## Abelian Duality

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## Contents

 $\Sigma$  a closed Riemann Surface with metric g. We have fields valued in  $S^1$ ,  $\phi: \Sigma \to S^1$ . The Lagrangian  $\mathcal{L}(\phi) := R^2/(4\pi) \int_{\Sigma} d\phi \wedge d * \phi$ . Where R is the radius of the circle  $S^1$ , which I think about as a connection.

Plan of todays talk: First we will take out the gauge invariance, then we introduce new fields in a bigger theory, integrate out the new field to recover the old theory, then integrate out the old theory to get a new theory, which is the dual.

## 1 setup

Let  $\phi$  be our field, we are going to couple it to electromagnetism:  $A \in mathcal A \cong \Omega^1_{\Sigma}(\mathbb{R})$ , then we have the covariant derivative  $D_A : \phi \mapsto d\phi + A$ .

I introduce another field  $\sigma: \Sigma \mapsto S^1$ . So we have introduce U(1) connection together with another field. Notes that  $d\sigma$  is well defined globally.

The new Langrangian is:

$$\mathcal{L}(A,\phi,\sigma) = R^2/(4\phi) \int_{\Sigma} D_A \phi \wedge *D_A \phi - i/(2\pi) \int_{Siama} \sigma \wedge F_A$$

, the last part is equal to  $i/(2\pi) \int_{\Sigma} d\sigma \wedge A$ .

The Partition function is

$$Z = 1/vol(\mathcal{G}) \int D\phi DAD\sigma \exp \mathcal{L}(A, \phi, \sigma)$$

We are going to show that this is equivalent to the old theory. Notice that  $\sigma$  looks like a Lagrange multiplier term, introduce to be integrate out to give a delta function of the connection at the trivial connection.

Choose  $\sigma_n : \Sigma \to S^1$  s.t.  $d\sigma_n \in [d\sigma] \subset H^1(\sigma; \mathbb{R})$ , and  $\sigma_n(P) = 0 \mod 2\pi\mathbb{Z}$ . We see that  $\sigma = \sigma_n + \sigma_{\mathbb{R}}$ , where  $\sigma_{\mathbb{R}} \in \Omega^0_{\Sigma}(\mathbb{R})$ . Note that  $1/(2\pi) \int_C d\sigma = 1/(2\pi) \int_C d\sigma_n \in \mathbb{Z}$ .

Choose  $(\gamma_j)$  basis for  $H^1(\Sigma; \mathbb{Z})$ , then we have that  $1/(2\pi)d\sigma_n = \sum_j m_j \gamma_j$ , then if we integrate out  $\sigma$ , we get

$$\int D\sigma \exp{-i/2\pi} \int (\sigma wedgeF_A) = \int D\sigma \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})} \exp{-i/2\pi} \int (\sigma_n + \sigma_{\mathbb{R}} wedgeF_A) = \sum_{[d\sigma_k] \in H^1(\Sigma; 2\pi\mathbb{Z})$$

 $\int D\sigma_{\mathbb{R}} \exp -i/2\pi \int \sigma_{\mathbb{R}} \wedge F_A(1)$ Note that the first piece become

$$\prod_{j} \sum_{m_j \in \mathbb{Z}} \exp^{j \int \gamma_j \wedge A}$$

Where  $\int \gamma_j \wedge A$  is the holonomy of A around  $C_j$ , where  $C_j$  is the corresponding curve to the harmonic 1-form. Notes that  $\prod_j \sum_{m_j} \exp(-im_j h_j) = \prod_j \delta(hol_{C_j}(A)) = 0$ 

The second part, we are just going to integrate  $\sigma_{\mathbb{R}}$  out and we get  $\delta F_A = 0$ , aka the curvature is zero.

So we see that the connection can have no curvature, nor holonomy. Thus it is a trivial (flat) connection.

Gauge fix and integrate out the gauge symmetry we see that A=0, and we recover the old theory.

## 2 Integrate out $\phi$ and A

Gauge fix: we can gauge fix such that  $\phi = 0$ . And the  $1\{vol\mathcal{G} \text{ goes away.}\}$ So we are left with

$$Z = \int DAD\sigma \exp{-R^2/4\pi} \int_{\Sigma} (A \wedge Awedge * A - i/2\pi) \int_{\Sigma} A \wedge d\sigma$$

Now we want to complete the square:

$$A' = A = i/R^2 * d\sigma = \int DA' D\sigma \exp(-R^2/2\pi \int |A|^2) \exp(-1/4\pi R^2 \int |d\sigma|^2))$$

We see that the out theory went inverted the radius. This is the duality that Leon discussed last time.

Another question we can ask is where did the operators go?

 $d\phi(x) \mapsto D_A\phi(x) \mapsto ?$ , where did it go? We need to integrate out  $\phi$  and A.

The vacuum expectation value of  $D_A\phi(x)$  is

$$\int DAD\phi D\sigma D_A\phi(X) exp(-S)$$

, when we integrate out  $D\phi$  and  $D\sigma$ , we see that  $D_A\phi(x)$  becomes  $d\phi(x)$ . Now let's see where  $D_A\phi(x)$  lands on the other side?

It is not very hard to show that it becomes  $-i/R^2*d\sigma$ . Basically set  $\phi = 0$  and  $A' = A + i/R^2*d\sigma = 0$  when we integrate it out.

Another operator  $exp(i\phi)(p)$