

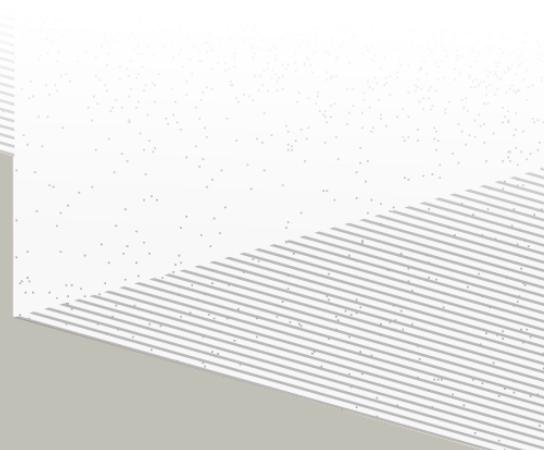
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Par

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**Precision measurement of solar neutrino oscillation parameters
with the JUNO small PMTs system and test of the unitarity of the
PMNS matrix**

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⁶⁵ Remerciements

66 Introduction

⁶⁷ **Chapter 1**

⁶⁸ **Neutrino physics**

⁶⁹ *The neutrino, or ν for the close friends, a fascinating and invisible particle. Some will say that dark matter also have those property but at least we are pretty confident that neutrinos exists.*

⁷⁰ **1.1 Standard model**

⁷¹ **1.1.1 Limits of the standard model**

⁷² **1.2 Historic of the neutrino**

⁷³ **First theories**

⁷⁴ **Discovery**

⁷⁵ **Milestones and anomalies**

⁷⁶ **1.3 Oscillation**

⁷⁷ **1.3.1 Phenomologies**

⁷⁸ **1.4 Open questions**

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⁷⁹ **Chapter 2**

⁸⁰ **The JUNO experiment**

⁸¹ “*Ave Juno, rosae rosam, et spiritus rex*”. It means nothing but I found it in tone.

⁸² The first idea of a medium baseline (~ 52 km) experiment, was explored in 2008 [1] where it was
⁸³ demonstrated that the Neutrino Mass Ordering (NMO) could be determined by a medium baseline
⁸⁴ experiment if $\sin^2(2\theta_{13}) > 0.005$ without the requirements of accurate knowledge of the reactor
⁸⁵ antineutrino spectra and the value of Δm_{32}^2 . From this idea is born the Jiangmen Underground
⁸⁶ Neutrino Observatory (JUNO) experiment.

⁸⁷ JUNO is a neutrino detection experiment under construction located in China, in Guangdong prov-
⁸⁸ ing, near the city of Kaiping. Its main objectives are the determination of the mass ordering at the
⁸⁹ $3\text{-}4\sigma$ level in 6 years of data taking and the measurement at the sub-percent precision of the oscillation
⁹⁰ parameters Δm_{21}^2 , $\sin^2 \theta_{12}$, Δm_{32}^2 and with less precision $\sin^2 \theta_{13}$ [2].



FIGURE 2.1 – On the left: Location of the JUNO experiment and its reactor sources in southern China. On the right: Aerial view of the experimental site

⁹¹ For this JUNO will measure the electronic anti-neutrinos ($\bar{\nu}_e$) flux coming from the nuclear reactors
⁹² of Taishan, Yangjiang, for a total power of 26.6 GW_{th} , and the Daya Bay power plant to a lesser
⁹³ extent. All of those cores are the second-generation pressurized water reactors CPR1000, which is a
⁹⁴ derivative of Framatome M310. Details about the power plants characteristics and their expected flux
⁹⁵ of $\bar{\nu}_e$ can be found in the table 2.1. The distance of 53 km has been specifically chosen to maximize
⁹⁶ the disappearance probability of the $\bar{\nu}_e$. The data taking is scheduled to start early 2025.

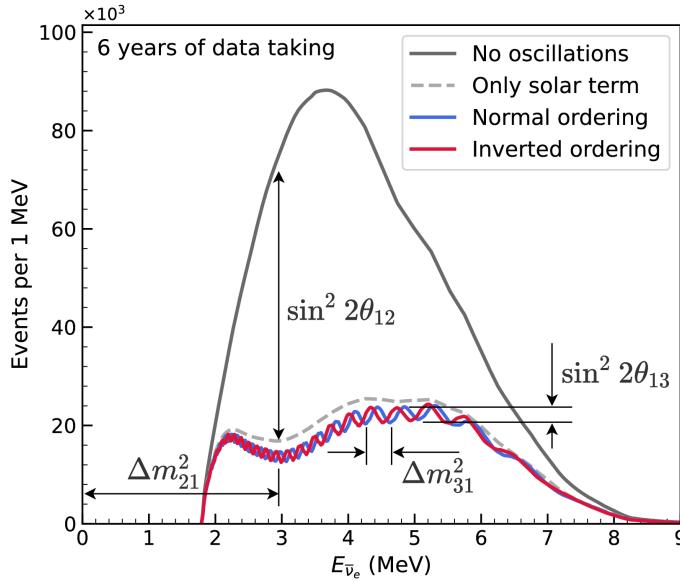


FIGURE 2.2 – Expected number of neutrinos event per MeV in JUNO after 6 years of data taking. The black curve shows the flux if there was no oscillation. The light gray curve shows the oscillation if only the solar terms are taken in account (θ_{12} , Δm_{21}^2). The blue and red curve shows the spectrum in the case of, respectively, NO and IO. The dependency of the oscillation to the different parameters are schematized by the double sided arrows. We can see the NMO sensitivity by looking at the fine phase shift between the red and the blue curve.

97 2.1 Neutrinos physics in JUNO

98 Even if the JUNO design detailed in section 2.2 was optimized for the measurement of the NMO, its
 99 large detection volume, excellent energy resolution and background level and understanding make it
 100 also an excellent detector to measure the flux coming from other neutrino sources. Thus the scientific
 101 program of JUNO extends way over reactor antineutrinos. The following section is an overview of
 102 the different physics topic JUNO will contribute in the coming years.

103 2.1.1 Reactor neutrino oscillation for NMO and precise measurements

Previous works [1, 3] shows that oscillation parameters and the NMO can be observed by looking at the $\bar{\nu}_e$ disappearance energy spectrum coming from medium baseline nuclear reactor. This disappearance probability can be expressed as [2] :

$$P(\bar{\nu}_e \rightarrow \bar{\nu}_e) = 1 - \sin^2 2\theta_{12} c_{13}^4 \sin^2 \frac{\Delta m_{21}^2 L}{4E} - \sin^2 2\theta_{13} \left[c_{12}^2 \sin^2 \frac{\Delta m_{31}^2 L}{4E} + s_{12}^2 \sin^2 \frac{\Delta m_{32}^2 L}{4E} \right]$$

104 Where $s_{ij} = \sin \theta_{ij}$, $c_{ij} = \cos \theta_{ij}$, E is the $\bar{\nu}_e$ energy and L is the baseline. We can see the sensitivity
 105 to the NMO in the dependency to Δm_{32}^2 and Δm_{31}^2 causing a phase shift of the spectrum as we can
 106 see in the figure 2.2. By carefully adjusting a theoretical spectrum to the data, one can extract the
 107 NMO and the oscillation parameters. The statistic procedure used to adjust the theoretical spectrum
 108 is reviewed in more details in the section 2.7. To reach the desired sensitivity, JUNO must meet
 109 multiple requirements but most notably:

- 110 1. An energy resolution of $3\%/\sqrt{E(\text{MeV})}$ to be able to distinguish the fine structure of the fast
oscillation.
- 111 2. An energy precision of 1% in order to not err on the location of the oscillation pattern.
- 112 3. A baseline of 53 ± 0.5 km to maximise the $\bar{\nu}_e$ oscillation probability.
- 113 4. At least $\approx 100,000$ events to limit the spectrum distortion due to statistical uncertainties.
- 114

115 **$\bar{\nu}_e$ flux coming from nuclear power plants**

116 To get such high measurements precision, it is necessary to have a very good understanding of the
117 sources characteristics. For its NMO and precise measurement studies, JUNO will observe the energy
118 spectrum of neutrinos coming from the nuclear power plants Taishan and Yangjiang's cores, located
119 at 53 km of the detector to maximise the disappearance probability of the $\bar{\nu}_e$.

Reactor	Power (GW _{th})	Baseline (km)	IBD Rate (day ⁻¹)	Relative Flux (%)
Taishan	9.2	52.71	15.1	32.1
Core 1	4.6	52.77	7.5	16.0
Core 2	4.6	52.64	7.6	16.1
Yangjiang	17.4	52.46	29.0	61.5
Core 1	2.9	52.74	4.8	10.1
Core 2	2.9	52.82	4.7	10.1
Core 3	2.9	52.41	4.8	10.3
Core 4	2.9	52.49	4.8	10.2
Core 5	2.9	52.11	4.9	10.4
Core 6	2.9	52.19	4.9	10.4
Daya Bay	17.4	215	3.0	6.4

TABLE 2.1 – Characteristics of the nuclear power plants observed by JUNO. The IBD rate are estimated from the baselines, the reactors full thermal power, selection efficiency and the current knowledge of the oscillation parameters

120 The $\bar{\nu}_e$ coming from reactors are emitted from β -decay of unstable fission fragments. The Taishan
121 and Yangjiang reactors are Pressurised Water Reactor (PWR), the same type as Daya Bay. In those
122 type of reactor more the 99.7 % and $\bar{\nu}_e$ are produced by the fissions of four fuel isotopes ^{235}U , ^{238}U ,
123 ^{239}Pu and ^{241}Pu . The neutrino flux per fission of each isotope is determined by the inversion of the
124 measured β spectra of fission product [4–8] or by calculation using the nuclear databases [9, 10].

125 The neutrino flux coming from a reactor at a time t can be predicted using

$$\phi(E_\nu, t)_r = \frac{W_{th}(t)}{\sum_i f_i(t) e_i} \sum_i f_i(t) S_i(E_\nu) \quad (2.1)$$

126 where $W_{th}(t)$ is the thermal power of the reactor, $f_i(t)$ is the fraction fission of the i th isotope, e_i its
127 thermal energy released in each fission and $S_i(e_\nu)$ the neutrino flux per fission for this isotope. Using
128 this method, the flux uncertainty is expected to be of an order of 2-3 % [11].

129 In addition to those prediction, a satellite experiment named TAO[12] will be setup near the reactor
130 core Taishan-1 to measure with an energy resolution of 2% at 1 MeV the neutrino flux coming from
131 the core, more details can be found in section 2.4.1. It will help identifying unknown fine structure
132 and give more insight on the $\bar{\nu}_e$ flux coming from this reactor.

133 One the open issue about reactor anti-neutrinos flux is the so-called neutrino anomaly [13], an
134 unexpected surplus of neutrino emission in the spectra around 5 MeV. Multiples scientists are trying
135 to explain this surplus by advanced recalculation of the nuclei model during beta decay [14, 15] but
136 no consensus on this issue has been reached yet.

¹³⁷ **Background in the neutrinos reactor spectrum**

¹³⁸ Considering the close reactor neutrinos flux as the main signal, the signals that are considered as
¹³⁹ background are:

- The geoneutrinos producing background in the $0.511 \sim 2.7$ MeV region.
- The neutrinos coming from the other nuclear reactors around Earth.

¹⁴⁰ In addition to all those physics signal, non-neutrinos signal that would mimic an IBD will also be
¹⁴¹ present. It is composed of:

- The signal coming from radioactive decay (α , γ , β) from natural radioactive isotopes in the material of the detector.
- Cosmogenic event such as fast neutrons and activated isotopes induced by muons passing through the detector, most notably the spallation on ^{12}C .

¹⁴² All those events represent a non-negligable part of the spectrum as shown in figure 2.3.

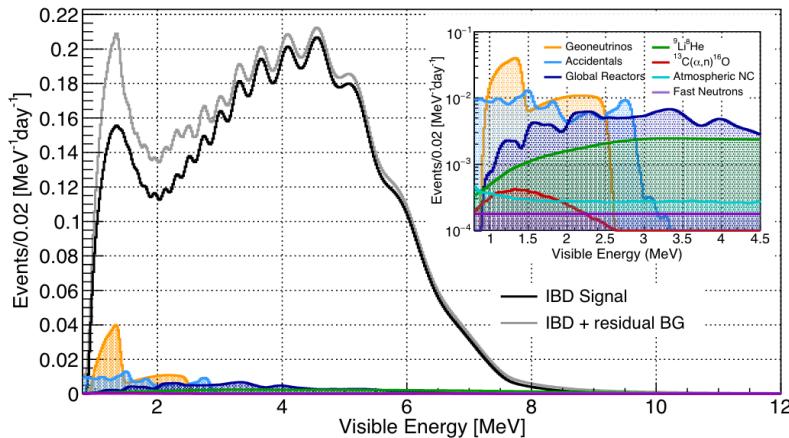


FIGURE 2.3 – Expected visible energy spectrum measured with the LPMT system with (grey) and without (black) backgrounds. The background amount for about 7% of the IBD candidate and are mostly localized below 3 MeV [11]

¹⁴⁹ **Identification of the mass ordering**

¹⁵⁰ To identify the mass ordering, we adjust the theoretical neutrino energy spectrum under the two hypothesis of NO and IO. Those give us two χ^2 , respectively χ^2_{NO} and χ^2_{IO} . By computing the difference $\Delta\chi^2 = \chi^2_{\text{NO}} - \chi^2_{\text{IO}}$ we can determine the most probable mass ordering and the confidence interval: NO if $\Delta\chi^2 > 0$ and IO if $\Delta\chi^2 < 0$. Current studies shows that the expected sensitivity the mass ordering would be of 3.4σ after 6 years of data taking in nominal setup[2]. More detailed explanations about the procedure can be found in the section 2.7.

¹⁵⁶ **Precise measurement of the oscillations parameters**

¹⁵⁷ The oscillations parameters θ_{12} , θ_{13} , Δm_{21}^2 , Δm_{31}^2 are free parameters in the fit of the oscillation spectrum. The precision on those parameters have been estimated and are shown in table 2.2. Wee see that for θ_{12} , Δm_{21}^2 , Δm_{31}^2 , precision at 6 years is better than the reference precision by an order of magnitude [11]

	Central Value	PDG 2020	100 days	6 years	20 years
$\Delta m_{31}^2 (\times 10^{-3} \text{ eV}^2)$	2.5283	± 0.034 (1.3%)	± 0.021 (0.8%)	± 0.0047 (0.2%)	± 0.0029 (0.1%)
$\Delta m_{21}^2 (\times 10^{-3} \text{ eV}^2)$	7.53	± 0.18 (2.4%)	± 0.074 (1.0%)	± 0.024 (0.3%)	± 0.017 (0.2%)
$\sin^2 \theta_{12}$	0.307	± 0.013 (4.2%)	± 0.0058 (1.9%)	± 0.0016 (0.5%)	± 0.0010 (0.3%)
$\sin^2 \theta_{13}$	0.0218	± 0.0007 (3.2%)	± 0.010 (47.9%)	± 0.0026 (12.1%)	± 0.0016 (7.3%)

TABLE 2.2 – A summary of precision levels for the oscillation parameters. The reference value (PDG 2020 [16]) is compared with 100 days, 6 years and 20 years of JUNO data taking.

2.1.2 Other physics

While the design of JUNO is tailored to measure $\bar{\nu}_e$ coming from nuclear reactor, JUNO will be able to detect neutrinos coming from other sources thus allowing for a wide range of physics studies as detailed in the table 2.3 and in the following sub-sections.

Research	Expected signal	Energy region	Major backgrounds
Reactor antineutrino	60 IBDs/day	0–12 MeV	Radioactivity, cosmic muon
Supernova burst	5000 IBDs at 10 kpc	0–80 MeV	Negligible
DSNB (w/o PSD)	2300 elastic scattering		
Solar neutrino	2–4 IBDs/year	10–40 MeV	Atmospheric ν
Atmospheric neutrino	hundreds per year for ${}^8\text{B}$	0–16 MeV	Radioactivity
Geoneutrino	hundreds per year	0.1–100 GeV	Negligible
	≈ 400 per year	0–3 MeV	Reactor ν

TABLE 2.3 – Detectable neutrino signal in JUNO and the expected signal rates and major background sources

Geoneutrinos

Geoneutrinos designate the antineutrinos coming from the decay of long-lived radioactive elements inside the Earth. The 1.8 MeV threshold necessary for the IBD makes it possible to measure geoneutrinos from ${}^{238}\text{U}$ and ${}^{232}\text{Th}$ decay chains. The studies of geoneutrinos can help refine the Earth crust models but is also necessary to characterise their signal, as they are a background to the mass ordering and oscillations parameters studies.

Atmospheric neutrinos

Atmospheric neutrinos are neutrinos originating from the decay of π and K particles that are produced in extensive air showers initiated by the interactions of cosmic rays with the Earth atmosphere. Earth is mostly transparent to neutrinos below the PeV energy, thus JUNO will be able to see neutrinos coming from all directions. Their baseline range is large (15km \sim 13000km), they can have energy between 0.1 GeV and 10 TeV and will contain all neutrino and antineutrinos flavour. Their studies is complementary to the reactor antineutrinos and can help refine the constraints on the NMO [2].

Supernovae burst neutrinos

Neutrinos are crucial component during all stages of stellar collapse and explosion. Detection of neutrinos coming for core collapse supernovae will provide us important informations on the mech-

182 mechanisms at play in those events. Thanks to its 20 kt sensible volume, JUNO has excellent capabilities
 183 to detect all flavour of the $\mathcal{O}(10 \text{ MeV})$ postshock neutrinos, and using neutrinos of the $\mathcal{O}(1 \text{ MeV})$
 184 will give informations about the pre-supernovae neutrinos. All those informations will allow to
 185 disentangle between the multiple hydro-dynamic models that are currently used to describe the
 186 different stage of core-collapse supernovae.

187 Diffuse supernovae neutrinos background

188 Core-collapse supernovae in our galaxy are rare events, but they frequently occur throughout the
 189 visible Universe sending burst of neutrinos in direction of the Earth. All those events contributes to
 190 a low background flux of low-energy neutrinos called the Diffuse Supernovae Neutrino Background
 191 (DSNB). Its flux and spectrum contains informations about the red-shift dependent supernovae rate,
 192 the average supernovae neutrino energy and the fraction of black-hole formation in core-collapse su-
 193 pernovae. Depending of the DSNB model, we can expect 2-4 IBD events per year in the energy range
 194 above the reactor $\bar{\nu}_e$ signal, which is competitive with the current Super-Kamiokande+Gadolinium
 195 phase [17].

196 Beyond standard model neutrinos interactions

197 JUNO will also be able to probe for beyond standard model neutrinos interactions. After the main
 198 physics topics have been accomplished, JUNO could be upgraded to probe for neutrinoless beta
 199 decay ($0\nu\beta\beta$). The detection of such event would give critical informations about the nature of
 200 neutrinos, is it a majorana or a dirac particle. JUNO will also be able to probe for neutrinos that
 201 would come for the decay or annihilation of Dark Matter inside the sun and neutrinos from putative
 202 primordial black hole. Through the unitary test of the mixing matrix, JUNO will be able to search
 203 for light sterile neutrinos. Thanks to JUNO sensitivity, multiple other exotic can be performed on
 204 neutrino related beyond standard model interactions.

205 2.2 The JUNO detector

206 The JUNO detector is a scintillator detector buried 693.35 meters under the ground (1800 meters
 207 water equivalent). It consist of Central Detector (CD), a water pool and a Top Tracker (TT) as showed
 208 in figure 2.4a. The CD is an acrylic vessel containing the 20 ktons of Liquid Scintillator (LS). It is
 209 supported by a stainless steel structure and is immersed in that water pool that is used as shielding
 210 from external radiation and as a cherenkov detector for the background. The top of the experiment
 211 is partially covered by the Top Tracker (TT), a plastic scintillator detector which is use to detect the
 212 atmospheric muons background and is acting as a veto detector.

213 The top of the experiment also host the LS purification system, a water purification system, a ven-
 214 tilation system to get rid of the potential radon in the air. The CD is observed by two system of
 215 Photo-Multipliers Tubes (PMT). They are attached to the steel structure and their electronic readout
 216 is submersed near them. A third system of PMT is also installed on the structure but are facing
 217 outward of the CD, instrumenting the water to be cherenkov detector. The CD and the cherenkov
 218 detector are optically separated by Tyvek sheet. A chimney for LS filling and purification and for
 219 calibration operations connects the CD to the experimental hall from the top.

220 The CD has been dimensioned to meet the requirements presented in section 2.1.1:

- 221 — Its 20 ktons monolithic LS provide a volume sizeable enough, in combination with the ex-
 222 pected $\bar{\nu}_e$ flux, to reach the desired statistic in 6 years. Its monolithic nature also allow for a
 223 full containment of most of the events, preventing the energy loss in non-instrumented parts
 224 that would arise from a segmented detector.

- 225 — Its large overburden shield it from most of the atmospheric background that would pollute
 226 the signal.
 227 — The localization of the experiment, chosen to maximize the disappearance with a 53km base-
 228 line and in a region that allow two nuclear power plant to be used as sources.

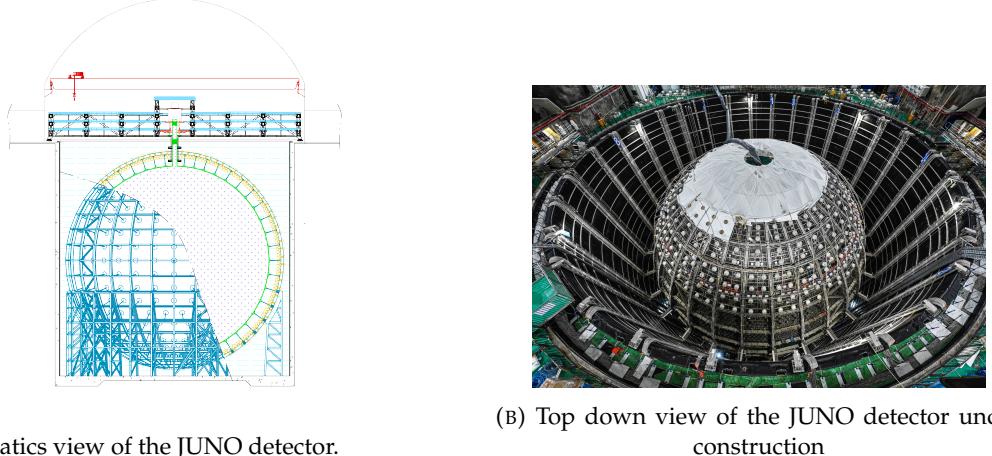


FIGURE 2.4

229 This section cover in details the different components of the detector and the detection systems.

230 2.2.1 Detection principle

The CD will detect the neutrino and measure their energy mainly via an Inverse Beta Decay (IBD) interaction with proton mainly from the ^{12}C and H nucleus in the LS:

$$\bar{\nu}_e + p \rightarrow n + e^+$$

231 Kinematics calculation shows that this interaction has an energy threshold for the $\bar{\nu}_e$ of $(m_n + m_e -$
 232 $m_p) \approx 1.806$ MeV [18] where m_λ is the mass of the λ particle. This threshold make the experiment
 233 blind to very low energy neutrinos. The residual energy $E_\nu - 1.806$ MeV is be distributed as kinetic
 234 energy between the positron and the neutron. The energy of the emitted positron E_e is given by [18]

$$E_e = \frac{(E_\nu - \delta)(1 + \epsilon_\nu) + \epsilon_\nu \cos \theta \sqrt{(E_\nu - \delta)^2 + \kappa m_e^2}}{\kappa} \quad (2.2)$$

235 where $\kappa = (1 + \epsilon_\nu)^2 - \epsilon_\nu^2 \cos^2 \theta \approx 1$, $\epsilon_\nu = \frac{E_\nu}{m_p} \ll 1$ and $\delta = \frac{m_n^2 - m_p^2 - m_e^2}{2m_p} \ll 1$. We can see from this
 236 equation that the positron energy is strongly correlated to the neutrino energy.

237 The positron and the neutron will then propagate in the detection medium, the Liquid Scintillator
 238 (LS), loosing their kinetic energy by exciting the molecule of the LS (more details in section 2.2.2).
 239 Once stopped, the positron will annihilate with an electron from the medium producing two 511
 240 KeV gamma. Those gamma will themselves interact with the LS, exciting it before being absorbed
 241 by photoelectrical effect. The neutron will be captured by an hydrogen, emitting a 2.2 MeV gamma
 242 in the process. This gamma will also deposit its energy before being absorbed by the LS.

243 The scintillation photons have frequency in the UV and will propagate in the LS, being re-absorbed
 244 and re-emitted by compton effect before finally be captured by PMTs instrumenting the acrylic
 245 sphere. The analog signal of the PMTs digitized by the electronic is the signal of our experiment.

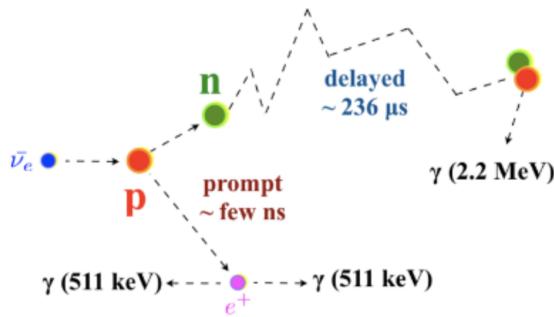


FIGURE 2.5 – Schematics of an IBD interaction in the central detector of JUNO

246 The signal produced by the positron is subsequently called the prompt signal, and the signal coming
 247 from the neutron the delayed signal. This naming convention come from the fact that the positron
 248 will deposit its energy rather quickly (few ns) where the neutron will take a bit more time ($\sim 236 \mu\text{s}$).

249 2.2.2 Central Detector (CD)

250 The central detector, composed of 20 ktons of Liquid Scintillator (LS), is the main part of JUNO. The
 251 LS is contained in a spherical acrylic vessel supported by a stainless steel structure. The CD and
 252 its structural support are submerged in a cylindrical water pool of 43.5m diameter and 44m height.
 253 We're confident that the water pool provide sufficient buffer protection in every direction against the
 254 rock radioactivity.

255 Acrylic vessel

256 The acrylic vessel is a spherical vessel of inner diameter of 35.4 m and a thickness of 120 mm. It is
 257 assembled from 265 acrylic panels, thermo bonded together. The acrylic recipes has been carefully
 258 tuned with extensive R&D to ensure it does not include plasticizer and anti-UV material that would
 259 stop the scintillation photons. Those panels requires to be pure of radioactive materials to not
 260 cause background. Current setup where the acrylic panels are molded in cleanrooms of class 10000,
 261 let us reach a uranium and thorium contamination of <0.5 ppt. The molding and thermoforming
 262 processes is optimized to increase the assemblage transparency in water to >96%. The acrylic vessel
 263 is supported by a stainless steel structure via supporting node (fig 2.6). The structure and the nodes
 264 are designed to be resilient to natural catastrophic events such as earthquake and can support many
 265 times the effective load of the acrylic vessel.

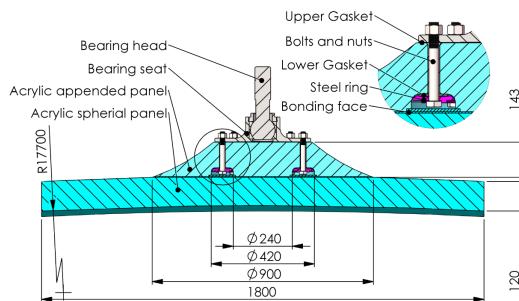


FIGURE 2.6 – Schematics of the supporting node for the acrylic vessel

266 **Liquid scintillator**

267 The Liquid Scintillator (LS) has a similar recipe as the one used in Daya Bay [19] but without gadolinium
 268 doping. It is made of three components, necessary to shift the wavelength of emitted photons to
 269 prevent their reabsorption:

- 270 1. The detection medium, the *linear alkylbenzene* (LAB). Selected because of its excellent transparency,
 271 high flash point, low chemical reactivity and good light yield. Accounting for \sim 98% of the LS, it is the main component with which ionizing particles and gamma interact.
 272 Charged particles will collide with its electronic cloud transferring energy to the molecules,
 273 gamma will interact via compton effect with the electronic cloud before finally be absorbed
 274 via photoelectric effect.
- 276 2. The second component of the LS is the *2,5-diphenyloxazole* (PPO). A fraction of the excitation
 277 energy of the LAB is transferred to the PPO, mainly via non radiative process [20]. The
 278 PPO molecules de-excites in the same way, transferring their energy to the bis-MSB. The PPO
 279 makes for 1.5 % of the LS.
- 280 3. The last component is the *p-bis(o-methylstyryl)-benzene* (bis-MSB). Once excited by the PPO, it
 281 will emit photon with an average wavelength of \sim 430 nm (full spectrum in figure 2.7) that
 282 can be detected by our photo-multipliers systems. It amount for \sim 0.5% of the LS.

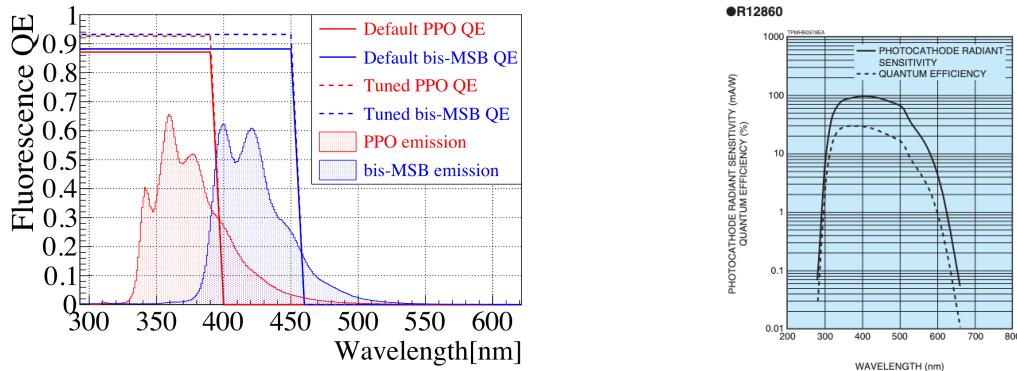


FIGURE 2.7 – On the left: Quantum efficiency (QE) and emission spectrum of the LAB and the bis-MSB [19]. On the right: Sensitivity of the Hamamatsu LPMT depending on the wavelength of the incident photons [21].

283 This formula has been optimized using dedicated studies with a Daya Bay detector [19, 22] to reach
 284 the requirements for the JUNO experiment:

- 285 — A light yield / MeV of the amount of 10^4 photons to maximize the statistic in the energy
 286 measurement.
- 287 — An attenuation length comparable to the size of the detector to prevent losing photons during
 288 their propagation in the LS. The final attenuation length is 25.8m [23] to compare with the CD
 289 diameter of 35.4m.
- 290 — Uranium/Thorium radiopurity to prevent background signal. The reactor neutrino program
 291 require a contamination fraction $F < 10^{-15}$ while the solar neutrino program require $F <$
 292 10^{-17} .

293 The LS will frequently be purified and tested in the Online Scintillator Internal Radioactivity In-
 294 vestigation System (OSIRIS) [24] to ensure that the requirements are kept during the lifetime of the
 295 experiment, more details to be found in section 2.4.2.

²⁹⁶ **Large Photo-Multipliers Tubes (LPMTs)**

²⁹⁷ The scintillation light produced by the LS is then collected by Photo-Multipliers Tubes (PMT) that
²⁹⁸ transform the incoming photon into an electric signal. As described in figure 2.8, the incident photons
²⁹⁹ interact with the photocathode via photoelectric effect producing an electron called a Photo-Electron
³⁰⁰ (PE). This PE is then focused on the dynodes where the high voltage will allow it to be multiplied.
³⁰¹ After multiple amplification the resulting charge - in coulomb [C] - is collected by the anode and
³⁰² the resulting electric signal can be digitalized by the readout electronics from which the charge and
³⁰³ timing can be extracted.

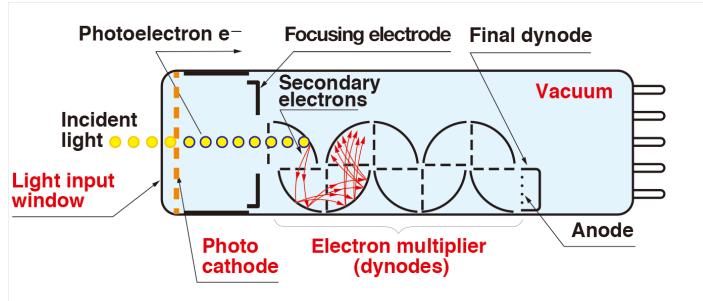


FIGURE 2.8 – Schematic of a PMT

³⁰⁴ The Large Photo-Multipliers Tubes (LPMT), used in the central detector and in the water pool, are
³⁰⁵ 20-inch (50.8 cm) radius PMTs. ~ 5000 dynode-PMTs [21] were produced by the Hamamatsu[®]
³⁰⁶ company and ~ 15000 Micro-Channel Plate (MCP) [25] by the NNVT[®] company. This system is
³⁰⁷ the one responsible for the energy measurement with a energy resolution of $3\%/\sqrt{E}$, resolution
³⁰⁸ necessary for the mass ordering measurement. To reach this precision, the system is composed of
³⁰⁹ 17612 PMTs quasi uniformly distributed over the detector for a coverage of 75.2% reaching ~ 1800
³¹⁰ PE/MeV or $\sim 2.3\%$ resolution due to statistic, leaving $\sim 0.7\%$ for the systematic uncertainties. They
³¹¹ are located outside the acrylic sphere in the water pool facing the center of the detector. To maintain
³¹² the resolution over the lifetime of the experiment, JUNO require a failure rate $< 1\%$ over 6 years.

³¹³ The LPMTs electronic are divided in two parts. One "near", located underwater, in proximity of the
³¹⁴ LPMT to reduce the cable length between the PMT and early electronic. A second one, outside of the
³¹⁵ detector that is responsible for higher level analysis before sending the data to the DAQ.

³¹⁶ The light yield per MeV induce that a LPMT can collect between 1 and 1000 PE per event, a wide
³¹⁷ dynamic range, causing non linearity in the PMT response that need to be understood and calibrated,
³¹⁸ see section 2.3 for more details.

³¹⁹ **Small Photo-Multipliers Tubes (SPMTs)**

³²⁰ The Small PMT (SPMTs) system is made of 3-inch (7.62 cm) PMTs. They will be used in the CD
³²¹ as a secondary detection system. Those 25600 SPMTs will observe the same events as the LPMTs,
³²² thus sharing the physics and detector systematics up until the photon conversion. With a detector
³²³ coverage of 2.7%, this system will collect ~ 43 PE/MeV for a final energy resolution of $\sim 17\%$.
³²⁴ This resolution is not enough to measure the NMO, θ_{13} , Δm_{31}^2 but will be sufficient to independently
³²⁵ measure θ_{12} and Δm_{21}^2 .

³²⁶ Due to the low PE rate, SPMTs will be running in photo-counting mode in the reactor range and thus
³²⁷ will be insensitive to non-linearity effect. Using this property, the intrinsic charge non linearity of
³²⁸ the LPMTs can be measured by comparing the PE count in the SPMTs and LPMTs [26]. Also, due
³²⁹ to their smaller size and electronics, SPMTs have a better timing resolutions than the LPMTs. At

330 higher energy range, like supernovae events, LPMTs will saturate where SPMTs due to their lower
 331 PE collection will to produce a reliable measure of the energy spectrum.

332 The Data Acquisition System (DAQ) is designed to support the event rate of IBD, background, dark
 333 noise and supplementary storage buffers are present in the LPMT electronics to withstand the event
 334 rate during supernovae burst.

335 2.2.3 Veto detector

336 The CD will be bathed in constant background noise coming from numerous sources : the radioac-
 337 tivity from surrounding rock and its own components or from the flux of cosmic muons. This
 338 background needs to be rejected to ensure the purity of the IBD spectrum. To prevent a big part
 339 of them, JUNO use two veto detector that will tag events as background before CD analysis.

340 Cherenkov in water pool

341 The Water Cherenkov Detector (WCD) is the instrumentation of the water buffer around the CD.
 342 When high speed charged particles will pass through the water, they will produced cherenkov
 343 photons. The light will be collected by 2400 MCP LPMTs installed on the outer surface of the CD
 344 structure. The muons veto strategy is based on a PMT multiplicity condition. WCD PMTs are
 345 grouped in ten zones: 5 in the top, 5 in the bottom. A veto is raised either when more than 19
 346 PMTs are triggered in one zone or when two adjacent zones simultaneously trigger more than 13
 347 PMTs. Using this trigger, we expect to reach a muon detection efficiency of 99.5% while keeping the
 348 noise at reasonable level.

349 Top tracker

350 The JUNO Top Tracker (TT) is a plastic scintillator detector located on the top of the experiment (see
 351 figure 2.9). Made from plastic scintillator from OPERA [27] layered horizontally in 3 layers on the
 352 top of the detector, the TT will be able to detect incoming atmospheric muons. With its coverage,
 353 about 1/3 of the of all atmospheric muons that passing through the CD will also pass through the 3
 354 layer of the detector. While it does not cover the majority of the CD, the TT is particularly effective
 355 to detect muons coming through the filling chimney region which might present difficulties from the
 other subsystems in some classes of events.

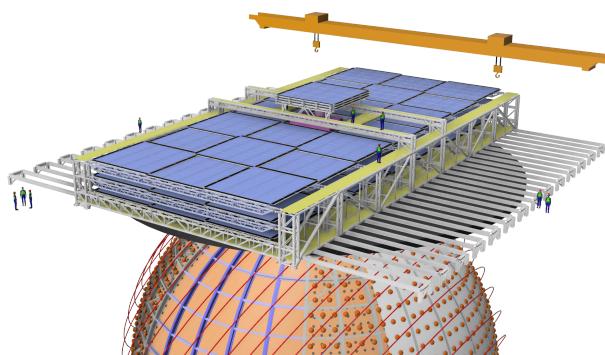


FIGURE 2.9 – The JUNO top tracker

357 2.3 Calibration strategy

358 The calibration is a crucial part of the JUNO experiment. Because we are looking at civil reactor
 359 neutrino it might be impossible to run measurement without signal, it would need to shut down
 360 every reactor from the Taishan and Yangjiang power plants which is realistically impossible. Because
 361 of this continuous rate, low frequency signal event, we need high frequency, recognisable sources in
 362 the energy range of interest : [0-12] MeV for the positron signal and 2.2 MeV for the neutron capture.
 363 It is expected that the CD response will be different depending on the type of particle, due to the
 364 interaction with LS, the position on the event and the optical response of the acrylic sphere (see
 365 section 2.6). We also expect a non-linear energy response of the CD due to the LS properties [19] but
 366 also due to the saturation of the LPMTs system when collecting a large amount of PE [26].

367 2.3.1 Energy scale calibration

368 While electrons and positrons sources would be ideal, for a large LS detector thin-walled electrons
 369 or positrons sources could lead to leakage of radionucleides causing radioactive contamination.
 370 Instead, we consider gamma sources in the range of the prompt energy of IBDs. The sources are
 371 reported in table 2.4.

Sources / Processes	Type	Radiation
^{137}Cs	γ	0.0662 MeV
^{54}Mn	γ	0.835 MeV
^{60}Co	γ	1.173 + 1.333 MeV
^{40}K	γ	1.461 MeV
^{68}Ge	e^+	annihilation 0.511 + 0.511 MeV
$^{241}\text{Am-Be}$	n, γ	neutron + 4.43 MeV ($^{12}\text{C}^*$)
$^{241}\text{Am-}^{13}\text{C}$	n, γ	neutron + 6.13 MeV ($^{16}\text{O}^*$)
$(n, \gamma)p$	γ	2.22 MeV
$(n, \gamma)^{12}\text{C}$	γ	4.94 MeV or 3.68 + 1.26 MeV

TABLE 2.4 – List of sources and their process considered for the energy scale calibration

372 For the ^{68}Ge source, it will decay in ^{68}Ga via electron capture, which will itself β^+ decay into ^{68}Zn .
 373 The positrons will be absorbed by the enclosure so only the annihilation gamma will be released. In
 374 addition, (α, n) sources like $^{241}\text{Am-Be}$ and $^{241}\text{Am-}^{13}\text{C}$ are used to provide both high energy gamma
 375 and neutrons, which will later be captured in the LS producing the 2.2 MeV gamma.

376 From this calibration we call E_{vis} the "visible energy" that is reconstructed by our current algorithms
 377 and we compare it to the true energy deposited by the calibration source. The results shown in figure
 378 2.10 show the expected response of the detector from calibration sources. The non-linearity is clearly
 379 visible from the $E_{\text{vis}}/E_{\text{true}}$ shape. See [28] for more details.

380 2.3.2 Calibration system

381 The non-uniformity due to the event position in the detector (more details in section 2.6) will be
 382 studied using multiples systems that are schematized in figure 2.11. They allow to position sources
 383 at different location in the CD.

- 384 — For a one-dimension vertical calibration, the Automatic Calibration Unit (ACU) will be able
 385 to deploy multiple radioactive sources or a pulse laser diffuser ball along the central axis of
 386 the CD through the top chimney. The source position precision is less than 1cm.

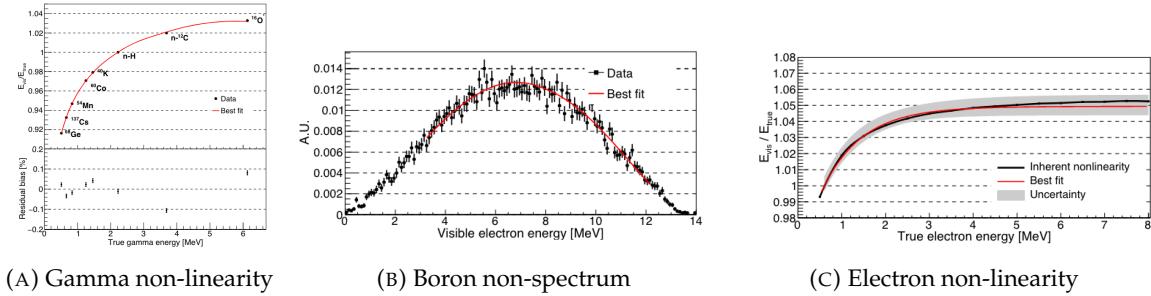


FIGURE 2.10 – Fitted and simulated non linearity of gamma, electron sources and from the ^{12}B spectrum. Black points are simulated data. Red curves are the best fits

- For off-axis calibration, a calibration source attached to a Cable Loop System (CLS) can be moved on a vertical half-plane by adjusting the length of two connection cable. Two set of CSL will be deployed to provide a 79% effective coverage of a vertical plane.
- A Guiding Tube (GT) will surround the CD to calibrate the non-uniformity of the response at the edge of the detector
- A Remotely Operated under-LS Vehicle (ROV) can be deployed to desired location inside LS for a more precise and comprehensive calibration. The ROV will also be equipped with a camera for inspection of the CD.

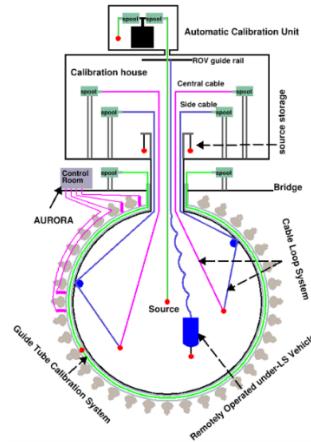


FIGURE 2.11 – Overview of the calibration system

- The preliminary calibration program is depicted in table 2.5.

2.4 Satellite detectors

- As introduced in section 2.1.1 and section 2.2.2, the precise knowledge and understanding of the detector condition is crucial for the measurements of the NMO and oscillation parameters. Thus two satellite detectors will be setup to monitor the experiment condition. TAO to monitor and understand the $\bar{\nu}_e$ flux and spectrum coming from the nuclear reactor and OSIRIS to monitor the LS response.

Program	Purpose	System	Duration [min]
Weekly calibration	Neutron (Am-C)	ACU	63
	Laser	ACU	78
Monthly calibration	Neutron (Am-C)	ACU	120
	Laser	ACU	147
	Neutron (Am-C)	CLS	333
	Neutron (Am-C)	GT	73
Comprehensive calibration	Neutron (Am-C)	ACU, CLS and GT	1942
	Neutron (Am-Be)	ACU	75
	Laser	ACU	391
	^{68}Ge	ACU	75
	^{137}Cs	ACU	75
	^{54}Mn	ACU	75
	^{60}Co	ACU	75
	^{40}K	ACU	158

TABLE 2.5 – Calibration program of the JUNO experiment

401 2.4.1 TAO

402 The Taishan Antineutrino Observatory (TAO) [12, 29] is a ton-level gadolinium doped liquid scin-
 403 tillator detector that will be located near the Taishan-1 reactor. It aim to measure the $\bar{\nu}_e$ spectrum at
 404 very low distance (< 30m) from the reactor to measure a quasi-unoscillated spectrum. TAO also aim
 405 to provide a major contribution to the so-called reactor anomaly [13]. Its requirement are to the level
 406 of 2 % energy resolution at 1 MeV.

407 Detector

408 The TAO detector is close, in concept, to the CD of JUNO. It is composed of an acrylic vessel
 409 containing 2.8 tons of gadolinium-loaded LS instrumented by an array of silicon photomultipliers
 410 (SiPM) reaching a 95% coverage. To efficiently reduce the dark count of those sensors, the detector
 411 is cooled to -50 °C. The $\bar{\nu}_e$ will interact with the LS via IBD, producing scintillation light, that will
 412 be detected by the SiPMs. From this signal the $\bar{\nu}_e$ energy and the full spectrum reconstructed. This
 413 spectrum will then be used by JUNO to calibrate the unoscillated spectrum, most notably the fission
 414 product fraction that impact the rate and shape of the spectrum. A schema of the detector is presented
 415 in figure 2.12a.

416 2.4.2 OSIRIS

417 The Online Scintillator Internal Radioactivity Investigation System (OSIRIS) [24] is an ultralow back-
 418 ground, 20 m³ LS detector that will be located in JUNO cavern. It aim to monitor the radioactive
 419 contamination, purity and overall response of the LS before it is injected in JUNO. OSIRIS will
 420 be located at the end of the purification chain of JUNO, monitoring that the purified LS meet the
 421 JUNO requirements. The setup is optimized to detect the fast coincidences decay of $^{214}\text{Bi} - ^{214}\text{Po}$
 422 and $^{212}\text{Bi} - ^{212}\text{Po}$, indicators of the decay chains of U and Th respectively.

423 Detector

424 OSIRIS is composed of an acrylic vessel that will contains 17t of LS. The LS is instrumented by
 425 a PMT array of 64 20 inch PMTs on the top and the side of the vessel. To reach the necessary

background level required by the LS purity measurements, in addition to being 700m underground in the experiment cavern, the acrylic vessel is immersed in a tank of ultra pure water. The water is itself instrumented by another array of 20 inch PMTs, acting as muon veto. A schema of the detector is presented in figure 2.12b.

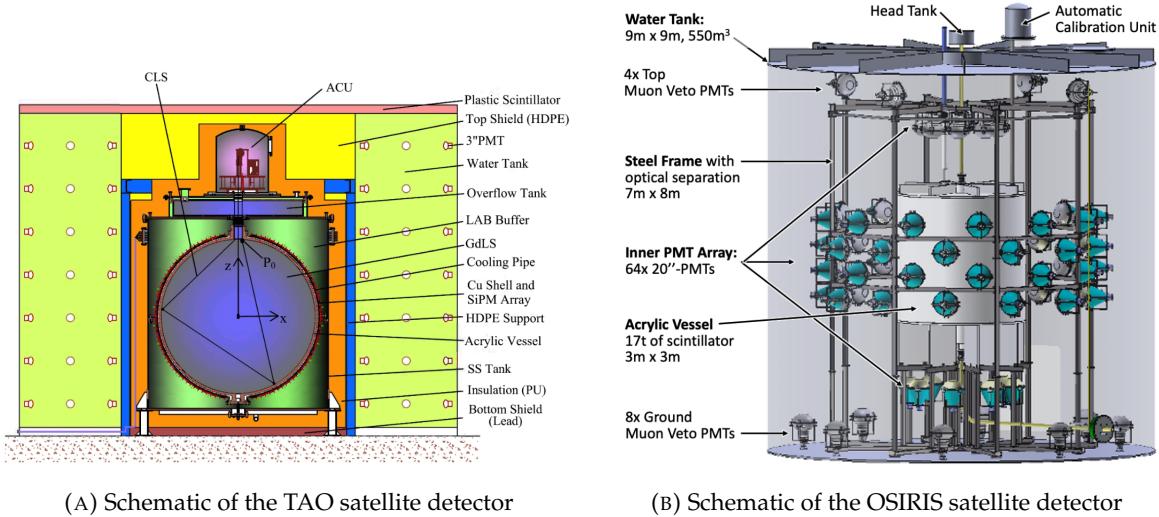


FIGURE 2.12

2.5 Software

The simulation, reconstruction and analysis algorithms are all packaged in the JUNO software, subsequently called the software. It is composed of multiple components integrated in the SNiPER [30] framework:

- Various primary particles simulators for the different kind of events, background and calibration sources.
- A Geant4 [31–33] Monte Carlo (MC) simulation containing the detectors geometries, a custom optical model for the LS and the supporting structures of the detectors. The Geant4 simulation integrate all relevant physics process for JUNO, validated by the collaboration. This step of the simulation is commonly called *Detsim* and compute up to the production of photo-electrons in the PMTs. The optics properties of the different materials and detector components have been measured beforehand to be used to define the material and surfaces in the simulation.
- An electronic simulation, simulating the response waveform of the PMTs, tracking it through the digitization process, accounting for effects such as non-linearity, dark noise, Time Transit Spread (TTS), pre-pulsing, after-pulsing and ringing if the waveform. It's also the step handling the event triggers and mixing. This step is commonly referenced as *Elecsim*.
- A waveform reconstruction where the digitized waveform are filtered to remove high-frequency white noise and then deconvoluted to yield time and charge informations of the photons hits on the PMTs. This step is commonly referenced as *Calib*.
- The charge and time informations are used by reconstruction algorithms to reconstruct the interaction vertex and the deposited energy. This step is commonly reported as *Reco*. See section 2.6 for more details on the reconstruction.
- Once the singular events are reconstructed, they go through event pairing and classification to select IBD events. This step is named Event Classification.

— The purified signal is then analysed by the analysis framework which depend of the physics topic of interest.

The steps Reco and Event Classification are divided into two category of algorithm. Fast but less accurate algorithms that are running during the data taking designated as the *Online* algorithms. Those algorithm are used to take the decision to save the event on tape or to throw it away. More accurate algorithms that run on batch of events designated *Offline* algorithms. They are used for the physics analysis. The Offline Reco will be one of the main topic of interest for this thesis.

2.6 State of the art of the Offline IBD reconstruction in JUNO

The main reconstruction method currently run in JUNO is a data-driven method based on a likelihood maximization [34, 35] using only the LPMTs. The first step is to reconstruct the interaction vertex from which the energy reconstruction is dependent. It is also necessary for event pairing and classification.

2.6.1 Interaction vertex reconstruction

To start the likelihood maximization, a rough estimation of the vertex and of the event timing is needed. We start by estimating the vertex position using a charge based algorithm.

Charge based algorithm

The charge-based algorithm is basically base on the charge-weighted average of the PMT position.

$$\vec{r}_{cb} = a \cdot \frac{\sum_i q_i \cdot \vec{r}_i}{\sum_i q_i} \quad (2.3)$$

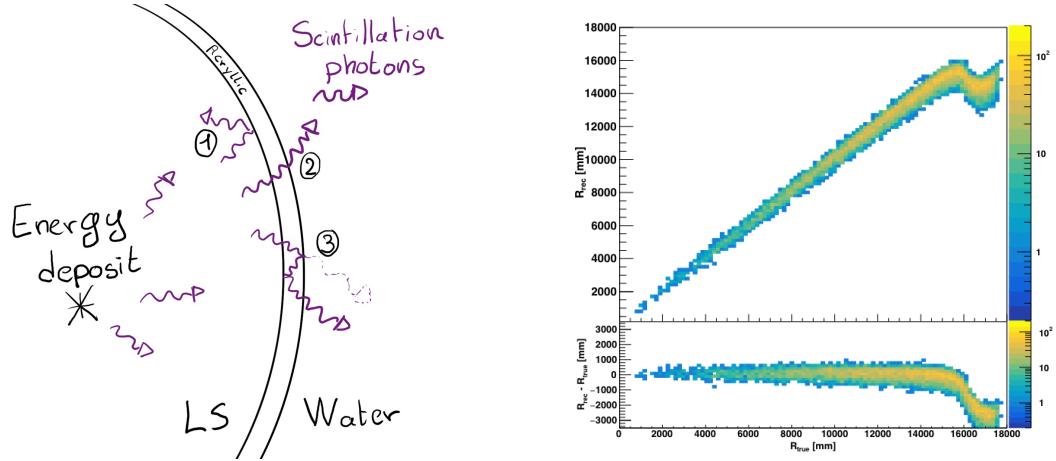
Where q_i is the reconstructed charge of the pulse of the i th PMT and \vec{r}_i is its position. \vec{r}_0 is the reconstructed interaction position. a is a scale factor introduced because a weighted average over a 3D sphere is inherently biased. Using calibration we can estimate $a \approx 1.3$ [36]. The results in figure 2.13b shows that the reconstruction is biased from around 15m and further. This is due to the phenomena called “total reflection area” or TR Area.

As depicted in the figure 2.13a the optical photons, given that they have a sufficiently large incidence angle, can be deviated of their trajectories when passing through the interfaces LS-acrylic and water-acrylic due to the optical index difference. This cause photons to be lost or to be detected by PMT further than anticipated if we consider their rectilinear trajectories. This cause the charge barycenter to be located closer to the center than the event really is.

It is to be noted that charge based algorithm, in addition to be biased near the edge of the detector, does not provide any information about the timing of the event. Therefore, a time based algorithm needs to be introduced to provide initial values.

Time based algorithm

The time based algorithm use the distribution of the time of flight corrections Δt (Eq 2.4) of an event to reconstruct its vertex and t_0 . It follow the following iterations:



(A) Illustration of the different optical photons reflection scenarios. 1 is the reflection of the photon at the interface LS-acrylic or acrylic-water. 2 is the transmission of the photons through the interfaces. 3 is the conduction of the photon in the acrylic.

(B) Heatmap of R_{rec} and $R_{rec} - R_{true}$ as a function of R_{true} for 4MeV prompt signals uniformly distributed in the detector calculated by the charge based algorithm

FIGURE 2.13

487 1. Use the charge based algorithm to get an initial vertex to start the iteration.

488 2. Calculate the time of flight correction for the i th PMT using

$$\Delta t_i(j) = t_i - \text{tof}_i(j) \quad (2.4)$$

489 where j is the iteration step, t_i is the timing of the i th PMT, and tof_i is the time-of-flight of the
490 photon considering an rectilinear trajectory and an effective velocity in the LS and water (see
491 [36] for detailed description of this effective velocity). Plot the Δt distribution and label the
492 peak position as Δt^{peak} (see fig 2.14a).

493 3. Calculate a correction vector $\vec{\delta}[\vec{r}(j)]$ as

$$\vec{\delta}[\vec{r}(j)] = \frac{\sum_i \left(\frac{\Delta t(j) - \Delta t^{\text{peak}}(j)}{\text{tof}_i(j)} \right) \cdot (\vec{r}_0(j) - \vec{r}_i)}{N^{\text{peak}}(j)} \quad (2.5)$$

494 where \vec{r}_0 is the vertex position at the beginning of this iteration, \vec{r}_i is the position of the i th
495 PMT. To minimize the effect of scattering, dark noise and reflection, only the pulse happening
496 in a time window (-10 ns, +5 ns) around Δt^{peak} are considered. N^{peak} is the number of PE
497 collected in this time-window.

498 4. if $\vec{\delta}[\vec{r}(j)] < 1\text{mm}$ or $j \geq 100$, stop the iteration. Otherwise $\vec{r}_0(j+1) = \vec{r}_0(j) + \vec{\delta}[\vec{r}(j)]$ and go to
499 step 2.

500 However because the earliest arrival time is used, t_i is related to the number photoelectrons N_i^{pe}
501 detected by the PMT [37–39]. To reduce bias in the vertex reconstruction, the following equation is
502 used to correct t_i into t'_i :

$$t'_i = t_i - p_0 / \sqrt{N_i^{\text{pe}}} - p_1 - p_2 / N_i^{\text{pe}} \quad (2.6)$$

503 The parameters (p_0, p_1, p_2) were optimized to (9.42, 0.74, -4.60) for Hamamatsu PMTs and (41.31,
504 -12.04, -20.02) for NNVT PMTs [36]. The results presented in figure 2.14b shows that the time based

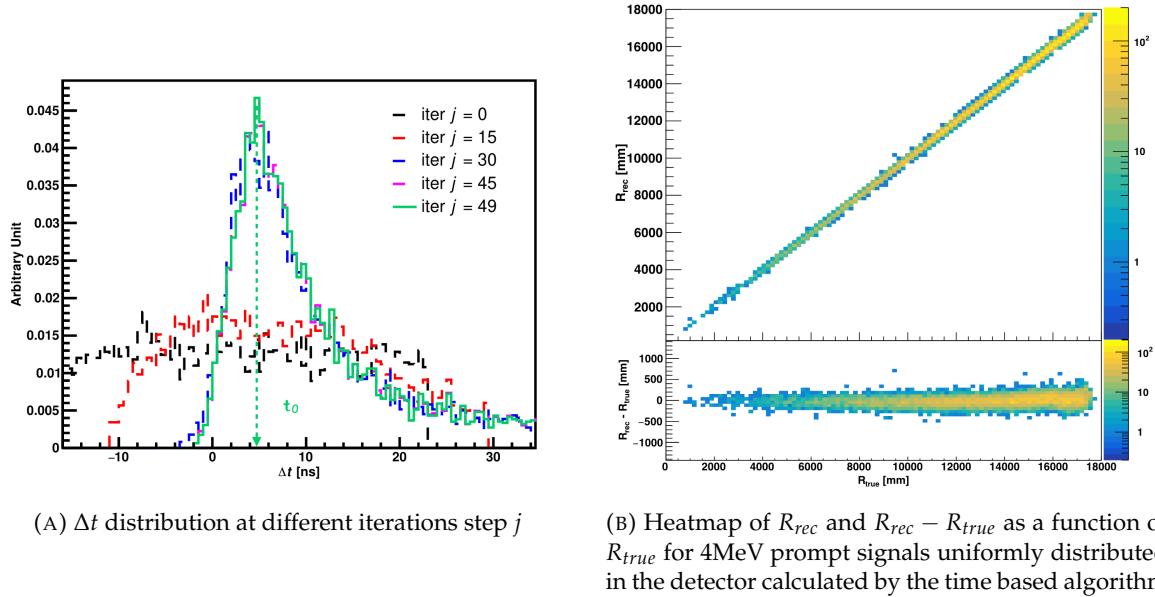
(A) Δt distribution at different iterations step j (B) Heatmap of R_{rec} and $R_{rec} - R_{true}$ as a function of R_{true} for 4MeV prompt signals uniformly distributed in the detector calculated by the time based algorithm

FIGURE 2.14

505 algorithm provide a more accurate vertex and is unbiased even in the TR area. This results (\vec{r}_0, t_0) is
506 used as initial value for the likelihood algorithm.

507 Time likelihood algorithm

508 The time likelihood algorithm use the residual time expressed as follow

$$t_{\text{res}}^i(\vec{r}_0, t_0) = t_i - \text{tof}_i - t_0 \quad (2.7)$$

509 In a first order approximation, the scintillator time response Probability Density Function (PDF) can
510 be described as the emission time profile of the scintillation photons, the Time Transit Spread (TTS)
511 and the dark noise of the PMTs. The emission time profile $f(t_{\text{res}})$ is described like

$$f(t_{\text{res}}) = \sum_k \frac{\rho_k}{\tau_k} e^{-\frac{t_{\text{res}}}{\tau_k}}, \sum_k \rho_k = 1 \quad (2.8)$$

512 as the sum of the k component that emit light in the LS each one characterised by it's decay time τ_k
513 and intensity fraction ρ_k . The TTS component is expressed as a gaussian convolution

$$g(t_{\text{res}}) = \frac{1}{\sqrt{2\pi}\sigma} e^{-\frac{(t_{\text{res}}-\nu)^2}{2\sigma^2}} \cdot f(t_{\text{res}}) \quad (2.9)$$

514 where σ is the TTS of PMTs and ν is the average transit time. The dark noise is not correlated with any
515 physical events and considered as constant rate over the time window considered T . By normalizing
516 the dark noise probability $\epsilon(t_{\text{res}})$ as $\int_T \epsilon(t_{\text{res}}) dt_{\text{res}} = \epsilon_{\text{dn}}$, it can be integrated in the PDF as

$$p(t_{\text{res}}) = (1 - \epsilon_{\text{dn}}) \cdot g(t_{\text{res}}) + \epsilon(t_{\text{res}}) \quad (2.10)$$

517 The distribution of the residual time t_{res} of an event can then be compared to $p(t_{\text{res}})$ and the best

518 fitting vertex \vec{r}_0 and t_0 can be chosen by minimizing

$$\mathcal{L}(\vec{r}_0, t_0) = -\ln \left(\prod_i p(t_{\text{res}}^i) \right) \quad (2.11)$$

519 The parameter of Eq. 2.10 can be measured experimentally. The results shown in figure 2.15 used
 520 PDF from monte carlo simulation. The results shows that $R_{\text{rec}} - R_{\text{true}}$ is biased depending on the
 521 energy. While this could be corrected using calibration, another algorithm based on charge likelihood
 522 was developed to correct this problem.

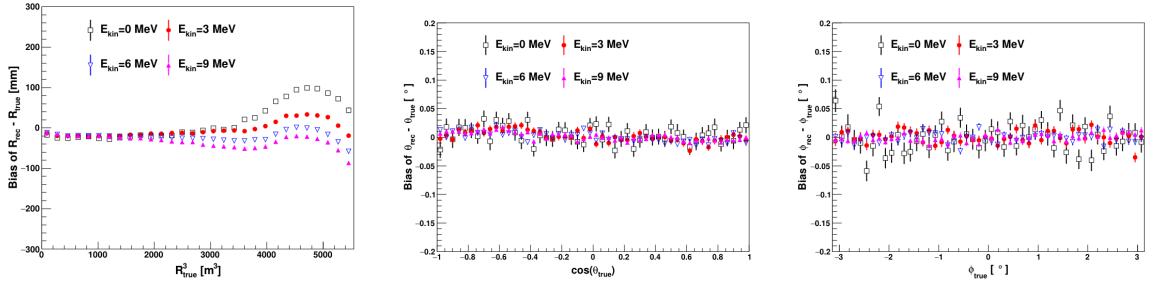


FIGURE 2.15 – Bias of the reconstructed radius R (left), θ (middle) and ϕ (right) for multiple energies by the time likelihood algorithm

523 Charge likelihood algorithm

524 Similarly to the time likelihood algorithms that use a timing PDF, the charge likelihood algorithm
 525 use a PE PDF for each PMT depending on the energy and position of the event. With $\mu(\vec{r}_0, E)$ the
 526 mean expected number of PE detected by each PMT, the probability to observe N_{pe} in a PMT follow
 527 a Poisson distribution. Thus

- 528 — The probability to observe no hit ($N_{pe} = 0$) in the j th PMT is $P_{\text{nohit}}^j(\vec{r}_0, E) = e^{-\mu_j}$
- 529 — The probability to observe $N_{pe} \neq 0$ in the i th PMT is $P_{\text{hit}}^i(\vec{r}_0, E) = \frac{\mu^{N_{pe}} e^{-\mu_i}}{N_{pe}!}$

530 Therefore, the probability to observe a specific hit pattern can be expressed as

$$P(\vec{r}_0, E) = \prod_j P_{\text{nohit}}^j(\vec{r}_0, E) \cdot \prod_i P_{\text{hit}}^i(\vec{r}_0, E) \quad (2.12)$$

531 The best fit values of \vec{R}_0 and E can then be calculated by minimizing the negative log-likelihood

$$\mathcal{L}(\vec{r}_0, E) = -\ln(P(\vec{r}_0, E)) \quad (2.13)$$

532 In principle, $\mu_i(\vec{r}_0, E)$ could be expressed

$$\mu_i(\vec{r}_0, E) = Y \cdot \frac{\Omega(\vec{r}_0, r_i)}{4\pi} \cdot \epsilon_i \cdot f(\theta_i) \cdot e^{-\sum_m \frac{d_m}{\zeta_m}} \cdot E + \delta_i \quad (2.14)$$

533 where Y is the energy scale factor, $\Omega(\vec{r}_0, r_i)$ is the solid angle of the i th PMT, ϵ_i is its detection
 534 efficiency, $f(\theta_i)$ its angular response, ζ_m is the attenuation length in the materials and δ_i the expected
 535 number of dark noise.

536 However Eq. 2.14 assume that the scintillation light yield is linear with energy and describe poorly
 537 the contribution of indirect light, shadow effect due to the supporting structure and the total reflec-

tion effects. The solution is to use data driven methods to produce the pdf by using the calibrations sources and position described in section 2.3. In the results presented in figures 2.16, the PDF was produced using MC simulation and 29 specific calibrations position [36] along the Z-axis of the detector. We see that the charge likelihood algorithm show little bias in the TR area and a better

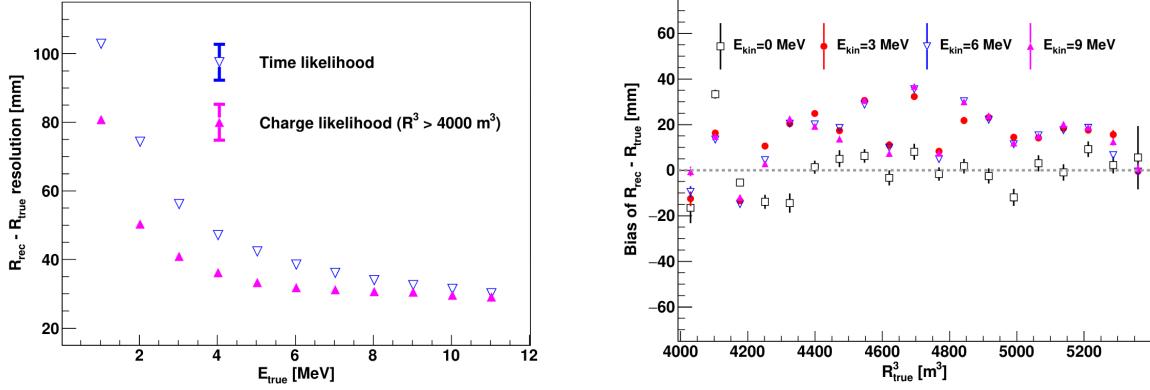


FIGURE 2.16 – On the left: Resolution of the reconstructed R as a function of the energy in the TR area ($R^3 > 4000 \text{ m}^3 \equiv R > 16 \text{ m}$) by the charge and time likelihood algorithms. On the right: Bias of the reconstructed R in the TR area for different energies by the charge likelihood algorithm

resolution than the time likelihood. The figure 2.17 shows the radial resolution of the different algorithm presented for this section, we can see the refinement at each step and that the charge likelihood yield the best results.

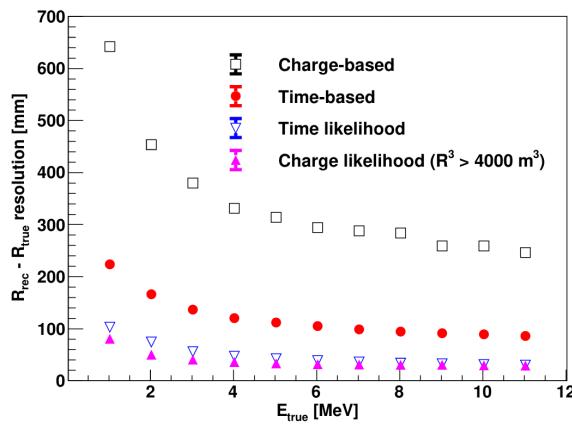


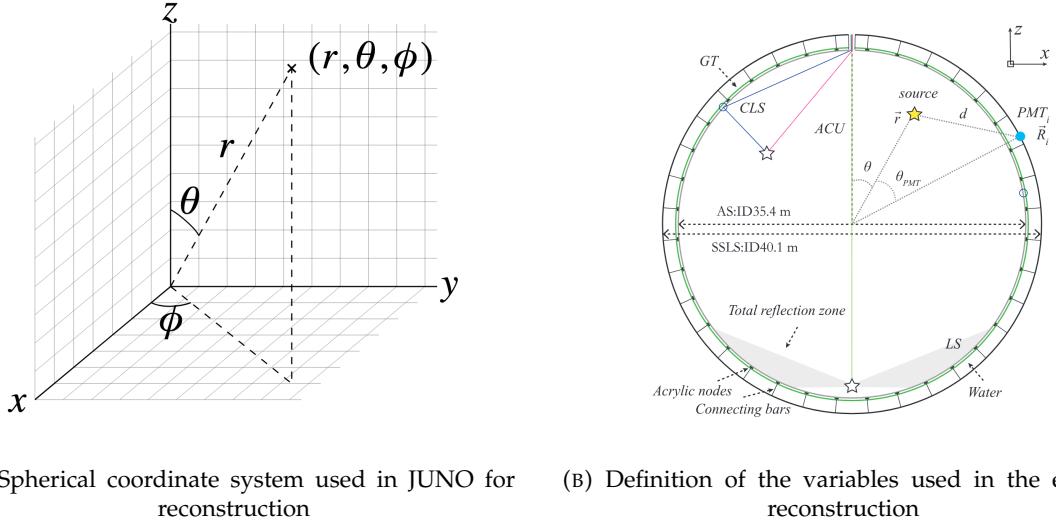
FIGURE 2.17 – Radial resolution of the different vertex reconstruction algorithms as a function of the energy

The charge based likelihood algorithms already give use some information on the energy as Eq. 2.13 is minimized but the energy can be further refined as shown in the next section.

2.6.2 Energy reconstruction

As explained in section 2.1.1, energy resolution is crucial for the NMO and oscillation parameters measurements. Thus the energy reconstruction algorithm should take into consideration as much

detector effect as possible. The following method is a data driven method based on calibration samples inspired by the charge likelihood algorithm described above [40].



(A) Spherical coordinate system used in JUNO for reconstruction

(B) Definition of the variables used in the energy reconstruction

FIGURE 2.18

Charge estimation

The most important element in the energy reconstruction is $\mu_i(\vec{r}_0, E)$ described in Eq. 2.14. For realistic cases, we also need to take into account the electronics effect that were omitted in the previous section. Those effect will cause a charge smearing due to the uncertainties in the N_{pe} reconstruction. Thus we define $\hat{\mu}^L(\vec{r}_0, E)$ which is the expected N_{pe}/E in the whole detector for an event with visible energy E_{vis} and position \vec{r}_0 . The position of the event and PMTs are now defined using $(r, \theta, \theta_{pmt})$ as defined in figure 2.18b.

$$\hat{\mu}(r, \theta, \theta_{pmt}, E_{vis}) = \frac{1}{E_{vis}} \frac{1}{M} \sum_i^M \frac{\bar{q}_i - \mu_i^D}{\text{DE}_i}, \quad \mu_i^D = \text{DNR}_i \cdot L \quad (2.15)$$

where i runs over the PMTs with the same θ_{pmt} , DE_i is the detection efficiency of the i th PMT. μ_i^D is the expected number of dark noise photoelectrons in the time window L . The time window have been optimized to $L = 280$ ns [40]. \bar{q}_i is the average recorded photoelectrons in the time window and \hat{Q}_i is the expected average charge for 1 photoelectron. The N_{pe} map is constructed following the procedure described in [35].

Time estimation

The second important observable is the hit time of photons that was previously defined in Eq. 2.7. It is here refined as

$$t_r = t_h - \text{tof} - t_0 = t_{LS} + t_{TT} \quad (2.16)$$

where t_h is the time of hit, t_{LS} is the scintillation time and t_{TT} the transit time of PMTs that is described by a gaussian

$$t_{TT} = \mathcal{N}(\overline{\mu_{TT} + t_d}, \sigma_{TT}) \quad (2.17)$$

569 where μ_{TT} is the mean transit time in PMTs, σ_{TT} is the Transit Time Spread (TTS) of the PMTs and t_d
 570 is the delay time in the electronics. The effective refraction index of the LS is also corrected to take
 571 into account the propagation distance in the detector.

572 The timing PDF $P_T(t_r|r, d, \mu_l, \mu_d, k)$ can now be generated using calibration sources [40]. This PDF
 573 describe the probability that the residual time of the first photon hit is in $[t_r, t_r + \delta]$ with r the radius
 574 of the event vertex, $d = |\vec{r} - \vec{r}_{PMT}|$ the propagation distance, μ_l and μ_d the expected number of PE
 575 and dark noise in the electronic reading window and k is the detected number of PE.

576 Now let denote $f(t, r, d)$ the probability density function of "photoelectron hit a time t" for an event
 577 happening at r where the photons traveled the distance d in the LS

$$F(t, r, d) = \int_t^L f(t', r, d) dt' \quad (2.18)$$

578 Based on the PDF for one photon $k = 1$, one can define

$$P_T^l(t|k = n) = I_n^l [f_l(t) F_l^{n-1}(t)] \quad (2.19)$$

579 where the indicator l means that the photons comes from the LS and I_n^l a normalisation factor. To this
 580 pdf we add the probability to have photons coming from the dark noise indicated by the indicator d
 581 using

$$f_d(t) = 1/L, F_d(t) = 1 - \frac{t}{L} \quad (2.20)$$

582 and so for the case where only one photon is detected by the PMT ($k = 1$)

$$P_T(t|\mu_l, \mu_d, k = 1) = I_1[P(1, \mu_l)P(0, \mu_d)f_l(t) + P(0, \mu_l)P(1, \mu_d)f_d(t)] \quad (2.21)$$

583 where $P(k_\alpha, \mu_\alpha)$ is the Poisson probability to detect k_α PE from $\alpha \in \{l, d\}$ with the condition $k_l + k_d = k$.
 584

585 Now that we have the individual timing and charge probability we can construct the charge likeli-
 586 hood referred as QMLE:

$$\mathcal{L}(q_1, q_2, \dots, q_N | \vec{r}, E_{vis}) = \prod_{j \in \text{unfired}} e^{-\mu_j} \prod_{i \in \text{fired}} \left(\sum_{k=1}^K P_Q(q_i|k) \cdot P(k, \mu_i) \right) \quad (2.22)$$

587 where $\mu_i = E_{vis}\hat{\mu}_i^L + \mu_i^D$ and $P(k, \mu_i)$ is the Poisson probability of observing k PE. $P_Q(q_i|k)$ is the
 588 charge pdf for k PE. And we can also construct the time likelihood referred as TMLE:

$$\mathcal{L}(t_{1,r}, t_{2,r}, \dots, t_{N,r} | \vec{r}, t_0) = \prod_{i \in \text{hit}} \frac{\sum_{k=1}^K P_T(t_{i,r}|r, d, \mu_i^l, \mu_i^d, k) \cdot P(k, \mu_i^l + \mu_i^d)}{\sum_{k=1}^K P(k, \mu_i^l + \mu_i^d)} \quad (2.23)$$

589 where K is cut to 20 PE and hit is the set of hits satisfying $-100 < t_{i,r} < 500$ ns.

590 Merging those two likelihood give the charge-time likelihood QTMLLE

$$\mathcal{L}(q_1, q_2, \dots, q_N; t_{1,r}, t_{2,r}, \dots, t_{N,r} | \vec{r}, t_0, E_{vis}) = \mathcal{L}(q_1, q_2, \dots, q_N | \vec{r}, E_{vis}) \cdot \mathcal{L}(t_{1,r}, t_{2,r}, \dots, t_{N,r} | \vec{r}, t_0) \quad (2.24)$$

591 The radial and energy resolutions of the different likelihood are presented in figure 2.19 (from [40]).
 592 We can see the improvement of adding the time information to the vertex reconstruction and that
 593 an increase in vertex precision can bring improvement in the energy resolution, especially at low
 594 energies.

595 Data driven methods prove to be performant in the energy and vertex reconstruction given that we
 596 have enough calibrations sources to produce the PDF. In the next section, we'll see another type of

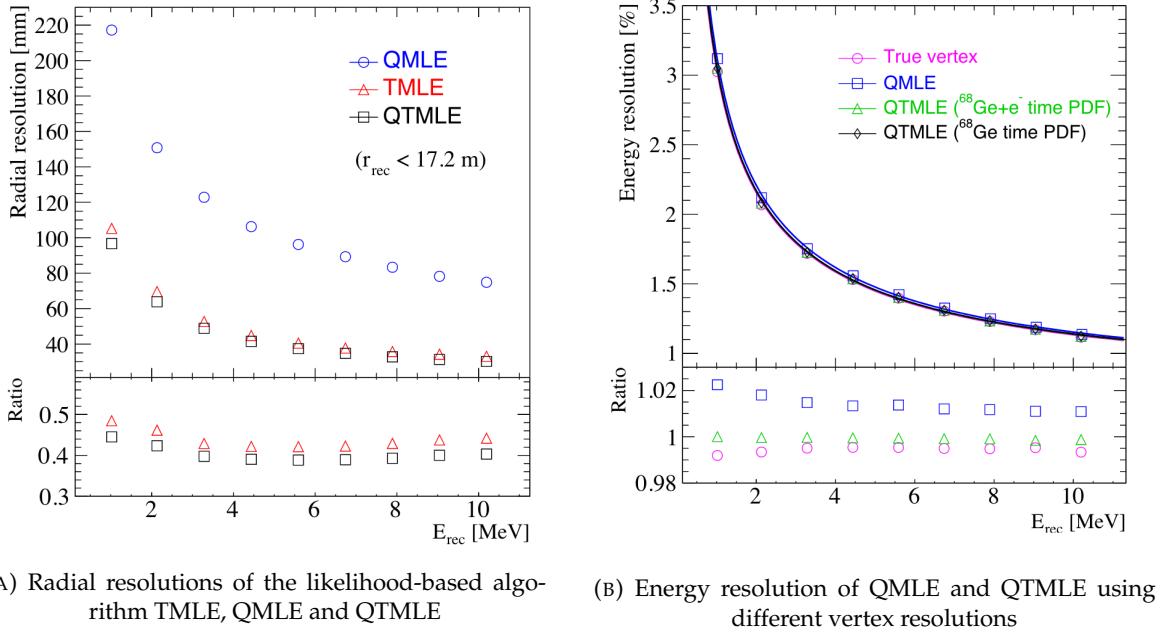


FIGURE 2.19

597 data-driven method based on machine learning.

598 2.6.3 Machine learning for reconstruction

599 Machine learning (ML) is family of data-driven algorithms that are inferring behavior and results
600 from a training dataset. A overview of methods and detailed explanation of the Neural Network
601 (NN) subfamily can be found in Chapter 3.

602 The power of ML is the ability to model complex response to a specific problem. In JUNO the
603 reconstruction problematic can be expressed as follow: knowing that each PMT, large or small,
604 detected a given number of PE Q at a given time t and their position is x, y, z where did the energy
605 was deposited and how much energy was it, modeling a function that naively goes:

$$\mathbb{R}^{5 \times N_{\text{pmt}}} \mapsto \mathbb{R}^4 \quad (2.25)$$

606 It is worth pointing that while this is already a lot in informations, this is not the rawest representa-
607 tion of the experiment. We could indeed replace the charge and time by the waveform in the time
608 window of the event but that would lead to an input representation size that would exceed our
609 computational limits. Also, due to those computational limits, most of the ML algorithm reduce this
610 input phase space either by structurally encoding the information (pictures, graph), by aggregating
611 it (mean, variance, ...) or by exploiting invariance and equivariance of the experiment (rotational
612 invariance due to the sphericity, ...).

613 For machine learning to converge to performant algorithm, a large dataset exploring all the phase
614 space of interest is needed. For the following studies, data from the monte carlo simulation presented
615 in section 2.5 are used for training. When the detector will be finished calibrations sources will be
616 complementarily be used.

617 **Boosted Decision Tree (BDT)**

618 On of the most classic ML method used in physics in last years is the Boosted Decision Tree (see
 619 chapter 3.1). They have been explored for vertex reconstruction [41] et for energy reconstruction [41,
 620 42].

621 For vertex and energy reconstruction a BDT was developed using the aggregated informations pre-
 622 sented in 2.6.

Parameter	description
$nHits$	Total number of hits
$x_{cc}, y_{cc}, z_{cc}, R_{cc}$	Coordinates of the center of charge
ht_{mean}, ht_{std}	Hit time mean and standard deviation

TABLE 2.6 – Features used by the BDT for vertex reconstruction

623 Its reconstruction performances are presented in figure 2.21.

624 A second and more advanced BDT, subsequently named BDTE, that only reconstruct energy use a
 625 different set of features [42]. They are presented in the table 2.7

626 **Neural Network (NN)**

627 The physics have shown a rising for Neural Network (NN) in the past years for event reconstruction,
 628 notably in the neutrino community [43–46]. Three type of neural networks have explored for event
 629 reconstruction in JUNO Deep Neural Network (DNN), Convolutional Neural Network (CNN) and
 630 Graph Network (GNN). More explanation about those neural network can be found in chapter 3.

631 The CNN are using 2D projection of the detector representing it as an image with two channel, one
 632 for the charge Q and one for the time t . The position of the PMTs is structurally encoded in the pixel
 633 containing the information of this PMT. In [41], the pixel is chosen based on a transformation of θ
 634 and ϕ coordinates to the 2D plane and rounded to the nearest pixel. A sufficiently large image has
 635 been chosen to prevent two PMT to be located in the same pixel. An example of this projection can
 636 be found in figure 2.20. The performances of the CNN can be found in figure 2.21.

637 Using 2D have the upside of encoding a large part of the informations structurally but loose the rota-
 638 tional invariance of the detector. It also give undefined information to the neural network (what is a
 639 pixel without PMT ? What should be its charge and time ?), cause deformation in the representation
 640 of the detector (sides of projection) and loose topological informations.

641 One of the way to present structurally the sphericity of JUNO to a NN is to use a graph: A collection
 642 of objects V called nodes and relations E called edges, each relation associated to a couple v_1, v_2

AccumCharge	$ht_{5\%-2\%}$
R_{cht}	pe_{mean}
z_{cc}	J_{cht}
pe_{std}	ϕ_{cc}
nPMTs	$ht_{35\%-30\%}$
$ht_{kurtosis}$	$ht_{20\%-15\%}$
$ht_{25\%-20\%}$	$pe_{35\%}$
R_{cc}	$ht_{30\%-25\%}$

TABLE 2.7 – Features used by the BDTE algorithm. pe and ht reference the charge
 and hit-time distribution respectively and the percentages are the quantiles of those
 distributions. cht and cc reference the barycenters of hit time and charge respectively

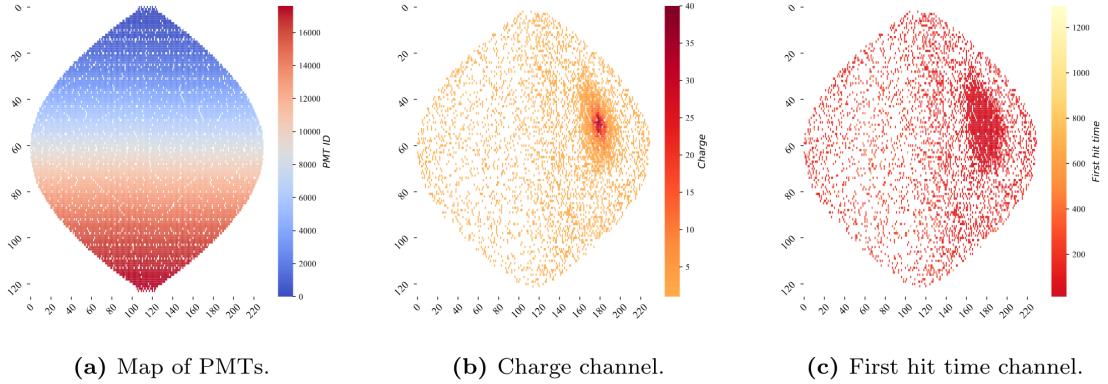


FIGURE 2.20 – Projection of the LPMTs in JUNO on a 2D plane. (a) Show the distribution of all PMTs and (b) and (c) are example of what the charge and time channel looks like respectively

643 forming the graph $G(E, V)$. Nodes and edges can hold informations or features. In [41] the nodes,
 644 are geometrical region of the detector as defined by the HealPix [47]. The features of the nodes are
 645 aggregated informations from the PMTs it contains. The edges contains geographic informations of
 646 the nodes relative positions.

647 This data representation has the advantages to keep the topology of the detector intact. It also permit
 648 the use of rotational invariant algorithms for the NN, thus taking advantage of the symmetries of the
 649 detector.

650 The neural network then process the graph using Chebyshev Convolutions [48]. The performances
 651 of the GNN are presented in figure 2.21.

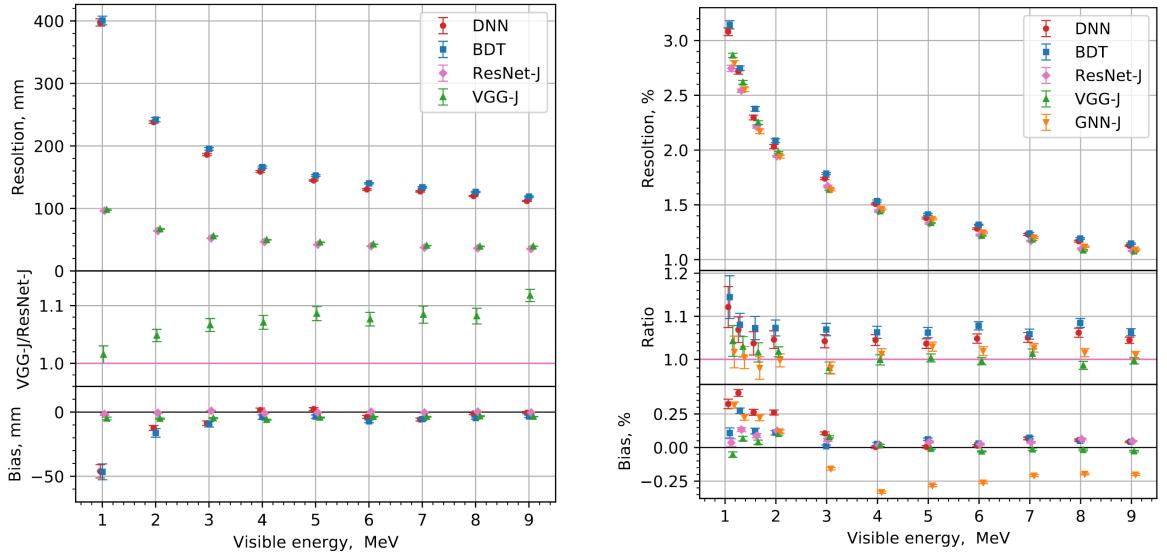


FIGURE 2.21 – Radial (left) and energy (right) resolutions of different ML algorithms. The results presented here are from [41]. DNN is a deep neural network, BDT is a BDT, ResNet-J and VGG-J are CNN and GNN-J is a GNN.

652 Overall ML algorithms show similar performances as classical algorithms in term of energy recon-
 653 structions with the more complex structure CNN and GNN showing better performances than BDT

and DNN. For vertex reconstruction, the BDT and DNN show poor performance while CNN are on the level of the classical algorithms.

2.7 JUNO sensitivity to NMO and precise measurements

Now that the event have been reconstructed, selected and that the non-IBD background have been rejected, we have access to the measured energy flux from JUNO. We consider two spectra, the one measured by the LPMT system and the one measured by the SPMT system. This give rise to three possible analysis: A LPMT only analysis, a SPMT only analysis and a joint analysis. This joint analysis is the subject of the chapter 7 of this thesis.

The following details about JUNO measurement is common to the three analysis. The details and specific of the joint analysis are detailed in chapter 7.

2.7.1 Theoretical spectrum

To extract the oscillation parameters and the NMO from the measured spectrum, it is compared to a theoretical spectrum. This theoretical spectrum is produced based on the theory of the three flavour oscillation (see section 1.3), the measurements of the calibration and satellite experiments and Monte Carlo simulation:

- The absolute flux and the fission product fraction calibrated by TAO.
- The estimation of the neutrinos flux from other sources, such as the geoneutrinos, by theoretical model.
- The computed cross-section of $\bar{\nu}_e$ and the LS.
- The estimation of mislabelled event, such as fast neutron events from cosmic muons, using Monte Carlo simulation.
- The measured bias and resolution of the LPMT and SPMT system by the calibration.
- The time dependent reactor parameters (age of fuel, instantaneous power of the reactors, etc...)

These systematics parameters come with their uncertainties that need to be taken into account by the fitting framework. This theoretical spectrum will, in the end, depend of the oscillation parameters of interest θ_{13} , θ_{12} , Δm_{21}^2 , Δm_{31}^2 . Noise parameters can be included in the parameters spectrum such as the earth density ρ between the power plants and JUNO.

2.7.2 Fitting procedure

The theoretical and measured spectra are represented as two histograms depending on the energy. The theoretical spectrum is adjusted with the data using a χ^2 minimization where χ^2 is naively defined as

$$\chi^2 = \sum_i \frac{(N_{th}^i - N_{data}^i)^2}{\sigma_i^2} \quad (2.26)$$

where N_{th}^i is the number event in the i th bin of the theoretical spectrum, N_{data}^i is the number of event in the i th bin of the measured spectrum and σ_i is the uncertainty of this bin. Two classic statistic test exist Pearson and Neyman where the difference is the estimation of σ_i parameters.

This σ_i is composed of the systematics uncertainties discussed above but also from the statistic uncertainty of the spectrum. Considering a Poisson process, the statistic uncertainty is estimated as $\sigma_{stat}^i = \sqrt{N^i}$. In a Pearson test, $N^i \equiv N_{th}^i$ whereas in a Neyman test $N^i \equiv N_{data}^i$. Under the assumption that the content of each bin follow a Gaussian distribution (a Poisson with high enough statistic), the two test are equivalent. But studies on Monte Carlo spectrum showed that the Pearson

and Neyman statistic are biased in opposite direction. It is easily visible where, for the same data, Pearson will prefer a higher N_{th}^i to reduce the ratio $\frac{1}{N_{th}^i}$ whereas Neyman will prefer a lower N_{th}^i to reduce the $(N_{th}^i - N_{data}^i)$ term.

This problematic can be circumvented by summing the two test, yielding the CNP statistic test and/or by adding a term

$$\chi^2 = \sum_i \frac{(N_{th}^i - N_{data}^i)^2}{\sigma_i^2} - \ln |\mathbf{V}| \quad (2.27)$$

where V is the covariance matrix of the theoretical spectrum yielding the PearsonV and CNPV statistic test.

The χ^2 is minimized by exploring the parameter phase space via gradient descent.

2.7.3 Physics results

The oscillation parameters are directly extracted from the minimization procedure and the error can be estimated directly from the procedure. For the NMO, the data are fitted under the two assumption of NO and IO. The difference in χ^2 give us the preferred ordering and the significance of our test. Latest studies show that the precision on oscillation parameters after six year of data taking will be of 0.2%, 0.3%, 0.5% and 12.1% for Δm_{31}^2 , Δm_{21}^2 , $\sin^2 \theta_{12}$ and $\sin^2 \theta_{13}$ respectively [11]. The expected sensitivity to mass ordering is 3σ after 6 years [49].

2.8 Summary

JUNO is one the biggest new generation neutrino experiment. Its goal, the measurements of oscillation parameters with unprecedeted precision and an NMO preference at the 3 sigma confidence level, needs an in depth knowledge and understanding of the detector and the physics at hand. The characterisation and calibration of the detector are of the utmost importance and the understanding of the detector response in its resolution and bias is capital to be able to correctly carry the high precision physics analysis of the neutrino oscillation.

In this thesis, I explore the usage of data-driven reconstruction methods to validate and optimize the reconstruction of IBD events in JUNO in the chapters 4, 5 and 6 and the usage of the dual calorimetry in the detection of possible mis-modelisation in the theoretical spectrum 7.

⁷¹⁸ **Chapter 3**

⁷¹⁹ **Machine learning and Artificial
Neural Network**

⁷²¹ *"I have the shape of a human being and organs equivalent to those of a human being. My organs, in fact, are identical to some of those in a prostheticized human being. I have contributed artistically, literally, and scientifically to human culture as much as any human being now alive. What more can one ask?"*

Isaac Asimov, *The Complete Robot*

⁷²² Machine Learning (ML) and more specifically Neural Network (NN) are families of data-driven ⁷²³ algorithm. They are used to model complex distributions from a finite dataset to extract a generalist ⁷²⁴ behavior. They learn, adapt their intrinsic parameters, interactively by computing its performance ⁷²⁵ or loss on those dataset. They take advantage of simple microscopic operation such as *if condition* or ⁷²⁶ non-continuous but differentiable function like *ReLU*. Through optimizers and the combination of a ⁷²⁷ lot of those microscopic operations, they can obtain complex and precise behaviours.

⁷²⁸ They are now widely used in a wide variety of domain including natural language processing, ⁷²⁹ computer vision, speech recognition and, the subject of this thesis, scientific studies.

⁷³⁰ We found them in particle physics, either as the main algorithm or as secondary algorithm, for event ⁷³¹ reconstruction, event classification, waveform reconstruction, etc..., domains where the underlying ⁷³² physic and detector process is complex and highly dimensional. Physicists have traditionally been ⁷³³ forced to use simplifications or assumptions to ease the development of algorithms or equations ⁷³⁴ (a good example is the algorithm presented in section 2.6) where machine learning could refine and ⁷³⁵ take into account those effects, provided that they have enough data and computing power.

⁷³⁶ This chapter present an overview of the different kind of machine learning methods and neural ⁷³⁷ networks that will be discussed in this thesis.

⁷³⁸ **3.1 Boosted Decision Tree (BDT)**

⁷³⁹ One of the most classic machine learning algorithm used in particle physics is Boosted Decision Tree ⁷⁴⁰ (BDT) [50] (or more recently Gradient Boosting Machine [51]). The principle of a BDT is fairly simple ⁷⁴¹ : based on a set of observables, a serie of decisions, represented as node in a tree, are taken by the ⁷⁴² algorithm. Each decision point, or node, takes its decision based on a set of trainable parameters ⁷⁴³ leading to a subtree of decision. The process is repeated until it reach the final node, yielding the ⁷⁴⁴ prediction. A simplistic example is given in figure 3.1.

⁷⁴⁵ The training procedure follow a simple score reward procedure. During the training phase the ⁷⁴⁶ prediction of the BDT is compared to a known truth about the data. The score is then used to

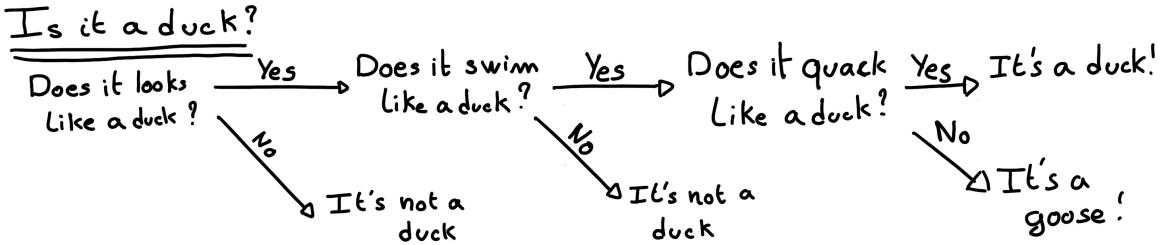


FIGURE 3.1 – Example of a BDT that determine if the given object is a duck

747 backpropagate corrections to the parameters of the tree. Modern BDT use gradient boosting where
 748 the gradient of the loss is calculated for each of the BDT parameters. Following the gradient descent,
 749 we can reach the, hopefully, global minima of the loss for our set of parameters.

750 3.2 Artificial Neural Network (NN)

751 One other big family of machine learning algorithm is the artificial Neural Networks (NN). The idea
 752 of developing automates which component mimic, in a simplistic way, the behavior of biological
 753 neurons emerge in 1959 with the paper “*What the Frog’s Eye Tells the Frog’s Brain*” [52]. They develop
 754 an automate where each component possess an *activation function*. Each one of those component then
 755 transmit its information to the other following a certain efficiency or *weight*. Those works influenced
 756 scientist and notably Frank Rosenblatt who published in 1958 what is considered the first neural
 757 network model the Perceptron [53].

758 Modern neural network still nowadays use the neuron metaphor to represent neural network, but
 759 approach them as a graph where the nodes are neurons possessing an activation function and edges
 760 holding the weights, or *parameters* in modern literature, between those nodes. Most of the modern
 761 neural network work with the principle of neurons layers. Each neurons belong to a layer and takes
 762 input from the preceding layer and forward it result to next layer. For example the most basic set
 763 layer is the fully connected layer where each of its neurons is connected to every other neurons of
 764 the precessing layer. All the neurons posses the same activation function F . The connection between
 765 two the two layers is expressed as a tensor T_j^i where i is the index of the precedent layer and j the
 766 index of the current layer. The propagation from the layer I to J is then described as

$$J_j = F_j(T_j^i I_i + B_j) \quad (3.1)$$

767 where the learning parameters are the tensor T_j^i and the bias tensor B_j . This is the fundamental
 768 component of the Fully Connected Deep NN (FCDNN) family presented in section 3.2.1. Most of the
 769 modern neural networks use gradient descent to optimize their parameters, i.e. the gradient of the
 770 parameter θ in respect of the loss function \mathcal{L} is subtracted to it

$$\theta_{i+1} = \theta_i - \frac{\partial \mathcal{L}}{\partial \theta} \quad (3.2)$$

771 i being the training iteration index. This needs the expression of \mathcal{L} dependent of θ to be differentiable,
 772 thus the layer and their activation function also need to be differentiable. This simple gradient
 773 descent, designated as Stochastic Gradient Descent (SGD), can be completed with first and second
 774 order momentum like with the Adam optimizer [54] (more details in section 3.2.5).

775 This description of neural networks as layer introduced the principle of *depth* and *width*, the number
 776 of layers in the NN and the number of neurons in each layer respectively. Those quantities that not

⁷⁷⁷ directly used for the computation of the results but describe the NN or its training are designated as
⁷⁷⁸ *hyperparameters*.

⁷⁷⁹ 3.2.1 Fully Connected Deep Neural Network (FCDNN)

⁷⁸⁰ Fully Connected Deep Neural Network (FCDNN) architecture is the natural evolution of the Perceptron.
⁷⁸¹ The input data is represented as a first order tensor I_j and then fed forward to multiple fully
⁷⁸² connected layers (Eq 3.1) as presented in the figure 3.2a. Most of the time, the classic ReLU function

$$\text{ReLU}(x) = \begin{cases} x & \text{if } x \geq 0 \\ 0 & \text{otherwise} \end{cases} \quad (3.3)$$

⁷⁸³ is used as activation function. PReLU and Sigmoid are also popular choices:

$$\text{Sigmoid}(x) = \frac{1}{1 + e^{-x}} \quad (3.4) \quad \text{PReLU}(x) = \begin{cases} x & \text{if } x \geq 0 \\ \alpha x & \text{otherwise} \end{cases} \quad (3.5)$$

⁷⁸⁵ The reasoning behind ReLU and PReLU is that with enough of them, you can mimic any continuous
⁷⁸⁶ function as illustrated in figure 3.2b. Sigmoid is more used in case of classification, its behavior going
⁷⁸⁷ hand in hand with the Cross Entropy loss function used in classification problems.

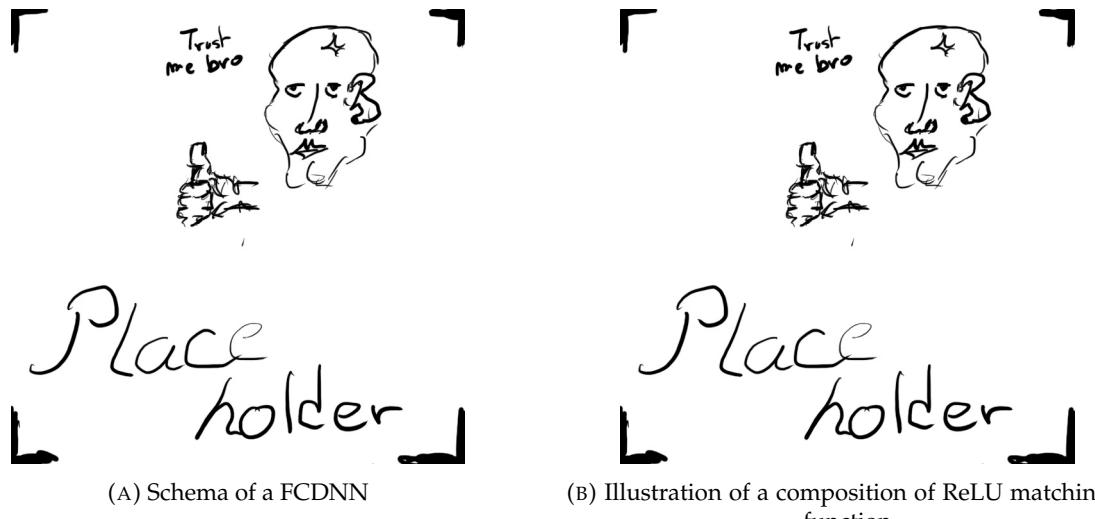


FIGURE 3.2

⁷⁸⁸ Due to its simplicity, FCDNN are also used as basic pieces for more complex architectures such as
⁷⁸⁹ the CNN and GNN that will be presented in the next section.

⁷⁹⁰ 3.2.2 Convolutional Neural Network (CNN)

⁷⁹¹ Convolutional Neural Networks are a family of neural networks that use discrete convolution filters,
⁷⁹² as illustrated in an example in figure 3.3, to process the input data, often images. They have the
⁷⁹³ advantage to be translation invariant by construction, this mean that they are capable of detecting
⁷⁹⁴ oriented features independently of their location on the image. The learning parameters are located
⁷⁹⁵ in the filters, the network thus learn the optimal filters to extract the desired features. 2D CNN,

796 where the filters are second order tensors that span over third order tensors, are commonly used in
 797 image recognition [55] for classification or regression problematics.

798 The convolution layers are commonly chained [56], reducing the input dimension while increasing
 799 the number of filters. The idea behind is that the first layers will process local informations and the
 800 latest layers will process more global informations. To try to preserve the amount of information, we
 801 tend to double the numbers of filters for each division of the input data. The results of the convolution
 802 filters is commonly then flattened and feed to a smaller FCDNN which will process the filters results
 803 to yield the desired output.

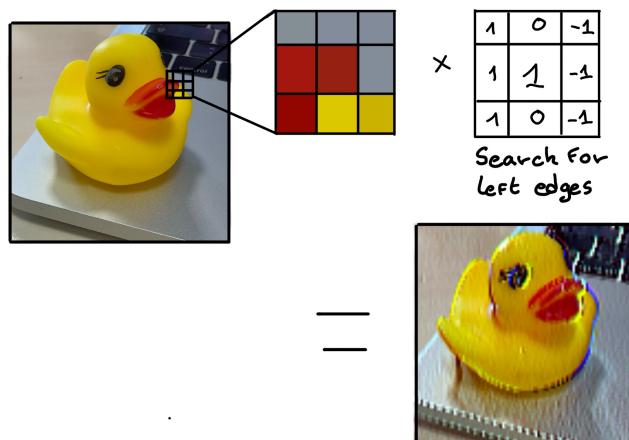


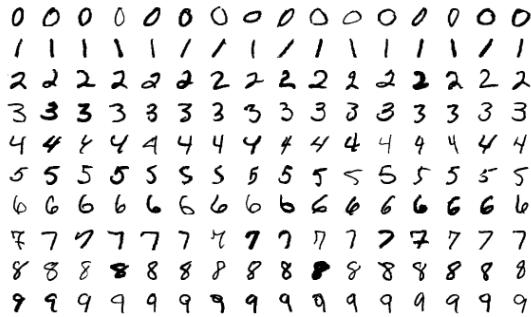
FIGURE 3.3 – Illustration of the effect of a convolution filter. Here we apply a filter with the aim do detect left edges. We see in the resulting image that the left edges of the duck are bright yellow where the right edges are dark blue indicating the contour of the object. The convolution was calculated using [57].

804 As an example, let's take the Pytorch [58] example for the MNIST [59], a dataset of black and white
 805 images of handwritten digits. Those images are 28×28 pixels with only one channel corresponding
 806 to the grey level of the pixel. Example of images from this dataset are presented in figure 3.4a

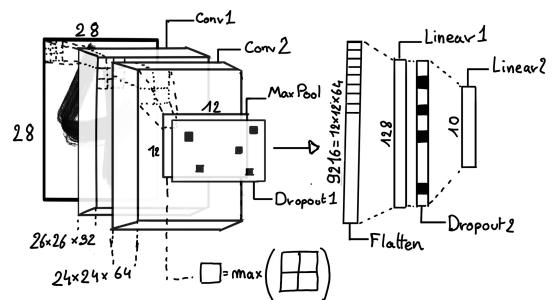
807 A schema of the CNN used in the Pytorch example is presented in figure 3.4b. Using this schema as
 808 a reference, the trained network is made of:

- 809 1. A convolutional layer of (3×3) filters yielding 32 channels. A bias parameter is applied
 810 to each channel for a total of $(32 \cdot (3 \times 3) + 32) = 320$ parameters. The resulting image is
 811 $(26 \times 26 \times 32)$ (26 per 26 pixels with 32 channels). The ReLU activation function is applied to
 812 each pixel.
- 813 2. A second convolutional layer of (3×3) filters yielding 64 channels. This channel also posses
 814 a bias parameter for a total of $(64 \cdot (3 \times 3) + 64) = 640$ parameters. Resulting image is $(24 \times$
 815 $24 \times 64)$. Also with with a ReLU activation function.
- 816 3. Then comes a (2×2) max pool layer with a stride of 1 meaning that for each channel the max
 817 value of pixels in a (2×2) block is condensed in a single resulting pixel. The resulting image
 818 is $(12 \times 12 \times 64)$.
- 819 4. This image goes through a dropout layer which will set the pixel to 0 with a probability of
 820 0.25. This help prevent overtraining of the neural network (see section 3.2.6 for more details).
- 821 5. The data is the flattened i.e. condensed into a vector of $(12 \times 12 \times 64) = 9216$ values.
- 822 6. Then comes a fully connected linear layer (Eq. 3.1) with a ReLU activation that output 128
 823 feature. It needs $(9216 \cdot 128) + 128 = 1'179'776$ parameters.
- 824 7. This 128 item vector goes through another dropout layer with a probability of 0.5

- 825 8. The vector is then transformed through a linear layer with ReLU activation. It output 10
 826 values, one for each digit class (0, 1, 2, ..., 9). It need $(128 \cdot 10) + 128 = 1408$ parameters.
- 827 9. Finally the 10 values are normalized using a log softmax function $\text{LogSoftmax}(x_i) = \log \left(\frac{\exp(x_i)}{\sum_j \exp(x_j)} \right)$
 828 to give the probability of the input image to be a certain digit.



(A) Example of images in the MNIST dataset



(B) Schema of the CNN used in Pytorch example to process the MNIST dataset

FIGURE 3.4

829 The final network needs 1'182'144 parameters or, if we consider each parameters to be a double
 830 precision floating point, 9.45 MB of data. To gives a order of magnitude, such neural network is
 831 considered "simple", train in a matter of minutes on T4 GPU [60] (14 epochs) and reach an accuracy
 832 in its prediction of 99%.

833 3.2.3 Graph Neural Network (GNN)

834 Graph neural network is a family of neural network where the data is represented as a graph $G(\mathcal{N}, \mathcal{E})$
 835 composed of vertex or node $n \in \mathcal{N}$ and edges $e \in \mathcal{E}$. The edges are associated to two nodes $(u, v) \in$
 836 \mathcal{N}^2 , "connecting" them. The node and the edges can hold features, commonly represented as vector
 837 $n \in \mathbb{R}^{k_n}$, $e \in \mathbb{R}^{k_e}$. We can thus define a graph using two tensors A_{ϵ}^{ij} the adjacency tensors that hold
 838 the features ϵ of the edge connecting the node i and j and the tensor N_v^i that hold the features v of a
 839 node i .

840 To efficiently manipulate such object we need to structurally encode their property in the neural
 841 network architecture: each node is equivalent (as opposite to ordered data in a vector), each node has
 842 a set of neighbours, ... One of this method is the message passing algorithm presented historically
 843 in "Neural Message Passing for Quantum Chemistry" [61]. In this algorithm, with each layer of
 844 message passing a new set of features is computed for each node following

$$n_i^{k+1} = \phi_u(n_i^k, \square_j \phi_m(n_i^k, n_j^k, e_{ij}^k)); n_j \in \mathcal{N}'_i \quad (3.6)$$

845 where ϕ_u is a differentiable update function, \square_j is a differentiable aggregation function and ϕ_m is a
 846 differentiable message function. $\mathcal{N}'_i = \{n_j \in \mathcal{N} | (n_i, n_j) \in \mathcal{E}\}$ is the set of neighbours of n_i , i.e. the
 847 nodes n_j from which it exist an edge $e_{ij} \rightarrow (n_i, n_j)$. k is the layer on which the message passing
 848 algorithm is applied. \square need also a few other property if we want to keep the graph property, most
 849 notably the permutational invariance of its parameters (example: mean, std, sum, ...).

850 The edges features can also be updated, either by directly taking the results of ϕ_m or by using another
 851 message function ϕ_e .

852 Message passing is a very generic way of describing the process of GNN and it can be specialized
 853 for convolutional filtering [48], diffusion [62] and many other specific operation. GNN are used in a
 854 wide variety of application such as regression problematics, node classification, edge classification,
 855 node and edge prediction, ...
 856 It is a very versatile but complex tool.

857 3.2.4 Adversarial Neural Network (ANN)

858 The adversarial machine learning, Adversarial Neural Networks (ANN) in the case of neural net-
 859 work, is a family of unsupervised machine learning algorithms where the learning algorithm (gen-
 860 erator) is competing against another algorithm (discriminator). Taking the example of Generative
 861 Adversarial Networks, concept initially developed by Goodfellow et al. [63], the discriminator goal
 862 is to discriminate between data coming from a reference dataset and data produced by the generator.
 863 The generator goal, on the other hand, is to produce data that the discriminator would not be able to
 864 differentiate from data from the reference dataset. The expression of duality between the two models
 865 is represented in the loss where, at least a part of it, is driven by the results of the discriminator.

866 3.2.5 Training procedure

867 A neural network without the adequate training is like an empty shell. If the parameters are not
 868 optimized they are, most of the time, initialized to random number and so the output will just be
 869 random. The training is a key step in the production of a solid and reliable NN. This section aim to
 870 give an overview of the different concept and tools used in the training of our neural networks.

871 Training lifecycle

872 The training of NN does not follow strict rules, you could imagine totally different lifecycle but I will
 873 describe here the one used in this thesis, the most common one.

874 The training is split into *epochs* during which the NN will train on a set of subsamples called *batch*.
 875 The size of those batch is called *batch size*, a.k.a. the number of data it contains (how many images,
 876 how many events,...). Each process of a batch is called a *step*. At the end of each epochs, the neural
 877 network is evaluated over a validation dataset. This validation dataset is not used for training (no
 878 gradient of the loss is computed) and is used as reference for the network performance and monitor
 879 overtraining (see section 3.2.6). Most of the time, the parameters are updated at each step using the
 880 mean loss over the batch and the optimizer hyperparameters are updated at each epochs.

881 The optimizer

882 As briefly introduced section 3.2, the parameters of the neural network are optimized using the
 883 gradient descent method. We calculate the gradient of the mean loss over the batch with respect
 884 of each parameters and we update the parameters in accord to minimize the loss. The gradient is
 885 computed backward from the loss up to the first layer parameters using the chain rule:

$$\frac{\partial \mathcal{L}}{\partial \theta_1} = \frac{\partial \theta_2}{\partial \theta_1} \frac{\partial \mathcal{L}}{\partial \theta_2} = \frac{\partial \theta_2}{\partial \theta_1} \frac{\partial \theta_3}{\partial \theta_2} \frac{\partial \mathcal{L}}{\partial \theta_3} = \frac{\partial \theta_2}{\partial \theta_1} \prod_{i=2}^{N-1} \frac{\partial \theta_{i+1}}{\partial \theta_i} \frac{\partial \mathcal{L}}{\partial \theta_N} \quad (3.7)$$

886 where θ is a parameter, i is the layer index. We see here that the gradient of the first layer is dependent
 887 of the gradient of all the following layers. We thus need to compute the gradient closest to loss first
 888 before computing the gradient of the earlier layers. This is called the *backward propagation*.

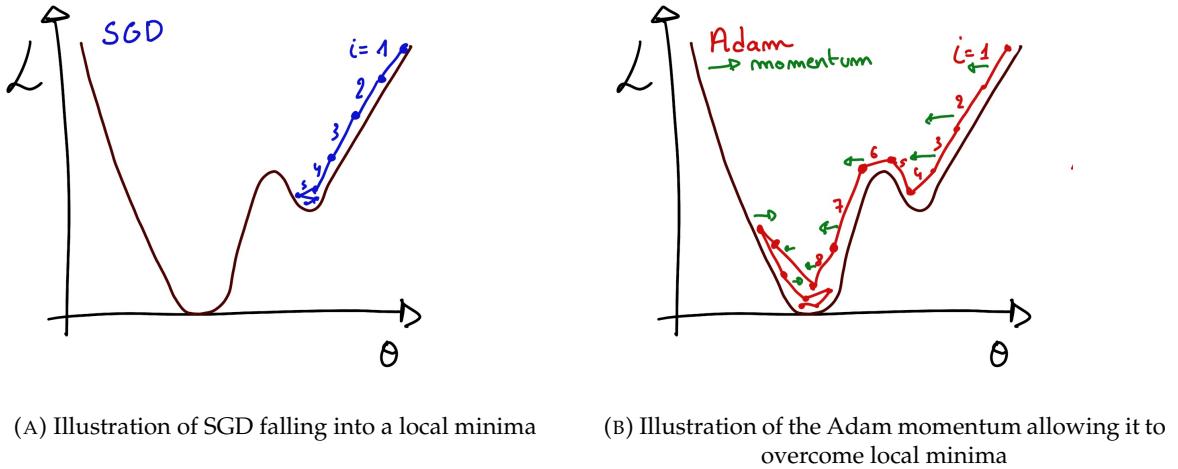


FIGURE 3.5

889 This update of the parameters is done following an optimizer policy. Those optimizers depends on
890 hyperparameters. The ones used in this thesis are:

- 891 1. SGD (Stochastic Gradient Descent). This is the simplest optimizer, it depend on only one
892 hyperparameter, the learning rate λ (LR) and update the parameters θ following

$$\theta_{t+1} = \theta_t - \lambda \frac{\partial \mathcal{L}}{\partial \theta} \Big|_{\theta_t} \quad (3.8)$$

893 where t is the step index. It is a powerful optimizer but is very sensible to local minima of the
894 loss in the parameters phase space as illustrated in figure 3.5a.

2. Adam [54]. The concept is, in short, to have and SGD but with momentum. Adam possess two momentum $m(\beta_1)$ and $v(\beta_2)$ which are respectively proportional to $\frac{\partial \mathcal{L}}{\partial \theta}$ and $(\frac{\partial \mathcal{L}}{\partial \theta})^2$. β_1 and β_2 are hyperparameters that dictate the moment update at each optimization step. The parameters are then upgraded following

$$m_{t+1} = \beta_1 m_t + (1 - \beta_1) \frac{\partial \mathcal{L}}{\partial \theta} \quad (3.9)$$

$$v_{t+1} = \beta_2 v_t + (1 - \beta_2) \left(\frac{\partial \mathcal{L}}{\partial \theta} \right)^2 \quad (3.10)$$

$$\theta_{t+1} = \theta_t - \lambda \frac{m_{t+1}}{\sqrt{v_{t+1}} + \epsilon} \quad (3.11)$$

895 where ϵ is a small number to prevent divergence when v is close to 0. These momentums
896 allow to overcome small local minima in the parameters phase space as illustrated in figure
897 3.5a.

898 Scheduler policies

899 Sometimes we want to update our hyperparameters or take a set of action during the training
900 procedure. We use for this scheduler policies, for example a common policy is a decrease of the
901 learning rate after each epochs. The reasoning is that if the learning rate is too high, the optimizer

902 will continuously miss the minimum and oscillate around it. By reducing the learning rate, we allow
 903 it to make more fine steps in the parameters phase space, hopefully converging to the true minima.

904 Another policy that is often used is the save of the best model. In some situations, the loss value after
 905 each epoch will strongly oscillate or even worsen. This policy allows us to keep the best version
 906 of the model attained during the training phase.

907 3.2.6 Potential pitfalls

908 Apart from being stuck in local minima, there are also other behaviors and effects we want to prevent
 909 during training.

910 Overtraining

911 This happens when the network learns the specificities of the training dataset instead of a more general
 912 representation of the underlying data distribution. This can happen if there is not enough data
 913 in comparison to the number of learning parameters, if the data contains some specific signatures
 914 specific to the training dataset or if it trains for too long on the same dataset. This behavior is illustrated
 915 in figure 3.6a. Overtraining can be fought in multiple ways, for example:

- 916 — **More data.** By having more data in the training dataset, the network will not be able to learn the
 917 specificities of every data.
- 918 — **Less parameters.** By reducing the number of parameters, we reduce the computing and
 919 learning capacities of the network. This will force it to fallback to generalist behaviors.
- 920 — **Dropout.** This technique implies to randomly set part of the neural network to 0. By doing
 921 this, we force the redundancy in its computing capability and, in a way, modify the data
 922 decreasing the possibility for specific learning.
- 923 — **Early stopping.** During the training we monitor the network performance over a validation
 924 dataset. The network does not train on this dataset and thus cannot learn its specificities. If
 925 the loss on the training dataset diverges too much from the loss on the validation dataset, we
 926 can stop the training earlier to prevent it from overtraining.

927 Gradient vanishing

928 Gradient vanishing is the effect of the gradient being so small for the upper layer that the parameters
 929 are barely updated after each step. This causes the network to be unable to converge to the minima.

930 This comes from the way the gradient descent is calculated. Imagine a simple network composed of
 931 three fully connected layers: the input layer, an intermediate layer and the output layer. Let L be the
 932 loss, θ_1 the parameter between the input and the intermediate layer and θ_2 the parameter between
 933 the intermediate and output layer. This network is schematized in figure 3.6b.

934 The gradient for θ_1 will be computed using the chain rule presented in equation 3.7. Because θ_1
 935 depends on θ_2 , if the gradient of θ_2 is small, so will be the gradient of θ_1 . Now if we have many
 936 more layers, we can see how the subsequent multiplication of small gradients would lead to
 937 very small updates of the parameters thus “vanishing gradient”.

938 Multiple actions can be taken to prevent this effect such as:

- 939 — **Batch normalization:** In this case we apply a normalization layer that will normalize the data
 940 so that, let D be the data, $\langle D \rangle = 0$ and $\sigma_D = 1$. This helps the weights of the network to
 941 maintain an appropriate scale.
- 942 — **Residual Network (ResNet) [64]:** Residual network is a technique for neural network in
 943 which, instead of just sequentially feeding the results of each layer to the next one, you ask
 944 each layer to calculate the residual of the input data. This technique is illustrated in figure 3.7.

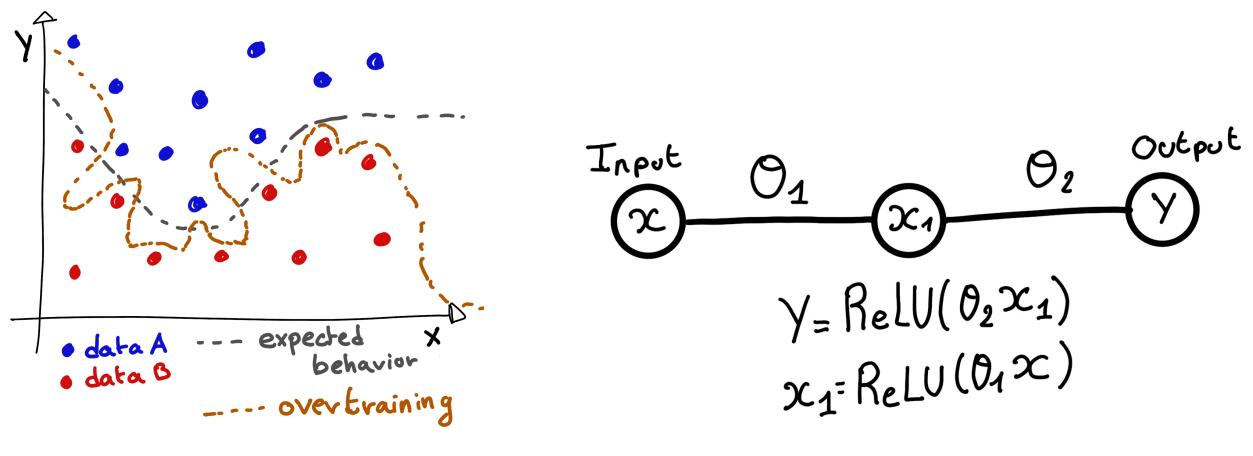


FIGURE 3.6

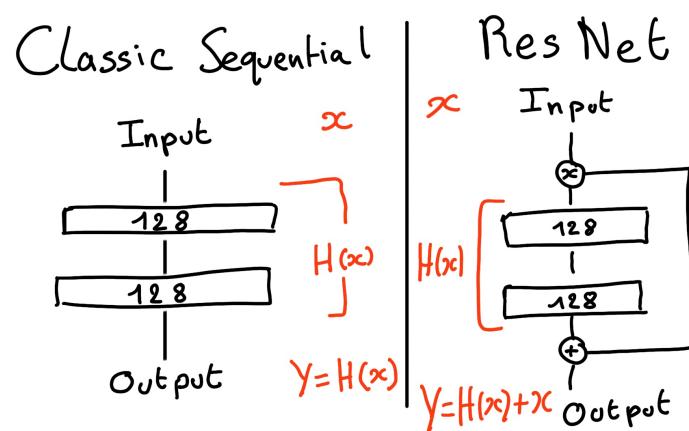


FIGURE 3.7 – Illustration of the ResNet framework

945 **Chapter 4**

946 **Image recognition for IBD
reconstruction with the SPMT system**

947 *Dave - Give me the position and momentum, HAL.*

HAL - I'm afraid I can't do that Dave.

Dave - What's the problem ?

HAL - I think you know what the problem is just as well as I do.

Dave - What are you talking about, HAL?

HAL - $\sigma_x \sigma_p \geq \frac{\hbar}{2}$

949 As explained in chapter 2, JUNO is an experiment composed of two systems, the Large Photomultiplier (LPMT) and the Small Photomultiplier (SPMT). Both of the system observe the same physics event inside of the same medium but they differ in their photo-coverage, respectively 75.2% and 950 2.7%, their dynamic range (see section 2.2.2, a thousands versus a few dozen, and their front-end 951 electronics (see section 2.2.2).

952 They are complementary in their strengths and weaknesses and support each other. One important 953 point is their differences in expected resolution, the LPMT system outperform largely the SPMT 954 system but is subject to effects such as charge non linearity [28] that could bias the reconstruction, 955 effect that the SPMT system is impervious to. This topic will be studied in more detail in chapter 7. 956 Also, due to the dynamic range of the LPMT, in case of high energy and high density event such as 957 core-collapse supernova, the LPMT system could saturate and the lower photo-coverage become a 958 benefit.

959 Thus, although event reconstruction algorithm and physics analysis combines both LPMT and SPMT 960 systems, individual approach are key studies to understand the detector and ensure their reliability. 961 This topic will also be studied in more details in chapter 7. The subject of this chapter is to propose 962 a machine learning algorithm for the SPMT reconstruction based on Convolutional Neural Network 963 (CNN).

964 **4.1 Motivations**

965 As explained in chapter 3, Machine Learning (ML) algorithms shine when modeling highly dimensional 966 data from a given dataset. In our case, we have access to complete monte-carlo simulation of 967 our detector to produce arbitrary large datasets that could represent multiple years of data taking. 968 Ideally ML algorithms would be able to consider the entirety of the information in the detector 969 and converge on the best parameters to yield optimal results, while classical methods where the 970 algorithms could be biased by the prior knowledge of the detector and physics processes. To study 971

973 this potential phenomena, we will compare our machine algorithm to a classical reconstruction
 974 method developed for energy and vertex reconstruction [65].

975 We have access to a very detailed simulation of the detector (section 2.5) that will allow us to simulate
 976 arbitrary large dataset of data while giving access to the all the physics parameters of the event. Those
 977 parameters include the target of our reconstruction algorithms: the vertex and position at with the
 978 event happened. As introduced above, we hope that the ML algorithm will be able to used all the
 979 informations in the event, meaning that potential mismodelings in our simulation could be exploited
 980 by the algorithm. This specific subject will be studied in chapter 6.

981 4.2 Method and model

982 One of simplest way to look at JUNO data is to consider the detector as an array of geometrically
 983 distributed sensors on a sphere. Their repartition is almost homogeneous, on this sphere surface
 984 providing an almost equal amount of information per unit surface on this sphere. It is then tempting
 985 to represent the detector as a spherical image with the PMT in place of pixel. Two events with two
 986 different energy or position would produce two different images.

987 The most common approach in machine learning for image processing and image recognition is the
 988 Convolutional Neural Network (CNN). It is widely used in research and industry [56, 66–68] due to
 989 its strengths (see section 3.2.2) and has proven its relevance in image processing.

990 Some CNN are developed to process spherical images [69] but for the sake of simplicity and as a
 991 first approach we decided to go with a planar projection of the detector, approach that has proven its
 992 efficiency using the LPMT system (see section 2.6.3). The details about this planar projection will be
 993 discussed in section 4.2.2.

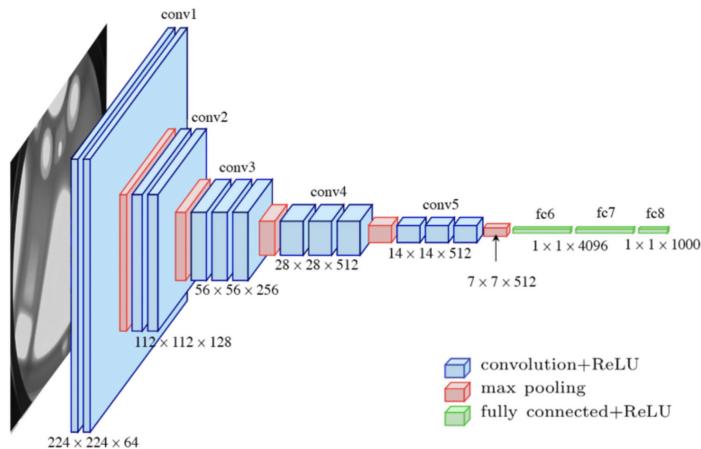


FIGURE 4.1 – Graphic representation of the VGG-16 architecture, presenting the different kind of layer composing the architecture.

994 4.2.1 Model

995 The architecture we use is derived from the VGG-16 architecture [56] illustrated in figure 4.1. We
 996 define a set of hyperparameters that will define the size, complexity and computational power of the
 997 NN. The chose hyperparameters are detailed below and their values are presented in table 4.1.

- 998 — N_{blocks} : the number of convolution blocks, a block being composed of two convolutional
999 layers with 3×3 filters using ReLU activation function, a 3×3 max-pooling layer (except for
1000 the last block) and a dropout layer.
- 1001 — $N_{channels}$: The number of channels in the first block. The number of channels in the subse-
1002 quent blocks are calculated using $N_{channels}^i = 2^i * N_{channels}$, $i \in [1..N_{blocks}]$.
- 1003 — **FCDNN configuration:** The result of the last convolution layer is flattened then fed to a
1004 FCDNN. Its configuration is expressed as a sequence of fully connected linear layer using
1005 the PReLU activation function. For example $2 * 1024 + 2 * 512$ is the sequence of 2 layers
1006 with a width of 1024 followed by 2 other layers with a width of 512. Finally the last layer
1007 is a 4 neurons wide linear layers without activation function. Each neurons of the last layer
1008 represent a component of the interaction vertex: Energy, X, Y, Z.
- 1009 — **Loss:** The loss function. In this work we study two different loss function $(E + V)$ and $(E_r +$
1010 $V_r)$ detailed below.

We explore in this work two different activation functions

$$(E + V)(E, x, y, z) = \left\langle (E - E_{true})^2 + 0.85 \sum_{\lambda \in [x, y, z]} (\lambda - \lambda_{true})^2 \right\rangle \quad (4.1)$$

$$(E_r + V_r)(E, x, y, z) = \left\langle \frac{(E - E_{true})^2}{E_{true}} + \frac{10}{R} \sum_{\lambda \in [x, y, z]} (\lambda - \lambda_{true})^2 \right\rangle \quad (4.2)$$

1011 where R is the radius of the CD. With the energy in MeV and the distance in meters, we use the factor
1012 0.85 and 10 to equilibrate the two term of the loss function so they have the same magnitude.

- 1013 — The loss function $(E + V)$ is close to a simple Mean Squared Error (MSE). MSE is one of the
1014 most basic loss function, the derivative is simple and continuous in every point. It is a strong
1015 starting point to explore the possibility of CNNs.
- 1016 — $(E_r + V_r)$ can be see as a relative MSE.

1017 The idea is that: due to the inherent statistic uncertainty over the number of collected Number of
1018 Photo Electrons (NPE), the absolute resolution $\sigma(E - E_{true})$ will be larger at higher energy than at
1019 low energy. But we expect the *relative* energy resolution $\frac{\sigma(E - E_{true})}{E_{true}}$ to be smaller at high energy than
1020 lower energy as illustrated in figure 2.19. Because of this, by using simple MSE the most important
1021 part in the loss come from the high energy part of the dataset whereas with a relative MSE, the most
1022 important become the low energy events in the dataset. We hope that by using a relative MSE, the
1023 neural network will focus on low energy events where the reconstruction is considered the hardest
1024 part of the dataset.

N_{blocks}	{2, 3, 4}
$N_{channels}$	{32, 64, 128}
FCDNN configurations	$2 * 1024$ $2 * 2048 + 2 * 1024$ $3 * 2048 + 3 * 512$ $2 * 4096$
Loss	{ $E + V$, $E_r + V_r$ }

TABLE 4.1 – Sets of hyperparameters values considered in this study

1025 Each combination of those hyperparameters (for example ($N_{blocks} = 2$, $N_{channels} = 32$, FCDNN =
1026 $(2 * 1024)$, Loss = $(E + V)$)), subsequently designated as configurations, is then tested and compared
1027 to each other over an analysis sample. We cannot use the mean loss because we consider multiple
1028 loss functions, there is no guarantee that comparison of their numerical value will be meaningful.
1029 We use multiple observables to rank the performances of each configuration:

- 1030 — The mean absolute energy error $\langle E \rangle = \langle |E - E_{true}| \rangle$. It is an indicator of the energy bias of our
1031 reconstruction.

- The standard deviation of the energy error $\sigma E = \sigma(E - E_{true})$. This is the indicator on our precision in energy reconstruction.
- The mean distance between the reconstructed vertex and the true vertex $\langle V \rangle = \langle |\vec{V} - \vec{V}_{true}| \rangle$. This is an indicator of the bias and precision of our vertex reconstruction.
- The standard deviation of the distance between the true and reconstructed vertex $\sigma V = \sigma |\vec{V} - \vec{V}_{true}|$. This is an indicator of the precision in our vertex reconstruction.

4.2.2 Data representation

This data is represented as 240×240 images, equivalent to third order tensor, with a charge Q channel and a time t channel. The SPMTs are then projected on the plane as illustrated in figure 4.2a. The x position is proportional to θ and the y position is defined by $\phi \sin \theta$ in spherical coordinates. $\theta = 0$ is defined as being the top of the detector and $\phi = 0$ is defined as an arbitrary direction in the detector. In practice, this is the $\phi = 0$ given by the MC simulation.

$$x = \left\lfloor \frac{\theta \cdot H}{\pi} \right\rfloor, \theta \in [0, \pi] \quad (4.3)$$

$$y = \left\lfloor \frac{(\phi + \pi) \sin \theta \cdot W}{2\pi} \right\rfloor, \phi \in [-\pi, \pi], \theta \in [0, \pi] \quad (4.4)$$

where H is the height of the image, W the width of the image and $(0, 0)$ the top left corner of the image.

When two SPMTs are in the same pixel, the charges are summed and the lowest of the hit-time is chosen. The SPMTs being located close to each other, we expect the time difference between two successive physics signals, two photons being collected, to be small. The first hit time is chosen because it can be considered as the relative propagation time of the photons that went the "straightest", i.e. that went under the less perturbation of the two. The only potential problem in using this first time come from the Dark Noise (DN). Its time distribution is uniform over the signal and could come before a signal hit on the other SPMT in the pixel. In that case, the time information in the pixel become irrelevant and we lose the timing information for this part of the detector. As illustrated in figure 4.2a the dimension have been chosen optimized so that at most two SPMTs are in the same pixel while keeping the number of empty pixels relatively low to prevent this kind of issue.

While it could be possible to use larger images (more pixel) to prevent overlapping, keeping image small images gives multiple advantages:

- As presented in section 4.2.1, the convolution filter we use are 3×3 convolution filter, meaning that if SPMTs would be separated by more than one pixel, the first filter would only see one SPMT per filter. This behavior would be kind of counterproductive as the first convolution block would basically be a transmission layer and would just induce noise in the data.
- It keep the network relatively small, while this do not impact the convolution layers, the flatten operation just before the FCDNN make the number parameters in the first layer of it dependent on the size of the image.
- It reduce the number of empty pixel in the image.

The question of empty pixel is an important question in this data representation. There is two kind of empty pixel in the data.

The first kind is pixel that contain a SPMT but the SPMT did not get hit nor registered any dark noise during the event. In this case, the charge channel is zero, which have a physical meaning but then come the question of the time layer. One could argue that the correct time would be infinity (or the largest number our memory allows us) because the hit "never" happened, so extremely far from the time of the event. This cause numerical problem as large number, in the linear operation that are

happening in the convolution layers, are more significant than smaller value. We could try to encode this feature in another way but no number have any significance due to our time being relative to the trigger of the experiment so -1 for example is out of question. Float and Double gives us access to special value such as NaN (Not a Number) [70] but the behavior is to propagate the NaN which leaves us with NaN for energy and position. We choose to keep the value 0 because it's the absorbing element of multiplication, absorbing the "information" of the parameter it would be multiplied by. It also can be thought as no activation in the ReLU activation function.

The second kind of pixel is pixel that do not represent parts of the detector such as the corners of the images. The question is basically the same, what to put in the charge and the time channel. The decision is to set the charge and time at 0 following the reasoning presented above. It's important to keep in mind that the fact that a part of the detector that has not been hit is also an information: There is no signal in this part of the detector. This problematic will be explored in more details in chapter 5.

Another problematic that happens with this representation, and this is not dependent of the chosen projection, is the deformation in the edges of the image and the loss of the neighbouring information in the for the SPMTs at the edge of the image $\phi \sim 180^\circ$. This deformation and neighbouring loss could be partially circumvented as explained in section 4.4

4.2.3 Dataset

In this study we use one million events coming from the full JUNO official monte-carlo simulations J23.0.1-rc8.dc1 (released the 7th January 2024). To put in perspective, the expected IBD rate in JUNO is 47 / days. Taking into account the calibration time, and the source reactor shutdown, it amounts to $\sim 94'000$ IBD events in 6 years. With this million of event, we are training the equivalent of ~ 10 years of data. With this amount we reach a density of $4783 \frac{\text{event}}{\text{m}^3 \cdot \text{MeV}}$, meaning our dataset is representative of the multiple event scenarios that could be happening in the detector.

While we expect and hope the monte-carlo simulation to give a realistic representation of the detector, there could be effects, even after the fine-tuning on calibration data, that the simulation cannot handle. Thus, once the calibration will be available, we will need to evaluate, and if needed retrain, the network on calibration data to establish definitive performances.

The data used during this analysis is monte carlo data using the official JUNO simulation software (see section 2.5 for details). The simulated data is composed of positron events, uniformly distributed in the CD volume and in kinetic energy over $E_k \in [0; 9]$ MeV producing a deposited energy $E_{dep} \in [1.022; 10.022]$ MeV. This is done to mimic the signal produced by the IBD prompt signal. Uniform distribution are used so that the CNN does not learn a potential energy distribution, favoring some part of the energy spectrum instead of others.

Those events can be considered as "optimistic" as there is no pile-up with potential background or other IBD.

The dark noise and time window come from the monte-carlo simulation of the electronic and trigger system. The dark noise rate used in this study is coming from studies and calibration of the PMTs outside of the experimental setup [25, 71].

4.2.4 Data characteristics

4.3 Results

Before presenting the results, let's discuss the different observables.

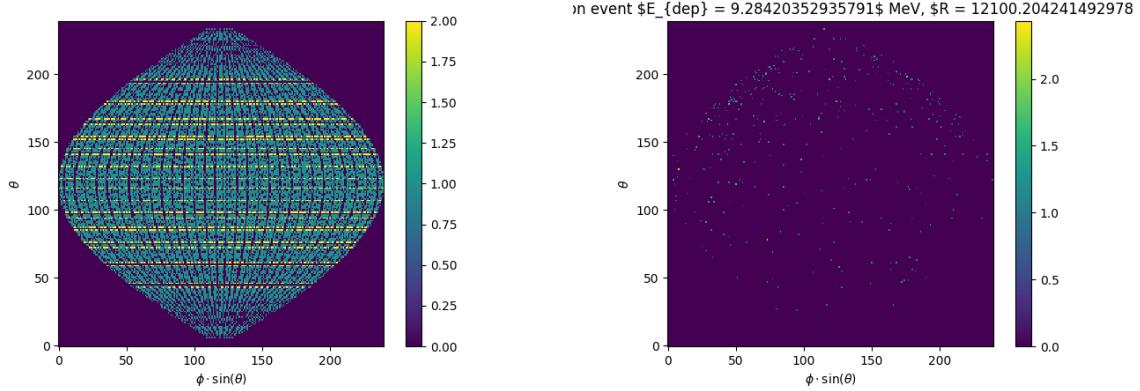


FIGURE 4.2

1115 The event are considered point like in this study. The target truth position, or vertex, is the mean po-
 1116 sition of the energy deposits of the positron and the two annihilation gammas. Due to the symmetries
 1117 of the detector, we mainly consider and discuss the bias and precision evolution depending of the
 1118 radius R but we will still monitor the performances depending of the spheric angle θ and ϕ . From the
 1119 detector construction and effect we expect relative important dependencies in radius thanks to the
 1120 TR area effect presented in section 2.6 and the possibility for the positron or the gammas to escape
 1121 from the CD for near the edge events. We also expect dependence in θ , the top of the experiment
 1122 being non-instrumented due to the filling chimney. It is also to be noted that the events in the dataset
 1123 are uniformly distributed in the CD, and so are uniformly distributed in R^3 and ϕ . The θ distribution
 1124 is not uniform and we will have more event for $\theta \sim 90^\circ$ than $\theta \sim 0^\circ$ or $\theta \sim 180^\circ$.

1125 We define multiple energy in JUNO:

- 1126 — E_ν : The energy of the neutrino.
- 1127 — E_k : The kinetic energy of the resulting positron from the IBD.
- 1128 — E_{dep} : The deposited energy of the positron and the two annihilation gammas.
- 1129 — E_{vis} : The equivalent visible energy, so E_{dep} after the detector effect such as the absorption of
 scintillation photons by the LS and the LS response non-linearity.
- 1131 — E_{rec} : The reconstructed energy by the reconstruction algorithm. The expected value depend
 on the algorithm we discuss about. For example the algorithm presented in section 2.6 is
 reconstructing E_{rec} while the ones presented in section 2.6.3 reconstruct E_{dep} .

1134 In this study, we will set E_{rec} as our target for energy reconstruction. This choice is motivated by the
 1135 ease with which we can retrieve this information in the monte-carlo data while E_{vis} is less trivial to
 1136 retrieve.

1137 4.4 Prospect

1138 4.5 Conclusion

1139 Introduction next chapter

1140 **Chapter 5**

1141 **Graph representation of JUNO for IBD
reconstruction with the LPMT system**

¹¹⁴³ **Chapter 6**¹¹⁴⁴ **Reliability of machine learning
methods**
¹¹⁴⁵

¹¹⁴⁶ “*Psychohistory was the quintessence of sociology; it was the science of human behavior reduced to mathematical equations. The individual human being is unpredictable, but the reactions of human mobs, Seldon found, could be treated statistically*”

Isaac Asimov, Second Foundation

¹¹⁴⁷ **Chapter 7**

¹¹⁴⁸ **Joint fit between the SPMT and LPMT
spectra**

¹¹⁴⁹

1150 Chapter 8

1151 Conclusion

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¹²⁵⁷ List of Abbreviations

ACU	Automatic Calibration Unit
BDT	Boosted Decision Tree
CD	Central Detector
CLS	Cable Loop System
CNN	Convolutional NN
DNN	Deep NN
DN	Dark Noise
FCDNN	Fully Connected Deep NN
GNN	Graph NN
GT	Guiding Tube
IBD	Inverse Beta Decay
IO	Inverse Ordering
JUNO	Jiangmen Underground Neutrino Observatory
LPMT	Large PMT
LR	Learning Rate
LS	Liquid Scintillator
MC	Monte Carlo simulation
ML	Machine Learning
MSE	Mean Squared Error
NMO	Neutrino Mass Ordering
NN	Neural Network
NO	Normal Ordering
NPE	Number of Photo Electron
OSIRIS	Online Scintillator Internal Radioactivity Investigation System
PE	Photo Electron
PMT	Photo-Multipliers Tubes
PReLU	Parametrized Rectified Linear Unit
ROV	Remotely Operated under-LS Vehicle
ReLU	Rectified Linear Unit
ResNet	Residual Network
SGD	Stochastic Gradient Descent
SPMT	Small PMT
TAO	Taishan Antineutrino Oservatory
TR Area	Total Reflexion Area
TTS	Time Transit Spread
TT	Top Tracker
WCD	Water Cherenkov Detector

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