# efficiently\_solfving\_stiff\_differential\_equations

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# 1 Efficiently Solving Stiff Differential Equations

#### 1.1 An Introduction at the Hand of the Brusselator Problem

In the following we will introduce some ideas around efficiently solving stiff Ordinary Differential Equations (ODEs) or rather systems of them. Let us make the sidenote that an equation is not per se stiff but a particular initial value problem for that equation may be stiff in some regions (C. W. Gear, 1982).

### 1.2 Imports

```
[]: # Mathematical tools
     using LinearAlgebra
     using DifferentialEquations
     using Symbolics
     using ForwardDiff
     using SparseDiffTools
     using LSODA
     using ODEInterfaceDiffEq
     using IterativeSolvers
     using LinearMaps
     # Graphs
     using Graphs
     using NetworkLayout
     # General tools
     using BenchmarkTools
     # Plotting
     using Colors, ColorSchemes
     using Plots
     # Plotting graphs
     using CairoMakie
     using GraphMakie
```

### 1.3 Example stiff ODE system | Brusselator

In this Notebook we consider a simplified "Brusselator" (as discussed in "Solving Ordinar Differential Equations I" by Hairer, Norsett and Wanner) (the Brusselator was oriinally introduced by Lefefer and Nicolis, 1971). Consider six substances A, B, D, E, X, Y undergoing chemical the following reactions

$$\begin{array}{cccc} A & \xrightarrow{k_1} & X \\ B+X & \xrightarrow{k_2} & Y+D & \text{(bimolecular reaction)} \\ 2X+Y & \xrightarrow{k_3} & 3X & \text{(autocatalytic trimol. reaction)} \\ X & \xrightarrow{k_4} & E \end{array}$$

where A(x), B(x), C(x), ... denote concentrations with respect to the place x. As of the mass action law (which you may know as the "Massenwirkungsgesetz" from school) (the rate of a chemical reaction is proportional to the product of the activities or concentrations of the reactants) one follows the following system of differential equations

$$\begin{split} A' &= -k_1 A \\ B' &= -k_2 B X \\ D' &= k_2 B X \\ E' &= k_4 X \\ X' &= k_1 A - k_2 B X + k_3 X^2 Y - k_4 X \\ Y' &= k_2 B X - k_3 X^2 Y \end{split}$$

where for simplicity, we \* leave out the equations for D and E as they do not influence the rest \* assume A and B to be constant (maintained) concentrations \* assume all reaction rates to be set to 1 \* assume a one-dimensiona setting with  $0 \le x \le 1$  and constant boundary conditions \* set u(x) := X(x) and v(x) := Y(x) \* assume diffusion terms term  $\alpha \frac{\partial^2 u}{\partial x^2}$  and  $\alpha \frac{\partial^2 v}{\partial x^2}$  and get

$$\begin{split} \frac{\partial u}{\partial t} &= A + u^2 v - (B+1)u + \alpha \frac{\partial^2 u}{\partial x^2} \\ \frac{\partial v}{\partial t} &= Bu - u^2 v + \alpha \frac{\partial^2 v}{\partial x^2} \end{split}$$

where from dicretizing the differentiation in space (i. e. using method of lines for approaching this partial differential equations) we follow (with  $x_i = \frac{i}{N+1} (1 \le i \le N), \Delta x = \frac{1}{N+1}, A = 1, B = 3, \alpha = \frac{1}{50}$ )

$$\begin{split} u_i' &= 1 + u_i^2 v_i - 4 u_i + \frac{\alpha}{(\Delta x)^2} \left( u_{i-1} - 2 u_i + u_{i+1} \right), \\ v_i' &= 3 u_i - u_i^2 v_i + \frac{\alpha}{(\Delta x)^2} \left( v_{i-1} - 2 v_i + v_{i+1} \right) \\ u_0(t) &= u_{N+1}(t) = 1, \quad v_0(t) = v_{N+1}(t) = 3 \\ u_i(0) &= 1 + \sin\left( 2\pi x_i \right), \quad v_i(0) = 3, \quad i = 1, \dots, N. \end{split}$$

where some boundary conditions have been chosen (constant boundary values enforced using ghost cells). A less simplified Brusselator occurs for instance in nature in the form of the "Oregonator" describing the reactions between  $HBrO_2$ ,  $Br^-$  and Ce(IV).

#### 1.3.1 Implementation of the Brusselator

```
[]: function brusselator!(dwdt, w, N, t)
        This function returns the right hand side of
        the ODE system of the Brusselator problem as given
        in "Solving Ordinary Differential Equations II" by
        Hairer and Wanner.
        # Extract u and v from w, a 1D array of length 2(N + 2)
        u = w[1:N+2]; v = w[N+3:end]
        # Brusselator parameter
        alpha = 1/50
        # Cell size
        h = 1/(N+1)
        # Right hand side of the ODE system, constant boundary conditions
        offset = N + 2
        \# dudt[1] = 0; dudt[end] = 0; dvdt[1] = 0; dvdt[end] = 0 not needed
        h^2 .* (u[1:N] .- 2 .* u[2:N+1] .+ u[3:N+2])
        dwdt[2 + offset:N+1 + offset] .= 3 .* u[2:N+1] .- u[2:N+1].^2 .* v[2:N+1] .
     \rightarrow+ alpha ./ h^2 .* (v[1:N] .- 2 .* v[2:N+1] .+ v[3:N+2])
        dwdt[1] = 0; dwdt[end] = 0; dwdt[1 + offset] = 0; dwdt[offset] = 0
    end
```

brusselator! (generic function with 1 method)

#### 1.3.2 Initialization of the Brusselator

```
[]: # Init brusselator
function init_brusselator(N)

"""

This function returns the initial values for the
Brusselator problem.

"""

x = range(0, 1, length=N+2)

u0 = 1 .+ sin.(2 .* pi .* x)
v0 = zeros(N+2) .+ 3
u0[1] = 1; u0[end] = 1; v0[1] = 3

w = vcat(u0, v0)

return x, w
```

```
end
```

init\_brusselator (generic function with 1 method)

### 1.3.3 Helpers for the Brusselator system

```
function sol_to_u_v(sol, N)

"""

This function returns the solution of the Brusselator
problem as a tuple of u and v.

"""

# sol is sadly Vector{Vector{Float64}}

u = [sol.u[i][j] for i = 1:length(sol.u), j = 1:N+2]
v = [sol.u[i][j] for i = 1:length(sol.u), j = N+3:length(sol.u[1])]
return u, v
end
```

sol\_to\_u\_v (generic function with 1 method)

```
[]: function wbrusselator!(dwdt, w)
    """
    Wrapped brusselator function that takes a vector w as input.
    N = (length(w) - 2) / 2
    """
    N = Int32(length(w) / 2 - 2)
    brusselator!(dwdt, w, N, 0)
end
```

wbrusselator! (generic function with 1 method)

```
[]: function init_ode_prob(N, t_0, t_end, num_saved_steps)
    """

    Returns an ODEProblem for the Brusselator problem.
    """

    tc = range(t_0, t_end, length = num_saved_steps) # saved time steps
    x, w = init_brusselator(N) # initial conditions
    return ODEProblem(brusselator!, w, (t_0, t_end), N), tc, x, w
end
```

init\_ode\_prob (generic function with 1 method)

### 1.4 Sparse matrix and graph functionality

```
[]: function sparsity_matrix_to_graph(sp_mat)
    """
    Each column of the sparsity matrix is a node in the graph.
    Two nodes are connected if the corresponding columns overlap, i. e.
    if there is a non-zero entry in the same row of both columns.
    """
```

```
# Get number of columns
    n = size(sp_mat, 2)
    # Create empty graph
    g = SimpleGraph(n)
    # Iterate over columns
    for i in 1:n
        # Get indices of non-zero entries in column i
        nz = findall(!iszero, sp_mat[:, i])
        for j in i + 1:n
            # Get indices of non-zero entries in column j
            nz2 = findall(!iszero, sp_mat[:, j])
            # Check if there is a non-zero entry in the same row
            if length(intersect(nz, nz2)) > 0
                # Add edge between column i and j
                add_edge!(g, i, j)
            end
        end
    end
    return g
end
function condense_sparsity_pattern(mat, coloring)
    This function condenses the sparsity pattern of a matrix.
    # number of colors
   n_colors = coloring.num_colors
    # condenset sparsity pattern
    cm = zeros(size(mat, 1), n_colors)
    for i in 1:n_colors
        cm[:, i] = sum(mat[:, coloring.colors .== i], dims=2)
    end
    return cm
end
```

condense\_sparsity\_pattern (generic function with 1 method)

### 1.5 Plot functionality

#### 1.5.1 Plotting the solution of the ODE system

```
[]: function plot_as_2D_surface(u, t, x, method, N)
         This function plots the solution of the Brusselator
         problem as a 2D surface.
         0.00
         # Plot
         f = Plots.surface(x, t, u, xlabel="spatial coordinate x", ylabel="time t", u
      ⇒zlabel="concentration u \n of substance X", color=:coolwarm, title = ∪
      _{\circlearrowleft}"Simplified Brusselator approached \\n via $method \\n number of gridpoints \\N_{\sqcup}
      \Rightarrow= $N", legend = :none)
         return f
     end
     function anim_plot_1D(u, N)
         This function animates the solution of the Brusselator
         problem as a 1D line plot.
         0.00
         x = range(0, stop=1, length=N+2)
         anim = @animate for i in 1:100
             Plots.plot(x, u[i, :], xlims=(0, 1), ylims=(0, 4), xlabel="x",\Box
      ⇔ylabel="u", title="Brusselator", legend=false, grid=true)
         end
         gif(anim, "figures/bruss_ani.gif", fps=10)
     end
```

anim\_plot\_1D (generic function with 1 method)

#### 1.5.2 Plotting sparse matrices

```
function scatter_matrix_plot(mat, cm, title)
    """
    Plot the sparsity pattern of a matrix as a scatter plot.
    """
    x = [i for i in 1:size(mat, 2) for j in 1:size(mat, 1)]
    y = [j for i in 1:size(mat, 2) for j in 1:size(mat, 1)]
    c = [cm[i, j] for j in 1:size(mat, 2) for i in 1:size(mat, 1)]
```

```
CairoMakie.scatter(x, y, color = c, marker = :rect, markersize = 50, axis =_{\sqcup}

→(aspect = DataAspect(), xgridvisible = false, ygridvisible = false, xticks = 

□
 41:size(mat, 2), yticks = 1:size(mat, 1), xlabel = "column-index", ylabel =

¬"row-index", title = title))
    # invert y axis
    CairoMakie.ylims!(size(mat, 1) + 0.8, 0)
    CairoMakie.xlims!(0, size(mat, 2) + 1)
    # show plot
    return CairoMakie.current_figure()
end
function plot_sparsity_pattern_black(mat, title = "Sparsity pattern")
    This function plots the sparsity pattern of a matrix.
    cm = zeros(size(mat, 1), size(mat, 2))
    cm[mat .== 0] .= 0.9
    cm[mat .== 1] .= 0
    cm = RGB.(cm)
    f = scatter_matrix_plot(mat, cm, title);
    display(f)
    return f
end
function plot sparsity pattern(mat, colors, title = "Sparsity Pattern")
    This function plots the sparsity pattern of a matrix.
    11 11 11
    cm = RGB.(mat)
    cm[mat .== 0] .= 0.9
    # loop through columns and assign color to all non-zero entries
    for i in 1:size(mat, 2)
        cm[mat[:, i] .!= 0, i] .= colors[i]
    end
    cm = RGB.(cm)
    RGB.(cm)
    f = scatter_matrix_plot(mat, cm, title);
    display(f)
    return f
end
```

plot\_sparsity\_pattern (generic function with 2 methods)

#### 1.5.3 Graph plotting

```
[]: # Use CairoMakie with SVG output
     CairoMakie.activate!(type = "svg")
     function plot_graph(g, coloring)
         This function plots the graph of the sparsity pattern, colored according to \Box
      →the coloring.
         0.00
         layout = Stress(Ptype=Float32)
         f, ax, p = graphplot(g, layout = layout, node_color = node_colors(coloring,_
      →g), nlabels = repr.(1:nv(g)), nlabels_distance = 0, nlabels_fontsize=25, u
      elabels_color=:grey, node_size = 60, nlabels_align=(:center,:center))
         hidedecorations!(ax); hidespines!(ax)
         # update_limits!(ax)
         # show the plot
         display(f)
         # save as svq
         save("figures/bruss_graph_colored.svg", f)
     end
     function plot_graph(g)
         This function plots the graph of the sparsity pattern.
         layout = Stress(Ptype=Float32)
         f, ax, p = graphplot(g, layout = layout, node_color = :black, nlabels = __
      repr.(1:nv(g)), nlabels_distance = 0, nlabels_fontsize=25, nlabels_color = :
      white, elabels_color=:grey, node_size = 60, nlabels_align=(:center,:center))
         # add a title, title!("Graph of the sparsity pattern") does not work
         # text!(ax, "Graph representation of \n Jacobian sparsity matrix", __
      \Rightarrowposition=(-1.6, 1), fontsize = 30)
         hidedecorations!(ax); hidespines!(ax)
         # update_limits!(ax)
         # show the plot
         display(f)
         # save as svq
         save("figures/bruss_graph.svg", f)
     end
     # Fix for GraphMakie labels being displayed outside of the plot
```

```
# taken from https://discourse.julialang.org/t/
 \neg textbox-qraph-with-networklayout-with-labels-cut-off-usinq-qraphmakie/96566/4
# by @hexaeder
function update limits!(ax)
    p = only(ax.scene.plots)
    @assert p isa GraphPlot
    nodep = get_node_plot(p)
    textp = get nlabel plot(p)
    to_corners = r -> (Point2f(r.origin[1:2]), Point2f(r.origin[1:2] + r.
 \hookrightarrow widths[1:2]))
    to_px = pts -> Makie.project(ax.scene, pts)
    # get lower left and upper right node bounding box in data and px space
    n1_dat, n2_dat = to_corners(boundingbox(nodep))
    n1, n2
                  = to_px(n1_dat), to_px(n2_dat)
    # get text bounding box in pixel sapce
    t1, t2 = to_corners(boundingbox(textp))
    # get axis limits in data and pixel space
    ax1_dat, ax2_dat = to_corners(ax.finallimits[])
    ax1, ax2
                    = to_px(ax1_dat), to_px(ax2_dat)
    # check difference between nodelabel and node bounding box in px
    px_diff1 = map(x->min(0, x), t1-n1)
    px diff2 = map(x->max(0, x), t2-n2)
    # define "targets" in pixelspace where n1 and n2 bounding should end up
    n1_target = ax1-px_diff1 + Point2f(20, 20)
    n2_target = ax2-px_diff2 - Point2f(20, 20)
    # calculate axis limits such as bounding box of node plot goes to target
    scale = (n2 dat - n1 dat) ./ (n2 target - n1 target)
    origin = n1_dat - scale .* (n1_target - ax1)
    width = ((n2_dat - n1_dat) ./ (n2_target - n1_target)) .* (ax2 - ax1)
    CairoMakie.xlims!(ax, origin[1], origin[1]+width[1] )
    CairoMakie.ylims!(ax, origin[2], origin[2]+width[2] )
end
```

update\_limits! (generic function with 1 method)

#### 1.5.4 Stepping on a 2D heatmap

plot\_heat\_stepper (generic function with 1 method)

### 1.6 Further helpers

```
[]: # Plot g with coloring and column indices as labels
function node_colors(coloring, g)
    return [ColorSchemes.tol_light.colors[coloring.colors[i]] for i in 1:nv(g)]
end

function unique_colors(coloring)
    return [ColorSchemes.tol_light.colors[i] for i in 1:coloring.num_colors]
end
```

unique\_colors (generic function with 1 method)

#### 1.7 Trying Explicit Euler on the Brusselator

Let us try explicit Euler on the Brusselator.

```
[]: # Explicit Euler method
function euler(f, u0, dt, num_steps)
    """
    Explicit Euler method for solving an ODE.
    """
    # Initialize array to store solution
    u = zeros(length(u0), num_steps)
    u[:, 1] .= u0
    # Initialize derivative
    du = zeros(size(u0))
```

```
for i in 2:num_steps
        # Compute derivative
        f(du, u[:, i - 1])
        # Update solution
        u[:, i] = u[:, i - 1] + dt * du
    end
    return u
end
# Explicit Euler method
function fast_euler(f!, u0, dt, num_steps)
    Explicit Euler method for solving an ODE.
    # Initialize array to store solution
    u = u0
    # Initialize derivative
    du = similar(u0)
    for i in 2:num_steps
        # Compute derivative
        f!(du, u)
        # Update solution
        u += dt * du
    end
    return u
end
function init_explicit_problem(N, t_0, t_end, dt)
    Initialize the explicit problem.
    # Initial conditions
   x, w = init_brusselator(N)
    num_steps = Int(round((t_end - t_0) / dt))
    t = range(t_0, t_end, length = num_steps)
    return x, w, t, num_steps
end
```

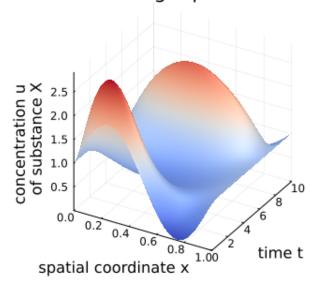
init\_explicit\_problem (generic function with 1 method)

```
[]: N = 40; dt = 0.01; t_0 = 0.0; t_end = 10.0;
x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)

sol = euler(wbrusselator!, w, dt, num_steps)
u_own = sol[1:N+2, :]';
```

```
[]: f = plot_as_2D_surface(u_own, t, x, "Explicit Euler with dt = $dt", N);
save("figures/bruss_ex_eu40.svg", f)
display(f)
```

# Simplified Brusselator approached via Explicit Euler with dt = 0.01 number of gridpoints N = 40



### 1.7.1 The explicit method has to take excessively small timesteps for large N $\mid$ stiffness

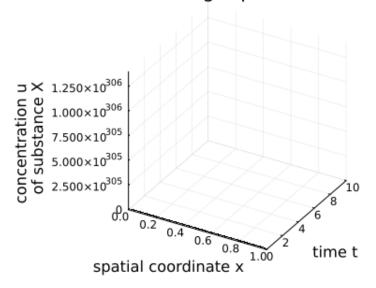
This does not look too bad, right? But what if we change N to say 400?

```
[]: N = 400; dt = 0.01; t_0 = 0.0; t_end = 10.0;
x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)

sol = @btime euler(wbrusselator!, w, dt, num_steps)

u_own = sol[1:N+2, :]';
f = plot_as_2D_surface(u_own, t, x, "Explicit Euler with dt = $dt", N)
save("figures/bruss_ex_eu400A.svg", f)
display(f)
```

# Simplified Brusselator approached via Explicit Euler with dt = 0.01 number of gridpoints N = 400



#### 10.579 ms (17985 allocations: 75.77 MiB)

Ok, that is not what we want. It turns out that we can avoid this instability by lowering the size of the timestep - but for say N=4000 we quickly run into issues. Let us make the previous a bit more concise using a physical argument:

The transport process at play is diffusion, where for a diffusive process we know for the spreading of some concentration to follow  $\sigma = \sqrt{2\alpha t}$ . Now in each Euler step we do, only neighboring cells have an effect on each other (compare the discretized ODE we introduced in the beginning). Therefore - in the style of a Courant-Friedrichs-Levy criterium - we can propose stability the constraint

$$\Delta x > \sigma(\Delta t) = \sqrt{2\alpha t}$$

so  $\Delta t < \frac{\Delta x^2}{2\alpha}$ . If we want to double N (cut in half  $\Delta x$ ) we need  $\mathcal{O}(N^2)$  more timesteps with the complexity of a function evaluation scaling with  $\mathcal{O}(N)$  resulting in a  $\mathcal{O}(N^3)$  scaling - calculations quickly become unfeasible. As a task to the reader, you can think about implications in higher dimensions.

```
[]: N = 400; dt = 0.8 * (1 / (N + 1)) ^ 2 / (2 * 1 / 50); t_0 = 0.0; t_end = 10.0;
print(dt)
x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)

sol = @btime euler(wbrusselator!, w, dt, num_steps)
end_state = @btime fast_euler(wbrusselator!, w, dt, num_steps)
```

```
display(norm(end_state .- sol[:, end]))
u_own = sol[1:N+2, :]';

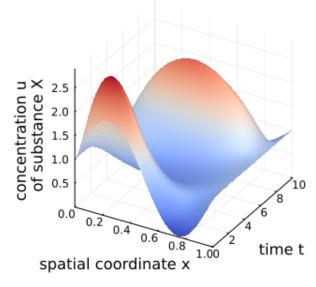
0.00012437733596184103  2.417 s (1447185 allocations: 5.95 GiB)

0.0

1.992 s (1286385 allocations: 4.49 GiB)

[]: dtr = round(dt, digits = 4)
f = plot_as_2D_surface(u_own, t, x, "Explicit Euler with dt = $dtr", N)
save("figures/bruss_ex_eu400B.svg", f)
display(f)
```

# Simplified Brusselator approached via Explicit Euler with dt = 0.0001 number of gridpoints N = 400



#### 1.7.2 Checking the normal stiffness indicators I: Eigenvalues of the Jacobian

The result is correct but has taken us lots of time and function evaluations. Let us also analyze the initial Jacobian.

```
[]: # ======== Analysis of the jacobian at initial conditions ============
function min_max_abs_ev(N)
```

```
Analysis of the jacobian at initial conditions for N = 40 min(abs(re(eigenvalues > 0))): 0.31408801361169625 max(abs(re(eigenvalues))): 136.9824405012886 Analysis of the jacobian at initial conditions for N = 400 min(abs(re(eigenvalues > 0))): 0.321033307345784 max(abs(re(eigenvalues))): 12866.578819965474
```

# 1.7.3 Checking the normal stiffness indicators II: Coefficients of vastly different magnitude in the ODE system

The larger we choose N the larger the stiffness ratio. We can also see this in the frame of the classical "factors of largely different magnitude" indicator for stiffness. Consider the ODE system we solve

$$\begin{split} u_i' &= 1 + u_i^2 v_i - 4u_i + \frac{\alpha}{(\Delta x)^2} \left( u_{i-1} - 2u_i + u_{i+1} \right), \\ v_i' &= 3u_i - u_i^2 v_i + \frac{\alpha}{(\Delta x)^2} \left( v_{i-1} - 2v_i + v_{i+1} \right) \\ u_0(t) &= u_{N+1}(t) = 1, \quad v_0(t) = v_{N+1}(t) = 3 \\ u_i(0) &= 1 + \sin\left( 2\pi x_i \right), \quad v_i(0) = 3, \quad i = 1, \dots, N. \end{split}$$

where  $\Delta x = \frac{1}{N+1}$  - so there is a factor decreasing quadratically with N in our equation. In addition the system becomes larger and larger with increasing N increasing the computational load.

#### 1.7.4 Checking the normal stiffness indicators III: Quickly varying solutions nearby

```
[]:  # to do
```

#### 1.8 Our Own Simple Implementation of Implicit Euler

We have already seen, that implicit Euler is pretty much unconditionally stable - even for large timesteps. Let us implement implicit Euler in the following and see if we can increase the overall efficiency regarding solving the stiff ODE.

#### 1.8.1 A first approach without a gain in efficiency

```
[]: # Mind that the brusselator is defined in place, which could be used to our
     \hookrightarrow computational
    # advantage but is rather bypassed here.
    function g_zero(f!, u, y_n, dt)
        Finding the root of this function is equivalent
        to finding the next time step in the implicit Euler
        scheme.
        0.00
        dudt = zeros(length(u))
        f!(dudt, u)
        return u - y_n - dt * dudt
    end
    function Jg(f!, u, dt)
        Jacobian of g_zero
        0.00
        N = Int32(length(u) / 2 - 2)
        dwdt0 = zeros(2(N + 2))
        J = ForwardDiff.jacobian(f!, dwdt0, u)
        return I - dt * J
    end
    function newtons_method(f!, u0, dt, tol, max_iter)
        The root of g_zero (and therefore the next step in time in
        the implicit Euler scheme) is found using Newton's method,
        which is implemented here.
        0.00
        u = copy(u0)
        for i in 1:max_iter
            # Here the LU decomposition is performed
            # the full Jacobian is constructed in every step
            luJ = lu(Jg(f!, u, dt)) # Jg(f!, u, dt) constructs
                                                   the Jacobian
            a = luJ \setminus g_zero(f!, u, u0, dt)
            u \cdot = u \cdot - a
            if norm(g_zero(f!, u, u0, dt)) < tol</pre>
                return u
```

```
end
end
return u
end

function implicit_euler(f!, g, zero_finder, dt, w, num_steps)

"""

A stiff differential equation can be
solved using the implicit Euler scheme.

"""

N = Int32(length(w) / 2 - 2)
u = zeros(2(N + 2), num_steps)
u[:, 1] = w
for i in 2:num_steps
u[:, i] = zero_finder(f!, u[:, i - 1], dt, 1e-4, 1000)
end
return u
end
```

implicit\_euler (generic function with 1 method)

```
Test of our first implementation
```

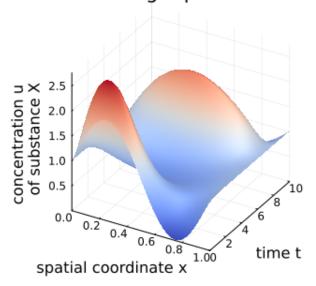
```
[]: N = 400; dt = 0.1; x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)

# Our solution
sol = Obtime implicit_euler(wbrusselator!, g_zero, newtons_method, dt, w,u_num_steps)
u_own = sol[1:N+2, :]';

8.218 s (396055 allocations: 11.48 GiB)
```

```
[]: plot_as_2D_surface(u_own, t, x, "Implicit Euler with dt = $dt", N)
```

# Simplified Brusselator approached via Implicit Euler with dt = 0.1 number of gridpoints N = 400



```
[]: sol = @btime implicit_euler(wbrusselator!, g_zero, newtons_method, dt, w,u onum_steps)
```

6.737 s (396055 allocations: 11.48 GiB)

#### 804×100 Matrix{Float64}:

1.0	1.0	1.0	1.0	1.0	•••	1.0	1.0	1.0
1.01567	1.01982	1.02416	1.02751	1.02926		0.994123	0.994099	0.994102
1.03133	1.03963	1.04832	1.05502	1.05852		0.988248	0.9882	0.988205
1.04699	1.05946	1.07249	1.08252	1.08776		0.982375	0.982304	0.982311
1.06263	1.07929	1.09667	1.11003	1.11699		0.976508	0.976412	0.976422
1.07826	1.09912	1.12085	1.13752	1.1462		0.970646	0.970526	0.970538
1.09387	1.11896	1.14504	1.16501	1.17538		0.964792	0.964648	0.964661
1.10946	1.13882	1.16924	1.19249	1.20452		0.958948	0.958779	0.958793
1.12502	1.15867	1.19344	1.21996	1.23361		0.953114	0.95292	0.952935
1.14055	1.17854	1.21765	1.2474	1.26265		0.947292	0.947072	0.947088
3.0	3.01803	3.02994	3.0388	3.04593		3.0514	3.05522	3.05875
3.0	3.01587	3.02632	3.03408	3.04031		3.04501	3.04835	3.05144
3.0	3.01368	3.02265	3.0293	3.03464		3.03861	3.04147	3.04411
3.0	3.01145	3.01894	3.02449	3.02893		3.03219	3.03457	3.03678
3.0	3.0092	3.0152	3.01963	3.02319		3.02576	3.02767	3.02943
3.0	3.00692	3.01142	3.01475	3.01742	•••	3.01933	3.02076	3.02208
3.0	3.00462	3.00763	3.00985	3.01162		3.01289	3.01384	3.01472

```
3.0 3.00231 3.00382 3.00493 3.00582 3.00645 3.00692 3.00736
3.0 3.0 3.0 3.0 3.0 3.0 3.0 3.0
```

#### 1.8.2 Quasi-Newton as a first improvement

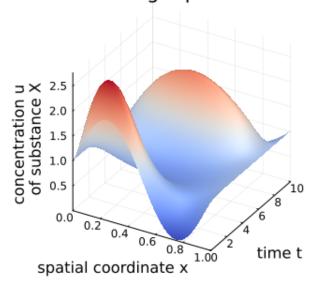
We see that even for  $\Delta t = 0.1$  at N = 400 we have no stability issues. However, the LU decomposition takes  $\mathcal{O}(N^3)$  operations and we calculate this multiple times in each step where building the Jacobian without optimizations takes  $\mathcal{O}(N)$  function evaluations (of a function  $f: \mathbb{R}^{2(N+2)} \to \mathbb{R}^{2(N+2)}$ ). We can improve by using a factorization more suited for sparse matrices (which is automatically selected) and quasi-Newton instead of Newton.

```
[]: function quasi_newtons_method(f!, u0, dt, tol, max_iter)
       The root of g_zero (and therefore the next step in time in
       the implicit Euler scheme) is found using Newton's method,
       which is implemented here.
       u = copy(u0)
       # Here a matrix decomposition automatically
       # chosen is performed the full Jacobian is
       # constructed in every step
       Fac = factorize(Jg(f!, u, dt))
        for i in 1:max_iter
           a = Fac \setminus g_zero(f!, u, u0, dt)
           u .= u .- a
           if norm(g_zero(f!, u, u0, dt)) < tol</pre>
              return u
           end
       end
       return u
    end
```

quasi\_newtons\_method (generic function with 1 method)

```
[]: f = plot_as_2D_surface(u_own, t, x, "Implicit Euler dt = $dt", N)
save("figures/bruss_impl_quasi_newton.svg", f)
display(f)
```

# Simplified Brusselator approached via Implicit Euler dt = 0.1 number of gridpoints N = 400



Notice that for each timestep we still calculate the full Jacobian - so  $(2(N+2))^2$  entries - which in a naive implementation not using sparsity takes 2(N+2) calls of f.

#### 1.8.3 Efficiently constructing the sparse Jacobian

Let us take a look at the sparsity pattern of the Jacobian at the hand of the case N=2 to make things a bit easier to understand and visualize.

#### Initialization of the simplified Brusselator with N=5

```
[]: # Brusselator parameters
N = 5 # number of grid points
t_0 = 0.0 # initial time
t_end = 10.0 # end time
t = range(t_0, t_end, length=100) # saved time steps
x, w = init_brusselator(N) # initial conditions

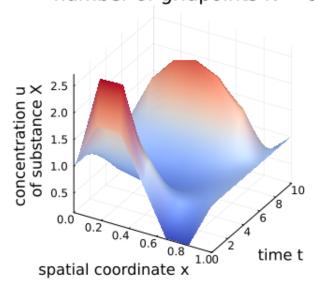
# Define the ODE problem
ode_prob = ODEProblem(brusselator!, w, (0, 10), N)

# Solve the ODE problem using a stiff solver
# from the DifferentialEquations package
sol = Obtime solve(ode_prob, TRBDF2(), saveat=t);
```

418.800 s (7752 allocations: 797.94 KiB)

```
[]: # Plot the solution as a surface over space and time
u, v = sol_to_u_v(sol, N)
f = plot_as_2D_surface(u, t, x, "TRBDF2", N)
display(f)
save("figures/bruss_sparse.svg", f)
```

# Simplified Brusselator approached via TRBDF2 number of gridpoints N = 5



```
[]: # Animate the the change of the solution over time anim_plot_1D(u, N)
```

Info: Saved animation to c:\Users\leona\OneDrive\Dokumente\Studium\Master\2.
Semester\Scientific Machine Learning\Code\figures\bruss\_ani.gif
@ Plots C:\Users\leona\.julia\packages\Plots\p3KMq\src\animation.jl:156

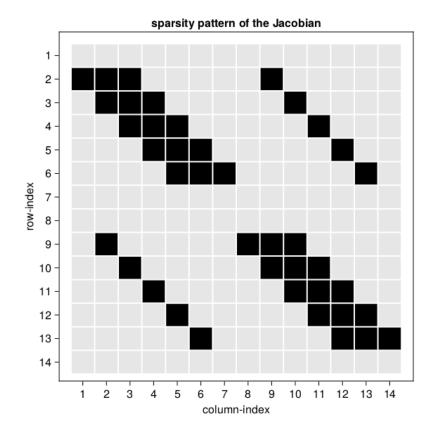
#### The Jacobian is sparse

```
[]: # Find the sparsity pattern of the Jacobian of the ODE considered dwdt0 = zeros(2(N + 2)) # initialization of the derivative sparsity = Symbolics.jacobian_sparsity((du, u) -> brusselator!(du, u, N, 0), dwdt0, w) # find the sparsity pattern

J = ForwardDiff.jacobian(wbrusselator!, dwdt0, w) # find the Jacobian display(J)
```

```
# Plot the sparsity pattern of the Jacobian
f = plot_sparsity_pattern_black(sparsity, "sparsity pattern of the Jacobian")
save("figures/black_sparsity.svg", f)
```

14×14 Matrix{Float64}:											
0.0	0.0	0.0	0.0		0.0	0.0	0.0	0.0			
0.72	5.75615	0.72	0.0		0.0	0.0	0.0	0.0			
0.0	0.72	5.75615	0.72		0.0	0.0	0.0	0.0			
0.0	0.0	0.72	0.56		1.0	0.0	0.0	0.0			
0.0	0.0	0.0	0.72		0.0	0.0179492	0.0	0.0			
0.0	0.0	0.0	0.0		0.0	0.0	0.0179492	0.0			
0.0	0.0	0.0	0.0		0.0	0.0	0.0	0.0			
0.0	0.0	0.0	0.0		0.0	0.0	0.0	0.0			
0.0	-8.19615	0.0	0.0		0.0	0.0	0.0	0.0			
0.0	0.0	-8.19615	0.0		0.72	0.0	0.0	0.0			
0.0	0.0	0.0	-3.0		-2.44	0.72	0.0	0.0			
0.0	0.0	0.0	0.0		0.72	-1.45795	0.72	0.0			
0.0	0.0	0.0	0.0		0.0	0.72	-1.45795	0.72			
0.0	0.0	0.0	0.0		0.0	0.0	0.0	0.0			



```
[]: # Compute the ratio of non-zero elements in the Jacobian
ratio = count(!iszero, sparsity) / length(sparsity)
println("The ratio of non-zero elements in the Jacobian is $ratio for N = $N")

# Compute the sparsity for N = 400
x400, w400 = init_brusselator(400); dwdt0400 = zeros(2(400 + 2)); sparsity400 = Symbolics.jacobian_sparsity((du, u) -> brusselator!(du, u, N, 0), dwdt0400, w400)
ratio400 = count(!iszero, sparsity400) / length(sparsity400)
println("The ratio of non-zero elements in the Jacobian is $ratio400 for N = 400")
```

The ratio of non-zero elements in the Jacobian is 0.20408163265306123 for N = 5 The ratio of non-zero elements in the Jacobian is 6.187965644414742e-5 for N = 400

Only a few elements in the Jacobian are non-zero - but how can be use this to our advantage in constructing the Jacobian?

Using Automatic Differentiation we need one function call of f to calculate one Jacobian-Vector product (where f is applied to a dual number) and if we choose the unit vectors as directions N function evaluation for the whole Jacobian.

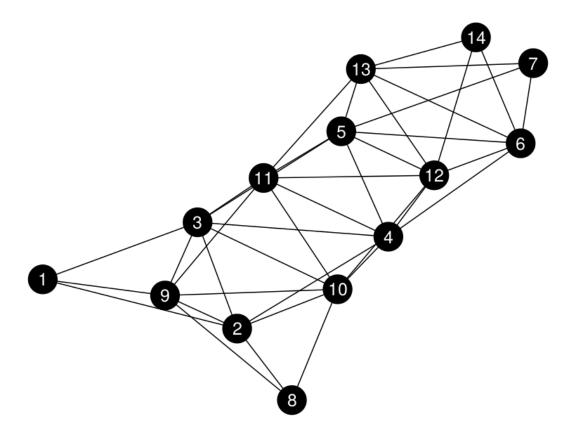
However, for non-overlapping columns (those not having a 1 in the same row), we can group multiple unit vectors into one direction we take the directional derivative along and can construct the whole Jacobian with far fewer function calls.

Finding the grouping with the least number of groups grouping the columns of the Jacobian into non-overlapping groups is equivalent to the distance-1 graph coloring problem.

#### Converting the Sparsity Pattern of the Jacobian into a Graph

```
[]: # Convert the sparsity pattern to a graph
g = sparsity_matrix_to_graph(sparsity)

# Plot the graph
plot_graph(g)
```



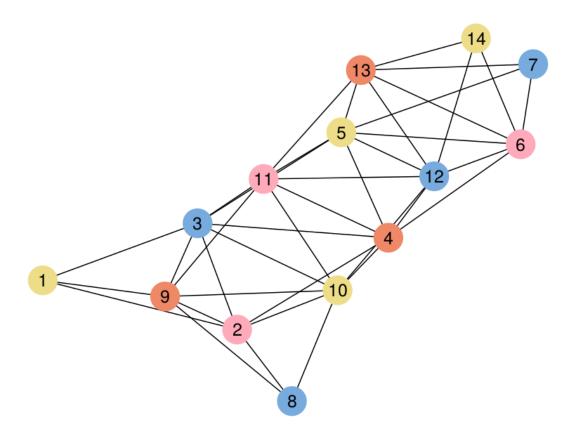
Solving the distance-1 Graph coloring problem Consider the graph where each node represents a column of S. Two nodes are connected by an edge if the respective columns are overlapping. Finding optimal directions for the directional derivatives then corresponds to finding the optimal distance-1 coloring of the graph.

Graph coloring is NP-complete (perfect coloring takes exponentially long) but we only need to do it once for our Jacobian and we use an approximate non-ideal greedy algorithms for finding groups.

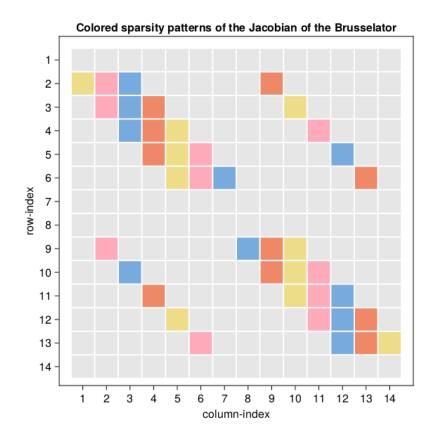
TODO: Maybe present Greedy algorithm

```
[]: # Find the distance-1 coloring of the graph
coloring = greedy_color(g; sort_degree = true)

# Plot the graph with the distance-1 coloring
plot_graph(g, coloring)
```



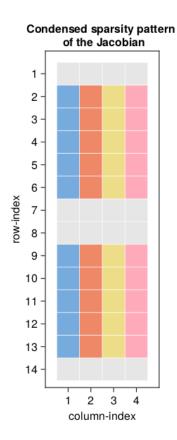
```
[]: # Plot the sparsity pattern with the coloring found previously
f = plot_sparsity_pattern(sparsity, node_colors(coloring, g), "Colored sparsity_
patterns of the Jacobian of the Brusselator")
save("figures/colored_sparsity_pattern.svg", f)
```



We can combine the directional derivatives of each group given by the same color and later construct the full Jacobian using the sparsity matrix.

```
[]: # Find the condensed pattern, i. e. the matrix resulting from
# overlapping the non-overlapping columns as found by the coloring
cm = condense_sparsity_pattern(sparsity, coloring)

# Plot the condensed pattern
f = plot_sparsity_pattern(cm, unique_colors(coloring), "Condensed sparsity_
→pattern \n of the Jacobian")
save("figures/colored_condensed_sparsity.svg", f)
```



Using this technique, we can again increase the efficiency of our implicit method.

```
forwarddiff_color_jacobian!(J, f!, u, colorvec = colors, sparsity = __
 ⇔sparsity)
   return I - dt * J
end
function quasi_newtons_method_c(f!, u0, dt, tol, max_iter, J, sparsity, colors)
   The root of g_zero (and therefore the next step in time in
   the implicit Euler scheme) is found using Newton's method,
   which is implemented here.
   11 11 11
   u = copy(u0)
   # Here a matrix decomposition automatically
   # chosen is performed the full Jacobian is
   # constructed in every step
   Fac = factorize(Jg2(f!, u, dt, J, sparsity, colors))
   for i in 1:max iter
       a = Fac \setminus g_zero(f!, u, u0, dt)
       u .= u .- a
       if norm(g_zero(f!, u, u0, dt)) < tol</pre>
           return u
       end
   end
   return u
end
function implicit_euler_c(f!, g, zero_finder, dt, w, num_steps, J, sparsity, u
 ⇔colors)
   0.00
   A stiff differential equation can be
   solved using the implicit Euler scheme.
   N = Int32(length(w) / 2 - 2)
   u = zeros(2(N + 2), num_steps)
   u[:, 1] = w
   for i in 2:num_steps
       u[:, i] = zero_finder(f!, u[:, i - 1], dt, 1e-4, 1000, J, sparsity, u[:, i]
 ⇔colors)
   end
   return u
end
```

implicit\_euler\_c (generic function with 1 method)

```
[]: N = 400; dt = 0.1; x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)
J, sparsity, colors = do_coloring(N, w)

# Our solution
sol = @btime implicit_euler_c(wbrusselator!, g_zero, quasi_newtons_method_c,_u
odt, w, num_steps, J, sparsity, colors);

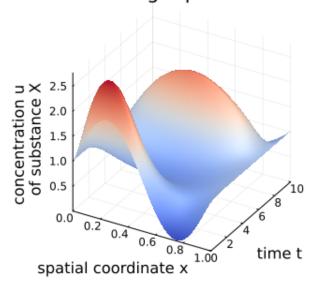
877.994 ms (17404 allocations: 1.51 GiB)

[]: u_own = sol[1:N+2, :]';
f = plot_as_2D_surface(u_own, t, x, "Implicit Euler", N)
```

save("figures/implicit\_euler\_colored.svg", f)

display(f)

# Simplified Brusselator approached via Implicit Euler number of gridpoints N = 400



### ${\bf 1.8.4}\quad {\bf Alternative:\ Jacobian\ free\ Newton\ -\ Krylov\ subspace\ method}$

In the Newton steps, we solve a system of linear equations. Using Krylov subspace methods this can be done using only matrix (here jacobian) vector products where we can use that one JVP (Jacobian vector product) can be calculated in one forward AD (automatic differentiation) pass.

Solving Linear Systems only based on matrix vector products I: Gradient descent

```
[]:  # Consider the matrix
A = [1 0.1; 0.1 10] # ill-conditioned
```

```
b = [1, 1]
# solve the system Ax = b
Af = factorize(A)
x = Af \ b
```

2-element Vector{Float64}: 0.990990990990991 0.09009009009009009

Solving this system is equivalent to finding the minimum of the quadratic form  $0.5 \text{ x}^T\text{Ax}$  -  $\text{x}^T\text{D}$  for A positive definite and symmetric.

```
[]: qf = x -> 0.5 * x' * A * x - b' * x

# plot the quadratic form as a heatmap for x1 and x2 in [-2, 2]

x1 = range(-100, 100, length=100)

x2 = range(-100, 100, length=100)

# plot heatmap

# fig, ax, pl = CairoMakie.contourf(x1, x2, (x1, x2) -> qf([x1, x2]), levels = 20, title="Quadratic form", xlabel="x1", ylabel="x2")

# CairoMakie.scatter!([x[1]], [x[2]], markersize=20, markercolor=:red)

# display(fig)

# Plots.surface(x1, x2, (x1, x2) -> qf([x1, x2]), color = :viridis, 40)

**Alabel="x1", ylabel="x2", zlabel="qf(x1, x2)")
```

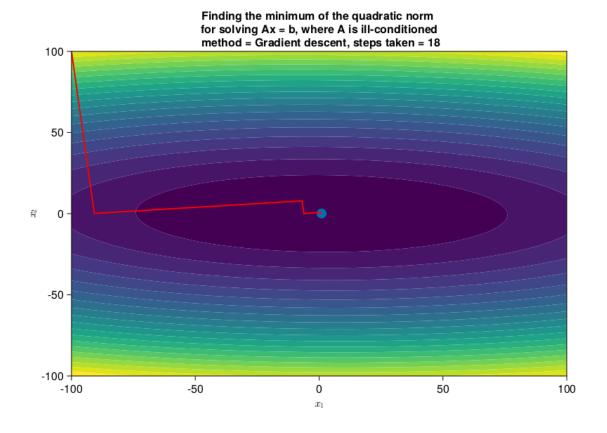
#### -100.0:2.0202020202020203:100.0

```
[]: # Implement and visualize a simple gradient descent
     # method for the quadratic form above
     function gradient_descent(A, b, x0, tol, max_iter)
         Gradient descent method for the quadratic form qf
         x = copy(x0)
         r = A * x - b
         x_history = zeros(length(x), max_iter)
         x_history[:, 1] = x
         for i in 2:max_iter
             gamma = (r' * r) / (r' * A * r)
             x = x - gamma * r
             x_history[:, i] .= x
             if norm(r) < tol</pre>
                 return x, x_history[:, 1:i]
             end
             r = A * x - b
         end
         return x, x_history
     end
```

```
# Initialize the gradient descent method
x0 = [-100, 100]
tol = 1e-6
max_iter = 1000
x, xh = gradient_descent(A, b, x0, tol, max_iter)
length(xh)
```

38

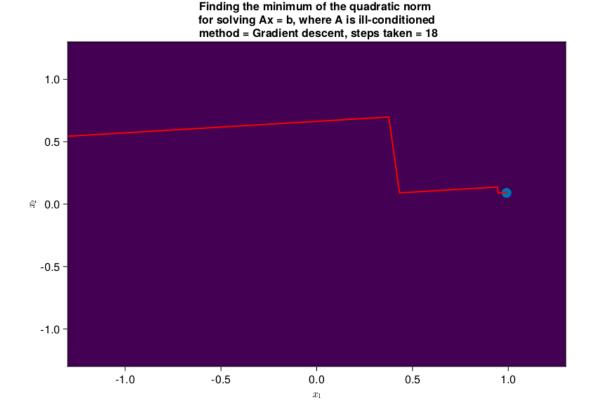
```
[]: f, ax = plot_heat_stepper(x1, x2, qf, xh, "Gradient descent"); save("figures/gradient_descent.svg", f)
```



### CairoMakie.Screen{SVG}

```
[]: # set x limits to [-10, 10], [-10, 10]
1 = 1.3
CairoMakie.xlims!(ax, -1, 1)
CairoMakie.ylims!(ax, -1, 1)
save("figures/gradient_descent_lims.svg", f)
```

display(f)



#### CairoMakie.Screen{IMAGE}

It turns out that the worse-conditioned the matrix is the more steps are needed to find the minimum of its quadratic form in gradient descent. Krylow subspace methods where consecutive step-directions are conjugate with respect to  $A((p^i)^T A p^j = 0, i \neq j)$  can be advantageous.

Solving Linear Systems only based on matrix vector products I: Krylow subspace methods

```
function simple_krylov_solver(A, b, x0, tol, max_iter)
    """

A simple Krylov subspace solver for Ax = b
    """

x = copy(x0)
p = b - A * x

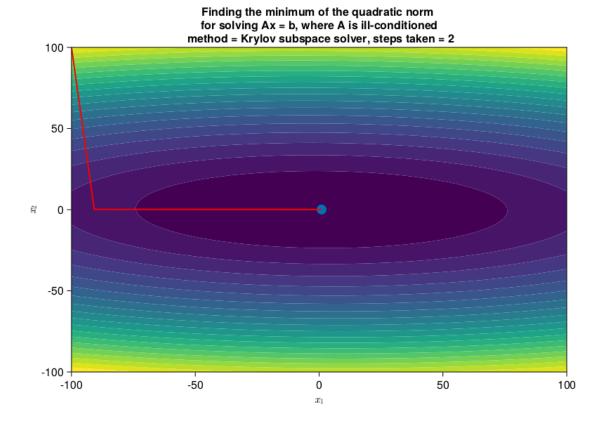
x_history = zeros(length(x), max_iter)
x_history[:, 1] .= x

for n in 2:max_iter
```

```
r = b - A * x
        alpha = (r' * p) / (p' * A * p)
        x = x + alpha * p
        x_history[:, n] .= x
        rn = b - A * x
        if norm(rn) < tol</pre>
            return x, x_history[:, 1:n]
        end
        pn = rn + (rn' * rn) / (r' * r) * p
        \# println("Directions are conjugate: pn' * A * p = ", pn' * A * p)
        # println("Residuals and steps are mutually orthogonal: rn' * p = ", \_
 \rightarrow rn' * p)
        # println("Residuals are orthogonal: rn' * p = ", rn' * r)
        p = pn
    return x, x_history
end
```

simple\_krylov\_solver (generic function with 1 method)

```
[]: x, xh = simple_krylov_solver(A, b, x0, tol, max_iter)
f, ax = plot_heat_stepper(x1, x2, qf, xh, "Krylov subspace solver");
save("figures/conjugate_gradient_descent.svg", f)
```



The problem with the Method shown above is that it only works for symmetric pdf matrices. We therefore resort to GMRES.

Wikipedia: In mathematics, the generalized minimal residual method (GMRES) is an iterative method for the numerical solution of an indefinite nonsymmetric system of linear equations. The method approximates the solution by the vector in a Krylov subspace with minimal residual. The Arnoldi iteration is used to find this vector.

```
[]: # Test the Krylov solver, with bigger A
A = rand(100, 100)
A = A'*A

println(isposdef(A))

b = rand(100)

x0 = rand(100)

tol = 1e-6

max_iter = 1000

AF = factorize(A)
xc = AF \ b

xk, xkh = simple_krylov_solver(A, b, x0, 1e-5, 1000)

norm(xk - xc)
```

true

#### 0.052517729505130506

Krylov subspace methods applied to solving stiff ODEs

```
function implicit_euler_k(f!, g, zero_finder, dt, w, num_steps)
    """

A stiff differential equation can be
    solved using the implicit Euler scheme.
    """

N = Int32(length(w) / 2 - 2)
    u = zeros(2(N + 2), num_steps)
    u[:, 1] = w
    for i in 2:num_steps
        u[:, i] = zero_finder(f!, u[:, i - 1], dt, 1e-2, 10)
    end
    return u
```

```
end
#
#
                             calls
#
function newtons_method_k(f!, u0, dt, tol, max_iter)
    The root of g_zero (and therefore the next step in time in
    the implicit Euler scheme) is found using Newton's method,
    which is implemented here.
    0.00
    u = copy(u0)
    function g!(dudt, u)
        f!(dudt, u)
        dudt[:] = u - u0 - dt * dudt
    end
    # Jm = Jg(f!, u, dt)
    J = JacVec(g!, u0, autodiff = true)
    function mat_vec!(res, w)
        mul!(res,J,w)
    end
    Jm = LinearMap(mat_vec!, 2*(N+2); issymmetric = false, ismutating = true)
    for i in 1:max_iter
        a = idrs(Jm, g_zero(f!, u, u0, dt); s = 8, log = false)
        u \cdot = u \cdot - a
        if norm(g_zero(f!, u, u0, dt)) < tol</pre>
            return u
        end
    end
    return u
end
                             calls
#
#
                 idrs or gmres or ...
```

```
# Note: The conjugate gradient method does not work here as the Jacobian is not_{\sqcup}
 ⇔positive definite
# function krylov_solve(f!, b, dt, x0, tol, max_iter, Jg)
#
      A simple Krylov subspace solver for Jx = b
#
#
      x = copy(x0)
     Jg = Jg' * Jg
#
      b = Jg' * b
#
      \# J = JacVec(f!, x0, autodiff = true)
      # Jg = Jg(f!, x0, dt)
#
      \# res = similar(x)
      # mul!(res, J, x)
      \# p = b - (x - dt * res)
#
     p = b - Jg * x
     println(issymmetric(Jg[2:end-1, 2:end-1]))
#
      for n in 2:max_iter
#
          # mul!(res, J, x)
          \# r = b - (x - dt * res)
#
          r = b - Jq * x
#
          # mul!(res, J, p)
#
          \# \ alpha = (r' * p) / (p' * (p - dt * res))
          alpha = (r' * p) / (p' * Jg * p)
#
          x = x + alpha * p
#
          # mul!(res, J, x)
          \# rn = b - (x - dt * res)
#
          rn = b - Jg * x
          if \ norm(rn) < tol
#
#
              return x
#
          p = rn + (rn' * rn) / (r' * r) * p
#
#
      end
#
      return x
# end
```

newtons\_method\_k (generic function with 1 method)

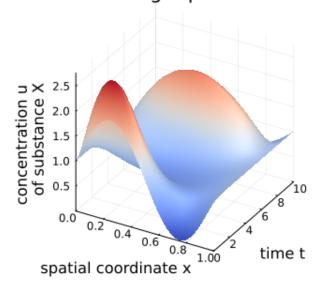
```
[]: N = 400; dt = 0.1; x, w, t, num_steps = init_explicit_problem(N, t_0, t_end, dt)
# Our solution
sol = @btime implicit_euler_k(wbrusselator!, g_zero, newtons_method_k, dt, w,___
num_steps);

1.961 s (935735 allocations: 4.95 GiB)

[]: u_own = sol[1:N+2, :]';
f = plot_as_2D_surface(u_own, t, x, "JFNK", N)
save("figures/jnfk.svg", f)
```

display(f)

# Simplified Brusselator approached via JFNK number of gridpoints N = 400



#### 1.9 Tests of Methods from the Differential Equations.jl package

```
ode_prob = ODEProblem(f!, w, (t_0, t_end), N);
[]: sol_def = Obtime solve(ode_prob)
     print(sol_def.destats)
     print(sol_def.alg)
      14.961 ms (33879 allocations: 27.41 MiB)
    DiffEqBase.Stats
    Number of function 1 evaluations:
                                                        502
    Number of function 2 evaluations:
                                                        0
    Number of W matrix evaluations:
                                                        28
    Number of linear solves:
                                                        388
    Number of Jacobians created:
                                                        5
    Number of nonlinear solver iterations:
                                                        315
    Number of nonlinear solver convergence failures:
                                                        1
    Number of rootfind condition calls:
                                                        0
    Number of accepted steps:
                                                        46
    Number of rejected steps:
                                                        5
    Maximum eigenvalue recorded:
    588955CompositeAlgorithm{Tuple{Tsit5{typeof(OrdinaryDiffEq.trivial_limiter!),
    typeof(OrdinaryDiffEq.trivial_limiter!), Static.False}, TRBDF2{12, false,
    KLUFactorization, NLNewton{Rational{Int64}, Rational{Int64}, Rational{Int64},
    Rational{Int64}}, typeof(OrdinaryDiffEq.DEFAULT_PRECS), Val{:forward}, true,
    nothing}},
    OrdinaryDiffEq.AutoSwitchCache{Tsit5{typeof(OrdinaryDiffEq.trivial_limiter!),
    typeof(OrdinaryDiffEq.trivial limiter!), Static.False}, TRBDF2{0, false,
    Nothing, NLNewton{Rational{Int64}, Rational{Int64}, Rational{Int64},
    Rational{Int64}}, typeof(OrdinaryDiffEq.DEFAULT_PRECS), Val{:forward}, true,
    nothing}, Rational{Int64}, Int64}}((Tsit5(stage_limiter! = trivial_limiter!,
    step_limiter! = trivial_limiter!, thread = static(false)), TRBDF2{12, false,
    KLUFactorization, NLNewton{Rational{Int64}, Rational{Int64}, Rational{Int64},
    Rational{Int64}}, typeof(OrdinaryDiffEq.DEFAULT PRECS), Val{:forward}, true,
    nothing}(KLUFactorization(true, true), NLNewton{Rational{Int64},
    Rational{Int64}, Rational{Int64}, Rational{Int64}, (1//100, 10, 1//5, 1//5,
    false, true, 0//1), OrdinaryDiffEq.DEFAULT_PRECS, true, :linear, :PI)),
    OrdinaryDiffEq.AutoSwitchCache{Tsit5{typeof(OrdinaryDiffEq.trivial_limiter!),
    typeof(OrdinaryDiffEq.trivial_limiter!), Static.False}, TRBDF2{0, false,
    Nothing, NLNewton{Rational{Int64}, Rational{Int64}, Rational{Int64},
    Rational{Int64}}, typeof(OrdinaryDiffEq.DEFAULT_PRECS), Val{:forward}, true,
    nothing}, Rational{Int64}, Int64}(44, 0, Tsit5(stage_limiter! =
    trivial_limiter!, step_limiter! = trivial_limiter!, thread = static(false)),
    TRBDF2{0, false, Nothing, NLNewton{Rational{Int64}, Rational{Int64},
    Rational{Int64}, Rational{Int64}}, typeof(OrdinaryDiffEq.DEFAULT_PRECS),
    Val{:forward}, true, nothing}(nothing, NLNewton{Rational{Int64},
    Rational{Int64}, Rational{Int64}, Rational{Int64}, (1//100, 10, 1//5, 1//5,
    false, true, 0//1), OrdinaryDiffEq.DEFAULT_PRECS, true, :linear, :PI), true, 10,
    3, 9//10, 9//10, 2, false, 5))
```

```
[]: \# ===================== Test some methods from the Differential Equations.jlu
     print("\n============ Tsitouras 5/4 Runge-Kutta method∪
     ω=======\n")
    print(@btime solve(ode_prob, Tsit5(), saveat=tc).destats)
   ======= Tsitouras 5/4 Runge-Kutta method ==============
     4.623 s (3080807 allocations: 9.58 GiB)
   DiffEqBase.Stats
   Number of function 1 evaluations:
                                              220047
   Number of function 2 evaluations:
   Number of W matrix evaluations:
   Number of linear solves:
   Number of Jacobians created:
   Number of nonlinear solver iterations:
   Number of nonlinear solver convergence failures:
   Number of rootfind condition calls:
   Number of accepted steps:
                                              36671
   Number of rejected steps:
[]: # print("\n=========== Jacobian-Free Newton Krylovu
    ⇒========\n")
    # print(solve(ode_prob, KenCarp47(linsolve = KrylovJL_GMRES()), save_everystepu
     ⇒= false).destats);
[]:  # print("\n========\n")
    # print(@btime solve(ode\_prob, Rosenbrock23(autodiff = true), abstol = 1e-2, ____
     \hookrightarrow reltol = 1e-2).destats)
[]: | # print("\n========\n")
    # print(@btime solve(ode_prob, Rodas5P()).destats)
[]:  # print("\n=======\n")
    # print(@btime solve(ode_prob, lsoda()).destats)
[]:|# print("\n========\n")
    # print(@btime solve(ode_prob, TRBDF2()).destats)
```