$ and $M_v = 1.8015 \times 10^{-2} \text{ kg mol}^{-1}$ are the molar mass of CO_2 and water vapor, respectively.

1.3.5 Marginal water use

The marginal water use efficiency (WUE) λ or the cost-parameter for the cost of water lost from leaves. The marginal WUE should change over time depending on the water availability Manzoni et al. (2011). The marginal WUE is estimated from photosynthesis, transpiration and stomatal conductance measurement, obtained simply from the gradient of WUE:

$$\text{WUE} = \frac{f_c}{f_v} \quad (1.97)$$

The observations derive a marginal WUE as a function of leaf water potential ψ_L :

$$\lambda(\psi_L) = \lambda_{\max}^* \frac{c_a}{c_a^*} \exp \left\{ -\beta \left(\langle \psi_L \rangle_{24h} - \psi_{L,\max} \right)^2 \right\} \quad (1.98)$$

where ψ_L is assumed to vary slowly such that $\langle \psi_L \rangle_{24h}$ is fixed within the secant iteration, λ_{\max}^* is the marginal WUE under well-watered soil condition at reference CO_2 concentration $c_a^* = 400 \text{ } \mu\text{mol mol}^{-1}$ or parts-per-million (ppm), β is the plant-specific sensitivity parameter (Huang et al. 2017).

1.3.6 Numerical method for determining leaf water potential

The water transport through the plant from soil to root, from root to xylem, through the xylem, and finally, from leaf stomata to air is a closed-problem once the leaf water potential is known. The leaf water potential is determined from the following constraint:

$$G_{v,\text{leaf}}(\psi_L) = G_{v,\text{root}}(\psi_L) \quad (1.99)$$

or as an optimization problem, it is defined as:

$$\arg \min_{\psi_L} \mathcal{G}(\psi_L) = \left| G_{v,\text{leaf}} - G_{v,\text{root}} \right| \quad (1.100)$$

As this is a non-linear closure problem (Manoli et al. 2014), a secant method is employed to iteratively converge to the leaf water potential. The $j+1^{\text{th}}$ leaf water potential estimate is determined as:

$$\psi_L^{j+1} = \psi_L^j - G(\psi_L^j) \frac{\psi_L^j - \psi_L^{j-1}}{G(\psi_L^j) - G(\psi_L^{j-1})} \quad (1.101)$$

where the initial estimate of $\psi_L^{j=0} = 0 \text{ MPa}$ and $\psi_L^{j=1} = -10 \text{ MPa}$ and with the additional constraint that $-10 \leq \psi_L \leq 0 \text{ MPa}$, enforcing that leaf water potential is negative

and not larger than -10 MPa (typically known to lower). The detailed solution strategy for determining for coupling all the models is detailed in next section.

1.3.7 Solution strategy for coupling

The numerical model for air domain, solid domains (soil, ground, building), the radiation model and the vegetation model is implemented into OpenFOAM. The solid and air domains are coupled at regular intervals t^m defined as exchange timesteps or air time steps (Saneinejad et al. 2014; Kubilay et al. 2018). The fluxes between air and solid domain consisting for thermal, moisture and radiative transfers are coupled at this step, chosen to be 10 min. At each t^m , the air domain is assumed to the quasi-steady and solving using steady-state RANS approach converged when residuals of ρ , u , h , k , ε are below threshold. During the steady-state computation, the leaf energy balance is evaluated periodically to correct the heat and mass fluxes, $q_{sen,leaf}$ and $g_{v,leaf}$, respectively.

The algorithm of the air domain $t^m \rightarrow t^{m+1}$ with Δt^m the air domain pseudo-timesteps of 10 min, is as follows:

1. Update the radiation fields in the air domain using q_{rad} from building surfaces and determined $q_{r,lw}$ and $q_{r,sw}$.
2. Solve the energy balance at the leaf surface:
 - a) Determine the radiative flux $q_{rad,leaf}$ using ??.
 - b) Calculate the stomatal and aerodynamic resistances r_a and r_s using ?? and ??, respectively. Note that $r_s = (k_{st,v})^{-1}$ and $\lambda(\psi_L)$ is constant. Therefore, r_s is only dependent on assimilation rate f_c , and the VPD.
 - c) Perform an initial estimate of leaf temperature $T_{leaf} = T$.
 - d) Calculate the saturated vapor pressure at the leaf surface $p_{vsat,leaf} = f(T_{leaf})$.
 - e) Calculate the latent heat flux $q_{lat,leaf}$ using ??.
 - f) Correct the leaf temperature T_{leaf} using ??.
 - g) Repeat steps (d) to (f) until the leaf temperature has converged with a convergence criterion of 10^{-8} .
3. Calculate all vegetation source terms s_ρ , s_u , s_T , s_w , s_k and s_ε using Eqs. ???????????.
4. Solve for the steady-state flow field for t^{m+1} , ??-??.
5. Repeat steps (2) to (4) until residuals of ??-?? have reached the convergence limit of 10^{-8} , $\delta f^{m+1} \leq \epsilon$

For every m exchange timesteps, the solid domains are solved using a transient approach with n adaptive solid timesteps (Janssen 2002; Kubilay et al. 2018). The solution from the air domain consisting of thermal and moisture fluxes are taken as boundary conditions for the solid domain equations. For each internal n iterations,

the thermal and radiative fluxes from solid domain are corrected until converges. For each of n solid timesteps, the linearized heat and mass transport equation are solved using k Picard iterations. Finally, for each k Picard iteration, the root water uptake is determined through j secant iterations minimizing the cost function \mathcal{G} .

The algorithm of the solid domain $t^m \leq t^n \leq t^{m+1}$, for $n \in 0, \dots, N$ with $t^{n=0} = t^m$ and $t^N = t^{m+1}$ such that $\Delta t^n < \Delta t^m$, is as follows:

1. Solve the linearized heat and mass transport equation using k Picard iteration such that $t^n \leq t^{n,k} \leq t^{n+1}$ with $k \in \{0, \dots, K\}$ where $t^{n,k=0} = t^n$ and $t^{n,k=K} = t^{n+1}$.
 - (a) Determine marginal WUE $\lambda (\langle \psi_L \rangle_{24h})$. As ψ_L is assumed to be slowly varying, λ is constant in secant iteration.
 - (b) Determine the stomatal conductance k_{st} , constant in the secant iteration.
 - (c) Determine the assimilation rate f_c and transpiration rate f_v , constant in the secant iteration.
 - (d) Determine the net transpiration rate $G_{v,leaf}$.
 - (e) Determine the effective soil-root conductance k_{sr}^* , constant in the secant iteration.
 - (f) Solve for leaf water potential $\psi_L^{n,k}$ of the k^{th} Picard iteration using $j \in \{0, \dots, J\}$ secant iterations.
 - (a) Initial guess of leaf water potential, $\psi_L^{j=0} = 0$ MPa, $\psi_L^{j=1} = -10$ MPa.
 - (b) Calculate the effective xylem conductance k_x^* .
 - (c) Calculate root water potential ψ_R^j .
 - (d) Calculate the root uptake $g_{v,root}$ and the net root uptake $G_{v,root}$.
 - (e) Determine the cost function \mathcal{G} .
 - (f) Correct the leaf water potential using secant method $\psi_L^j \rightarrow \psi_L^{j+1}$.
 - (g) Repeat till leaf water potential converged, $\delta \psi_L \leq \epsilon$.
 - (g) Calculate the sink in soil moisture due to root water uptake s_r .
 - (h) Solve linearized form of heat and mass equation using PCG until $\delta p_c = \delta T = 10^{-2}$, repeating steps before.
2. Use the final surface temperature T_s^N to update the radiation model updating $q_{r,lw}$ fluxes from all surfaces.
3. Final surface temperature T_s^N and moisture fluxes g_v^N are boundary condition for the air domain for $t^m \rightarrow t^{m+1}$.

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COLOPHON

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