

The BSSN evolution equations. Pt.1

These are the equations as given in the Phys.Rev.D (62) 044034 paper.

$$\partial_t \bar{g}_{ij} = -2N \bar{A}_{ij} \tag{eq09.01}$$

$$\partial_t \phi = -\frac{1}{6} N \text{tr} K \tag{eq10.01}$$

$$\partial_t \text{tr} K = -g^{ij} D_{ij} N + N \left(\bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} \text{tr} K^2 \right) \tag{eq11.01}$$

$$\partial_t \bar{A}_{ij} = N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a{}_j \right) + \exp(-4\phi) \left(N R_{ij} - D_{ij} N - \frac{1}{3} g_{ij} (N R_{ab} - D_{ab} N) g^{ab} \right) \tag{eq12.01}$$

$$\partial_t \bar{\Gamma}^i = -2 \bar{A}^{ij} \partial_j N + 2N \left(\bar{\Gamma}^i{}_{jk} \bar{A}^{kj} - \frac{2}{3} \bar{g}^{ij} \partial_j \text{tr} K + 6 \bar{A}^{ij} \partial_j \phi \right) \tag{eq20.01}$$

$$R_{ab} = \bar{R}_{ab} + R^\phi_{ab} \tag{eq14.01}$$

$$\begin{aligned} &= -2 \bar{D}_{ab} \phi - 2 \bar{g}_{ab} \bar{g}^{cd} \bar{D}_{cd} \phi + 4 \bar{D}_a \phi \bar{D}_b \phi - 4 \bar{g}_{ab} \bar{g}^{cd} \bar{D}_c \phi \bar{D}_d \phi - \frac{1}{2} \bar{g}^{lm} \partial_{lm} \bar{g}_{ab} + \frac{1}{2} \bar{g}_{ka} \partial_b \bar{\Gamma}^k + \frac{1}{2} \bar{g}_{kb} \partial_a \bar{\Gamma}^k + \frac{1}{2} \bar{\Gamma}^k \bar{\Gamma}_{abk} + \frac{1}{2} \bar{\Gamma}^k \bar{\Gamma}_{bak} \\ &\quad + \bar{g}^{lm} \bar{g}^{ke} (\bar{\Gamma}_{ela} \bar{\Gamma}_{bkm} + \bar{\Gamma}_{elb} \bar{\Gamma}_{akm} + \bar{\Gamma}_{kam} \bar{\Gamma}_{elb}) \end{aligned} \tag{eq14.02}$$

The next task is to rewrite the right hand sides of this set of equations in terms of the BSSN variables and their partial derivatives.

The results are summarised on the following page.

The BSSN evolution equations. Pt.2

Notice that equation (14) contains $\partial\bar{\Gamma}^i$ but not $\bar{\Gamma}^i$. This is a consequence of the advice given by Alcubierre, Bruggmann et al. (Phys Rev D (67) 084023, 2nd-3rd paragraph on pg. 084023-4)

... if one wants to achieve numerical stability. In the computer code we do not use the numerically evolved $\bar{\Gamma}^i$ in all places, but we follow this rule:

Partial derivatives $\partial_j\bar{\Gamma}^i$ are computed as finite differences of the independent variables $\bar{\Gamma}^i$ that are evolved using ...

Compare equation (14) on the previous page and that on this page.

$$\partial_t\bar{g}_{ij} = -2N\bar{A}_{ij} \quad (\text{eq09.99})$$

$$\partial_t\phi = -\frac{1}{6}N\text{tr}K \quad (\text{eq10.99})$$

$$\partial_t\text{tr}K = N\bar{A}_{ab}\bar{A}^{ab} + \frac{1}{3}N\text{tr}K^2 + \exp(-4\phi) \left(-\bar{g}^{ab}\partial_{ab}N - \partial_aN\partial_b\bar{g}^{ab} - 2\bar{g}^{ab}\partial_a\phi\partial_bN \right) \quad (\text{eq11.99})$$

$$\begin{aligned} \partial_t\bar{A}_{ij} = N\text{tr}K\bar{A}_{ij} - 2N\bar{A}_{ia}\bar{A}_{jb}\bar{g}^{ab} + \exp(-4\phi) \left(NR_{ij} - \partial_{ij}N + \frac{1}{2}\bar{g}^{ab}\partial_aN\partial_j\bar{g}_{ib} + \frac{1}{2}\bar{g}^{ab}\partial_aN\partial_j\bar{g}_{ib} - \frac{1}{2}\bar{g}^{ab}\partial_aN\partial_b\bar{g}_{ij} + 2\partial_i\phi\partial_jN + 2\partial_j\phi\partial_iN \right. \\ \left. - \frac{4}{3}\bar{g}_{ij}\bar{g}^{ab}\partial_a\phi\partial_bN - \frac{1}{3}N\bar{g}_{ij}\bar{g}^{ab}R_{ab} + \frac{1}{3}\bar{g}_{ij}\bar{g}^{ab}\partial_{ab}N + \frac{1}{3}\bar{g}_{ij}\partial_aN\partial_b\bar{g}^{ab} \right) \end{aligned} \quad (\text{eq12.99})$$

$$\partial_t\bar{\Gamma}^i = -2\bar{A}^{ia}\partial_aN + 2N\bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - N\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{4}{3}N\bar{g}^{ia}\partial_a\text{tr}K + 12N\bar{A}^{ia}\partial_a\phi \quad (\text{eq20.99})$$

$$\begin{aligned} R_{ab} = -2\partial_{ab}\phi + \bar{g}^{cd}\partial_a\phi\partial_b\bar{g}_{cd} + \bar{g}^{cd}\partial_d\phi\partial_b\bar{g}_{ac} - \bar{g}^{cd}\partial_d\phi\partial_c\bar{g}_{ab} + 12\partial_a\phi\partial_b\phi - 2\bar{g}_{ab}\bar{g}^{cd}\partial_{cd}\phi - 2\bar{g}_{ab}\partial_d\phi\partial_c\bar{g}^{cd} - 12\bar{g}_{ab}\bar{g}^{cd}\partial_c\phi\partial_d\phi - \frac{1}{2}\bar{g}^{cd}\partial_{cd}\bar{g}_{ab} + \frac{1}{2}\bar{g}_{ac}\partial_b\bar{\Gamma}^c \\ + \frac{1}{2}\bar{g}_{bc}\partial_a\bar{\Gamma}^c - \frac{1}{2}\partial_c\bar{g}_{ab}\partial_d\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{bc}\partial_a\bar{g}^{dc} - \frac{1}{2}\bar{g}^{ed}\bar{g}^{cf}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{2}\bar{g}^{ef}\bar{g}^{cd}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{4}\partial_b\bar{g}_{cd}\partial_a\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{ac}\partial_b\bar{g}^{dc} \end{aligned} \quad (\text{eq14.99})$$

The full details of the above computations can be found on the following pages.

PhysRevD.62.044034 equation (9)

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-09.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  DgBarDt := \partial_{t}{gBar_{ij}}.          # cdb(eq09.00,DgBarDt)
10 DgBarDt := -2 N ABar_{i j}.                # cdb(eq09.01,DgBarDt)
11
12 canonicalise (DgBarDt)                      # cdb(eq09.02,DgBarDt)  # no change
13                                              # cdb(eq09.99,DgBarDt)  # no change
14
15 cdblib.put ('DgBarDt',DgBarDt,jsonfile)
```

$$\partial_t \bar{g}_{ij} = -2N \bar{A}_{ij} \quad (\text{eq09.02})$$

PhysRevD.62.044034 equation (10)

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-10.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  DphiDt := \partial_{t}\{\phi\}.          # cdb(eq10.00,DphiDt)
10 DphiDt := -(1/6) N trK.                # cdb(eq10.01,DphiDt)
11
12 canonicalise (DphiDt)                  # cdb(eq10.02,DphiDt) # no change
13                                         # cdb(eq10.99,DphiDt) # no change
14
15 cdblib.put ('DphiDt',DphiDt,jsonfile)
```

$$\partial_t \phi = -\frac{1}{6} N \text{tr} K \quad (\text{eq10.02})$$

PhysRevD.62.044034 equation (11)

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-11.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  DtrKDt := \partial_{t}{trK}.          # cdb(eq11.00,DtrKDt)
10 DtrKDt := - g^{i j} D_{i j}{N}
11         + N (ABar_{i j} ABar^{i j} + (1/3) trK**2).
12
13                                     # cdb(eq11.01,DtrKDt)
14
15 substitute (DtrKDt, defD2)          # cdb(eq11.02,DtrKDt)
16 substitute (DtrKDt, defGamma2GammaBar) # cdb(eq11.03,DtrKDt)
17
18 foo := g^{a b} -> \exp(-4\phi) gBar^{a b}.
19
20 substitute (DtrKDt, foo)            # cdb(eq11.04,DtrKDt)
21 distribute (DtrKDt)                 # cdb(eq11.05,DtrKDt)
22 eliminate_kronecker (DtrKDt)       # cdb(eq11.06,DtrKDt)
23 canonicalise (DtrKDt)               # cdb(eq11.07,DtrKDt)
24 substitute (DtrKDt, defGBarSq)      # cdb(eq11.08,DtrKDt)
25
26 DtrKDt = product_sort (DtrKDt)
27
28 factor_out (DtrKDt, $\exp(-4\phi)$) # cdb(eq11.08,DtrKDt)
29 substitute (DtrKDt, defGammaBarU)   # cdb(eq11.09,DtrKDt)
30 distribute (DtrKDt)                 # cdb(eq11.10,DtrKDt)
31
32 DtrKDt = product_sort (DtrKDt)      # cdb(eq11.11,DtrKDt)
33
34 canonicalise (DtrKDt)               # cdb(eq11.12,DtrKDt)
35
36 foo := gBar^{b c} \partial_{a}{gBar_{b c}} -> 0. # follows from det(g) = 1
```

```

37 bah := gBar^{e b} gBar^{f c} \partial_{a}{gBar_{b c}} -> - \partial_{a}{gBar^{e f}}.
38
39 substitute (DtrKDt, foo)           # cdb(eq11.13,DtrKDt)
40 substitute (DtrKDt, bah)           # cdb(eq11.14,DtrKDt)
41
42 DtrKDt = product_sort (DtrKDt)
43
44 canonicalise (DtrKDt)               # cdb(eq11.15,DtrKDt)
45 factor_out   (DtrKDt, $\exp(-4\phi)$) # cdb(eq11.16,DtrKDt)
46                                     # cdb(eq11.99,DtrKDt)
47
48 cdblib.put ('DtrKDt',DtrKDt,jsonfile)

```

$$\partial_t \text{tr} K = -g^{ij} D_{ij} N + N \left(\bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} \text{tr} K^2 \right) \quad (\text{eq11.01})$$

$$= -g^{ij} (\partial_{ij} N - \Gamma^c_{ij} \partial_c N) + N \left(\bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} \text{tr} K^2 \right) \quad (\text{eq11.02})$$

$$= -g^{ij} (\partial_{ij} N - (\bar{\Gamma}^c_{ij} + 2\bar{g}^c_j \partial_i \phi + 2\bar{g}^c_i \partial_j \phi - 2\bar{g}^{ce} \bar{g}_{ij} \partial_e \phi) \partial_c N) + N \left(\bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} \text{tr} K^2 \right) \quad (\text{eq11.03})$$

$$= -\exp(-4\phi) \bar{g}^{ij} (\partial_{ij} N - (\bar{\Gamma}^c_{ij} + 2\bar{g}^c_j \partial_i \phi + 2\bar{g}^c_i \partial_j \phi - 2\bar{g}^{ce} \bar{g}_{ij} \partial_e \phi) \partial_c N) + N \left(\bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} \text{tr} K^2 \right) \quad (\text{eq11.04})$$

$$= -\exp(-4\phi) \bar{g}^{ij} \partial_{ij} N + \exp(-4\phi) \bar{g}^{ij} \bar{\Gamma}^c_{ij} \partial_c N + 2 \exp(-4\phi) \bar{g}^{ij} \bar{g}^c_j \partial_i \phi \partial_c N + 2 \exp(-4\phi) \bar{g}^{ij} \bar{g}^c_i \partial_j \phi \partial_c N - 2 \exp(-4\phi) \bar{g}^{ij} \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi \partial_c N \\ + N \bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} N \text{tr} K^2 \quad (\text{eq11.05})$$

$$= -\exp(-4\phi) \bar{g}^{ij} \partial_{ij} N + \exp(-4\phi) \bar{g}^{ij} \bar{\Gamma}^c_{ij} \partial_c N + 2 \exp(-4\phi) \bar{g}^{ic} \partial_i \phi \partial_c N + 2 \exp(-4\phi) \bar{g}^{cj} \partial_j \phi \partial_c N - 2 \exp(-4\phi) \bar{g}^{ij} \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi \partial_c N \\ + N \bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} N \text{tr} K^2 \quad (\text{eq11.06})$$

$$= -\exp(-4\phi) \bar{g}^{ij} \partial_{ij} N + \exp(-4\phi) \bar{g}^{ci} \bar{\Gamma}^j_{cj} \partial_j N + 2 \exp(-4\phi) \bar{g}^{ci} \partial_c \phi \partial_i N + 2 \exp(-4\phi) \bar{g}^{cj} \partial_c \phi \partial_j N - 2 \exp(-4\phi) \bar{g}^{ce} \bar{g}^{ij} \bar{g}_{ce} \partial_i \phi \partial_j N \\ + N \bar{A}_{ij} \bar{A}^{ij} + \frac{1}{3} N \text{tr} K^2 \quad (\text{eq11.07})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 + \exp(-4\phi) (-\bar{g}^{ab} \partial_{ab} N + \bar{g}^{ab} \bar{\Gamma}^c_{ab} \partial_c N - 2\bar{g}^{ab} \partial_a \phi \partial_b N) \quad (\text{eq11.08})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 + \exp(-4\phi) \left(-\bar{g}^{ab} \partial_{ab} N + \frac{1}{2} \bar{g}^{ab} \bar{g}^{ce} (\partial_a \bar{g}_{eb} + \partial_b \bar{g}_{ae} - \partial_e \bar{g}_{ab}) \partial_c N - 2\bar{g}^{ab} \partial_a \phi \partial_b N \right) \quad (\text{eq11.09})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \exp(-4\phi) \bar{g}^{ab} \partial_{ab} N + \frac{1}{2} \exp(-4\phi) \bar{g}^{ab} \bar{g}^{ce} \partial_a \bar{g}_{eb} \partial_c N + \frac{1}{2} \exp(-4\phi) \bar{g}^{ab} \bar{g}^{ce} \partial_b \bar{g}_{ae} \partial_c N - \frac{1}{2} \exp(-4\phi) \bar{g}^{ab} \bar{g}^{ce} \partial_e \bar{g}_{ab} \partial_c N \\ - 2 \exp(-4\phi) \bar{g}^{ab} \partial_a \phi \partial_b N \quad (\text{eq11.10})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N + \frac{1}{2} \bar{g}^{cb} \bar{g}^{da} \exp(-4\phi) \partial_d N \partial_c \bar{g}_{ab} + \frac{1}{2} \bar{g}^{ac} \bar{g}^{db} \exp(-4\phi) \partial_d N \partial_c \bar{g}_{ab} - \frac{1}{2} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_c N \partial_d \bar{g}_{ab} \\ - 2\bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N \quad (\text{eq11.11})$$

$$\partial_t \text{tr} K = N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N + \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_c \bar{g}_{bd} - \frac{1}{2} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{cd} - 2 \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N \quad (\text{eq11.12})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N + \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_c \bar{g}_{bd} - 2 \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N \quad (\text{eq11.13})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \partial_c \bar{g}^{ac} \exp(-4\phi) \partial_a N - 2 \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N \quad (\text{eq11.14})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 - \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \exp(-4\phi) \partial_a N \partial_b \bar{g}^{ab} - 2 \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N \quad (\text{eq11.15})$$

$$= N \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3} N \text{tr} K^2 + \exp(-4\phi) (-\bar{g}^{ab} \partial_{ab} N - \partial_a N \partial_b \bar{g}^{ab} - 2 \bar{g}^{ab} \partial_a \phi \partial_b N) \quad (\text{eq11.16})$$

PhysRevD.62.044034 equation (12)

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-12.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  DABarDt := \partial_{t}{ABar_{i j}}.          # cdb(eq12.00,DABarDt)
10 DABarDt := N (trK ABar_{i j} - 2 ABar_{i a} ABar^{a}_{j})
11           + \exp(-4\phi) (N R_{i j} - D_{i j}{N}
12                       - (1/3) g_{i j} (N R_{a b} - D_{a b}{N}) g^{a b}).
13
14                                     # cdb(eq12.01,DABarDt)
15
16  # -----
17
18  substitute (DABarDt, defD2)                # cdb(eq12.02,DABarDt)
19  substitute (DABarDt, defGamma2GammaBar)    # cdb(eq12.03,DABarDt)
20
21  foo := g_{a b} -> \exp(4\phi) gBar_{a b}.
22  bah := g^{a b} -> \exp(-4\phi) gBar^{a b}.
23
24  substitute (DABarDt, foo)                  # cdb(eq12.04,DABarDt)
25  substitute (DABarDt, bah)                  # cdb(eq12.05,DABarDt)
26  distribute (DABarDt)                      # cdb(eq12.06,DABarDt)
27  eliminate_kronecker (DABarDt)             # cdb(eq12.07,DABarDt)
28  substitute (DABarDt, defGBarSq)            # cdb(eq12.08,DABarDt)
29
30  DABarDt = product_sort (DABarDt)           # cdb(eq12.09,DABarDt)
31
32  rename_dummies (DABarDt)                  # cdb(eq12.10,DABarDt)
33  canonicalise (DABarDt)                    # cdb(eq12.11,DABarDt)
34
35  map_sympy (DABarDt, "simplify")           # cdb(eq12.12,DABarDt)
36  factor_out (DABarDt, $\exp(-4\phi)$)      # cdb(eq12.13,DABarDt)
```

```

37
38 foo := ABar^{a}_{b} -> gBar^{a c} ABar_{c b}.
39
40 substitute (DABarDt, foo)
41
42 DABarDt = product_sort (DABarDt)           # cdb(eq12.14,DABarDt)
43
44 substitute (DABarDt,defGammaBarU)         # cdb(eq12.15,DABarDt)
45 distribute (DABarDt)
46
47 DABarDt = product_sort (DABarDt)           # cdb(eq12.16,DABarDt)
48
49 canonicalise (DABarDt)                     # cdb(eq12.17,DABarDt)
50
51 foo := gBar^{b c} \partial_{a}{gBar_{b c}} -> 0.   # follows from det(g) = 1
52 bah := gBar^{e b} gBar^{f c} \partial_{a}{gBar_{b c}} -> - \partial_{a}{gBar^{e f}}.
53
54 substitute (DABarDt,foo)                   # cdb(eq12.18,DABarDt)
55 substitute (DABarDt,bah)                   # cdb(eq12.19,DABarDt)
56
57 DABarDt = product_sort (DABarDt)
58
59 canonicalise (DABarDt)                     # cdb(eq12.20,DABarDt)
60 factor_out (DABarDt, $\exp(-4\phi)$)       # cdb(eq12.21,DABarDt)
61
62                                           # cdb(eq12.99,DABarDt)
63
64 cdblib.put ('DABarDt',DABarDt,jsonfile)

```

$$\partial_t \bar{A}_{ij} = N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a_j \right) + \exp(-4\phi) \left(N R_{ij} - D_{ij} N - \frac{1}{3} g_{ij} (N R_{ab} - D_{ab} N) g^{ab} \right) \quad (\text{eq12.01})$$

$$= N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a_j \right) + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + \Gamma^c_{ij} \partial_c N - \frac{1}{3} g_{ij} (N R_{ab} - \partial_{ab} N + \Gamma^c_{ab} \partial_c N) g^{ab} \right) \quad (\text{eq12.02})$$

$$= N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a_j \right) + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + (\bar{\Gamma}^c_{ij} + 2 \bar{g}^c_j \partial_i \phi + 2 \bar{g}^c_i \partial_j \phi - 2 \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi) \partial_c N \right. \\ \left. - \frac{1}{3} g_{ij} (N R_{ab} - \partial_{ab} N + (\bar{\Gamma}^c_{ab} + 2 \bar{g}^c_b \partial_a \phi + 2 \bar{g}^c_a \partial_b \phi - 2 \bar{g}^{ce} \bar{g}_{ab} \partial_e \phi) \partial_c N) g^{ab} \right) \quad (\text{eq12.03})$$

$$= N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a_j \right) + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + (\bar{\Gamma}^c_{ij} + 2 \bar{g}^c_j \partial_i \phi + 2 \bar{g}^c_i \partial_j \phi - 2 \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi) \partial_c N \right. \\ \left. - \frac{1}{3} \exp(4\phi) \bar{g}_{ij} (N R_{ab} - \partial_{ab} N + (\bar{\Gamma}^c_{ab} + 2 \bar{g}^c_b \partial_a \phi + 2 \bar{g}^c_a \partial_b \phi - 2 \bar{g}^{ce} \bar{g}_{ab} \partial_e \phi) \partial_c N) g^{ab} \right) \quad (\text{eq12.04})$$

$$= N \left(\text{tr} K \bar{A}_{ij} - 2 \bar{A}_{ia} \bar{A}^a_j \right) + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + (\bar{\Gamma}^c_{ij} + 2 \bar{g}^c_j \partial_i \phi + 2 \bar{g}^c_i \partial_j \phi - 2 \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi) \partial_c N \right. \\ \left. - \frac{1}{3} \exp(4\phi) \bar{g}_{ij} (N R_{ab} - \partial_{ab} N + (\bar{\Gamma}^c_{ab} + 2 \bar{g}^c_b \partial_a \phi + 2 \bar{g}^c_a \partial_b \phi - 2 \bar{g}^{ce} \bar{g}_{ab} \partial_e \phi) \partial_c N) \exp(-4\phi) \bar{g}^{ab} \right) \quad (\text{eq12.05})$$

$$= N \text{tr} K \bar{A}_{ij} - 2 N \bar{A}_{ia} \bar{A}^a_j + \exp(-4\phi) N R_{ij} - \exp(-4\phi) \partial_{ij} N + \exp(-4\phi) \bar{\Gamma}^c_{ij} \partial_c N + 2 \exp(-4\phi) \bar{g}^c_j \partial_i \phi \partial_c N + 2 \exp(-4\phi) \bar{g}^c_i \partial_j \phi \partial_c N \\ - 2 \exp(-4\phi) \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi \partial_c N - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} N R_{ab} \exp(-4\phi) \bar{g}^{ab} + \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_{ab} N \exp(-4\phi) \bar{g}^{ab} \\ - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{\Gamma}^c_{ab} \partial_c N \exp(-4\phi) \bar{g}^{ab} - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{g}^c_b \partial_a \phi \partial_c N \exp(-4\phi) \bar{g}^{ab} \\ - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{g}^c_a \partial_b \phi \partial_c N \exp(-4\phi) \bar{g}^{ab} + \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{g}^{ce} \bar{g}_{ab} \partial_e \phi \partial_c N \exp(-4\phi) \bar{g}^{ab} \quad (\text{eq12.06})$$

$$= N \text{tr} K \bar{A}_{ij} - 2 N \bar{A}_{ia} \bar{A}^a_j + \exp(-4\phi) N R_{ij} - \exp(-4\phi) \partial_{ij} N + \exp(-4\phi) \bar{\Gamma}^c_{ij} \partial_c N + 2 \exp(-4\phi) \partial_i \phi \partial_j N + 2 \exp(-4\phi) \partial_j \phi \partial_i N \\ - 2 \exp(-4\phi) \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi \partial_c N - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} N R_{ab} \exp(-4\phi) \bar{g}^{ab} + \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_{ab} N \exp(-4\phi) \bar{g}^{ab} \\ - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{\Gamma}^c_{ab} \partial_c N \exp(-4\phi) \bar{g}^{ab} - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_a \phi \partial_b N \exp(-4\phi) \bar{g}^{ab} \\ - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_b \phi \partial_a N \exp(-4\phi) \bar{g}^{ab} + \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{g}^{ce} \bar{g}_{ab} \partial_e \phi \partial_c N \exp(-4\phi) \bar{g}^{ab} \quad (\text{eq12.07})$$

$$\begin{aligned}
\partial_t \bar{A}_{ij} = & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j + \exp(-4\phi) N R_{ij} - \exp(-4\phi) \partial_{ij} N + \exp(-4\phi) \bar{\Gamma}^c_{ij} \partial_c N + 2 \exp(-4\phi) \partial_i \phi \partial_j N + 2 \exp(-4\phi) \partial_j \phi \partial_i N \\
& - 2 \exp(-4\phi) \bar{g}^{ce} \bar{g}_{ij} \partial_e \phi \partial_c N - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} N R_{ab} \exp(-4\phi) \bar{g}^{ab} + \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_{ab} N \exp(-4\phi) \bar{g}^{ab} \\
& - \frac{1}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{\Gamma}^c_{ab} \partial_c N \exp(-4\phi) \bar{g}^{ab} - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_a \phi \partial_b N \exp(-4\phi) \bar{g}^{ab} \\
& - \frac{2}{3} \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \partial_b \phi \partial_a N \exp(-4\phi) \bar{g}^{ab} + 2 \exp(-4\phi) \exp(4\phi) \bar{g}_{ij} \bar{g}^{ce} \partial_e \phi \partial_c N \exp(-4\phi) \quad (\text{eq12.08})
\end{aligned}$$

$$\begin{aligned}
= & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \bar{\Gamma}^a_{ij} \exp(-4\phi) \partial_a N + 2 \partial_i \phi \exp(-4\phi) \partial_j N + 2 \partial_j \phi \exp(-4\phi) \partial_i N \\
& - 2 \bar{g}_{ij} \bar{g}^{ab} \partial_b \phi \exp(-4\phi) \partial_a N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_{ab} N \\
& - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{\Gamma}^c_{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_c N - \frac{2}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_b N \\
& + \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_b \phi \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_a N \quad (\text{eq12.09})
\end{aligned}$$

$$\begin{aligned}
= & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \bar{\Gamma}^a_{ij} \exp(-4\phi) \partial_a N + 2 \partial_i \phi \exp(-4\phi) \partial_j N + 2 \partial_j \phi \exp(-4\phi) \partial_i N \\
& - 2 \bar{g}_{ij} \bar{g}^{ba} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_{ab} N \\
& - \frac{1}{3} \bar{g}_{ij} \bar{g}^{bc} \bar{\Gamma}^a_{bc} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_a N - \frac{2}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_b N \\
& + \frac{4}{3} \bar{g}_{ij} \bar{g}^{ba} \partial_a \phi \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_b N \quad (\text{eq12.10})
\end{aligned}$$

$$\begin{aligned}
= & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \bar{\Gamma}^a_{ij} \exp(-4\phi) \partial_a N + 2 \partial_i \phi \exp(-4\phi) \partial_j N + 2 \partial_j \phi \exp(-4\phi) \partial_i N \\
& - 2 \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_{ab} N \\
& - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{\Gamma}^c_{ab} \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_c N + \frac{2}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \exp(-4\phi) \exp(4\phi) \partial_b N \quad (\text{eq12.11})
\end{aligned}$$

$$\begin{aligned}
= & \text{tr} K N \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j + N \exp(-4\phi) R_{ij} - \exp(-4\phi) \partial_{ij} N + \bar{\Gamma}^a_{ij} \exp(-4\phi) \partial_a N + 2 \partial_i \phi \exp(-4\phi) \partial_j N + 2 \partial_j \phi \exp(-4\phi) \partial_i N \\
& - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \exp(-4\phi) \bar{g}_{ij} \bar{g}^{ab} R_{ab} + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{\Gamma}^c_{ab} \exp(-4\phi) \partial_c N \quad (\text{eq12.12})
\end{aligned}$$

$$\begin{aligned}
= & \text{tr} K N \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}^a_j \\
& + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + \bar{\Gamma}^a_{ij} \partial_a N + 2 \partial_i \phi \partial_j N + 2 \partial_j \phi \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_{ab} N - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{\Gamma}^c_{ab} \partial_c N \right) \quad (\text{eq12.13})
\end{aligned}$$

$$\begin{aligned}\partial_t \bar{A}_{ij} &= N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{aj} \bar{A}_{ib} \bar{g}^{ba} \\ &+ \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + \bar{\Gamma}^a_{ij} \partial_a N + 2\partial_i \phi \partial_j N + 2\partial_j \phi \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_{ab} N - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{\Gamma}^c_{ab} \partial_c N \right) \quad (\text{eq12.14})\end{aligned}$$

$$\begin{aligned}&= N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{aj} \bar{A}_{ib} \bar{g}^{ba} + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + \frac{1}{2} \bar{g}^{ae} (\partial_i \bar{g}_{ej} + \partial_j \bar{g}_{ie} - \partial_e \bar{g}_{ij}) \partial_a N + 2\partial_i \phi \partial_j N + 2\partial_j \phi \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \partial_b N \right. \\ &\quad \left. - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_{ab} N - \frac{1}{6} \bar{g}_{ij} \bar{g}^{ab} \bar{g}^{ce} (\partial_a \bar{g}_{eb} + \partial_b \bar{g}_{ae} - \partial_e \bar{g}_{ab}) \partial_c N \right) \quad (\text{eq12.15})\end{aligned}$$

$$\begin{aligned}&= N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{aj} \bar{A}_{ib} \bar{g}^{ba} + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \frac{1}{2} \bar{g}^{ba} \exp(-4\phi) \partial_b N \partial_i \bar{g}_{aj} + \frac{1}{2} \bar{g}^{ba} \exp(-4\phi) \partial_b N \partial_j \bar{g}_{ia} \\ &\quad - \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \exp(-4\phi) \partial_j N + 2\partial_j \phi \exp(-4\phi) \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \\ &\quad + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \frac{1}{6} \bar{g}_{ij} \bar{g}^{cb} \bar{g}^{da} \exp(-4\phi) \partial_d N \partial_c \bar{g}_{ab} - \frac{1}{6} \bar{g}_{ij} \bar{g}^{ac} \bar{g}^{db} \exp(-4\phi) \partial_d N \partial_c \bar{g}_{ab} + \frac{1}{6} \bar{g}_{ij} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_c N \partial_d \bar{g}_{ab} \quad (\text{eq12.16})\end{aligned}$$

$$\begin{aligned}&= N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}_{jb} \bar{g}^{ab} + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_i \bar{g}_{jb} + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_j \bar{g}_{ib} \\ &\quad - \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \exp(-4\phi) \partial_j N + 2\partial_j \phi \exp(-4\phi) \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \\ &\quad + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_c \bar{g}_{bd} + \frac{1}{6} \bar{g}_{ij} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{cd} \quad (\text{eq12.17})\end{aligned}$$

$$\begin{aligned}&= N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}_{jb} \bar{g}^{ab} + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_i \bar{g}_{jb} + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_j \bar{g}_{ib} \\ &\quad - \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \exp(-4\phi) \partial_j N + 2\partial_j \phi \exp(-4\phi) \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \\ &\quad + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N - \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \bar{g}^{cd} \exp(-4\phi) \partial_a N \partial_c \bar{g}_{bd} \quad (\text{eq12.18})\end{aligned}$$

$$\begin{aligned}
\partial_t \bar{A}_{ij} = & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}_{jb} \bar{g}^{ab} + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_i \bar{g}_{jb} + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_j \bar{g}_{ib} \\
& - \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \exp(-4\phi) \partial_j N + 2\partial_j \phi \exp(-4\phi) \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \\
& + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N + \frac{1}{3} \bar{g}_{ij} \partial_c \bar{g}^{ac} \exp(-4\phi) \partial_a N
\end{aligned} \tag{eq12.19}$$

$$\begin{aligned}
= & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}_{jb} \bar{g}^{ab} + N R_{ij} \exp(-4\phi) - \exp(-4\phi) \partial_{ij} N + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_i \bar{g}_{jb} + \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_j \bar{g}_{ib} \\
& - \frac{1}{2} \bar{g}^{ab} \exp(-4\phi) \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \exp(-4\phi) \partial_j N + 2\partial_j \phi \exp(-4\phi) \partial_i N - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \exp(-4\phi) \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} \exp(-4\phi) \\
& + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \exp(-4\phi) \partial_{ab} N + \frac{1}{3} \bar{g}_{ij} \exp(-4\phi) \partial_a N \partial_b \bar{g}^{ab}
\end{aligned} \tag{eq12.20}$$

$$\begin{aligned}
= & N \text{tr} K \bar{A}_{ij} - 2N \bar{A}_{ia} \bar{A}_{jb} \bar{g}^{ab} + \exp(-4\phi) \left(N R_{ij} - \partial_{ij} N + \frac{1}{2} \bar{g}^{ab} \partial_a N \partial_i \bar{g}_{jb} + \frac{1}{2} \bar{g}^{ab} \partial_a N \partial_j \bar{g}_{ib} - \frac{1}{2} \bar{g}^{ab} \partial_a N \partial_b \bar{g}_{ij} + 2\partial_i \phi \partial_j N + 2\partial_j \phi \partial_i N \right. \\
& \left. - \frac{4}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_a \phi \partial_b N - \frac{1}{3} N \bar{g}_{ij} \bar{g}^{ab} R_{ab} + \frac{1}{3} \bar{g}_{ij} \bar{g}^{ab} \partial_{ab} N + \frac{1}{3} \bar{g}_{ij} \partial_a N \partial_b \bar{g}^{ab} \right)
\end{aligned} \tag{eq12.21}$$

PhysRevD.62.044034 equation (14)

The advice given by Miguel Alcubierre, Bernd Bruggmann et al (Phys Rev D (67) 084023, 2nd-3rd paragraph on pg. 084023-4)

... if one wants to achieve numerical stability. In the computer code we do not use the numerically evolved $\bar{\Gamma}^i$ in all places, but we follow this rule:

Partial derivatives $\partial_j \bar{\Gamma}^i$ are computed as finite differences of the independent variables $\bar{\Gamma}^i$ that are evolved using ...

The Einstein Toolkit code uses the same rule – the only place where the *evolved* $\bar{\Gamma}^i$ are used is in computing the $\partial_j \bar{\Gamma}^i$ terms in the equation for \bar{R}_{ij} , that is equation (18) of the Phys Rev D (62) 044034 paper.

```

1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-14.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  Rphi := -2 DBar_{a b}{\phi} - 2 gBar_{a b} gBar^{c d} DBar_{c d}{\phi}
10         +4 DBar_{a}{\phi} DBar_{b}{\phi} - 4 gBar_{a b} gBar^{c d} DBar_{c}{\phi} DBar_{d}{\phi}.
11
12                                     # cdb(eq15.prd,Rphi)
13
14  RBar := - (1/2) gBar^{l m} \partial_{l m}{gBar_{a b}}
15         + (1/2) gBar_{k a} \partial_{b}{GammaBar^{k}}
16         + (1/2) gBar_{k b} \partial_{a}{GammaBar^{k}}
17         + (1/2) GammaBar^{k} GammaBar_{a b k}
18         + (1/2) GammaBar^{k} GammaBar_{b a k}
19         + gBar^{l m} gBar^{k e} ( GammaBar_{e l a} GammaBar_{b k m}
20                                 + GammaBar_{e l b} GammaBar_{a k m}
21                                 + GammaBar_{k a m} GammaBar_{e l b}).
22
23                                     # cdb(eq18.prd,RBar)
24
25  defRab := R_{a b} -> @(Rphi) + @(RBar).
26
27  Rab := RBar_{a b} + Rphi_{a b}.                                     # cdb(eq14.01,Rab)
28  Rab := R_{a b}.                                                  # cdb(eq14.00,Rab)
29
30  substitute (Rab, defRab)                                           # cdb(eq14.02,Rab)
31  substitute (Rab, defDBar1)                                         # cdb(eq14.03,Rab)
32  substitute (Rab, defDBar2)                                         # cdb(eq14.04,Rab)
33  substitute (Rab, defGamma2GammaBar)                               # cdb(eq14.05,Rab)
34  distribute (Rab)                                                  # cdb(eq14.06,Rab)
35  eliminate_kronecker (Rab)                                         # cdb(eq14.07,Rab)
36
37  Rab = product_sort (Rab)                                           # cdb(eq14.08,Rab)
38

```



```

39 rename_dummies (Rab) # cdb(eq14.09,Rab)
40 canonicalise (Rab) # cdb(eq14.10,Rab)
41
42 foo := GammaBar^{a} GammaBar_{b c a} -> gBar^{d e} GammaBar^{a}_{d e} GammaBar_{b c a}.
43
44 substitute (Rab, foo) # cdb(eq14.11,Rab)
45 substitute (Rab, defGBarSq) # cdb(eq14.12,Rab)
46 substitute (Rab, defGammaBarD) # cdb(eq14.13,Rab)
47 substitute (Rab, defGammaBarU) # cdb(eq14.14,Rab)
48 distribute (Rab) # cdb(eq14.15,Rab)
49
50 foo := \partial_{a}\{gBar_{b c}\} gBar^{b c} -> 0. # follows from det(g) = 1
51
52 substitute (Rab,foo) # cdb(eq14.16,Rab)
53 canonicalise (Rab) # cdb(eq14.17,Rab)
54
55 foo := gBar^{b e} gBar^{c f} \partial_{a}\{gBar_{b c}\} -> - \partial_{a}\{gBar^{e f}\}.
56 bah := gBar^{e b} gBar^{f c} \partial_{a}\{gBar_{b c}\} -> - \partial_{a}\{gBar^{e f}\}.
57 moo := gBar^{e b} gBar^{c f} \partial_{a}\{gBar_{b c}\} -> - \partial_{a}\{gBar^{e f}\}.
58
59 substitute (Rab,foo) # cdb(eq14.18,Rab)
60 substitute (Rab,bah) # cdb(eq14.19,Rab)
61 substitute (Rab,moo) # cdb(eq14.20,Rab)
62
63 Rab = product_sort (Rab) # cdb(eq14.21,Rab)
64 # cdb(eq14.99,Rab)
65
66 defRab := R_{a b} -> @(Rab). # used later in bssn-ricci-scalar.tex
67
68 cdblib.put ('Rab',Rab,jsonfile)
69 cdblib.put ('defRab',defRab,jsonfile)

```


There is a single term in this final expression that appears to be neither symmetric in ab nor part of a symmetric pair, namely

$$\partial_b \bar{g}_{cd} \partial_a \bar{g}^{cd}$$

It is, however, easy to show that this term is symmetric in ab . Start by noting that, for any \bar{g}_{ab} ,

$$\partial_a \bar{g}^{cd} = -\bar{g}^{ce} \bar{g}^{df} \partial_a \bar{g}_{ef}$$

Now contract both sides with $\partial_b \bar{g}_{cd}$ to obtain

$$\partial_a \bar{g}^{cd} \partial_b \bar{g}_{cd} = -\bar{g}^{ce} \bar{g}^{df} \partial_a \bar{g}_{ef} \partial_b \bar{g}_{cd}$$

The right hand side is clearly symmetric in ab and thus the left hand must also be symmetric in ab .

PhysRevD.62.044034 equation (20)

```

1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-eqtns-20.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  DGiBarDt := \partial_{t}{GammaBar^{i}}. # cdb(eq20.00,DGiBarDt)
10 DGiBarDt := - 2 ABar^{i j} \partial_{j}{N}
11           + 2 N ( GammaBar^{i}_{j k} ABar^{k j}
12                 - (2/3) gBar^{i j} \partial_{j}{trK}
13                 + 6 ABar^{i j} \partial_{j}{\phi}). # cdb(eq20.01,DGiBarDt)
14
15 substitute (DGiBarDt,defGammaBarU) # cdb(eq20.02,DGiBarDt)
16
17 distribute (DGiBarDt)
18 DGiBarDt = product_sort (DGiBarDt) # cdb(eq20.03,DGiBarDt)
19
20 canonicalise (DGiBarDt) # cdb(eq20.04,DGiBarDt)
21 # cdb(eq20.99,DGiBarDt)
22
23 cdblib.put ('DGiBarDt',DGiBarDt,jsonfile)

```

$$\partial_t \bar{\Gamma}^i = -2\bar{A}^{ij} \partial_j N + 2N \left(\bar{\Gamma}^i_{jk} \bar{A}^{kj} - \frac{2}{3} \bar{g}^{ij} \partial_j \text{tr} K + 6\bar{A}^{ij} \partial_j \phi \right) \quad (\text{eq20.01})$$

$$= -2\bar{A}^{ij} \partial_j N + 2N \left(\frac{1}{2} \bar{g}^{ie} (\partial_j \bar{g}_{ek} + \partial_k \bar{g}_{je} - \partial_e \bar{g}_{jk}) \bar{A}^{kj} - \frac{2}{3} \bar{g}^{ij} \partial_j \text{tr} K + 6\bar{A}^{ij} \partial_j \phi \right) \quad (\text{eq20.02})$$

$$= -2\bar{A}^{ia} \partial_a N + N \bar{A}^{ab} \bar{g}^{ic} \partial_b \bar{g}_{ca} + N \bar{A}^{ab} \bar{g}^{ic} \partial_a \bar{g}_{bc} - N \bar{A}^{ab} \bar{g}^{ic} \partial_c \bar{g}_{ba} - \frac{4}{3} N \bar{g}^{ia} \partial_a \text{tr} K + 12N \bar{A}^{ia} \partial_a \phi \quad (\text{eq20.03})$$

Lapse function

This is about as easy it gets – choose a static lapse (which is fine for the Kasner spacetime).

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-lapse.json'
5  cdblib.create (jsonfile)
6
7  # -----
8
9  # DNDt := -2 N trK.  # 1+log
10 # DNDt := -N*N trK.  # Harmonic
11 DNDt := 0.           # Static lapse
12
13 cdblib.put ('DNDt',DNDt,jsonfile)
```

Ricci scalar

Here we compute the Ricci scalar R in terms of the BSSN data.

Note that this expression for R will only be used when evaluating the constraints. It will *not* be used in the evolution equations so the advice that the evolved $\bar{\Gamma}^i$ should be expressed in terms of \bar{g}_{ij} does not apply here.

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-ricci-scalar.json'
5  cdblib.create (jsonfile)
6
7  defRab = cdblib.get ('defRab','bssn-eqtns-14.json')
8
9  # -----
10
11  defG2GBarU := g^{a b} -> \exp(-4\phi) gBar^{a b}.
12
13  Rscalar := R. # cdb(Rscalar.00,Rscalar)
14  Rscalar := g^{a b} R_{a b}. # cdb(Rscalar.01,Rscalar)
15
16  substitute (Rscalar, defRab) # cdb(Rscalar.02,Rscalar)
17  substitute (Rscalar, defG2GBarU) # cdb(Rscalar.03,Rscalar)
18  distribute (Rscalar) # cdb(Rscalar.04,Rscalar)
19
20  Rscalar = product_sort (Rscalar) # cdb(Rscalar.05,Rscalar)
21
22  rename_dummies (Rscalar) # cdb(Rscalar.06,Rscalar)
23  canonicalise (Rscalar) # cdb(Rscalar.07,Rscalar)
24
25  foo := gBar^{b c} \partial_{a}{gBar_{b c}} -> 0. # follows from det(g) = 1
26
27  substitute (Rscalar, foo) # cdb(Rscalar.08,Rscalar)
28
29  foo := gBar_{a b} gBar^{a b} -> 3.
30  bah := gBar_{a b} gBar^{a c} -> gBar_{b}^{c}.
31  moo := gBar^{c d} gBar^{e f} \partial_{a}{gBar_{c e}} -> - \partial_{a}{gBar^{d f}}.
```



```

32
33 substitute (Rscalar, foo)           # cdb(Rscalar.09,Rscalar)
34 substitute (Rscalar, bah)           # cdb(Rscalar.10,Rscalar)
35 substitute (Rscalar, moo)           # cdb(Rscalar.11,Rscalar)
36 eliminate_kronecker (Rscalar)       # cdb(Rscalar.12,Rscalar)
37 rename_dummies (Rscalar)            # cdb(Rscalar.13,Rscalar)
38 canonicalise (Rscalar)              # cdb(Rscalar.14,Rscalar)
39
40 foo := gBar^{a b} gBar^{c d} \partial_{c}{gBar_{b d}} -> - \partial_{c}{gBar^{a c}}.
41 bah := \partial_{b}{gBar^{a b}} -> - GammaBar^{a}. # prd62.eqn17
42
43 substitute (Rscalar, foo)           # cdb(Rscalar.15,Rscalar)
44 substitute (Rscalar, bah)           # cdb(Rscalar.16,Rscalar)
45
46 Rscalar = product_sort (Rscalar)     # cdb(Rscalar.17,Rscalar)
47
48 rename_dummies (Rscalar)            # cdb(Rscalar.18,Rscalar)
49 canonicalise (Rscalar)              # cdb(Rscalar.19,Rscalar)
50
51 foo := gBar^{a b} gBar^{c d} \partial_{a b}{gBar_{c d}} ->
52     - gBar^{a b} \partial_{a}{gBar_{c d}} \partial_{b}{gBar^{c d}}. # follows from det(g) = 1
53
54 substitute (Rscalar, foo)           # cdb(Rscalar.20,Rscalar)
55 factor_out (Rscalar, $\exp(-4\phi)$) # cdb(Rscalar.21,Rscalar)
56
57 cdblib.put ('Rscalar',Rscalar,jsonfile)

```

$$R = g^{ab} R_{ab} \quad (\text{Rscalar.01})$$

$$= g^{ab} \left(-2\partial_{ab}\phi + \bar{g}^{cd}\partial_d\phi\partial_a\bar{g}_{bc} + \bar{g}^{cd}\partial_d\phi\partial_b\bar{g}_{ac} - \bar{g}^{cd}\partial_d\phi\partial_c\bar{g}_{ab} + 12\partial_a\phi\partial_b\phi - 2\bar{g}_{ab}\bar{g}^{cd}\partial_{cd}\phi - 2\bar{g}_{ab}\partial_d\phi\partial_c\bar{g}^{cd} - 12\bar{g}_{ab}\bar{g}^{cd}\partial_c\phi\partial_d\phi - \frac{1}{2}\bar{g}^{cd}\partial_{cd}\bar{g}_{ab} + \frac{1}{2}\bar{g}_{ac}\partial_b\bar{\Gamma}^c \right. \\ \left. + \frac{1}{2}\bar{g}_{bc}\partial_a\bar{\Gamma}^c - \frac{1}{2}\partial_c\bar{g}_{ab}\partial_d\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{bc}\partial_a\bar{g}^{dc} - \frac{1}{2}\bar{g}^{ed}\bar{g}^{cf}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{2}\bar{g}^{ef}\bar{g}^{cd}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{4}\partial_b\bar{g}_{cd}\partial_a\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{ac}\partial_b\bar{g}^{dc} \right) \quad (\text{Rscalar.02})$$

$$= \exp(-4\phi) \bar{g}^{ab} \left(-2\partial_{ab}\phi + \bar{g}^{cd}\partial_d\phi\partial_a\bar{g}_{bc} + \bar{g}^{cd}\partial_d\phi\partial_b\bar{g}_{ac} - \bar{g}^{cd}\partial_d\phi\partial_c\bar{g}_{ab} + 12\partial_a\phi\partial_b\phi - 2\bar{g}_{ab}\bar{g}^{cd}\partial_{cd}\phi - 2\bar{g}_{ab}\partial_d\phi\partial_c\bar{g}^{cd} - 12\bar{g}_{ab}\bar{g}^{cd}\partial_c\phi\partial_d\phi - \frac{1}{2}\bar{g}^{cd}\partial_{cd}\bar{g}_{ab} \right. \\ \left. + \frac{1}{2}\bar{g}_{ac}\partial_b\bar{\Gamma}^c + \frac{1}{2}\bar{g}_{bc}\partial_a\bar{\Gamma}^c - \frac{1}{2}\partial_c\bar{g}_{ab}\partial_d\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{bc}\partial_a\bar{g}^{dc} - \frac{1}{2}\bar{g}^{ed}\bar{g}^{cf}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{2}\bar{g}^{ef}\bar{g}^{cd}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{4}\partial_b\bar{g}_{cd}\partial_a\bar{g}^{cd} - \frac{1}{2}\partial_d\bar{g}_{ac}\partial_b\bar{g}^{dc} \right) \quad (\text{Rscalar.03})$$

$$= -2\exp(-4\phi) \bar{g}^{ab}\partial_{ab}\phi + \exp(-4\phi) \bar{g}^{ab}\bar{g}^{cd}\partial_d\phi\partial_a\bar{g}_{bc} + \exp(-4\phi) \bar{g}^{ab}\bar{g}^{cd}\partial_d\phi\partial_b\bar{g}_{ac} - \exp(-4\phi) \bar{g}^{ab}\bar{g}^{cd}\partial_d\phi\partial_c\bar{g}_{ab} + 12\exp(-4\phi) \bar{g}^{ab}\partial_a\phi\partial_b\phi \\ - 2\exp(-4\phi) \bar{g}^{ab}\bar{g}_{ab}\bar{g}^{cd}\partial_{cd}\phi - 2\exp(-4\phi) \bar{g}^{ab}\bar{g}_{ab}\partial_d\phi\partial_c\bar{g}^{cd} - 12\exp(-4\phi) \bar{g}^{ab}\bar{g}_{ab}\bar{g}^{cd}\partial_c\phi\partial_d\phi - \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\bar{g}^{cd}\partial_{cd}\bar{g}_{ab} \\ + \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\bar{g}_{ac}\partial_b\bar{\Gamma}^c + \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\bar{g}_{bc}\partial_a\bar{\Gamma}^c - \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\partial_c\bar{g}_{ab}\partial_d\bar{g}^{cd} - \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\partial_d\bar{g}_{bc}\partial_a\bar{g}^{dc} \\ - \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\bar{g}^{ed}\bar{g}^{cf}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\bar{g}^{ef}\bar{g}^{cd}\partial_e\bar{g}_{ac}\partial_f\bar{g}_{bd} + \frac{1}{4}\exp(-4\phi) \bar{g}^{ab}\partial_b\bar{g}_{cd}\partial_a\bar{g}^{cd} - \frac{1}{2}\exp(-4\phi) \bar{g}^{ab}\partial_d\bar{g}_{ac}\partial_b\bar{g}^{dc} \quad (\text{Rscalar.04})$$

$$= -2\bar{g}^{ab}\partial_{ab}\phi \exp(-4\phi) + \bar{g}^{ca}\bar{g}^{bd}\partial_d\phi \exp(-4\phi) \partial_c\bar{g}_{ab} + \bar{g}^{ac}\bar{g}^{bd}\partial_d\phi \exp(-4\phi) \partial_c\bar{g}_{ab} - \bar{g}^{ab}\bar{g}^{cd}\partial_d\phi \exp(-4\phi) \partial_c\bar{g}_{ab} + 12\bar{g}^{ab}\partial_a\phi\partial_b\phi \exp(-4\phi) \\ - 2\bar{g}_{ab}\bar{g}^{ab}\bar{g}^{cd}\partial_{cd}\phi \exp(-4\phi) - 2\bar{g}_{ab}\bar{g}^{ab}\partial_d\phi \exp(-4\phi) \partial_c\bar{g}^{cd} - 12\bar{g}_{ab}\bar{g}^{ab}\bar{g}^{cd}\partial_c\phi\partial_d\phi \exp(-4\phi) - \frac{1}{2}\bar{g}^{ab}\bar{g}^{cd} \exp(-4\phi) \partial_{cd}\bar{g}_{ab} + \frac{1}{2}\bar{g}_{ab}\bar{g}^{ac} \exp(-4\phi) \partial_c\bar{\Gamma}^b \\ + \frac{1}{2}\bar{g}_{ab}\bar{g}^{ca} \exp(-4\phi) \partial_c\bar{\Gamma}^b - \frac{1}{2}\bar{g}^{ab} \exp(-4\phi) \partial_c\bar{g}_{ab}\partial_d\bar{g}^{cd} - \frac{1}{2}\bar{g}^{ca} \exp(-4\phi) \partial_d\bar{g}_{ab}\partial_c\bar{g}^{db} - \frac{1}{2}\bar{g}^{ac}\bar{g}^{be}\bar{g}^{fd} \exp(-4\phi) \partial_f\bar{g}_{ab}\partial_e\bar{g}_{cd} \\ + \frac{1}{2}\bar{g}^{ac}\bar{g}^{bd}\bar{g}^{ef} \exp(-4\phi) \partial_e\bar{g}_{ab}\partial_f\bar{g}_{cd} + \frac{1}{4}\bar{g}^{cd} \exp(-4\phi) \partial_d\bar{g}_{ab}\partial_c\bar{g}^{ab} - \frac{1}{2}\bar{g}^{ac} \exp(-4\phi) \partial_d\bar{g}_{ab}\partial_c\bar{g}^{db} \quad (\text{Rscalar.05})$$

$$= -2\bar{g}^{ab}\partial_{ab}\phi \exp(-4\phi) + \bar{g}^{bc}\bar{g}^{da}\partial_a\phi \exp(-4\phi) \partial_b\bar{g}_{cd} + \bar{g}^{cb}\bar{g}^{da}\partial_a\phi \exp(-4\phi) \partial_b\bar{g}_{cd} - \bar{g}^{cd}\bar{g}^{ba}\partial_a\phi \exp(-4\phi) \partial_b\bar{g}_{cd} + 12\bar{g}^{ab}\partial_a\phi\partial_b\phi \exp(-4\phi) \\ - 2\bar{g}_{cd}\bar{g}^{cd}\bar{g}^{ab}\partial_{ab}\phi \exp(-4\phi) - 2\bar{g}_{cd}\bar{g}^{cd}\partial_a\phi \exp(-4\phi) \partial_b\bar{g}^{ba} - 12\bar{g}_{cd}\bar{g}^{cd}\bar{g}^{ab}\partial_a\phi\partial_b\phi \exp(-4\phi) - \frac{1}{2}\bar{g}^{cd}\bar{g}^{ab} \exp(-4\phi) \partial_{ab}\bar{g}_{cd} + \frac{1}{2}\bar{g}_{ca}\bar{g}^{cb} \exp(-4\phi) \partial_b\bar{\Gamma}^a \\ + \frac{1}{2}\bar{g}_{ca}\bar{g}^{bc} \exp(-4\phi) \partial_b\bar{\Gamma}^a - \frac{1}{2}\bar{g}^{cd} \exp(-4\phi) \partial_a\bar{g}_{cd}\partial_b\bar{g}^{ab} - \frac{1}{2}\bar{g}^{bc} \exp(-4\phi) \partial_a\bar{g}_{cd}\partial_b\bar{g}^{ad} - \frac{1}{2}\bar{g}^{cd}\bar{g}^{eb}\bar{g}^{af} \exp(-4\phi) \partial_a\bar{g}_{ce}\partial_b\bar{g}_{df} \\ + \frac{1}{2}\bar{g}^{cd}\bar{g}^{ef}\bar{g}^{ab} \exp(-4\phi) \partial_a\bar{g}_{ce}\partial_b\bar{g}_{df} + \frac{1}{4}\bar{g}^{ba} \exp(-4\phi) \partial_a\bar{g}_{cd}\partial_b\bar{g}^{cd} - \frac{1}{2}\bar{g}^{cb} \exp(-4\phi) \partial_a\bar{g}_{cd}\partial_b\bar{g}^{ad} \quad (\text{Rscalar.06})$$

[illegible]

PhysRevD.67.084023 equation (19)

```
1  from shared import *
2  import cdblib
3
4  jsonfile = 'bssn-constraints.json'
5  cdblib.create (jsonfile)
6
7  # -----
8  # Hamiltonian constraint
9
10 Ham := R + K_{a b} g^{a b} K_{c d} g^{c d} - K_{a b} K_{c d} g^{a c} g^{b d}.
11                                     # cdb(Ham.101,Ham)
12
13 Ham := R + (2/3) (trK)**2 - ABar_{a b} ABar^{a b}.          # cdb(Ham.102,Ham)
```

$$\mathcal{H} = R + K_{ab} g^{ab} K_{cd} g^{cd} - K_{ab} K_{cd} g^{ac} g^{bd} \quad (\text{Ham.101})$$

$$= R + \frac{2}{3} \text{tr} K^2 - \bar{A}_{ab} \bar{A}^{ab} \quad (\text{Ham.102})$$

PhysRevD.67.084023 equation (20)

```
1  # -----
2  # Momentum constraint
3
4  confMom := 6 ABar^{i a} \partial_{a}{\phi}
5           + \partial_{a}{ABar^{i a}}
6           + ABar^{a b} GammaBar^{i}_{a b}
7           - (2/3) gBar^{i a} \partial_{a}{trK}.
8
9  defGammaBar := GammaBar^{a}_{b c} ->
10              (1/2) gBar^{a e} ( \partial_{b}{gBar_{e c}}
11                                + \partial_{c}{gBar_{b e}}
12                                - \partial_{e}{gBar_{b c}}).
13
14  substitute (confMom, defGammaBar)           # cdb(confMom.101,confMom)
15  distribute (confMom)                       # cdb(confMom.102,confMom)
16
17  confMom = product_sort (confMom)            # cdb(confMom.103,confMom)
18
19  rename_dummies (confMom)                   # cdb(confMom.104,confMom)
20  canonicalise (confMom)                     # cdb(confMom.105,confMom)
21
22  foo := \partial_{a}{ABar^{i a}} -> \partial_{a}{gBar^{i c} gBar^{a d} ABar_{c d}}.
23
24  substitute (confMom, foo)                   # cdb(confMom.106,confMom)
25  product_rule (confMom)                     # cdb(confMom.107,confMom)
26
27  confMom = product_sort (confMom)            # cdb(confMom.108,confMom)
28
29  rename_dummies (confMom)                   # cdb(confMom.109,confMom)
30  canonicalise (confMom)                     # cdb(confMom.110,confMom)
31
32  cdblib.put ('Ham',Ham,jsonfile)
33  cdblib.put ('confMom',confMom,jsonfile)
```

$$\exp(4\phi)\mathcal{D}^j = 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{A}^{ia} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ie}(\partial_a\bar{g}_{eb} + \partial_b\bar{g}_{ae} - \partial_e\bar{g}_{ab}) - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.101})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{A}^{ia} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ie}\partial_a\bar{g}_{eb} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ie}\partial_b\bar{g}_{ae} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ie}\partial_e\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.102})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{A}^{ia} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{cb} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_b\bar{g}_{ac} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.103})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{A}^{ia} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{cb} + \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_b\bar{g}_{ac} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.104})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{A}^{ia} + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.105})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a(\bar{g}^{ic}\bar{g}^{ad}\bar{A}_{cd}) + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.106})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \partial_a\bar{g}^{ic}\bar{g}^{ad}\bar{A}_{cd} + \bar{g}^{ic}\partial_a\bar{g}^{ad}\bar{A}_{cd} + \bar{g}^{ic}\bar{g}^{ad}\partial_a\bar{A}_{cd} + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.107})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \bar{A}_{ab}\bar{g}^{cb}\partial_c\bar{g}^{ia} + \bar{A}_{ab}\bar{g}^{ia}\partial_c\bar{g}^{cb} + \bar{g}^{cb}\bar{g}^{ia}\partial_c\bar{A}_{ab} + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.108})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \bar{A}_{ab}\bar{g}^{cb}\partial_c\bar{g}^{ia} + \bar{A}_{ab}\bar{g}^{ia}\partial_c\bar{g}^{cb} + \bar{g}^{cb}\bar{g}^{ia}\partial_c\bar{A}_{ab} + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.109})$$

$$= 6\bar{A}^{ia}\partial_a\phi + \bar{A}_{ab}\bar{g}^{ac}\partial_c\bar{g}^{ib} + \bar{A}_{ab}\bar{g}^{ia}\partial_c\bar{g}^{bc} + \bar{g}^{ia}\bar{g}^{bc}\partial_b\bar{A}_{ac} + \bar{A}^{ab}\bar{g}^{ic}\partial_a\bar{g}_{bc} - \frac{1}{2}\bar{A}^{ab}\bar{g}^{ic}\partial_c\bar{g}_{ab} - \frac{2}{3}\bar{g}^{ia}\partial_a\text{tr}K \quad (\text{confMom.110})$$