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# Vibrational properties of bulk LaAlO<sub>3</sub> from Fourier-transform infrared ellipsometry

Travis Willett-Gies, Eric DeLong, Stefan Zollner\*

Department of Physics, New Mexico State University, MSC 3D, P.O. Box 30001, Las Cruces, NM 88003-8001, USA

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## ABSTRACT

We used Fourier-transform infrared spectroscopic ellipsometry to determine the dielectric function of twinned single-crystalline bulk lanthanum aluminate at 300 K in the region of lattice vibrations from 250 to 1000 cm<sup>−1</sup>. We fit the experimental data using a classical sum of Lorentz oscillators as well as a factorized model. We were able to determine the parameters of five infrared-active optical phonons within our spectral range. Transverse phonons appear as peaks in the imaginary part of the dielectric function which are clearly visible without fitting. By transforming the data to obtain the loss function, we are able to observe the longitudinal phonons as peaks in the imaginary part. The polar nature of LaAlO<sub>3</sub> causes a strong splitting between the transverse optical (TO) and longitudinal optical (LO) phonon energies. We report energies, amplitudes and broadenings of five TO/LO phonon pairs and compare the two models used to describe the data.

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## 1. Introduction

Lanthanum aluminate (LaAlO<sub>3</sub>) is widely used as a substrate material in oxide epitaxy [1]. Recently, thin films of LaAlO<sub>3</sub> have drawn considerable attention owing to the discovery of a two-dimensional electron gas in LaAlO<sub>3</sub>/SrTiO<sub>3</sub> heterostructures [2–4]. During the early development of high-*k* gate dielectrics, the semiconductor industry considered LaAlO<sub>3</sub> as a replacement for SiO<sub>2</sub> as a gate dielectric, mostly because of its close lattice match with Si [5–7].

LaAlO<sub>3</sub> is a polar perovskite with a distorted cubic crystal structure. It can be viewed as an alternating stack of positively charged LaO and negatively charged AlO<sub>2</sub> planes. The distortions lead to a rhombohedral structure with space group  $R\bar{3}c$  or  $D_{3d}^{5d}$  (space group 167) at room temperature [8–12]. The rhombohedral lattice constants of LaAlO<sub>3</sub> are found to be  $a = b = 5.365$  Å and  $c = 13.111$  Å using neutron powder diffraction [9,10]. The long-wavelength optical phonon modes expected for this crystal structure are given by the factor group [12–14]

$$\Gamma(D_{3d}^6) = 2A_{1u} + 3A_{2g} + A_{1g} + 3A_{2u} + 4E_g + 5E_u. \quad (1)$$

The relationship between the cubic  $Pm\bar{3}m$  or  $O_h^h$  perovskite phonons and the crystal field splittings due to the rhombohedral distortions are summarized in Table 1 [12–14].

Near 800 K, LaAlO<sub>3</sub> transforms from the rhombohedral to the cubic crystal structure [9,10,12,15]. Raman measurements of low-energy phonons with energy below 200 cm<sup>−1</sup> have revealed that there is a soft phonon [9,12] associated with this transition. The energy splitting of

the lowest-energy  $E_g/A_{1g}$  phonon pair (representing the rotation of an AlO<sub>6</sub> octahedron) decreases with increasing temperature. At 800 K, they merge into a zero-frequency phonon with  $F_{2u}$  symmetry [9,12].

The vibrational structure of LaAlO<sub>3</sub> has been studied both theoretically and experimentally. Abrashev et al. [14] calculated the long-wavelength optical phonon frequencies using a semi-empirical shell model and plotted diagrams showing the displacement patterns of the atoms involved in the vibration. They also included the splittings between transverse (TO) and longitudinal (LO) optical phonons due to the long-range polar interactions. More recently, Delugas et al. [16] presented an ab initio density-functional and self-interaction-corrected calculation of the optical phonon frequencies, in which the TO/LO splittings were ignored. None of these calculations are sufficiently accurate for comparison with experimental results for the phonon frequencies.

The factor group analysis in Eq. (1) describes the symmetry of the optical phonons and how they can be observed experimentally. The results are summarized in Table 1: One  $A_{1g}$  mode and four  $E_g$  modes are Raman active. Three  $A_{2u}$  and five  $E_u$  modes are infrared active and can be observed with Fourier-transform infrared (FTIR) ellipsometry. Finally, two  $A_{1u}$  modes and three  $A_{2g}$  modes are silent and not accessible experimentally using first-order Raman or infrared measurements.

While most Raman spectroscopy studies have focused on the temperature dependence of the  $A_{1g}$  and  $E_g$  soft modes and their implications about the phase transition from the rhombohedral to the cubic crystal structure [9,12,14], the higher-energy Raman modes have also been investigated [12,14]. Four Raman modes have clearly been identified, while the assignment of the fifth mode is controversial. See Table 1 for a summary of experimental results.

The IR-active modes have been investigated by infrared reflectance or transmittance measurements, usually followed by fitting of the dielectric function as a sum of Lorentz oscillators [13,17–19]. Three pairs

\* Corresponding author.

E-mail address: [zollner@nmsu.edu](mailto:zollner@nmsu.edu) (S. Zollner).

**Table 1**

Summary of transverse optical phonon modes for LaAlO<sub>3</sub> at 300 K. All energies are listed in units of cm<sup>-1</sup>.

R3c symmetry	Pm3m symmetry	Activity	Pattern	Energy (exp)
E <sub>g</sub>	F <sub>2u</sub> (R)	Raman	AlO <sub>6</sub>	33 <sup>a</sup>
A <sub>1g</sub>	F <sub>2u</sub> (R)	Raman	AlO <sub>6</sub>	123 <sup>b</sup>
E <sub>g</sub>	F <sub>1u</sub> (R)	Raman	La	152 <sup>b</sup>
E <sub>g</sub>	F <sub>1u</sub> (R)?	Raman	O	470 <sup>b</sup> (?)
E <sub>g</sub>	E <sub>u</sub> (R)?	Raman	O	487 <sup>b</sup>
A <sub>2u</sub>	F <sub>1u</sub> (I)	IR	AlO-La	188(1) <sup>c</sup>
A <sub>2u</sub>	F <sub>1u</sub> (I)	IR	O bend	427.0(1) <sup>c</sup>
A <sub>2u</sub>	F <sub>1u</sub> (I)	IR	O stretch	650.79(5) <sup>c</sup>
E <sub>u</sub>	F <sub>1u</sub> (I)	IR	AlO-La	188(1) <sup>c</sup>
E <sub>u</sub>	F <sub>2u</sub> (I)	IR	O	weak
E <sub>u</sub>	F <sub>1u</sub> (I)	IR	O bend	427(1) <sup>c</sup>
E <sub>u</sub>	F <sub>2g</sub> (R)	IR	Al	495.72(1) <sup>c</sup>
E <sub>u</sub>	F <sub>1u</sub> (I)	IR	O stretch	708.2(9) <sup>c</sup>
A <sub>1u</sub>	F <sub>2u</sub> (I)	silent		NA
A <sub>1u</sub>	F <sub>2g</sub> (R)	silent	Al	NA
A <sub>2g</sub>	F <sub>1u</sub> (R)	silent	La	NA
A <sub>2g</sub>	F <sub>1u</sub> (R)	silent	O	NA
A <sub>2g</sub>	A <sub>2u</sub> (R)	silent		NA

<sup>a</sup> Ref. [12].

<sup>b</sup> Ref. [14].

<sup>c</sup> This work using Eq. (4).

of A<sub>2u</sub>/E<sub>u</sub> modes, arising from the crystal field splitting of the three cubic F<sub>1u</sub>(I) modes, lead to strong features in the spectra [14]. One weak zone-folded E<sub>u</sub> mode is also seen clearly, while the second E<sub>u</sub> mode (derived from F<sub>2u</sub>, which is silent in the cubic parent structure) has been considered too weak to be observable [14]. There is considerable disagreement in the literature about the exact phonon energies, amplitudes, and broadenings used to fit experimental spectra.

In this work, we determined the energies, broadenings, and amplitudes of the infrared-active optical phonons with very high accuracy using Fourier-transform infrared ellipsometry. We also describe the symmetries of all phonons and their relationships, such as rhombohedral splittings and the relationship to their cubic phonon counterparts.

## 2. Experiment and models

Two single-sided polished, 2-inch LaAlO<sub>3</sub> wafers with 0.5 mm thickness and (100) surface orientation were obtained commercially [20]. To reduce the reflections from the backside, one wafer was roughened further using a bead blaster. Our substrates are twinned and the surface orientation refers to the pseudo-cubic structure. Additional information about the samples and their properties are given in our earlier work [21], where we also report the dielectric function of LaAlO<sub>3</sub> from 0.8 to 6.6 eV between 77 and 700 K. Most importantly for the present work, our earlier research found  $\epsilon_{\infty} = 4.12 \pm 0.01$  at 300 K.

Infrared ellipsometry measurements were performed on a J.A. Woollam FTIR-VASE variable angle of incidence ellipsometer at the Center for Integrated Nanotechnologies user facility. This instrument is based on a fixed analyzer (at 0° and 180°), a fixed polarizer (at  $\pm 45^\circ$ ), and a rotating compensator. To increase accuracy, two fixed positions for the analyzer and polarizer were chosen (four-zone measurements), as this cancels experimental errors to first order in the analyzer and polarizer position. We measured at four angles of incidence: 60°, 65°, 70°, and 75°. Nominally, the instrument reports data between 250 and 8000 cm<sup>-1</sup> but we restrict our analysis to the region of the lattice vibrations between 250 and 1000 cm<sup>-1</sup>. We did not observe features in the spectra above 1000 cm<sup>-1</sup> other than normal dispersion. In the spectral range around 6000 cm<sup>-1</sup>, our FTIR ellipsometry results are consistent with our earlier near-IR work [21].

For various reasons, FTIR ellipsometry measurements on a bulk LaAlO<sub>3</sub> wafer are quite challenging. On one hand, the reflectance of LaAlO<sub>3</sub> is quite low (about 11% at normal incidence in the mid-IR), much lower than for a bulk semiconductor or for a thin film on Si. In

the region of the lattice vibrations, the reflectance becomes very small [18,22] near 300, 620, and 800 cm<sup>-1</sup>. On the other hand, the phonon broadenings in LaAlO<sub>3</sub> are very low (near 4 cm<sup>-1</sup> or less). We therefore select a resolution of 2 cm<sup>-1</sup> for the FTIR spectrometer. These conditions create noise below 350 and near 800 cm<sup>-1</sup>. We are thus forced to select a data acquisition time of eight hours to improve the signal to noise ratio, using 20 FTIR scans per spectrum and 15 spectra for each revolution of the rotating compensator.

Spectroscopic ellipsometry measures the ellipsometric angles  $\psi$  and  $\Delta$  as a function of photon energy. These ellipsometric angles and the Fresnel reflectance ratio  $\rho = e^{i\Delta} \tan \psi$  are related to the pseudo-refractive index  $\hat{n}$  and the pseudo-dielectric function  $\epsilon^* = \hat{n}^2$  of the sample through [23,24]

$$\rho = \frac{(\hat{n} \cos \phi_0 - \cos \phi_1) (\cos \phi_0 + \hat{n} \cos \phi_1)}{(\hat{n} \cos \phi_0 + \cos \phi_1) (\cos \phi_0 - \hat{n} \cos \phi_1)} \quad (2)$$

where  $\phi_0$  is the angle of incidence and  $\phi_1$  the angle of refraction. For an ideal sample without surface overlayers,  $\hat{n}$  and  $\epsilon^*$  are equal to the refractive index  $n$  and the dielectric function  $\epsilon = n^2$ .  $\Delta$  equals zero or  $\pi$  for an ideal transparent substrate (for an insulator outside of the region of lattice vibrations), because all quantities in Eq. (2) are real. For our LaAlO<sub>3</sub> substrates, the surface overlayers are very thin (see Table I in Ref. [21]). The surface roughness is between 15 and 20 Å for our samples.

Normally, spectroscopic ellipsometry requires extensive data analysis to determine the optical constants (the complex dielectric function  $\epsilon = n^2$ ) from the ellipsometric angles. This is not the case for our analysis. We can simply convert the measured ellipsometric angles into the dielectric function using Eq. (2), because the effect of surface roughness is negligible. The transverse optical (TO) phonons appear as peaks in the dielectric function [25,26]. Longitudinal optical (LO) vibrations are possible at energies where  $\epsilon(\omega) = 0$ . LO phonons therefore appear as peaks in the loss function  $\text{Im}(-1/\epsilon)$ .

To determine accurate phonon energies, amplitudes, and broadenings, we write the dielectric function  $\epsilon$  as a function of photon energy  $\omega$  as a sum of uncoupled damped harmonic oscillators [25–27]

$$\epsilon(\omega) = \epsilon_{\infty} + \sum_i \frac{A_i \omega_i^2}{\omega_i^2 - \omega^2 - i\gamma_i \omega} \quad (3)$$

The first term  $\epsilon_{\infty} = 4.12 \pm 0.01$  (see Ref. [21]) describes the contributions of electronic transitions to the dielectric function. In principle, we expect eight terms in the sum, one for each infrared active phonon. In practice, some of these phonons may be very weak, while others may be degenerate. Normally, we use five oscillators as described below.

The Lorentz model (3) is derived for charges oscillating in an electric field. This model is classical and it assumes that the frictional force acting on the charge is proportional to their velocity [25]. This results in harmonic damping and a constant broadening term. (The LO broadening parameter is generally different from the TO broadening, most often larger, but not independent.)

The anharmonic coupling of phonons causes the decay of optical phonons into acoustic or other optical phonons with lower energy. Usually, the decay products are zone-edge phonons with a high density of states. If the splitting between the TO and LO phonons is large, then they will have different decay paths and their damping constants  $\gamma_{\text{TO}}$  and  $\gamma_{\text{LO}}$  may differ. Therefore, the uncoupled Lorentz oscillator model (3) often gives a good description of experimental data, which can be improved by assigning different damping parameters to LO and TO phonons yielding [26–29]

$$\epsilon(\omega) = \epsilon_{\infty} \prod_i \frac{\omega_{i,\text{LO}}^2 - \omega^2 - i\gamma_{i,\text{LO}} \omega}{\omega_{i,\text{TO}}^2 - \omega^2 - i\gamma_{i,\text{TO}} \omega} \quad (4)$$

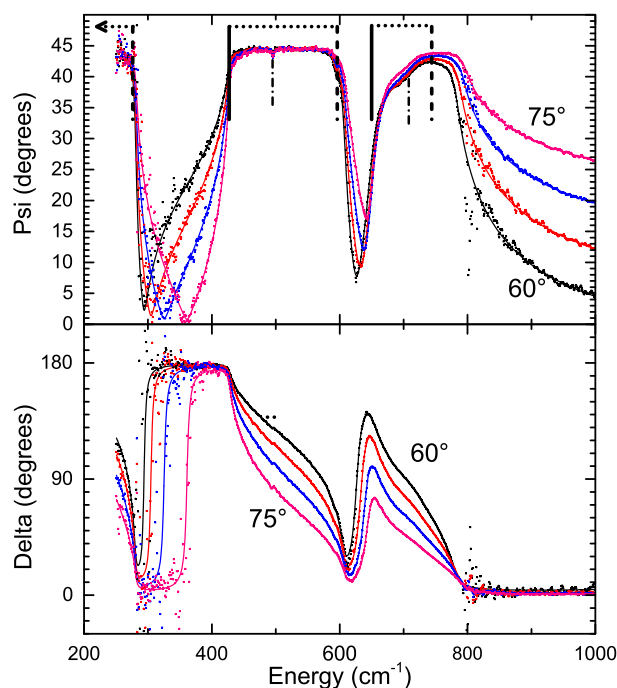
as the functional form for the infrared dielectric function of insulators. If we set  $\gamma_{i,\text{TO}} = \gamma_{i,\text{LO}} = \gamma_i$ , then both descriptions (3) and (4) become

equivalent (at least for a single oscillator or for clearly separated narrow absorption lines) [22,26,28].

### 3. Experimental results

The ellipsometric angles  $\psi$  and  $\Delta$  from 250 to 1000  $\text{cm}^{-1}$  for  $\text{LaAlO}_3$  at 300 K are shown in Fig. 1. The data were taken at four angles of incidence (60°, 65°, 70°, and 75°). The shape of the  $\psi$  spectra can be understood using arguments presented by Humlíček [22] for  $\alpha$ -quartz: The first restrahlen band with  $\psi = 45^\circ$  extends from the lowest TO phonon energy at 188  $\text{cm}^{-1}$  (below our spectral range) to the corresponding LO phonon near 280  $\text{cm}^{-1}$ .  $\psi$  then drops to a minimum located at an energy that increases with the angle of incidence. The second restrahlen band ( $\psi = 45^\circ$ ) extends from 427  $\text{cm}^{-1}$  to 596  $\text{cm}^{-1}$ . It is interrupted by a small dip at 496  $\text{cm}^{-1}$  due to the third phonon. The energy of the second minimum also increases with the angle of incidence. The third restrahlen band has an irregular shape with contributions from two TO phonons located between 650 and 710  $\text{cm}^{-1}$ . After the third restrahlen band,  $\psi$  drops again. For an angle of incidence of 60°,  $\psi$  reaches zero at 1200  $\text{cm}^{-1}$ .  $\psi$  converges to a steady state determined by  $\epsilon_\infty$  near 2000  $\text{cm}^{-1}$ .  $\psi$  never rises above 45°, which would be a sign of optical anisotropy [22].  $\psi$  increases and decreases smoothly at the restrahlen bands (another potential sign of crystal anisotropy found in SiC) [30].

The mid-IR Brewster angle given by  $\tan\phi_B = \sqrt{\epsilon_\infty}$  is 63.8°. Therefore, in the mid-IR,  $\Delta = \pi$  for our measurements at an angle of incidence of  $\phi = 60^\circ$ , while  $\Delta = 0$  for  $\phi = 65^\circ$  or larger.  $\epsilon$  drops below 3 at 1200  $\text{cm}^{-1}$  and thus  $\Delta = 0$  between 800 and 1000  $\text{cm}^{-1}$  for all our incidence angles. Just below the main phonon peak at 427  $\text{cm}^{-1}$ ,  $\epsilon_1$  is very large and therefore the Brewster angle is above 75°. Therefore,  $\Delta = \pi$  in this range for our incidence angles (60° to 75°).  $\epsilon_1$  drops below the main phonon resonance (427  $\text{cm}^{-1}$ ) and therefore  $\Delta$  drops to 0 as the Brewster angle crosses the angle of incidence. The weaker TO phonons between 495 and 710  $\text{cm}^{-1}$  cause more gradual changes in  $\Delta$ . Below 300  $\text{cm}^{-1}$ ,  $\Delta$  rises again due to the influence of the TO phonon at 188  $\text{cm}^{-1}$  (below our spectral range).



**Fig. 1.** (Color online) Ellipsometric angles  $\psi$  (top) and  $\Delta$  (bottom) for  $\text{LaAlO}_3$  at 300 K at four angles of incidence from 60° to 75°. Symbols: experimental data. Lines: fit with Eq. (4) and parameters in Table 3. Solid and dashed vertical lines indicate the location of transverse and longitudinal optical phonons, respectively.

An excellent description of the ellipsometric angles can be achieved using a fit with the Lorentz model in Eq. (3) and the parameters in Table 2. The average mean square deviation between our experimental data and the Lorentz model is only 40% larger than the experimental errors. The low-frequency dielectric constant is given by [25]

$$\epsilon_0 = \epsilon_\infty + \sum_i A_i = 23.81 \pm 0.06 \quad (5)$$

using the parameters in Table 2. The dominant contribution arises from the strong TO phonons at 188 and 427  $\text{cm}^{-1}$ . This infrared optical measurement is in excellent agreement with electrical measurements [31] at 145 GHz and with density functional calculations [16].

An even better description of our experimental data (where the mean square deviation is slightly smaller than the experimental errors) is achieved with Lowndes' model in Eq. (4) and the parameters in Table 3. At the scale in Fig. 1, it is not possible to pinpoint the improvements in the fit with the Lowndes model compared to the Lorentz oscillator model.

Therefore, we compare the dielectric functions and loss functions for both models in Figs. 2, 3, 4 and 5. As a reminder, TO phonons appear as peaks in  $\epsilon_2$  and LO phonons appear as peaks in  $\text{Im}(-1/\epsilon)$ . The Lorentz model has the following issues describing the experimental data, which are clearly improved by the Lowndes model: (1) The minimum and maximum of  $\epsilon_1$  between 600 and 700  $\text{cm}^{-1}$  are too shallow. (2) The asymmetry of the peak in  $\epsilon_2$  at 650  $\text{cm}^{-1}$  is not described well. See especially near 600  $\text{cm}^{-1}$ . (3) The minima and maxima in  $\text{Re}(-1/\epsilon)$  and the peaks in  $\text{Im}(-1/\epsilon)$  are too shallow.

The static dielectric constant can also be calculated from Eq. (4) by setting  $\omega = 0$ , leading to the common Lyddane-Sachs-Teller (LST) relation for multiphonon systems

$$\epsilon_0 = \epsilon_\infty \prod_i \frac{\omega_{i,LO}^2}{\omega_{i,TO}^2} = 22.3 \pm 0.3 \quad (6)$$

with parameters and errors taken from Table 3. The largest sources of error are uncertainties in the high-frequency dielectric constant  $\epsilon_\infty$  and the lowest TO phonon energy (which is a fit parameter, but below our spectral range). Since the Lowndes model (4) gives a better description than the Lorentz model (3), we believe that the LST relation (6) gives a more accurate value of  $\epsilon_0$  than Eq. (5), despite the larger error bar.

Finally, we note that our data show no evidence of a weak  $E_u$  phonon peak expected at 300  $\text{cm}^{-1}$ . Our data is very noisy at such long-wavelengths. We do not find any “ghost” peaks at other energies.

### 4. Discussion

Our infrared ellipsometry data yield very accurate phonon parameters for the Lorentz model (see Table 2) and the Lowndes model (see Table 3). We find that the Lowndes model gives a superior description of the experimental results, especially at peaks in the loss function related to LO phonons. This result is not surprising, since the main feature of

**Table 2**

Transverse optical phonon mode parameters (amplitude  $A$ , TO phonon energy  $\omega$ , and TO phonon broadening  $\gamma$ ) for  $\text{LaAlO}_3$  at 300 K using a Lorentz oscillator fit as in Eq. (3). 90% confidence limits for experimental results are given in parentheses. The additional parameter  $\epsilon_\infty = 4.12 \pm 0.01$  was fixed based on our earlier work [21].

Mode	$A_i$ (Exp)	$\omega_i$ (Exp)	$\gamma_i$ (Exp)
	(1)	$\text{cm}^{-1}$	$\text{cm}^{-1}$
AlO-La	15.24(5)	182(f)	4(f)
O bend	4.121(7)	426.94(6)	3.7(1)
Al	0.008(1)	495.8(3)	3.8(7)
O stretch	0.285(1)	652.9(1)	21.3(1)
O stretch	0.031(1)	688.6(3)	31.4(5)

**Table 3**

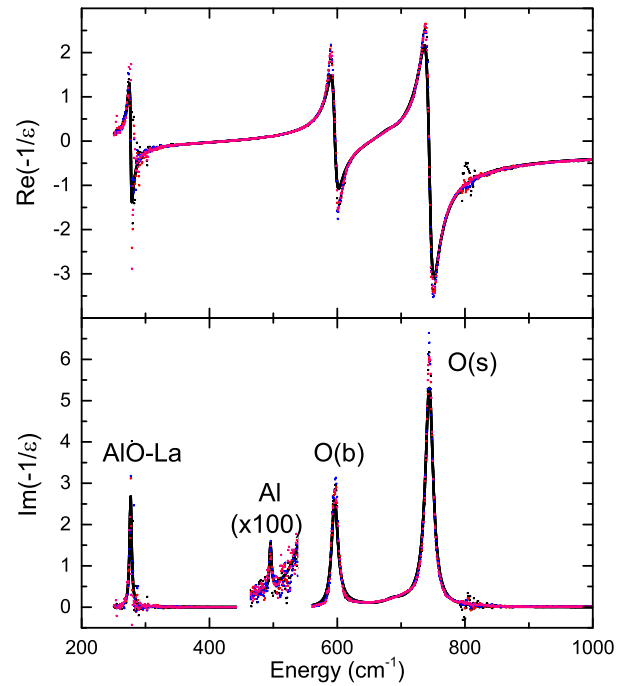
Transverse (TO) and longitudinal (LO) optical phonon energies and broadenings for  $\text{LaAlO}_3$  at 300 K derived from a Lowndes oscillator fit as in Eq. (4), in units of  $\text{cm}^{-1}$ . 90% confidence limits are given in parentheses. The additional parameter  $\epsilon_\infty = 4.12 \pm 0.01$  was fixed based on our earlier work [21].

mode	$\omega_{i,\text{TO}}$	$\gamma_{i,\text{TO}}$	$\omega_{i,\text{LO}}$	$\gamma_{i,\text{LO}}$
AlO-La	188(1)	0.4(1)	276.4(2)	3.7(7)
O bend	427.0(1)	5.0(1)	596.1(7)	7.2(1)
Al	495.72(1)	3.8(7)	495.5(3)	3.8(7)
O stretch	650.79(5)	22.5(7)	744.1(9)	12.1(1)
O stretch	708.2(9)	55.3(9)	702.2(9)	66(1)

the Lowndes model is the assignment of independent broadening parameters to the LO phonons, while the Lorentz model describes both TO and LO phonons with one broadening parameter.

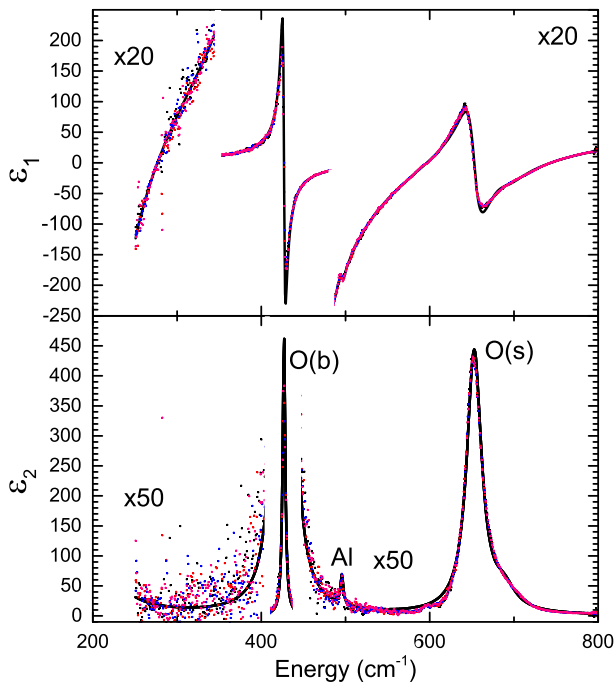
For isolated strong phonon modes, both models yield nearly identical TO phonon energies (compare Tables 2 and 3). However, if two TO phonon peaks overlap, then the two models find significantly differently TO phonon energies. This is true especially for the weak  $E_u$  oxygen stretch mode at  $709 \text{ cm}^{-1}$ . In general, the TO phonon energy from our model is lower than the LO phonon energy (as expected), but there is an exception for the very weak phonon located at  $497 \text{ cm}^{-1}$ , where the LO–TO separation is small and negative. In our analysis, the negative Coulomb splitting for the Al mode is only  $0.2 \text{ cm}^{-1}$ , much smaller than the resolution of our spectra ( $2 \text{ cm}^{-1}$ ). It is possible that a good fit to our raw data could also be found with a positive Coulomb splitting for the Al mode at  $496 \text{ cm}^{-1}$ . The weak phonon mode near  $705 \text{ cm}^{-1}$  is also inverted, but the splitting is ten times smaller than its phonon broadening. There has been some discussion in the literature about the sign of the TO/LO phonon splitting in quartz and complex metal oxides, which we did not find convincing [33–35].

Based on an argument presented by Lowndes [29] and a generalized condition for multi-phonon systems in Ref. [36], we expect that the LO broadening should be larger than the TO broadening for a material like GaAs with a single TO–LO phonon pair. This result is confirmed for most phonon modes, except for the anomalously small LO phonon broadening of the LO phonon with the largest energy. The broadenings

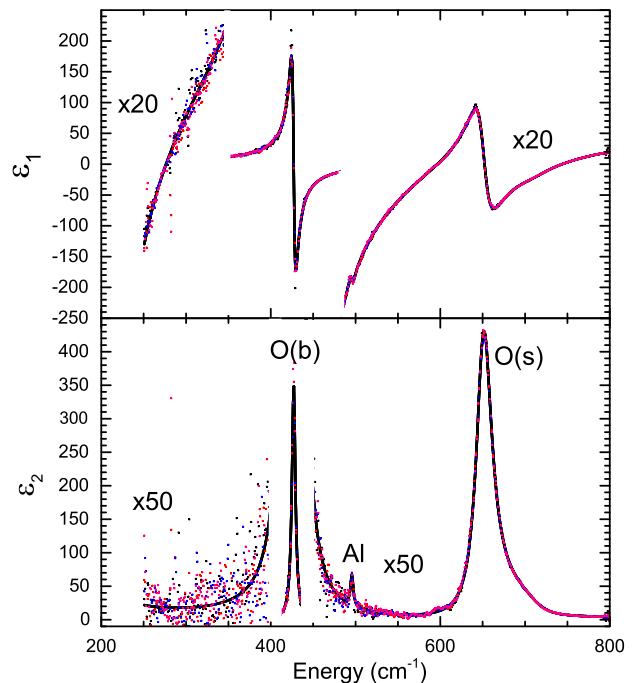


**Fig. 3.** (Color online) Loss function  $-1/\epsilon$  for  $\text{LaAlO}_3$  at 300 K calculated from the data in Fig. 1 using Eq. (2). Symbols: experimental data. Lines: fit with Eq. (3). Peaks corresponding to longitudinal optical phonons are found at 277, 596, 703, and  $744 \text{ cm}^{-1}$ .

of the two highest-energy phonons can also be found by visual inspection of  $\epsilon_2$  and  $\text{Im}(-1/\epsilon)$  in Figs. 4 and 5. The  $\epsilon_2$  spectrum shows a sharper TO phonon at a lower energy (near  $650 \text{ cm}^{-1}$ ) and a much broader TO phonon at a higher energy. Similarly, the loss function  $\text{Im}(-1/\epsilon)$  shows a broad TO phonon at about  $700 \text{ cm}^{-1}$  accompanied by a sharper LO phonon at a higher energy. Therefore, the broadenings of the TO and LO phonons at the two highest energies are a direct experimental observation, not a numerical artifact.

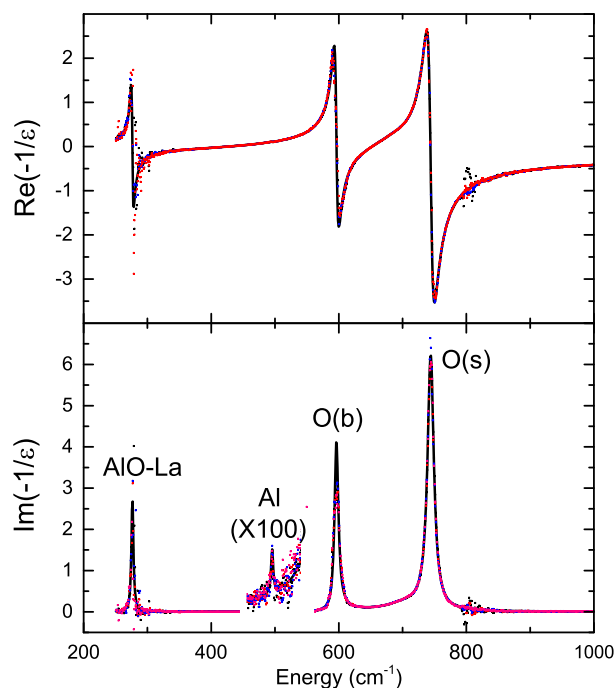


**Fig. 2.** (Color online) Pseudodielectric function for  $\text{LaAlO}_3$  at 300 K calculated from the data in Fig. 1 using Eq. (2). Symbols: experimental data. Lines: fit with Eq. (3). Peaks corresponding to transverse optical phonons are found at 427, 496, 653, and  $689 \text{ cm}^{-1}$ .



**Fig. 4.** (Color online) Pseudodielectric function for  $\text{LaAlO}_3$  at 300 K calculated from the data in Fig. 1 using Eq. (2). Symbols: experimental data. Lines: fit with Eq. (4).





**Fig. 5.** (Color online) Loss function  $-1/\epsilon$  for  $\text{LaAlO}_3$  at 300 K calculated from the data in Fig. 1 using Eq. (2). Symbols: experimental data. Lines: fit with Eq. (4).

**Table 4**

Comparison of our experimental FTIR ellipsometry results (last column) with theory and previous FTIR reflectance measurements. TO and LO phonon energies are in units of  $\text{cm}^{-1}$ . 90% confidence limits are given in parentheses.

Mode	Theory [14]	Theory [16]	Reflectance [17]	Ellipsometry (this work)
	TO/LO	TO	TO	TO/LO
$A_{2u}$	213/263	168	182	188(1)/276.4(2)
$E_u$	220/263	179		
$E_u$	270/270	297		
$A_{2u}$	366/496	409	429	427.0(1)/596.1(7)
$E_u$	371/475	411		
$E_u$	481/505	478	501	495.72(1)/495.5(3)
$A_{2u}$	706/712	627	657	650.79(5)/744.1(9)
$E_u$	707/712	637	695	708.2(9)/702.2(9)

A comparison of our FTIR ellipsometry results with previous FTIR reflectance measurements and theoretical results is given in Table 4.

## 5. Conclusions

We have characterized the parameters of the five IR-active optical phonons of  $\text{LaAlO}_3$ . We were able to find the energies, amplitudes, and broadenings not only of the reasonably well-documented TO phonons, but also of the LO phonons which have not previously been determined experimentally. Combining our Lorentz amplitudes with the  $\epsilon_\infty$  previously determined by Nelson et al. [21], we have found an experimental dielectric constant of about 22–24. This agrees well with published values for  $\text{LaAlO}_3$  which range from 23 to 25 [32]. By characterizing

both the TO and LO phonons, we have a more complete picture of the lattice dynamics of our material. This is vital when considering the possible applications of  $\text{LaAlO}_3$  in microelectronics.

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