



## *Escherichia coli* as a model active colloid: A practical introduction



Jana Schwarz-Linek<sup>a</sup>, Jochen Arlt<sup>a</sup>, Alys Jepson<sup>a</sup>, Angela Dawson<sup>a</sup>, Teun Vissers<sup>a</sup>,  
Dario Miroli<sup>a,b</sup>, Teuta Pilizota<sup>b</sup>, Vincent A. Martinez<sup>a,\*</sup>, Wilson C.K. Poon<sup>a,\*</sup>

<sup>a</sup> SUPA and School of Physics & Astronomy, The University of Edinburgh, James Clerk Maxwell Building, Peter Guthrie Tait Road, Edinburgh EH9 3FD, UK

<sup>b</sup> School of Biological Sciences and Centre for Synthetic and Systems Biology, The University of Edinburgh, Darwin Building, Max Born Crescent, Edinburgh EH9 3BF, UK

### ARTICLE INFO

#### Article history:

Received 1 June 2015

Received in revised form 17 July 2015

Accepted 19 July 2015

Available online 30 July 2015

#### Keywords:

*Escherichia coli*

Active colloids

Motility

Differential dynamic microscopy

Metabolism

Bioenergetics

Proton motive force

### ABSTRACT

The flagellated bacterium *Escherichia coli* is increasingly used experimentally as a self-propelled swimmer. To obtain meaningful, quantitative results that are comparable between different laboratories, reproducible protocols are needed to control, 'tune' and monitor the swimming behaviour of these motile cells. We critically review the knowledge needed to do so, explain methods for characterising the colloidal and motile properties of *E. coli* cells, and propose a protocol for keeping them swimming at constant speed at finite bulk concentrations. In the process of establishing this protocol, we use motility as a high-throughput probe of aspects of cellular physiology via the coupling between swimming speed and the proton motive force.

© 2015 Elsevier B.V. All rights reserved.

Some time ago, our lab wanted to culture motile bacteria as 'model active colloids'. We obtained a strain of *Escherichia coli* with the full complement of motility genes and a culturing protocol from a local microbiologist. For some time, we thought we were experimenting with motile *E. coli*, until one day we checked in the microscope. Few, if any, of the cells were swimming! So we set out to learn how to modify the standard protocol to optimise motility by collating literature, talking to other researchers and trial and error; we also implemented differential dynamic microscopy (DDM) to quantify motility.

This article reviews what we have learnt. Some of the material is previously known, but seldom critically discussed in one place. We have explained some basic bacterial bioenergetics and genetics, because physical scientists can use *E. coli* and collaborate with biologists more effectively if these topics are understood. Much of the material is new, arising from using DDM to quantify motility. While we aim primarily at researchers working on active colloids [1], this article should also be useful to others studying motility biophysics [2].

From the outset, we refer to various culture media (BMB, TB, LB) and protocols, and freely use terminology related to molecular biology (plasmid, gene names, etc.) and measurement techniques (OD, DDM, etc.). Readers should refer to Section 5 on matters of cell culture, Sections 3 and 4 for methodology, and Sections 10 and Appendix C for biological jargon. We also provide a table of symbols in Appendix F. Data for all the figures in this manuscript can be access on the Edinburgh DataShare repository [163].

### 1. *E. coli* as a model active colloid

'Active colloids' [1] are  $\approx 5\text{ nm} - 5\text{ }\mu\text{m}$  particles suspended in a fluid that consume 'fuel' to propel themselves, such as motile bacteria or various synthetic 'colloidal swimmers' [3,4]. Active colloids are out of thermal equilibrium even without external driving. Unusual phenomena displayed by such systems, such as 'negative viscosity increment' [5,6] and 'rectification' [7], pose a 'grand challenge' to statistical mechanics [8,9]. The self assembly of active colloids, possibly mixed with passive particles, may provide routes to new 'smart materials' [10].

Progress in this new area will be greatly facilitated by experimental data from 'model systems' that can most 'cleanly' confront theory and simulations. Historically, well-characterised model passive colloids have enabled progress at critical points, from sedimentation equilibrium [11,12] to hard-sphere crystallisation [13] and

\* Corresponding authors.

E-mail addresses: [vincent.martinez@ed.ac.uk](mailto:vincent.martinez@ed.ac.uk) (V.A. Martinez), [w.poon@ed.ac.uk](mailto:w.poon@ed.ac.uk) (W.C.K. Poon).

sticky-sphere glass transitions [14]. The minimal requirements for a 'model passive colloid' include reproducible synthesis/preparation, known particle size/shape distributions, quantifiable and 'tuneable' interparticle interactions and accurate particle volume fractions [15–18]. For model active colloids, we also need propulsion mechanisms that are understood and 'tuneable', and knowledge of motility-specific interparticle interactions, e.g., via hydrodynamics. By these criteria, there is as yet no ideal model active colloid.

*Escherichia coli* is widely used in active colloids research. Its self propulsion is understood in essence [19–21], while propulsion mechanisms in many synthetic swimmers are still debated [22,23]. However, for *E. coli* to become a fully-fledged model, reproducibility, control and characterisation need to be improved. We address these issues in this article.

## 2. *E. coli* as a bacterium

*Escherichia coli* [24] is the best understood living organism on earth today. It is often the model of choice for understanding molecular biological processes: 'Tout ce qui est vrai pour le Colibacille est vrai pour l'éléphant' [25].

The cytoplasm of *E. coli* is enclosed by two lipid bilayer membranes. Between them is a peptidoglycan 'cell wall' (or 'periplasmic space'), a network of poly-sugars ('glycans') linked by short peptides. Within the species known as *E. coli* there are many genetic variants, or 'strains' (so that 'bacterial species' is a problematic concept [26]). The genome of strain MG1655 was first to be sequenced [27]; today, the genome of 2085 strains are available in GenBank [28], although  $\approx 30\%$  of the genes remain of unknown function. The best single source of information on *E. coli* (and its close relative *Salmonella*) is a two-volume 'bible' [29] now updated digitally [30].

Each cell carries multiple helical flagella [31,32] powered by rotary motors [2,19] embedded in the membranes and cell wall [33]. When all flagella rotate counterclockwise (CCW) (viewed from behind), they bundle to propel the cell forward. When a motor reverses to clockwise (CW), its flagellum unbundles and the cell tumbles. Wild-type (WT) cells run ( $\approx 1$  s), tumble ( $\approx 0.1$  s), and run again in a more or less random new direction [34]. If successive sensings of the environment return an increasing concentration of an 'attractant', the rate of CCW  $\rightarrow$  CW  $\rightarrow$  CCW switching (i.e. the tumble rate) is reduced, giving a biased random walk, which constitutes chemotaxis.

Some *E. coli* strains used for active colloids and related research are listed in Table B.2, Appendix B. Many were derived from K-12, a 'debilitated' laboratory strain [35,36] that does not normally infect humans and is classified as a category 1 biohazard (=safe for laboratory use) [37]. Mutants in which one or more key chemotaxis genes have been deleted, e.g.,  $\Delta cheY$ , lose the ability to tumble. These 'smooth swimmers' simply swim forward, with a persistence time set by their inverse rotational diffusivity  $D_r^{-1} = 2\pi\eta L\sigma^2/k_B T \approx 5$  s, where we have modelled a cell body (length  $l$ ) with its flagella bundle as an ellipsoid of major and minor axes  $L \approx 10 \mu\text{m}$  ( $=l + \text{flagella}$ ) and  $\sigma \approx 1 \mu\text{m}$ .

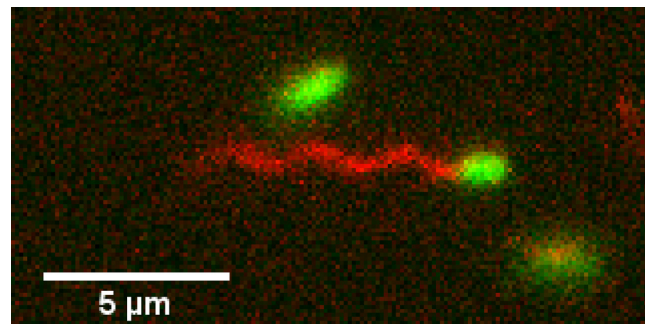
## 3. *E. coli* as a colloid

The cell body of an *E. coli* bacterium is a net negatively charged spherocylinder (diameter  $\sigma$ , pole-to-pole length  $l$ ).<sup>1</sup> For mostly unknown reasons [41,42], bacterial shapes are highly conserved<sup>2</sup>: shape mutations tend to be lethal. In this section, we review the

**Table 1**

Average length and standard deviation of 30 cell bodies of *E. coli* AB1157 grown in LB/TB at 30 °C/37 °C to late-exponential (OD = 0.5, see Footnote 5) or stationary (16 h growth) phase. Note that these cells have not been washed in BMB. Washing, however, makes little difference.

	LB <sub>30,stat</sub>	TB <sub>30,stat</sub>	LB <sub>37,stat</sub>	TB <sub>37,stat</sub>
$l$ ( $\mu\text{m}$ )	1.7	1.5	1.6	1.7
$\delta l$ ( $\mu\text{m}$ )	0.4	0.3	0.3	0.5
$\delta l/l$	0.24	0.2	0.19	0.29
	LB <sub>30,0.5</sub>	TB <sub>30,0.5</sub>	LB <sub>37,0.5</sub>	TB <sub>37,0.5</sub>
$l$ ( $\mu\text{m}$ )	4.4	2.3	4.4	2.4
$\delta l$ ( $\mu\text{m}$ )	1.2	0.6	1.0	0.6
$\delta l/l$	0.27	0.26	0.23	0.25



**Fig. 1.** Visualisation of both body and flagella bundle of swimming *E. coli* from a real-time movie. Flagella are fluorescently labelled using Alexa546 and the body with a plasmid expressing GFP (AD1 pHC60, Table B.2). Using epi-fluorescence microscopy these can be imaged into separate channels at high enough frame rates to avoid motion blur. This image is part of a 100 fps movie acquired using an Orca Flash 4.0 camera on a Nikon TI microscope with a PA60 $\times$ /1.4 OIL objective using a Cairn Optosplit II.

single-particle 'colloidal' properties of *E. coli* and how to estimate cell concentrations.

### 3.1. Shape and size

#### 3.1.1. Microscopy

The shape and size of *E. coli* cell bodies can be obtained using microscopy or scattering (compared in [44]). Phase-contrast microscopy of strain AB1157 harvested using our standard protocol and washed into Bergs motility buffer (BMB) gave  $\langle l \rangle = 2.4 \pm 0.6 \mu\text{m}$ , where the uncertainty is a polydispersity. At mid-cell,  $\langle \sigma \rangle = 0.86 \pm 0.07 \mu\text{m}$ , where the uncertainty represents measurement errors. Importantly, these 'vital statistics' depend on strain, growth conditions and the time of harvest, Table 1. The latter is important because bacteria grow in well defined stages: from rapid division in an exponential phase (cell number density  $n$  grows exponentially with time), through a stationary phase ( $n = \text{constant}$ ), to a death phase ( $dn/dt < 0$ ). We find that stationary-phase cells tend to be less polydisperse in length than mid-exponential cells, as found previously [45]. There may also be some dependence of  $\sigma$  on  $l$  [46].

It is often desirable to visualise the flagella bundle, which renders the cell highly anisotropic and introduces strong steric effects [47,48]. If the cell body and flagella are labelled with the same dye [49,50], fluorescence from the former will dominate. To visualise the flagella alone, the *fliC* gene<sup>3</sup> can be mutated to facilitate binding to Alexa Fluor C5 maleimide (e.g., Alexa546) [51]. If cells carry GFP-plasmid pHC60 (Table B.2), then the cell body and flagella can

<sup>1</sup> For physicochemical 'vital statistics' of *E. coli*, see [38–40].

<sup>2</sup> It is possible to turn bacteria into silica colloids [43], opening up a novel route to synthetic colloids with many shapes [41].

<sup>3</sup> This encodes flagellin, the protein units used to build up flagella filament.

visualised simultaneously, Fig. 1. Note that such mutants should always be checked for potential changes to their motility and other phenotypes.

### 3.1.2. Scattering

Cell shape and size can also be studied by static light scattering either as an average [52–57] or individually in flow cytometry [58]. Given the refractive index [59], models can be fitted to scattering data. Typically, an *E. coli* cell (cytoplasm + cell wall) is modelled as a core-shell ellipsoid [52,60,61], giving interference fringes that survive polydispersity in  $l$  because there is very little variation in either  $\sigma$  or periplasm thickness [53,62]. How optical inhomogeneities due to the chromosome affect dynamic light scattering, which also gives information on size and shape, has been studied [63]. Polarised scattering gives extra information [64,65].

### 3.2. Sedimentation

The density of *E. coli* cells is variable:  $1.08 \lesssim \rho \lesssim 1.16 \text{ g cm}^{-3}$  [66–68]. The density difference with the solvent,  $\Delta\rho \lesssim 0.1 \text{ g cm}^{-3}$ , causes cells to sediment. Sedimentation equilibrium defines the colloidal length scale [17]. In this dynamic equilibrium, concentration-driven diffusive flux from the bottom is balanced by a sedimentation flux from the top. In a dilute suspension, the number density profile with height is

$$n(z) = n(0)e^{-z/\ell_g}, \quad (1)$$

where the gravitational height  $\ell_g$  is given by

$$\ell_g = \frac{D_0}{v_s}. \quad (2)$$

Here,  $D_0$  is the free-particle diffusivity, and  $v_s$  is the single-particle sedimentation speed. A colloid is suspendable against gravity, so that  $\ell_g$  has to be larger than its characteristic size.

For an axially symmetric *E. coli* cell body, we can define an average, isotropic friction coefficient  $\bar{\xi}$  [69], from which we obtain the isotropic sedimentation velocity

$$\bar{\xi}v_s = \Delta\rho g V_{\text{cell}} \Rightarrow v_s = \frac{\Delta\rho g V_{\text{cell}}}{\bar{\xi}}, \quad (3)$$

and the isotropic diffusivity

$$D = \frac{k_B T}{\bar{\xi}}, \quad (4)$$

so that from Eq. (2),

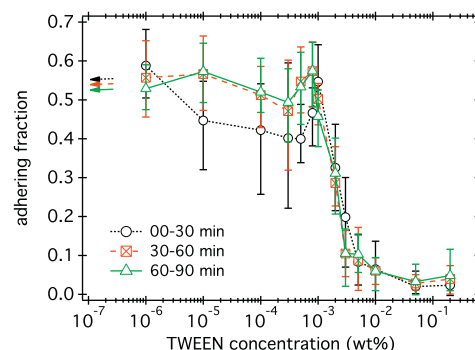
$$\ell_g^{(\text{coli})} = \frac{k_B T}{\Delta\rho g V_{\text{cell}}} \approx 4 \mu\text{m}, \quad (5)$$

for an *E. coli* with  $V_{\text{cell}} \gtrsim 1 \mu\text{m}^3$ , which is therefore colloidal.

Run-and-tumble *E. coli* has higher effective diffusivity  $D \sim v^2 \tau / 3 \approx 100 \mu\text{m}^2 \text{ s}^{-1}$ , with  $v$  and  $\tau$  being the swimming speed ( $\approx 15 \mu\text{m s}^{-1}$ ) and the average time between tumbles ( $\approx 1 \text{ s}$ ), respectively, so that  $\ell_g \sim 1 \text{ mm}$  [70]. This is  $\gtrsim 2 \times$  the thickness of our capillaries, so that gravity-induced inhomogeneities can be neglected in our experimental protocol. Note that motility coupled with growth can display rich phenomena [71].

### 3.3. Interactions: stability and surface adhesion

The surface properties of *E. coli* [72–74] are strain dependent [75]. Interpreting electrophoretic mobilities ( $\mu$ ) is non-trivial because surface macromolecules [73,75,74] and the periplasm are ion-permeable. Even with a ‘soft’ electrokinetic theory allowing for this [76,77], relating  $\mu$  to zeta potentials and surface structures is fraught.



**Fig. 2.** Fraction of bacteria sticking to the lower capillary wall in a  $\approx 4$  cell width layer at different concentrations of TWEEN 20 surfactant in BMB with 0.72 mM glucose for WT AB1157 at  $\approx 6 \times 10^7$  cells/ml. Points are averages over 0–30 min, 30–60 min and 60–90 min after loading. Arrows: adhering fractions with no surfactant.

Nevertheless, *E. coli* definitely carries a net negative charge under physiological conditions ( $-3 \lesssim \mu \lesssim -2 \mu\text{m cm V}^{-1} \text{ s}^{-1}$ ) largely due to ionised carboxylate and phosphate groups [72,75]. This charge is normally sufficient to confer colloidal stability in growth media or in BMB.<sup>4</sup> However, a strain that has behaved stably for many months may suddenly start aggregating due to ‘phase variation’ in the expression of the surface protein Ag43 (antigen 43) [78]. Each generation, there is an  $O(10^{-3})$  probability that a cell switches between ‘on’ and ‘off’ states of Ag43 expression. ‘On’ renders cells ‘sticky’ and leads to ‘autoaggregation’. The expression state is heritable and reversible.

When working with *E. coli* close to surfaces [49,79–81], it is important to minimise adhesion, which is far from understood. Various additives such as non-ionic surfactants or polymers prevent *E. coli* from sticking to untreated glass [49,81–83]. Fig. 2 quantifies this effect for a popular surfactant used for this purpose.

We imaged AB1157 cells in BMB within an optical slice of  $\approx 4$  cell width next to the bottom surface of an untreated glass capillary using a Mikrotrotron MC1362 camera and a Nikon Ti microscope with a PF60 $\times$ /0.7 objective. Automated software for identifying [84] and tracking rods was used to classify cells into ‘free’ (diffusing or swimming) and ‘adhering’ (stuck). The average fraction of adhering cells, extracted from many movies, dropped abruptly to nearly 0 at a TWEEN 20 concentration of  $\approx 10^{-2}$  wt%. Note that concentrations up to 0.2% have been used previously [81]. Preliminary DDM data (not shown) suggest that commercial TWEEN may contain impurities rapidly metabolisable by *E. coli*, so that high enough concentrations of TWEEN may significantly alter swimming behaviour.

### 3.4. Cell concentration

The definition of cell concentration and its determination are non-trivial. The cell body volume fraction,  $\phi$ , gives a better intuitive feel for the degree of ‘crowding’ than the number density of cells,  $n$ . The effective volume fraction, which takes into account the cell body (length  $l$ ) and flagella bundle with total length  $L \lesssim 10 \mu\text{m}$ , is  $\phi_{\text{eff}} \sim (L/l)^3 \phi \gtrsim 10^2 \phi$  if the cells are randomly oriented, so that they reach ‘overlap’ at  $\phi \lesssim 0.01$ .

The standard way to obtain  $n$  is from the spectrophotometric optical density (OD)<sup>5</sup> of cell suspensions, which is proportional to  $n$  for  $n \lesssim 10^9$  cells/ml (corresponding to  $\phi_{\text{eff}} \approx 0.2$ ). Calibration is by ‘viable plate counting’ (Section 5.3), and is found to depend on growth stage and other conditions.

<sup>4</sup> The screening length in BMB is  $\lesssim 1 \text{ nm}$ .

<sup>5</sup> The OD depends on the wavelength used; we use 600 nm throughout.

#### 4. Characterising *E. coli* motility

We turn next to characterising *E. coli* as self-propelled particles. The average swimming speed,  $\bar{v}$ , and the fraction,  $\beta$ , of non-swimmers that only diffuse,<sup>6</sup> can be measured using a high-throughput technique, differential dynamic microscopy (DDM) [86–88]. Experimental details as well as application to swimming algae [88] and passive magnetic rods [89] can be found elsewhere. Here we present a summary and focus on practical aspects of using DDM to characterize motile *E. coli*.

We take a low-magnification movie of a cell population, and calculate the squared modulus of the Fourier transform (F.T., reciprocal vector  $\mathbf{q}$ ) of the difference of two images separated by  $\tau$  in time, and averaged over starting times:

$$g(\mathbf{q}, \tau) = \langle |\text{F.T.}[\mathcal{I}(x, y; t + \tau) - \mathcal{I}(x, y; t)]|^2 \rangle. \quad (6)$$

If the image intensity  $\mathcal{I}(x, y; t)$  at pixel position  $(x, y)$  at time  $t$  is linearly related to the cell density at the corresponding sample position and time, then  $g(\mathbf{q}, \tau)$  is directly related to the density–density time correlation function,  $f(\mathbf{q}, \tau)$ , of the system, with  $q = |\mathbf{q}| \sim 2\pi(\text{image size})^{-1}$  to  $\pi(\text{pixel size})^{-1}$  being experimentally accessible. Indeed, under appropriate conditions,

$$g(q, \tau) = A(q)[1 - f(q, \tau)] + B(q). \quad (7)$$

Sometimes called the normalised dynamic structure factor or intermediate scattering function (ISF),  $f(q, \tau)$  is in principle (but not in practice [87]) also measurably by dynamic light scattering. We assume isotropy, so that  $g$  and  $f$  only depend on  $q$ . In Eq. (7),  $A(q)$  is related to the form and (static) structure factors of the bacterial suspension, and  $B(q)$  is instrument noise.

The ISF can be modelled once we know how cells move. For smooth-swimming *E. coli* in 3D mixed with non-motile cells,

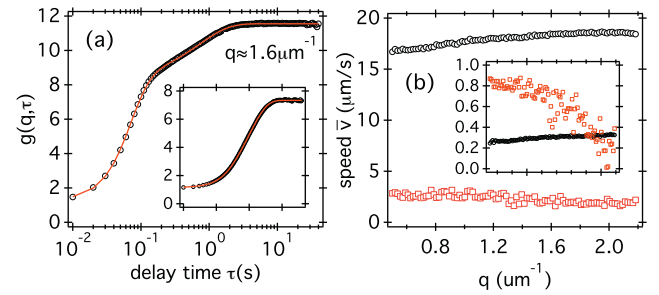
$$f(q, \tau) = \beta e^{-q^2 D \tau} + (1 - \beta) e^{-q^2 D \tau} \int_0^\infty P(v) \frac{\sin(qv\tau)}{qv\tau} dv, \quad (8)$$

where  $P(v)$  is the speed distribution of the swimmers, and  $D$  is the free translational diffusion coefficient, which is mainly determined by the diffusivity of the non-motile cells (but see footnote 6). We have ignored polydispersity in cell size (which induces a distribution in  $D$ ) for simplicity.

Self consistency demands that for straight swimmers mixed with diffusers, the dynamical quantities obtained by fitting Eq. (8) to data, in particular  $\bar{v} = \int_0^\infty v P(v) dv$ ,  $\beta$  and  $D$ , should be  $q$ -independent. Residual  $q$  dependence is nevertheless observed in practice, Fig. 3(b), e.g., due to tumbling<sup>7</sup> [88]; such  $q$  dependence is the main source of experimental uncertainties in the fitted parameters from DDM quoted in this work.

The average body rotation angular frequency,  $\bar{\Omega}$ , can be determined using another high-throughput technique, dark-field flicker microscopy (DFM) [91]. Essentially, the lowest peak in the Fourier transform of the time-dependent image intensity of a single cell in dark-field microscopy estimates  $\bar{\Omega}$ . In the low-Reynolds-number limit,  $\bar{v}/\bar{\Omega}$  depends solely on the geometry of the swimmers, and not on viscosity, provided that the medium is a Newtonian fluid [1]. Thus, non-constancy of  $\bar{v}/\bar{\Omega}$  indicates either changing cell geometry or non-Newtonian effects [91].

Note that the relative merits of DDM and real-space tracking for obtaining  $\bar{v}$  and  $\beta$  has been discussed in depth before [88].



**Fig. 3.** (a) Fitting DDM data  $g(q, \tau)$ , Eq. (6) and (7), using the ISF given by Eq. (8) at one  $q$  for WT AB1157 at  $5 \times 10^8$  cells/ml. Main plot: data for bacteria swimming fast enough for ballistic and diffusive motions to show up clearly as two processes, separated by a point of inflection at  $\tau \approx 0.4$  s. Inset (axis labels as main): the bacteria have slowed down to the extent that swimming and diffusion are no longer well separated. (b) Fitted speed  $v$  (main plot) and non-motile fraction  $\beta$  (inset, bottom axis as main) for the data shown in the main part of (a) (black) and in the inset of (a) (red). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

##### 4.1. Imaging protocol

We use  $d \approx 400 \mu\text{m}$  deep flat glass sample cells filled with  $\approx 150 \mu\text{l}$  of suspension and rendered air tight with petroleum jelly to prevent drift due to evaporation and stop replenishment of  $\text{O}_2$  after it is exhausted by respiring bacteria.<sup>8</sup> DDM shows that contact with petroleum jelly does not change ( $\bar{v}$ ,  $\beta$ ). When sealing samples into capillaries, it is important to avoid trapping air bubbles, which induce aerotactic motion in the cells.

Movies in phase-contrast illumination (PF10 $\times$ /0.3 at 100 frames/s,  $\sim 4000$  images at  $512 \times 512$  pixels  $\Rightarrow \sim 40$  s movies) are recorded in an inverted microscope (Nikon Ti) with a Mikrotron high-speed camera (MC 1362) and frame grabber (Inspecta 5, 1 Gb memory) at  $22 \pm 1^\circ\text{C}$ . Custom LabVIEW recording software controlling the microscope stage allows scanning of many samples (giving good averages) repeatedly (giving time resolution), giving data such as Figs. 4, 9 and 10.

##### 4.2. Observational complications

The majority of experiments to date using synthetic swimmers produced data at or near surfaces. *Escherichia coli* cells can swim in the bulk for extended time, although cells encountering a wall tend to be trapped there [92,82]. If performed sufficiently far from walls, DDM yields a genuine three dimensional speed distribution [87,88] via fitting an ISF to Eq. (7). Near a wall, the swimming is two-dimensional, and the  $\sin(qv\tau)/qv\tau$  in Eq. (7) is replaced by  $J_0(qv\tau)$ , the zeroth-order Bessel function. However, it may not be possible to distinguish between these functional forms unless  $P(v)$  is sufficiently narrow. Moreover, *E. coli* next to a wall swims in circles [93], giving  $q$ -dependent fitted dynamical quantities if the circle radii are comparable to the accessible range of  $2\pi/q$ . The complex motion of cells ‘tethered’ to walls further complicates data analysis. We therefore typically image at  $100 \mu\text{m}$  away from the bottom of our capillaries (cf. total cell length  $L \lesssim 10 \mu\text{m}$ ).

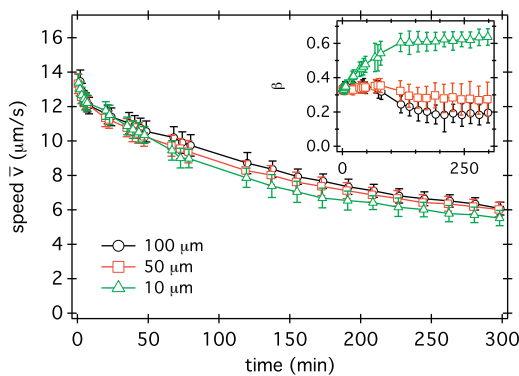
The sedimentation height of non-motile cells, Eq. (2), can be enhanced by 50% or more because swimmers enhance their diffusivity [85]. Nevertheless,  $l_g \approx \mathcal{O}(10 \mu\text{m})$ , so that non-swimmers will accumulate on the bottom on a time scale of  $d/v_s \lesssim 70$  min (using  $v_s \approx 0.1 \mu\text{m s}^{-1}$ , Eq. (3)). Measurement of WT swimmers  $10 \mu\text{m}$  from the bottom capillary surface, Fig. 4 inset, shows that  $\beta$  increases with time, until it saturates just over an hour ( $\approx d/v_s$ ) into the measurement. All non-motile cells have now sedimented into a

<sup>6</sup> Note that this diffusion is enhanced by the swimmers [85].

<sup>7</sup> Calculations of the effects of tumbling on the ISF exist [90].

<sup>8</sup> See Section 8, especially Fig. 10(b), to see why we want no  $\text{O}_2$  supply.





**Fig. 4.** Average swimming speed  $\bar{v}$  (main) and non-motile fraction  $\beta$  (inset, bottom axis as main) vs. time at three heights from bottom wall for WT AB1157 at  $5 \times 10^8$  cells/ml.

layer of thickness  $\approx l_g$ . Correspondingly,  $\beta$  is constant at any single bulk position (e.g., 50 or 100 μm from the bottom, Fig. 4 inset) until the sedimentation front passes (at some fraction of  $d/v_s$ ), whereupon  $\beta$  drops.

In a sealed capillary at our standard height of 100 μm from the bottom,  $\bar{v}$  decrease with time, Fig. 4. A priori, it is conceivable that this is because fast swimmers encounter surfaces and become trapped [92,82] more rapidly, thus leaving slower swimmers in the bulk. However, measurements at three heights, Fig. 4, give essentially identical results,<sup>9</sup> so that the decreasing  $\bar{v}(t)$  is *not* due to gradual kinematic accumulation at surfaces.

Drift gives ballistic motions that masquerade as swimming. Evaporative drift can be eliminated by proper sealing, but unavoidable drift occurs at short times due to loading. These transients decay on a time scale  $\gtrsim \rho L^2/\eta$  for a liquid with viscosity  $\eta$  and density  $\rho$ , where  $L$  is a characteristic length [94]. We used a 5 cm  $\times$  8 mm  $\times$  0.4 mm capillary throughout. The 8 mm dimension leads to resolvable speed changes on the  $\approx 1$  min time scale, which probably accounts for the initial fast decay in  $\bar{v}(t)$ , Fig. 4. Decreasing this dimension gives faster initial decay.

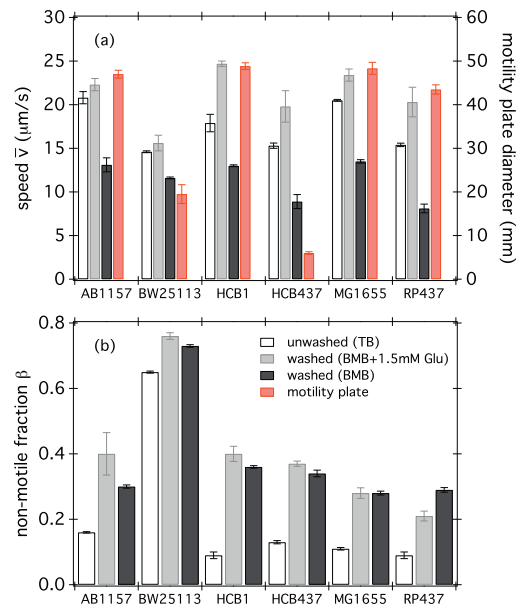
#### 4.3. Fitting lore

The decay of the ISF, Eq. (8), is due to diffusion (parameterised by  $D$ ) and ballistic motion (parameterised by  $v$ ). Satisfactory fitting of the data in the form of  $g(q, \tau)$  using Eqs. (7) and (8) to obtain  $D$  and  $v$  depends on convincingly decoupling these contributions, which is straightforward when the two decays have well-separated characteristic times, Fig. 3. Overlap of these two contributions occurs when their relaxation times approach, i.e.  $v \rightarrow qD$ , which, for WT *E. coli* in bulk ( $D \gtrsim 0.3 \mu\text{m}^2 \text{s}^{-1}$ ) and  $q \lesssim 2.2 \mu\text{m}^{-1}$  (using a 10 $\times$  objective), occurs at  $v \lesssim 1 \mu\text{m s}^{-1}$ . However, due to the distribution of speeds, this limit is underestimated. In practice, we can resolve  $\bar{v}$  down to 2–3  $\mu\text{m s}^{-1}$  in BMB, Fig. 3. Note also that if  $\beta \rightarrow 0$  or 1, the fitted value of  $D$  or  $v$ , respectively will be subject to large errors.

#### 4.4. Motility plates

When chemotactic strains are inoculated into ‘soft agar’ ( $\gtrsim 0.2$  wt%), they spread out in rings [95]. Non-chemotactic mutants [95], and the chemotactic WT when the agar is  $\gtrsim 0.3$  wt% [96], spread out instead as solid discs. The diameters of these rings or solid discs are often used to assay motility.

We inoculated 5  $\mu\text{l}$  drops of OD = 0.3 tryptone broth (TB)-grown *E. coli* into 0.3 wt% Luria broth (LB) agar plates, and observed solid



**Fig. 5.** (a) Average swimming speed  $\bar{v}$  (left axis) and (b) non-motile fraction  $\beta$  for selected *E. coli* strains, from one batch culture in each case, before washing (in TB, white) and after washing (in BMB + 1.5 mM glucose, grey or in BMB, black) at a cell density of  $5 \times 10^8$  cells/ml. Red bars in (a) are diameters of solid, circular colonies in 22-hour LB motility plates (right axis). Note that for HCB437, the cells did not penetrate the agar but stayed on the surface. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

discs (rather than rings), consistent with non-chemotactic spreading. We measured the diameter of the solid colony discs after 22 h of incubation at 30 °C, and compared these with  $\bar{v}$  measured from DDM in three different media, Fig. 5(a). For chemotactic WT cells, there is reasonable correlation between motility plates and DDM measurements in all three media, but especially in glucose-supplemented BMB.

The non-chemotactic spreading of growing cells can be modelled using the Fisher equation [97]. This predicts a front speed of  $u \sim 2\sqrt{D\alpha}$ , where  $\alpha$  is the growth rate, and  $D$  is the cells' effective diffusivity. The latter is  $\approx \bar{v}^2 \tau/3$ , with  $\tau$  being the average time between tumbles [98]. The colony diameter after time  $\Delta t$  is  $d = 2u\Delta t$  (ignoring the presence of a lag-time). Thus  $\bar{v} = (d/4\Delta t)\sqrt{3/\tau\alpha}$ . Using  $\tau \sim 1$  s [34] as an upper bound for  $\tau$  in agar [96],  $d \approx 48$  mm over  $\Delta t = 22$  h, Fig. 5(a), and taking  $\alpha \approx 0.7 \text{ h}^{-1}$  [96], we estimate  $\bar{v} \gtrsim 20 \mu\text{m s}^{-1}$ , which is not unreasonable in light of measured swimming speeds, Fig. 5(a).

Two caveats are in order. First, if the agar concentration allows chemotaxis, the dependence of the speed of the front (now a ‘ring’) on  $\bar{v}$  becomes more complex [96]. Secondly, we find that HCB437, a smooth swimmer, did not penetrate the agar; its minimal spread on the plate surface does *not* reflect  $\bar{v}$  in bulk media, especially since surface drying might play a role. Thus, motility plates must only be used to estimate  $\bar{v}$  with caution.

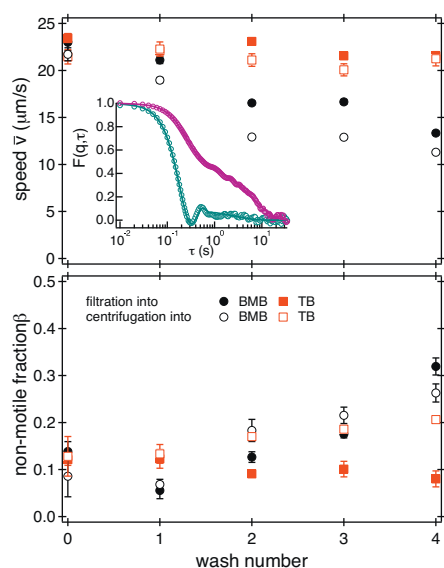
### 5. Preparing motile *E. coli*

We turn now to culturing bacterial swimmers. Readers needing to learn generic microbiological protocols should use standard manuals [99,100] and consult a microbiologist. We focus on extra procedures for obtaining *motile* cells.

#### 5.1. Culturing cells

Our protocol is based on that of H.C. Berg (Harvard), and involves Luria broth (LB, originally ‘lysogeny broth’ [101]) and tryptone broth (TB). TB is a mixture of amino acids from hydrolysing milk

<sup>9</sup> Note the longer time axis here than those shown later in Figs. 9 and 10.



**Fig. 6.** (a) Average swimming speed  $\bar{v}$  and (b) non-motile fraction  $\beta$  during a washing process from TB to BMB (black) and from TB to TB (red) of WT AB1157, using filtration (filled) or centrifugation (open) at final density of  $5 \times 10^8$  cells/ml. Inset: measured (symbols) and fitted (line)  $f(q, \tau)$  of these cells in TB before washing (turquoise, lower) and in BMB after 4 washing steps (purple, upper). Data sets are from one batch culture. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

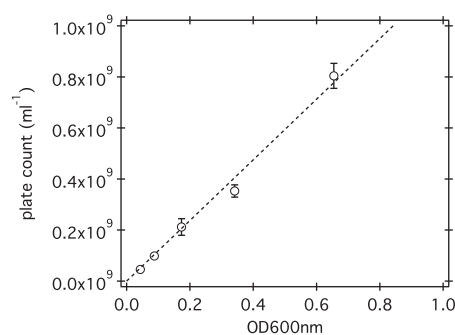
casein proteins; LB is tryptone broth plus yeast extract, the latter containing various carbohydrates not in TB.

We grow single colonies from frozen stocks on LB agar plates at 30 °C overnight. A single colony is transferred from a plate to 10 ml of liquid LB and incubated overnight ( $\approx 16$  h) at 30 °C<sup>10</sup>, shaken (for aeration) at 200 rpm. Finally, cells are diluted 100-fold into 35 ml of TB and grown for  $\sim 4$  h at 30 °C shaken at 200 rpm. DDM shows that transferring to TB from LB results in higher  $\bar{v}$ . As cell sizes (Table 1) and other properties change with growth phase, it is important that growth times for different experiments should be kept constant.

## 5.2. Washing cells

The next step is to transfer cells from TB to a ‘minimal medium’ with no exogenous nutrients. This prevents growth, which is an unwanted complication in the long run. More importantly, we find that cells use oxygen very quickly in TB, and once oxygen is exhausted,  $\bar{v}$  decreases dramatically. For WT AB1157 at OD = 0.3, this occurs within 10 min, which is not long enough for experiments. We therefore wash (by filtration) and transfer cells into Berg’s motility buffer (BMB), containing 6.2 mM  $K_2HPO_4$ , 3.8 mM  $KH_2PO_4$ , 67 mM NaCl, and 0.1 mM EDTA. Cells undergo 4 successive filtrations (=3 washes, see Appendix D) until the concentration of TB is diluted  $10^3$ -fold by BMB. After the last wash, the suspension on top of the filter is removed and transferred to a 50 ml centrifuge tube together with the filter, where cells are carefully suspended by rolling the tube. This yields 1–3 ml of highly concentrated cells (OD  $\sim 10$ ), allowing the preparation of many dilute samples.

Transfer from TB into BMB decreases  $\bar{v}$  and increases  $\beta$  for all strains studied, Fig. 5. This change in cellular motility is reflected in the ISF measured by DDM, Fig. 6(a) inset. In TB, the population of mostly motile cells (a single decay) has a narrow speed distribution (oscillations). In BMB, two processes are clearly visible: a sizeable fraction of non-motile (diffusive) cells coexists with



**Fig. 7.** Example of viable plate count vs. OD used for calibrating the latter. These values are typical for our culture protocol.

(ballistic) swimmers, whose speed distribution is now considerably wider (no oscillations).

Fig. 6 also shows that washing by filtration has non-trivial effects on  $\bar{v}$  and  $\beta$ , probably due to a combination of mechanical damage to flagella, regrowth (possible in TB) and change of medium. Pipetting may also have an effect. After 20 pipetting cycles in BMB using a 0.6 mm tip at normal expulsion speeds, we find that  $\bar{v}$  drops by  $\approx 50\%$  and  $\beta$  increases by a factor of  $\lesssim 3$ . It is not until the tip diameter has been increased to  $\approx 2$  mm that the effect on ( $\bar{v}$ ,  $\beta$ ) begins to saturate. Such flow-driven damage of *E. coli* flagella has been noted before [103].

If low cell concentrations are required or many strains need to be processed simultaneously, we use a bench-top centrifuge to wash small volumes ( $\lesssim 2$  ml) 3 times at moderate speeds (2 min at 8000 rpm) to produce  $\approx 1$  ml of OD  $\approx 0.3$  cell suspension. Fig. 6 shows that this procedure produces populations with higher  $\beta$  in TB, but it is considerably faster than filtration.

## 5.3. Counting cells

The most common protocol for relating OD to cell number is the ‘viable plate count’ [99,100]. Briefly, a sample is diluted to the relevant OD range (say, 0.05–1) using phosphate buffered saline (PBS).<sup>11</sup> One part of the sample is used to measure OD; another part is used to prepare 10-fold serial dilutions (down to  $10^{-6}$ ) by transferring 100  $\mu$ l sample to 900  $\mu$ l PBS with vigorous mixing between each dilution. 100  $\mu$ l of the final two dilutions are spread on replicate LB plates and incubated at 37 °C overnight. Next day, individual colonies can be seen. Each of these has arisen from a ‘colony forming unit’ (CFU), a shorthand for ‘probably but not certainly a single cell’. These are counted to arrive at the viable plate count in units of CFU/ml. The cell density of the original sample is then calculated taking into account the dilution. Plotting the OD vs. the cell density from plate counting, Fig. 7, gives the desired calibration.

## 6. E. coli energetics

Self-propulsion requires energy. The energetics of synthetic swimmers is partially understood [104,105]. Here we give a simplified introduction to *E. coli* bioenergetics. Details can be found in textbooks [106–111].<sup>12</sup>

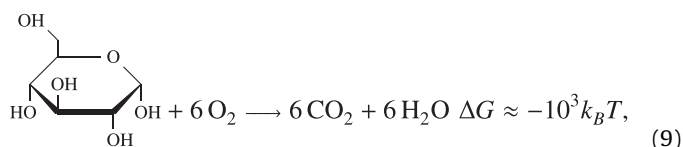
Bioenergy comes from transferring electrons from high-energy to low-energy bonds through a sequence of molecules. The energy is used to pump protons out of the cell, giving rise to a proton motive force (PMF), which powers swimming. In aerobic respiration, the

<sup>11</sup> PBS = 137 mM NaCl, 2 mM KCl, 10 mM  $Na_2HPO_4$  and 2 mM  $KH_2PO_4$ .

<sup>12</sup> An unusual, ‘coarse-grained’ introduction is given by Nobel laureate Christian de Duve [112]. See also tutorials at [113] under ‘The Microbe’ tab.

<sup>10</sup> Growth is optimal at  $T = 37$  °C, but lower  $T$  promotes swimming [19,102].

electron donors are various reduced foods rich in C–H bonds, while the final electron acceptor is typically oxygen. The electron pair in the C–H bond lowers its energy as it is transferred (notionally as  $\text{H}^-$ ) to an O–H bond (in water), where, because O is more electronegative than C, the electron pair interacts more strongly with positively charged nuclei than in C–H, and therefore has lower energy. The cell harnesses the energy released for chemical work. The ‘favourite’ supplier of high energy electrons for *E. coli* is  $\alpha$ -D-glucose ( $\text{C}_6\text{H}_{12}\text{O}_6$ ).<sup>13</sup> The full oxidation of this molecule in aerobic respiration generates  $\approx 2.8 \text{ MJ mol}^{-1}$  under physiological conditions

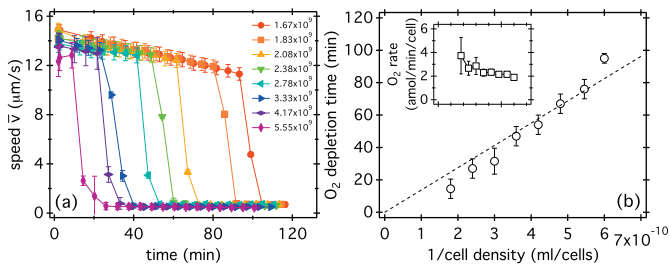


<sup>14</sup> When glucose is the sole carbon and energy source for *E. coli*, ~75% of it channels into the EMP pathway; the rest feeds into the hexose monophosphate (HMP) pathway, which mainly provides precursors for biosynthesis, i.e. anabolism [110]. However, the HMP pathway can also feed pyruvate into the TCA cycle [110], partly to provide anabolic precursors. Such ‘anabolic syphons’ together with overflow metabolism, reaction (12), mean that the stoichiometry in reaction (9) is, seldom, if ever, satisfied on average.

$$\text{C}_6\text{H}_{12}\text{O}_6 + 2\text{O}_2 \longrightarrow 2\text{CH}_3\overset{\text{O}}{\underset{\text{O}^-\text{H}^+}{\text{C}}} + 2\text{CO}_2 + 2\text{H}_2\text{O}. \quad (12)$$

Under starvation conditions (i.e. no exogenous nutrients), such as occurs in BMB, *E. coli* obtains energy by metabolising intracellular resources. Such 'endogenous metabolism' is complex [123,124] and poorly understood. For starved cells grown in TB, catabolism of free amino acids and free or RNA-derived ribose (a sugar) feeds into the TCA cycle [109], so that  $O_2$  or other terminal acceptors are needed.

*Escherichia coli* and other bacteria [125–127] can swim powered entirely by endogenous metabolism. We now probe the time dependence of  $\bar{v}$ . All data are obtained from DDM with WT AB1157 in BMB prepared using our standard protocol.



**Fig. 9.** (a) Average swimming speed  $\bar{v}$  vs. time for the capillary protocol in BMB at cell densities indicated in the legend (cells/ml). Lines guide the eye. The 'crash' in  $\bar{v}$  in each case is due to  $O_2$  exhaustion. Data sets from one batch culture. (b) Corresponding  $O_2$  exhaustion time as a function of inverse of cell density. Error bars relate to the accuracy of measuring the 'crash' time.

### 7.1. Observations and practical implications

Fig. 9(a) shows the average swimming speed as a function of time,  $\bar{v}(t)$ , in sealed capillaries at various cell densities,  $n$ . Cells slow down, approximately linearly with time. If  $\approx 10\%$  slowing down is tolerable, then there is a useable window of  $\approx 1$  h at  $n = 2 \times 10^9$  cells/ml, which ends in a 'crash'.

The time of this 'crash',  $t_c$ , drops with  $n$ ; we suggest that it is due to  $O_2$  depletion. Quantitatively,  $nt_c \approx \text{constant}$ , Fig. 9(b), as expected if the initial  $[O_2]$  and the  $O_2$  consumption rate per cell,  $Q$ , are invariant with  $n$ . Assuming that the BMB was initially saturated with  $O_2$ , and using literature solubilities [128], which are only very weakly dependent on ionicity [129], we estimate  $Q$  at each cell density, Fig. 9(b) inset, which is weakly dependent on  $n$ , rising from  $Q_{\text{endo}} \approx 2$  amol/min/cell at  $n = 1.67 \times 10^9$  cells  $\text{ml}^{-1}$  to  $Q_{\text{endo}} \lesssim 4$  amol/min/cell at  $n = 5.55 \times 10^9$  cells  $\text{ml}^{-1}$  ('a'  $\equiv$  'atto' =  $10^{-18}$ ). These values are comparable to those measured before using a different strain and under somewhat different conditions [123].<sup>15</sup>

The highest density data in Fig. 9 were obtained at  $n \approx 5 \times 10^9$  cells  $\text{ml}^{-1}$ , or  $\phi \approx 0.5\%$  of cell bodies. Extrapolating from  $Q$  at this  $n$ ,  $t_c \approx 1$  min at  $\phi \approx 5\%$ ; this is an upper bound, since  $Q$  probably increases with  $n$ . Oxygen exhaustion therefore precludes experimentation at  $\phi \gtrsim 1\%$  under these conditions.

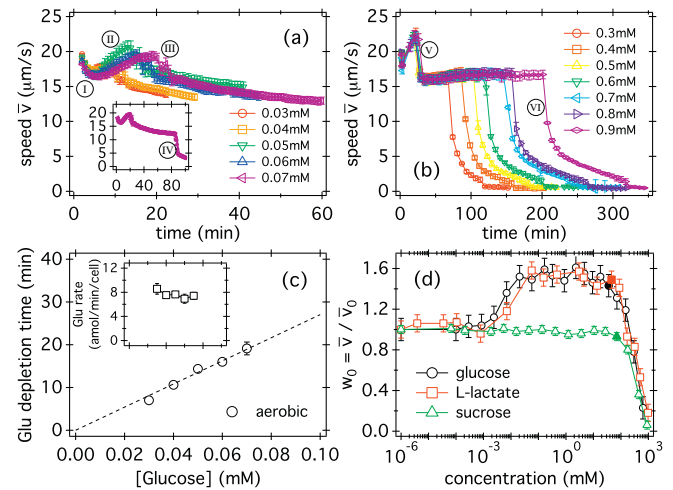
### 7.2. Physiological interpretation

The observed slowing down in  $\bar{v}$  at  $t \lesssim 1$  h is probably not due to substrate exhaustion, as extrapolated literature data [123,124] suggests this happens in  $\approx 10$  h. That  $d\bar{v}/dt$  is identical at different  $n$ , Fig. 9(a), supports this suggestion.

The ubiquitous signalling molecule cyclic di-guanylate (c-di-GMP) can slow down bacterial rotary motors by binding to the protein YcgR to form a molecular 'brake' [130]. However, we found that a  $\Delta YcgR$  mutant behaved identically to the WT in sealed capillaries; so c-di-GMP is unlikely implicated. Instead, we infer a slow decrease in the PMF driving the rotary motors. The adaptive value and mechanism of this putative starvation response is unknown, but it can affect many PMF-dependent cellular functions, e.g., ATP synthesis by  $F_1F_0$ -ATPase.

## 8. E. coli swimming powered by glucose

If glucose is present, *E. coli* utilises it first, and suppresses the expression of enzymes for processing other nutrients [131]. Glucose increases the swimming speed of *E. coli* [102] (see also Fig. 5(a))



**Fig. 10.** Effect of glucose on time-dependent average swimming speed  $\bar{v}(t)$  for WT AB1157 at  $5 \times 10^8$  cells/ml. (a) At low [Glu], glucose depletes first, causing the rapid drop (III) in  $\bar{v}$ . Inset: a longer run at [Glu] = 0.07 mM reveals a subsequent sharp drop (IV) due to  $O_2$  depletion. (b) At high [Glu],  $O_2$  depletes first, causing the first sharp drop (V) in  $\bar{v}$ . The second, larger sharp drop (VI) in  $\bar{v}$  is due to glucose depletion. Datasets are from one batch culture. (c) Depletion time of glucose in aerobic condition versus glucose concentration. Inset: aerobic glucose consumption rate versus [Glu] (with horizontal axis as main figure). (d) Speed  $\bar{v}$  versus molecule concentration normalised to the speed in BMB,  $\bar{v}_0$ , for three small molecules, measured  $\approx 5$  min after mixing and filling the capillary. Filled symbols are concentration markers for two and three independent batch cultures.

and other enteric bacteria [132]. We now demonstrate how to use glucose to enable *E. coli* AB1157 (cultured by standard protocol) to swim at a constant  $\bar{v}$  for extended time.

### 8.1. Observations and practical implications

Fig. 10(a) and (b) show  $\bar{v}$  at  $n = 5 \times 10^8$  cells  $\text{ml}^{-1}$  of WT AB1157 sealed in a capillary in BMB supplemented with low and high glucose concentrations, respectively. At  $0.03 \text{ mM} \leq [\text{Glu}] \leq 0.07 \text{ mM}$ , Fig. 10(a),  $\bar{v}$  rises gradually with time (II) after initial transients (I), until it drops suddenly (III). Thereafter, it decreases slowly with time. A longer run at [Glu] = 0.07 mM, Fig. 10(a) inset, reveals a second rapid drop (IV) that is sharper and of larger amplitude. The long, slow decrease with time and the final sudden drop are reminiscent of 'endogenous propulsion', Fig. 9. Thus, the data suggest that glucose is consumed aerobically until it is depleted, whereupon cells revert to endogenous metabolism until  $O_2$  is, in turn, depleted.

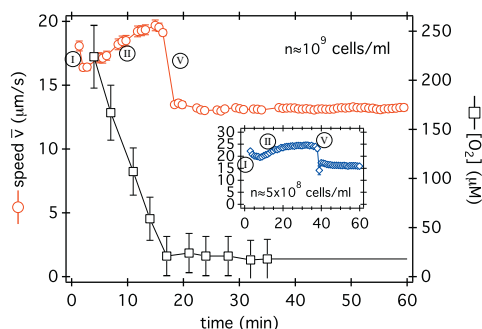
Consistent with this, the time of the first sharp drop (III) is proportional to [Glu], Fig. 10(c), and the deduced specific glucose consumption rate, inset, is nearly independent of [Glu], averaging to  $\bar{G} \approx 6.9 \pm 0.5$  amol/min/cell. We also measured  $\bar{G}$  over a range of [Glu] biochemically (Appendix E), and found a comparable value of  $\bar{G} \approx 5 \pm 0.6$  amol/min/cell.

Fig. 10(c) implies that there will be a [Glu] at which  $O_2$  depletion will occur before glucose depletion. This indeed happens at the higher [Glu] explored in Fig. 10(b). Now, the very sharp, first 'crash' (V) in  $\bar{v}$  is due to  $O_2$  depletion, at which point there is still plentiful glucose in the medium. Consistent with this interpretation, the time of the first abrupt speed decrease (V) is independent of [Glu]. Again using literature  $O_2$  solubilities [128,129], we find an average  $O_2$  utilisation rate of  $Q_{\text{Glu}} = 22.6 \pm 0.6$  amol/min/cell, which, as expected, is considerable higher than the  $Q_{\text{endo}} \approx 2$ –4 amol/min/cell reported by us (Section 7.1) and others [123] for endogenous metabolism.

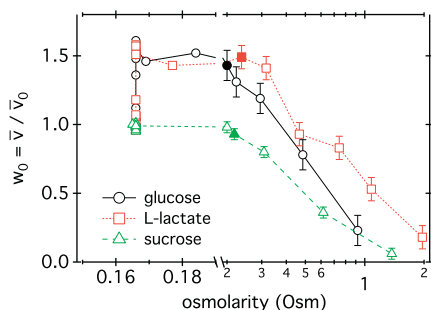
Independent evidence for  $O_2$  depletion comes from monitoring the fluorescence lifetime  $\tau$  of an  $O_2$ -sensitive dye [133]. By measuring  $\tau$  and using our experimentally derived calibration

<sup>15</sup>  $Q_{\text{endo}} \approx 10$ –20 ml/h/mg dry cell mass [123]  $\rightarrow$  2–4 amol/min/cell using  $O_2$  molar volume = 24 dm<sup>3</sup> and dry cell weight  $\approx 0.3 \times 10^{-9}$  mg cell [106].





**Fig. 11.** Simultaneous measurements of average swimming speed  $\bar{v}$  (○, left axis) and  $[O_2]$  (□, right axis), using fluorescence lifetime imaging microscopy (FLIM) vs. time in a sealed capillary for WT AB1157 at  $\approx 10^9$  cells  $ml^{-1}$  and  $[Glu]=0.5$  mM. For FLIM, 47 mM of the  $O_2$ -sensitive dye ruthenium tris(2,2'-dipyridyl) dichloride hydrate (RTDP) was added to BMB (which did not affect the cells' motility). RTDP was excited using sub-ps pulses ( $\lambda = 450$  nm, 1 mW at 1 MHz repetition rate) and its fluorescence imaged using a gated intensified CCD camera (Picostar HR-12QE, LaVision GmbH, Germany). The data was fitted with a single exponential decay yielding a (homogeneous) fluorescence lifetime map, which was averaged to estimate  $[O_2]$ . Inset:  $\bar{v}$  at  $5 \times 10^8$  cells/ml to better show the saturation before  $O_2$  depletion (horizontal axis as in main figure). Data sets are from independent batch cultures with speeds in BMB without glucose of  $\bar{v}(t=0) \approx 12 \mu m s^{-1}$  and  $\approx 15 \mu m s^{-1}$ , respectively.



**Fig. 12.** Normalised average swimming speed,  $\bar{v}/\bar{v}_0$ , from Fig. 10(b) as a function of the osmolarity. Note the bottom split axis with a linear scale from 150 to 190 mOsm (bottom left) and a log-scale from 200 to 2000 mOsm (bottom right). Filled symbols corresponds to markers in Fig. 10(d).

$[O_2]=568.1 \mu M$  ( $604 \text{ ns}/\tau - 1$ ) (consistent with [134]), we can assay  $[O_2]$  in our sealed samples, Fig. 11, which drops to  $\approx 0$  precisely at the point of the abrupt decrease (V) in  $\bar{v}$ .

Subsequent swimming powered by anaerobic glucose mixed-acid fermentation (Section 6) shows  $\bar{v}$  constant to  $\pm 1 \mu m s^{-1}$  until glucose depletion and  $\bar{v}$  abruptly drops by a factor of 3 or more, Fig. 10(c) (VI). Thus, *cells in initially aerated BMB supplemented with enough glucose will swim at approximately constant average speed after  $O_2$  depletion and before glucose depletion*. We note that *E. coli* is also able to utilise L-serine anaerobically, so that the speed of *E. coli* swimming powered by this amino acid shows a similar time course [135] to what we have seen for glucose, Fig. 10(b), before the final 'crash'.

The period of constant- $\bar{v}$  swimming can, at first sight, be arbitrarily extended by increasing  $[Glu]$ , Fig. 10(b), up to saturation ( $\approx 5$  M at  $25^\circ C$ ). However, a practical limit exists. Fig. 10(d) plots  $w_0$ , the average swimming speed at various  $[Glu]$  normalised to the average speed at  $[Glu]=0$ , both measured immediately after loading into sealed capillaries. At  $[Glu] \lesssim 10^{-3}$  mM, glucose depletion occurs before measurements could begin, so that  $w_0 = 1$ . At  $[Glu] \gtrsim 50$  mM,  $w_0$  catastrophically drops. Repeating using L-lactate, which *E. coli* can metabolise, and sucrose, which it cannot, also shows the same abrupt drop, which we believe is due to osmotic shock.

Fig. 12 replots our data against solution osmolarity (determined using an Osmomat30, Genotec, Germany), showing that  $\bar{v}$  always starts to decrease at an osmolarity between 0.2 and 0.4 Osm, comparable to the osmolarities needed to start decreasing cell volumes in previous studies [136]. Thus, there is a limit to the highest useable  $[Glu]$  for maintaining constant  $\bar{v}$ .<sup>16</sup>

## 8.2. Physiological interpretation

Reaction (9) for complete glucose oxidation requires an oxygen to glucose stoichiometry of  $Q_{Glu}/\dot{G} = 6$ . We measure  $Q_{Glu}/\dot{G} \approx 3.3 \pm 0.33$ , i.e. about 50% of the glucose consumed is fully oxidised. A combination of overflow metabolism and 'anabolic siphon' (footnote 14) probably explains this finding.

Fig. 11 shows the early-time portion of  $\bar{v}(t)$  on an expanded time axis. After initial transients (I),  $\bar{v}(t)$  rises (II) and then saturates (this is particularly clear in the inset) before the crash (V) due to  $O_2$  exhaustion. Immediately before these experiments, our cells have been harvested from TB, which contains little carbohydrates. While the operon of genes necessary for glucose transport (*pts*) is constitutively expressed irrespective of whether glucose is present, exposure to glucose can increase the expression of some *pts* genes by up to threefold [137,138], which probably causes the rise (II) of  $\bar{v}(t)$ .

Averaging over data sets, the ratio of the average speeds before and after the  $O_2$  'crash' (V) is in the range 1.3–1.6. The PMF of K-12 *E. coli* (strain AN387) grown aerobically on glycerol and fermentatively on glucose was previously measured to be  $-160$  mV and  $-117$  mV, respectively [139];  $160/117 = 1.36$  is remarkably close to our speed ratio.

## 8.3. How much energy does swimming use?

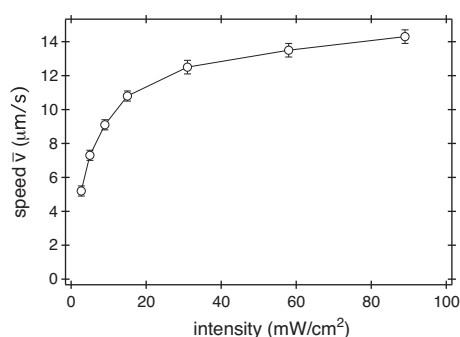
The detailed mechanics of a peritrichous swimmer (=many flagella distributed over the cell body) is not yet understood [50], so that the vast majority of the literature on *E. coli* is couched in terms of propulsion by a single 'effective flagellum'. The power developed by the 'effective motor' turning this 'effective flagellum' has been estimated to be  $P \approx 4 \times 10^{-16}$  W for swimming in minimal medium [140]. Each  $H^+$  driven into the cell by a PMF of  $\approx 150$  mV [139] can perform  $w = 2.4 \times 10^{-20}$  J of work. If we assume that flagella motors extract work with 100% efficiency from the protons, then  $P/w \lesssim 2 \times 10^4 H^+ s^{-1}$  is needed to power the 'effective motor'. In reality, of course, a somewhat higher flux would be needed, because the motor is not 100% efficient in harnessing energy from the PMF.

On the other hand, full oxidation of one glucose gives 12 NADH; each NADH can export up to  $8H^+$ . Thus,  $\dot{G}/2 \approx 4$  amol/min/cell  $\approx 4 \times 10^4$  /s/cell fully-oxidised glucose molecules export up to  $384 \times 10^4 H^+ /s/cell$ . This is in large excess of the  $\sim 10^4 H^+ /s/cell$  we estimate as needed to power the  $\approx 50\%$  increase in  $\bar{v}$  upon addition of glucose. Thus, swimming accounts for only a small part of a cell's energy budget.

## 8.4. Other small molecules

We have also studied swimming powered by acetate ( $C_2$ ), glycerol and lactate ( $C_3$ ), xylose ( $C_5$ ), galactose ( $C_6$ ), and maltose and lactose ( $C_{12}$ ), which can be expected to generate different amounts of energy per molecule. Interestingly, however, the maximum increase in speed in each case is approximately the same

<sup>16</sup> As expected from previous work [136], the results in Fig. 12 were time dependent: waiting longer produced partial recovery in the 'crashed' speed. Nevertheless, a definite upper limit exists at  $\approx 100$  mM small-molecule solutes.



**Fig. 13.** Average steady-state swimming speed  $\bar{v}$  of *E. coli* MG1655 ( $n = 10^9$  cells  $\text{ml}^{-1}$ ) expressing SAR86  $\gamma$ -proteorhodopsin at different intensities of incident green light ( $\lambda = 510 - 540$  nm) after  $\text{O}_2$  depletion; data obtained starting from the highest intensity.

as observed for glucose (about 50%, see Fig. 10(d) for L-lactate). This suggests that the flagellar motor speed saturates at high PMF (as observed for *Bacillus subtilis* [141]), and supports our conclusion that swimming uses only a small fraction of metabolic energy. At high enough concentration of these alternative carbon sources, we find that the post- $\text{O}_2$ -exhaustion speed drops to zero (for lactate, cf. [135]), rather than to some intermediate, constant level as observed for glucose. This is because *E. coli* cannot ferment these substrates.

## 9. *E. coli* swimming powered by light

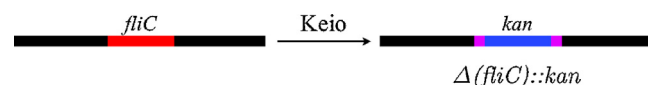
Some micro-organisms directly harvest energy from light. Thus, photosynthesizers use photons to split water to generate high energy electrons. Alternatively, light is harvested directly to pump protons using proton pumps such as proteorhodopsin (PR). PR, first discovered through the genomic sequencing of marine bacteria [142], has been incorporated into *E. coli* [143].

Most PRs identified to date can generate PMF up to or slightly above WT levels [144]. To enable illuminated PR to take over PMF generation, other sources need to be deactivated. This has been demonstrated on single cells using poisons for the respiratory enzymes (cf. Fig. 8) or by  $\text{O}_2$  exclusion. As the motor, and therefore flagellum rotation speed, is proportional to the PMF, illuminating poisoned cells expressing PR leads to a controlled increase in flagellum rotation speed [143].

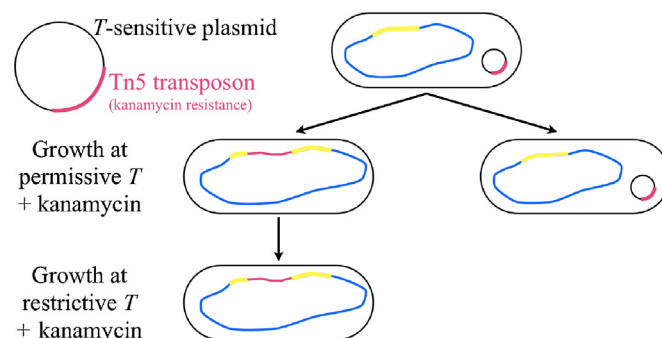
Here, we demonstrate this effect in a cell population using MG1655 transformed with plasmid pBAD-HisC-PR expressing PR when induced by arabinose (DM4). We already know that  $\text{O}_2$  is exhausted after a period when swimmers are sealed in capillaries, Fig. 9. After our transformed cells have exhausted  $\text{O}_2$ , their swimming is wholly powered by PR-generated PMF. The initial increase in speed with incident intensity, Fig. 13, saturates at high intensities, in agreement with previous measurements of the rotation rate of individual flagellar motors [143]. This and other PR mutants therefore enable simple external control of their swimming speed, making it feasible to change the samples activity both in time and space. They are living analogues of light-driven synthetic swimmers [10,145].

## 10. *E. coli* – the genetic toolbox

A great attraction of *E. coli* is the availability of many mutants and a versatile genetic ‘tool box’. Thus, e.g., the FliC flagella protein could be mutated to include a cysteine residue to make disulphide bonds with dyes, or the *cheY* chemotaxis gene could be deleted to give smooth (non-tumbling) swimmers. We briefly introduce two ‘off the shelf’ mutant libraries and two methods for mutagenesis



**Fig. 14.** A non-motile Keio knockout [147],  $\Delta(fliC)::kan$ , where the flagella protein gene *fliC* (red) is replaced by the kanamycin resistance gene *kan* (blue), flanked by DNA sequences (purple) that allow its easy excision. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)



**Fig. 15.** Schematic of transposon mutagenesis. See text for details.

to illustrate principles and introduce terminology. Details can be found in the textbooks, e.g., [146].

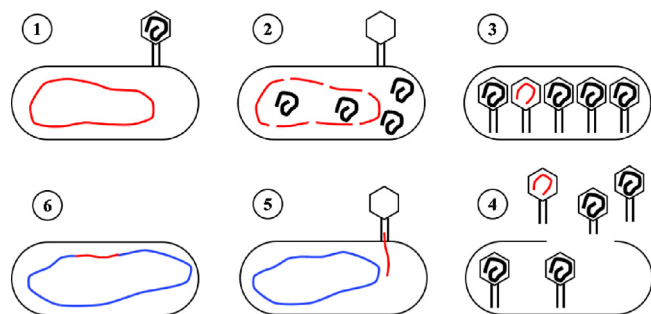
A major resource is the Keio collection [147] of ‘single knockout’ mutants of strain BW25113 (closely related to MG1655 but more amenable to molecular biology), available on a non-profit making basis to academics. Each non-essential gene was precisely deleted and replaced with an excisable kanamycin resistance gene, Fig. 14.

Many mutagenetic methods use plasmids, extrachromosomal pieces of DNA that can be ‘transformed’ into bacteria and maintained independently. Fig. 15 illustrates one strategy to carry out ‘transposon mutagenesis’. It relies on a temperature-sensitive plasmid, which is only maintained within the cell at the ‘permissive’ temperature, and is lost at a (higher) ‘restrictive’ temperature. This ‘suicide’ plasmid acts as a ‘vector’ (carrier) for a transposon, a piece of genetic code that can ‘jump’ between DNA locations. Here, the plasmid carries transposon Tn5, which confers resistance to the antibiotic kanamycin.

When the transposon jumps, or transposes, from the plasmid onto the bacterial chromosome, it may insert into a gene, generating a mutation. Following a shift of the *E. coli* from the permissive to the restrictive temperature, the plasmid will be lost. Then, all kanamycin-resistant cells carry a chromosomal copy of the Tn5 element transposed onto the chromosome, where it can be stably maintained. Crucially, the transposon jumps into random chromosomal locations, and so can generate mutations in any gene. A particular mutation can be selected by a change in phenotype and the insertion site mapped by sequencing.

P1 transduction is another popular mutagenesis technique (Fig. 16). Here DNA is transferred from a donor to a recipient cell using the P1 bacterial virus (a bacteriophage). When P1 infects a host cell, it not only packages its own genome into the virus particle, but also can package fragments of the host chromosome. When these virus particles are purified and used to infect the target strain of *E. coli*, DNA homology between the donor DNA packaged in the virus particle and the recipient chromosome can generate insertion of the donor DNA.

More generally, plasmids constitute a vital resource for introducing new genes into *E. coli*, e.g., transforming with a plasmid carrying the GFP gene will render cells fluorescent. Each member in another useful library, the ASKA collection [148], consists of an *E. coli* strain containing a plasmid in which a single gene from *E. coli* strain W3110 has been cloned. The DNA encoding each gene has



**Fig. 16.** Schematic of P1 phage transduction. ① Phage attaches to cell. ② Phage injects DNA (black) into cell and uses its machinery to replicate as well as fragment cell's genome (red). ③ DNA is packaged into viral particles, some of which contain fragments of the cell's DNA, giving a 'Trojan' phage. ④ Cell lyses, releasing all viral particles. ⑤ 'Trojan' phage infects target cell. ⑥ Donor DNA fragment is incorporated into target cell genome by recombination. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

been fused to sequences that permit precise temporal control of the expression of the associated protein and its fluorescent tagging with GFP.

We end with two 'health warnings'. First, genetic modification in motility-related work typically starts with parent strains selected for motility (e.g., HCBxx or RPxx, Table B.2). Many of these retain wild-type features not optimised for DNA manipulation. Moreover, the genome of many 'motility favourites' remains less well characterised compared to 'canonical' strains such as MG1655 [27]. Sequencing and annotating the genome of some of these strains is a desirable future step.

Secondly, care must be taken in obtaining a good reliable starter culture of the strain that one chooses to work with, either from the originating author or from a national culture collection, e.g., ATCC [149] and CGSC [150] in America; NCIMB [151] in the UK; DSMZ [152] in Germany; and CGMCC in China [153]. Bacterial genomes are changeable, so that strains deposited at the back of freezers 20 years ago whose provenance is uncertain are not good starting points.

## 11. Summary and conclusions

*Escherichia coli* is increasingly used in active colloids experiments as an alternative to or to contrast with synthetic colloidal swimmers. We have reviewed, explained and proposed a set of procedures that, together, should enable the collection of data from such experiments between different laboratories that can be quantitatively compared. At the heart of these procedures is the use of a new, high-throughput method to characterise *E. coli* motility, differential dynamic microscopy (DDM), which delivers the average swimming speed and motile fraction averaged over  $\approx 10^4$  cells in a matter of minutes [87,88]. Using DDM, we have established that *E. coli* swims at constant speed when motility is powered by mixed-acid fermentation, Fig. 10(b). Along the way, we have explored the use of DDM to interrogate aspects of cellular physiology via the dependence of swimming speed on PMF. Our results suggest that *E. coli* cells powered by endogenous metabolism actively reduce their PMF with time. It appears that about 50% of the glucose under our conditions is used for generating energy, and that a very small fraction (maybe  $\approx 1\%$ ) of this energy is used for motility.

Whether our proposed procedures, or something similar, are followed, it is important that sufficient information is recorded in publications using *E. coli* as a model active colloid to enable faithful reproduction of experiments. We therefore propose a 'checklist' of what needs to be reported in Appendix A.

More needs to be done before *E. coli* can become an ideal model for active matter experiments. Biologically, it would be desirable to sequence the genomes of some favourite motility strains. Armed with accurate genomic information, it should be possible to design bespoke swimmers that have ideal characteristics for different aspects of active matter research, e.g. a strain that does not depend on either poison, or  $O_2$  depletion, or light to display constant swimming speed over extended time. Theoretically, it would be useful to construct coarse-grained models that take into account the flagella bundle. Predictions from such theoretical models confronted with data from experiments using increasingly well-characterised *E. coli* suspensions should propel the subject of active colloids significantly forward.

## Acknowledgements

WCKP initiated work with *E. coli* under an EPSRC Senior Research Fellowship (EP/D071070). He and J.S.L., J.A., A.D. and V.A.M. are now funded by an EPSRC Programme Grant (EP/J007404/1) and an ERC Advanced Grant (ADG-PHYAPS). V.A.M.'s was also supported by Marie Curie Fellowship 'ActiveDynamics' FP7-PEOPLE (PIIF-GA-2010-276190). A.J. held an EPSRC studentship. T.V. is funded through Marie Curie fellowship 'Living Patchy Colloids' (LivPaC, 623364) under the FP7-PEOPLE-2013-IEF program. T.P. is supported by an Edinburgh Chancellor's Fellowship and the BBSRC/EPSRC/MRC Synthetic Biology Research Centre grant (BB/M018040/1).

Many past lab members contributed to our knowledge of *E. coli* motility, especially Otti Croze (who cultured our first *E. coli*) and Laurence Wilson (who performed the first DDM). We thank Howard Berg (Harvard), Ian Booth (Aberdeen) and Gail Ferguson (Edinburgh, then Aberdeen) for educational and illuminating discussions, and all the authors represented in Table B.2 for providing strains and plasmids.

## Appendix A. Protocol reporting checklist

To facilitate reproducing and comparing results, we suggest that published methods should always include:

- Strain name and details of any genetic modifications.
- Growth conditions (temperature, growth medium composition, cell density of harvest).
- Washing procedure (method: filtration or centrifugation, number of washing steps, exact composition of motility medium).
- Conditions for motility experiments (cell concentration and how determined, sample cell details: material and dimensions, sealing, position and conditions for image/video acquisition).
- Details for image analysis and data processing.

## Appendix B. Strains

Table B.2 summarises a number of strains of *E. coli* that have been used by us and others in active colloids and related motility work. Note that publication in a journal carries the obligation of curating and sharing the strains used in a published work. However, one should not give away a strain obtained from another laboratory; instead the protocol is to direct enquiries to the strain originator. Note that intra-strain variability can arise rapidly because certain genetic elements (called insertion sequences, IS) can 'jump' into regions of the genome controlling flagella gene expression, enhancing the motility [154].

**Table B.2**  
*E. coli* strains and plasmids.

Name	Motility genotype/relevant characteristics	Source
<i>E. coli</i> strains		
MG1655	Motility WT	Laboratory strain [27]
AB1157	Motility WT	Laboratory strain [155]
RP437	Motility WT	J. Parkinson [156]
HCB1	Motility WT	H.C. Berg [157]
BW25113	Parent strain of Keio collection, motile	[147]
HCB437	$\Delta$ ( <i>tsr</i> )7021 $\Delta$ ( <i>trg</i> )100 $\Delta$ ( <i>cheA-cheZ</i> )2209 Defective in chemotaxis, a smooth swimmer	H. C. Berg
DM4	MG1655 + proteorhodopsin-expressing pBAD-HisC-PR	This work
JSL1	AB1157 $\Delta$ ( <i>cheY::kan</i> ) Smooth swimmer	This work
AD1	AB1157 <i>fliC</i> (S353C) Motility WT with <i>FliC</i> mutation to bind to Alexa dye	This work
AD2	AD1 $\Delta$ ( <i>cheY::kan</i> ) Smooth swimmer with <i>FliC</i> mutation to bind to Alexa dye	This work
AD3	AD1 $\Delta$ ( <i>ycgR::kan</i> ) Motility WT, unable to 'brake'	This work
Plasmids		
pHC60	Tet <sup>R</sup> , GFP constitutive	[158]
pBAD-HisC-PR	Amp <sup>R</sup> , arabinose inducible proteorhodopsin	[144]

### Appendix C. Nomenclature for bacterial genetics

Researchers working with bacteria, especially if they collaborate with biologists, need some facility in the shorthand used to described genotypes [159,160], which in turn requires some knowledge of molecular genetics [146]. Here we give a very brief introduction based on 'worked examples' from Table B.2.

Each gene is given a three-letter, lower-case italicised name (possibly followed by an extra capital letter for different members of a gene family), often suggestive of its function; different versions of a gene (alleles) are distinguished by numbers. Thus, *tsr*7021 is allele 7021 of the *tsr* gene in *E. coli* coding for part of a receptor system for sensing L-serine and related amino acids. Deletion of this allele in HCB427 is indicated by  $\Delta$ (*tsr*)7021. The protein product of a gene is denoted by the same three letters unitalicised and capitalised. Thus, in the *fliC*(S353C) mutation carried by AD1, residue 353 of the *FliC* protein used to build flagella filaments has been mutated from serine (S) to cysteine (C).<sup>17</sup> Sometimes, a different gene is inserted in place of a deleted gene. Thus, in JSL1, a cassette of genes conferring resistance to the antibiotic kanamycin has been inserted in place of the deleted *cheY* gene for chemotaxis –  $\Delta$ (*cheY::kan*); the advantage is that such mutants can be easily selected by growth in the presence of kanamycin. Finally, many bacteria carry plasmids, such as plasmid pHC60, which carries tetracycline resistance (Tet<sup>R</sup>) and expresses green fluorescence protein (GFP) constitutively (i.e. continuously rather than subject to environmental control).

Genotype of strains can be found from various online databases, including the various culture collections [149–153]. Information

on individual genes are available from appropriate databases, e.g., [161].

### Appendix D. Filtration

To transfer the bacteria from TB growth media to BMB, we use a sterile Nalgene filtration unit consisting of two compartments separated by a Milipore 45  $\mu$ m HATF filter [162]. The filter is soaked in BMB for  $\approx$ 5 min before being placed centrally on the unit with sterile tweezers. After screwing tight the upper and lower halves of the unit, 35 ml of cell culture is slowly poured onto the filter. A tap-powered suction pump is used to enhance the flow rate, with a 70% ethanol bath between the unit and the suction tap to prevent water supply contamination.

When the level of TB has fallen to  $\approx$ 3 mm above the filter, leaving  $\approx$ 3 ml of liquid, 35 ml of BMB is pipetted into the unit. Ensuring that the filter does not run dry during the procedure minimises the number of non-motile cells. This washing step is repeated 3 times. The water flow is adjusted to maintain a constant filtration rate so that each wash step takes 10 min. When  $\sim$ 1–3 ml of filtrate remains after the third wash, a sterile cut pipette tip is used to transfer the filtrate to a plastic 50 ml polystyrene test tube (Greiner). (Smaller tubes are not wide enough for next step.) The filter is removed from the unit using sterile tweezers and deposited on the side of the tube. By gently rolling the liquid over the filter, bacteria are resuspended to reach a final OD between 5 and 15, depending on the volume of liquid left on-top of the filter after the final washing step.

### Appendix E. Glucose assay

We used a Sigma GAGO20 glucose assay kit. Hydrogen peroxide from the oxidation of glucose by glucose oxidase reacts with o-dianisidine in the presence of peroxidase to form a colored intermediate, which further reacts with sulfuric acid to form a more stable, pink final product. The OD measured at 540 nm is proportional to the original glucose concentration. Calibration runs showed that in our equipment, the linear range for OD<sub>540</sub> vs. glucose concentration extended to  $\approx$ 600  $\mu$ M.

Bacterial cells were grown using the standard motility protocol, washed and OD adjusted to 0.3 as for motility experiments, except that the buffer here lacks EDTA, which interfered with glucose oxidase. Cells and glucose were mixed and samples taken at different time points, put on ice and filtered (0.22  $\mu$ m pore size) as quickly as possible. The filtered supernatant was used according to manufacturer's protocol to measure glucose concentration. Glucose utilization rates were calculated using glucose and time difference as well as wet or dry weight of cells in 1 ml of OD = 0.3 suspension. Results in Table E.3 are reported as averages and standard deviations using different time periods measured in one set of experiments. There is no systematic trend. We therefore average the value determined at different [Glu] to give an estimated glucose utilisation rate of  $5.2 \pm 0.6$   $\mu$ mol/min/g wet cells, or  $\approx 5 \pm 0.6$  amol/min/cell.

**Table E.3**  
Glucose utilisation rate

Initial [Glu] (mM)	Rate ( $\mu$ mol/min/g wet cells)
0.50	$4.9 \pm 1.0$
0.20	$5.2 \pm 0.9$
0.10	$6.1 \pm 1.6$
0.05	$4.7 \pm 0.8$

<sup>17</sup> See the textbooks, e.g., [107], for single-letter abbreviations of amino acids.



### Appendix F. List of symbols

Symbol	Typical units	Brief definition
[...]	μM or mM	concentration of ...
$\bar{x}$ or $\langle x \rangle$	–	average of $x$
<b>Latin</b>		
$c$	μM or mM	small molecule conc.
$D$	μm <sup>2</sup> s <sup>−1</sup>	diffusion coefficient
$d$	mm	colony diameter
$\dot{G}$	amol/min/cell	glucose consumption rate
$g$	m s <sup>−2</sup>	gravitational acceleration
$L$	μm	cell body + flagellum length
$l$	μm	cell body length
$\ell_g$	μm	gravitational length
$n$	cells/ml	cell density
$n$	–	refractive index
$Q$	amol/min/cell	O <sub>2</sub> consumption rate
$q$	μm <sup>−1</sup>	scattering vector
$T$	K or °C	temperature
$t$	s or min	time
$t_c$	s or min	time of speed crashes
$u(t)$	–	$\bar{u}(t)/\bar{v}_0$
$V$	μm <sup>3</sup>	volume
$\bar{v}$	μms <sup>−1</sup>	average speed
$\bar{v}_0$	μms <sup>−1</sup>	$\bar{v}(t = 0) \equiv \bar{v}(0)$
$v_s$	μms <sup>−1</sup>	sedimentation speed
$w_0(c)$	–	$w_0(c) = \bar{v}_0(c)/\bar{v}_0(c = 0)$
<b>Greek</b>		
$\alpha$	s <sup>−1</sup> or min <sup>−1</sup>	growth rate
$\beta$	–	non-motile fraction
$\gamma$	–	$l/\sigma$
$\eta$	mPa s	viscosity
$\lambda$	μm	wavelength <i>in vacuo</i>
$\mu$	μm cm V <sup>−1</sup> s <sup>−1</sup>	electrophoretic motility
$\nu$	–	number of protons (H <sup>+</sup> )
$\xi$	pN s μm <sup>−1</sup>	friction coefficient
$\rho$	g cm <sup>−3</sup>	(mass) density
$\sigma$	μm	cell body diameter
$\tau$	s	delay time, run time, lifetime
$\phi$	–	cell body volume fraction

### References

[1] W.C.K. Poon, From Clarkia to *Escherichia* and *Janus*: the physics of natural and synthetic active colloids, in: C. Bechinger, F. Sciortino, P. Ziaherl (Eds.), *Physics of Complex Colloids*, Società Italiana di Fisica, Bologna, 2013, pp. 317–386 [ArXiv:1306.4799](#)

[2] Y. Sowa, R.M. Berry, Bacterial flagellar motor, *Q. Rev. Biophys.* **41** (2008) 103–132.

[3] S.J. Ebbens, J.R. Howse, In pursuit of propulsion at the nanoscale, *Soft Matter* **6** (2010) 726–738.

[4] Y. Hong, D. Velegol, N. Chaturvedi, A. Sen, Biomimetic behavior of synthetic particles: from microscopic randomness to macroscopic control, *Phys. Chem. Chem. Phys.* **12** (2010) 1423–1435.

[5] A. Sokolov, I.S. Aranson, *Phys. Rev. Lett.* **103** (2009) 148101.

[6] H.M. López, J. Gachelin, C. Douarche, H. Auradou, E. Clément, Turning bacteria suspensions into superfluids, *Phys. Rev. Lett.* **115** (2015) 028301.

[7] P. Galajda, J. Keymer, P. Chaikin, R. Austin, A wall of funnels concentrates swimming bacteria, *J. Bacteriol.* **189** (2007) 8704–8707.

[8] M.E. Cates, Diffusive transport without detailed balance in motile bacteria: does microbiology need statistical physics? *Rep. Prog. Phys.* **75** (2012) 042601.

[9] M.C. Marchetti, J.F. Joanny, S. Ramaswamy, T.B. Liverpool, J. Prost, M. Rao, R.A. Simha, Hydrodynamics of soft active matter, *Rev. Mod. Phys.* **85** (2014) 1143.

[10] J. Palacci, S. Sacanna, A.P. Steinberg, D.J. Pine, P.M. Chaikin, Living crystals of light-activated colloidal surfers, *Science* **339** (2013) 936–940.

[11] J. Perrin, *Atoms*, Ox Bow Press, Woodbridge, 1990 (Original: Les Atomes, Librairie Félix Alcan, Paris, 1913).

[12] A. Pais, ‘Subtle is the Lord . . .’ The Science and the Life of Albert Einstein, OUP, Oxford, 1982.

[13] P.N. Pusey, W. van Megen, Phase behaviour of concentrated suspensions of nearly hard colloidal spheres, *Nature* **320** (1985) 340–342.

[14] K.N. Pham, A.M. Puentes, J. Bergenholtz, S.U. Egelhaaf, A. Moussaïd, P.N. Pusey, A.B. Schofield, M.E. Cates, M. Fuchs, W.C.K. Poon, Multiple glassy states in a simple model system, *Science* **296** (2002) 104–106.

[15] P.N. Pusey, Colloidal suspensions, in: J.P. Hansen, D. Levesque, J. Zinn-Justin (Eds.), *Liquids, Freezing and the Glass Transition*, Elsevier, 1991, pp. 765–942.

[16] W.C.K. Poon, P.N. Pusey, *Phase Transition of Spherical Colloids*, Kluwer, Dordrecht, 1995.

[17] W.C.K. Poon, Chapter 1: Colloidal suspensions, in: E.M. Terentjev, D.A. Weitz (Eds.), *Oxford Handbook of Soft Condensed Matter*, Oxford University Press, Oxford, 2015.

[18] W.C.K. Poon, E.R. Weeks, C.P. Royall, On measuring colloidal volume fractions, *Soft Matter* **8** (2012) 21–30.

[19] H.C. Berg, *E. coli* in Motion, Springer, Berlin, 2003.

[20] Q. Liao, G. Subramanian, M.P. DeLisa, D.L. Koch, M. Wu, Pair velocity correlations among swimming *Escherichia coli* bacteria are determined by force-quadrupole hydrodynamic interactions, *Phys. Fluids* **18** (2007) 061701.

[21] K. Drescher, R.E. Goldstein, N. Michel, M. Polin, I. Tuval, Direct measurement of the flow field around swimming microorganisms, *Phys. Rev. Lett.* **105** (2010) 168101.

[22] A.T. Brown, W.C.K. Poon, Ionic effects in self-propelled Pt Janus swimmers, *Soft Matter* **10** (2014) 4016–4027.

[23] S. Ebbens, D.A. Gregory, G. Dunderdale, J.R. Howse, Y. Ibrahim, T.B. Liverpool, R. Golestanian, Electrokinetic effects in catalytic platinum-insulator Janus swimmers, *EPL* **106** (2014) 58003.

[24] S.T. Shulman, H.C. Friedmann, R.H. Sims, Theodor Escherich: The first pediatric infectious diseases physician? *Clin. Infect. Dis.* **45** (2007) 1025–1029.

[25] F. Jacob, *The Statue Within: An Autobiography*, Cold Spring Harbor Laboratory Press, Cold Spring Harbor, 1998.

[26] M. Achtman, M. Wagner, Microbial diversity and the genetic nature of microbial species, *Nat. Rev. Microbiol.* **6** (2008) 431–440.

[27] F. Blattner, G. Plunkett, C. Bloch, N. Perna, V. Burland, M. Riley, J. ColladoVides, J. Glasner, C. Rode, G. Mayhew, J. Gregor, N. Davis, H. Kirkpatrick, M. Goeden, D. Rose, B. Mau, Y. Shao, The complete genome sequence of *Escherichia coli* K-12, *Science* **277** (1997) 1453–1462.

[28] Genbank: The NH Genetic Sequence Database, 2015 <http://www.ncbi.nlm.nih.gov/genbank/> (last accessed on 17.07.15).

[29] F.C. Neihardt (Ed.), *Escherichia coli* and *Salmonella*: Cellular and Molecular Biology, 2nd ed., American Society for Microbiology, Washington D.C., 1996.

[30] EcoSal plus, Access by subscription, <http://www.ecosal.org>

[31] C.J. Calladine, Design requirements for the construction of bacterial flagella, *J. Theor. Biol.* **57** (1976) 469–489.

[32] C.J. Calladine, Change in waveform in bacterial flagella: the role of mechanics at the molecular level, *J. Mol. Biol.* **118** (1978) 457–379.

[33] R.M. McNab, Flagella and motility, in: F.C. Neidhardt (Ed.), *E. coli* and *Salmonella*, ASM Press, Washington, DC, 1996, pp. 123–145.

[34] H.C. Berg, D.A. Brown, Chemotaxis in *Escherichia coli* analysed by three-dimensional tracking, *Nature* **239** (1972) 500–504.

[35] J.L. Hobman, C.W. Penn, M.J. Pallen, Laboratory strains of *Escherichia coli*: model citizens or deceitful delinquents growing old disgracefully? *Mol. Microbiol.* **64** (2007) 881–885.

[36] T. Hayashi, K. Makino, M. Ohnishi, K. Kurokawa, K. Ishii, K. Yokoyama, C.-G. Han, E. Ohtsubo, K. Nakayama, T. Murata, M. Tanaka, T. Tobe, T. Iida, H. Takami, T. Honda, C. Sasakawa, N. Ogasawara, T. Yasunaga, S. Kuhara, T. Shiba, M. Hattori, H. Shinagawa, Complete genome sequence of enterohemorrhagic *Escherichia coli* O157:H7 and genomic comparison with a laboratory strain K-12, *DNA Res.* **8** (2001) 11–22.

[37] Recognizing the biosafety levels, n.d., <http://www.cdc.gov/training/QuickLearns/biosafety/> (accessed 17.07.15).

[38] U. Moran, R. Phillips, R. Milo, SnapShot. Key numbers in biology, *Cell* **141** (2010) 1262–1263.

[39] CyberCell Database CCDB, <http://ccdb.wishartlab.com/CCDB/index.html>

[40] Bionumbers, <http://www.bionumbers.hms.harvard.edu> (last accessed 17.07.15).

[41] K.D. Young, Bacterial morphology: why have different shapes? *Curr. Opin. Microbiol.* **10** (2007) 596–600.

[42] O. Hallatschek, P. Hersen, S. Ramanathan, D.R. Nelson, Genetic drift at expanding frontiers promotes gene segregation, *Proc. Natl. Acad. Sci. U. S. A.* **104** (2007) 19926–19930.

[43] T. Nomura, Y. Morimoto, M. Ishikawa, H. Tokumoto, Y. Konishi, Synthesis of hollow silica microparticles from bacterial templates, *Adv. Powder Technol.* **21** (2010) 8–12.

[44] P. Latimer, Light-scattering vs. microscopy for measuring average cell-size and shape, *Biophys. J.* **27** (1979) 117–126.

[45] H. Åkerlund, K. Nörtström, R. Bernander, Analysis of cell size and dna content in exponentially growing and stationary-phase batch cultures of *Escherichia coli*, *J. Bacteriol.* **177** (1995) 6791–6797.

[46] F.J. Trueba, C.L. Woldringh, Changes in cell diameter during the division cycle of *Escherichia coli*, *J. Bacteriol.* **142** (1980) 869–878.

[47] M.J. Solomon, P.T. Spicer, Microstructural regimes of colloidal rod suspensions, gels, and glasses, *Soft Matter* **6** (2010) 1391–1400.

[48] Y. Han, A.M. Ellsayed, M. Nobili, J. Zhang, T.C. Lubensky, A.G. Yodh, Brownian motion of an ellipsoid, *Science* **314** (2006) 626–630.

[49] L. Turner, W. Ryu, H. Berg, Real-time imaging of fluorescent flagellar filaments, *J. Bacteriol.* **182** (2000) 2793–2801.

[50] N.C. Darnton, L. Turner, S. Rojevsky, H.C. Berg, On torque and tumbling in swimming *Escherichia coli*, *J. Bacteriol.* **189** (2007) 1756–1764.

[51] L. Turner, R. Zhang, N.C. Darnton, H.C. Berg, Visualization of flagella during bacterial swimming, *J. Bacteriol.* **192** (2010) 3259–3267.

- [52] A.L. Koch, Theory of the angular dependence of light scattered by bacteria and similar-sized biological objects, *J. Theor. Biol.* 18 (1968) 133–156.
- [53] D.A. Cross, P. Latimer, Angular dependence of scattering from *Escherichia coli* cells, *Appl. Opt.* 11 (1972) 1225–1228.
- [54] V.J. Morris, B.R. Jennings, Light scattering by bacteria I. angular dependence of the scattered intensity, *Proc. R. Soc. Lond. A* 338 (1974) 197–208.
- [55] M. Holz, S.H. Chen, Structural effects in quasi-elastic light scattering from motile bacteria of *E. coli*, *Appl. Opt.* 17 (1978) 1930–1937.
- [56] P. Latimer, Light scattering and absorption as methods of studying cell population parameters, *Ann. Rev. Biophys. Bioeng.* 11 (1982) 129–150.
- [57] A. Katz, A. Alimova, M. Xu, E. Rudolph, M.K. Shah, H.E. Savage, R.B. Rosen, S.A. McCormick, R.R. Alfano, Bacteria size determination by elastic light scattering, *IEEE J. Sel. Top. Quantum Electron.* 9 (2003) 9277–9287.
- [58] A.I. Konokhova, A.A. Gelash, M.A. Yurkin, A.V. Chernyshev, V.P. Maltsev, High-precision characterization of individual *E. coli* cell morphology by scanning flow cytometry, *Cytometry* 83A (2013) 568–575.
- [59] J.B. Bateman, J. Wagman, E.L. Carstensen, Refraction and absorption of light in bacterial suspensions, *Kolloid Z. Z. Polym.* 208 (1965) 44–58.
- [60] P. Latimer, Light scattering by ellipsoids, *J. Colloid Interface Sci.* 53 (1975) 102–109.
- [61] D.S. Wang, P.W. Barber, Scattering by inhomogeneous nonspherical objects, *Appl. Opt.* 18 (1979) 1190–1197.
- [62] D.S. Wang, H.C.H. Chen, P.W. Barber, P.J. Wyatt, Light scattering by polydisperse suspensions of inhomogeneous nonspherical particles, *Appl. Opt.* 18 (1979) 2672–2678.
- [63] B.J. Berne, R. Nossal, Inelastic light scattering by large structured particles, *Biophys. J.* 14 (1974) 865–880.
- [64] B.V. Bronk, S.D. Druger, J. Czégé, W.P. Van De Merwe, Measuring diameters of rod-shaped bacteria in vivo with polarized light scattering, *Biophys. J.* 69 (1995) 1170–1177.
- [65] W.P. Van De Merwe, Z.Z. Li, B.V. Bronk, J. Czégé, Polarized light scattering for rapid observation of bacterial size changes, *Biophys. J.* 73 (1997) 500–506.
- [66] W.W. Baldwin, R. Myer, T. Kung, E. Anderson, A.L. Koch, Growth and buoyant density of *Escherichia coli* at very low osmolarities, *J. Bacteriol.* 177 (1995) 235–237.
- [67] W.W. Baldwin, R. Myer, N. Powell, E. Anderson, A.L. Koch, Buoyant density of *Escherichia coli* is determined solely by the osmolarity of the culture medium, *Arch. Microbiol.* 164 (1995) 155–157.
- [68] M. Godin, A.K. Bryan, T.P. Burg, K. Babcock, S.R. Manalis, Measuring the mass, density, and size of particles and cells using a suspended microchannel resonator, *Appl. Phys. Lett.* 91 (2007) 123121.
- [69] H.C. Berg, *Random Walks in Biology*, Princeton University Press, Princeton, NJ, 1993.
- [70] J. Palacci, C. Cottin-Bizonne, C. Ybert, L. Bocquet, Sedimentation and effective temperature of active colloidal suspensions, *Phys. Rev. Lett.* 105 (2010) 088304.
- [71] C. Barrett-Freeman, M.R. Evans, D. Marenduzzo, W.C.K. Poon, Nonequilibrium phase transition in the sedimentation of reproducing particles, *Phys. Rev. Lett.* 101 (2008) 100602.
- [72] J.T. Davies, D.A. Haydon, E. Rideal, Surface behaviour of *Bacillus coli* I. The nature of the surface, *Proc. R. Soc. Lond. B* 145 (1956) 375–383.
- [73] Y. Liu, The impact of bacterial surface polymers on bacterial adhesion, in: H. Zeng (Ed.), *Polymer Adhesion, Friction, and Lubrication*, John Wiley and Sons, Hoboken, NJ, 2013, pp. 575–616.
- [74] J.B. Kaper, J.P. Nataro, H.L.T. Mobley, Pathogenic *Escherichia coli*, *Nat. Rev. Microbiol.* 2 (2004) 123–140.
- [75] A.T. Poortinga, R. Bosa, W. Norde, H.J. Busscher, Electric double layer interactions in bacterial adhesion to surfaces, *Surf. Sci. Rep.* 47 (2002) 1–32.
- [76] A. van der Wal, M. Minor, W. Norde, A.J.B. Zehnder, J. Lyklema, Electrokinetic potential of bacterial cells, *Langmuir* 13 (1997) 165–171.
- [77] A.J. de Kerchove, M. Elimelech, Relevance of electrokinetic theory for “soft” particles to bacterial cells: Implications for bacterial adhesion, *Langmuir* 21 (2005) 6462–6472.
- [78] M.W. van der Wolde, I.R. Henderson, Regulation and function of Ag43 (flu), *Annu. Rev. Microbiol.* 62 (2008) 153–169.
- [79] A.T. Brown, I.D. Vladescu, A. Dawson, T. Vissers, J. Schwarz-Linek, J.S. Lintuvuori, W.C.K. Poon, Swimming in a Crystal, *Soft Matter* (2015), submitted for publication, arXiv:1411.6847v3, <http://arxiv.org/abs/1411.6847v3>.
- [80] C. Niu, J.D. Graves, F.O. Mokuolu, S.E. Gilbert, E.S. Gilbert, Enhanced swarming of bacteria on agar plates containing the surfactant tween 80, *J. Microbiol. Methods* 62 (2005) 129–132.
- [81] N. Koumakis, A. Lepore, C. Maggi, R. Di Leonardo, Targeted delivery of colloids by swimming bacteria, *Nat. Commun.* 4 (2013) 2588.
- [82] A.P. Berke, L. Turner, H.C. Berg, E. Lauga, Hydrodynamic attraction of swimming microorganisms by surfaces, *Phys. Rev. Lett.* 101 (2008) 038102.
- [83] N.-C. Choi, S.-J. Park, C.-G. Lee, J.-A. Park, S.-B. Kim, Influence of surfactants on bacterial adhesion to metal oxide-coated surfaces, *Environ. Eng. Res.* 16 (2011) 219–225.
- [84] T. Besseling, M. Hermes, A. Kuijk, B. de Nijs, T.-S. Deng, M. Dijkstra, A. Imhof, A. van Blaaderen, Determination of the positions and orientations of concentrated rod-like colloids from 3D microscopy data, *J. Phys.: Condens. Matter* 27 (2015) 194109, <http://dx.doi.org/10.1088/0953-8984/27/19/194109>, arXiv preprint arXiv:1406.4985.
- [85] A. Jepson, V.A. Martinez, J. Schwarz-Linek, A. Morozov, W.C.K. Poon, Enhanced diffusion of nonswimmers in a three-dimensional bath of motile bacteria, *Phys. Rev. E* 88 (2013), 041002(R).
- [86] R. Cerbino, V. Trappe, Differential dynamic microscopy: probing wave vector dependent dynamics with a microscope, *Phys. Rev. Lett.* 100 (2008) 188102.
- [87] L.G. Wilson, V.A. Martinez, J. Schwarz-Linek, J. Tailleur, G. Bryant, P.N. Pusey, W.C.K. Poon, Differential dynamic microscopy of bacterial motility, *Phys. Rev. Lett.* 106 (2011) 018101.
- [88] V.A. Martinez, R. Besseling, O.A. Croze, J. Tailleur, M. Reufer, J. Schwarz-Linek, L.G. Wilson, M.A. Bees, W.C.K. Poon, Differential dynamic microscopy: a high-throughput method for characterizing the motility of microorganisms, *Biophys. J.* 103 (2012) 1637–1647.
- [89] M. Reufer, V.A. Martinez, P. Schurtenberger, W.C.K. Poon, Differential dynamic microscopy for anisotropic colloidal dynamics, *Langmuir* 28 (2012) 4618–4624.
- [90] M. Paoluzzi, R. Di Leonardo, L. Angelani, Effective run-and-tumble dynamics of bacteria baths, *J. Phys.: Condens. Matter* 25 (2013) 415102.
- [91] V.A. Martinez, J. Schwarz-Linek, M. Reufer, L.G. Wilson, A.N. Morozov, W.C.K. Poon, Flagellated bacterial motility in polymer solutions, *Proc. Natl. Acad. Sci. U. S. A.* 111 (2014) 17771–17776.
- [92] G. Li, J.X. Tang, Accumulation of microswimmers near a surface mediated by collision and rotational brownian motion, *Phys. Rev. Lett.* 103 (2009) 078101.
- [93] E. Lauga, W.R. Di Luzio, G.M. Whitesides, H.A. Stone, Swimming in circles: motion of bacteria near solid boundaries, *Biophys. J.* 90 (2006) 400–412.
- [94] E. Guyon, J.P. Hulin, L. Petit, C.D. Matescu, *Phys. Hydrodyn.*, OUP, Oxford, 2001.
- [95] A.J. Wolfe, H.C. Berg, Migration of bacteria in semisolid agar, *Proc. Natl. Acad. Sci. U. S. A.* 86 (1989) 6973–6977.
- [96] O.A. Croze, G.P. Ferguson, M.E. Cates, W.C.K. Poon, Migration of chemotactic bacteria in soft agar: role of gel concentration, *Biophys. J.* 101 (2011) 525–534.
- [97] J.D. Murray, *Mathematical Biology I: An Introduction*, 3rd ed., Springer, Berlin, 2002.
- [98] P.S. Loveley, F.W. Dahlquist, *J. Theor. Biol.* 50 (1975) 477–496.
- [99] K. Barker, *At the Bench: A Laboratory Navigator*, Cold Spring Laboratory Publisher Press, New York, 1998.
- [100] P.L. Bonner, A.J. Hargreaves, *Basic Bioscience Laboratory Techniques: A Pocket Guide*, Wiley, Chichester, 2011.
- [101] G. Bertani, Lysogeny at mid-twentieth century: P1, P2, and other experimental systems, *J. Bacteriol.* 186 (2004) 595–600.
- [102] J. Adler, B. Templeton, The effect of environmental conditions on the motility of *Escherichia coli*, *J. Gen. Microbiol.* 46 (1967) 175–184.
- [103] L. Turner, A.S. Stern, H.C. Berg, Growth of flagellar filaments of *Escherichia coli* is independent of filament length, *J. Bacteriol.* 194 (2012) 2437–2442.
- [104] B. Sabass, U. Seifert, Dynamics and efficiency of a self-propelled, diffusiophoretic swimmer, *J. Chem. Phys.* 136 (2012) 064508.
- [105] W. Wang, T.Y. Chiang, D. Velegol, T.E. Mallouk, Understanding the efficiency of autonomous nano- and microscale motors, *J. Am. Chem. Soc.* 135 (2013) 10557–10565.
- [106] F.C. Neidhardt, J.L. Ingraham, M. Schaechter, *Physiology of the Bacterial Cell: A Molecular Approach*, Sinauer Associates, Sunderland, Massachusetts, 1990.
- [107] J.M. Berg, L. Stryer, J.L. Tymoczko, N.D. Clarke, *Biochemistry*, 5th revised ed., W.H. Freeman, New York, 2002.
- [108] A.G. Moat, J.W. Foster, M.P. Spector (Eds.), *Microbial Physiology*, 4th ed., Wiley, 2002.
- [109] J.E. McMuray, T.P. Begley, *The Organic Chemistry of Biological Pathways*, Roberts and Company, Englewood, Colorado, 2005.
- [110] B.H. Kim, G.M. Gadd, *Bacterial Physiology and Metabolism*, Cambridge University Press, Cambridge, 2008.
- [111] D.G. Nicholls, S. Ferguson, *Bioenergetics*, 4th ed., Elsevier, Amsterdam, 2013.
- [112] C. de Duve, *Singularities*, Cambridge University Press, Cambridge, 2005.
- [113] *E. coli Student Portal*, <http://ecolistudentportal.org/> (last accessed 17.07.15).
- [114] G. Uden, Chapter 11: Aerobic respiration and regulation of aerobic/anaerobic metabolism, in: J. Lengeler, G. Drews, H. Schlegel (Eds.), *Biology of the Prokaryotes*, Blackwell, Oxford, 1999, pp. 261–277.
- [115] G.N. Vemuri, E. Altman, D.P. Sangurdekar, A.B. Khodursky, M.A. Eiteman, Overflow metabolism in *Escherichia coli* during steady-state growth: transcriptional regulation and effect of the redox ratio, *Appl. Environ. Microbiol.* 72 (2006) 3653–3661.
- [116] Extensive, exometabolome analysis reveals extended overflow metabolism in various microorganisms, *Microb. Cell Factor.* 11 (2012) 122.
- [117] A.D. Thomas, H.W. Doelle, A.W. Westwood, G.L. Gordon, Effect of oxygen on several enzymes involved in the aerobic and anaerobic utilization of glucose in *Escherichia coli*, *J. Bacteriol.* 112 (1972) 1099–1105.
- [118] M.W. Smith, F.C. Neidhardt, Proteins induced by anaerobiosis in *Escherichia coli*, *J. Bacteriol.* 154 (1983) 336–343.
- [119] M.W. Smith, F.C. Neidhardt, Proteins induced by aerobiosis in *Escherichia coli*, *J. Bacteriol.* 154 (1983) 344–350.
- [120] A.M. Sánchez, G.N. Bennett, K.Y. San, Novel pathway engineering design of the anaerobic central metabolic pathway in *Escherichia coli* to increase succinate yield and productivity, *Metab. Eng.* 7 (2005) 229–239.
- [121] C.V. Gabel, H. Berg, The speed of the flagellar rotary motor of *Escherichia coli* varies linearly with protonmotive force, *Proc. Natl. Acad. Sci. U. S. A.* 100 (2003) 8748–8751.

- [122] M.D. Manson, P. Tedesco, H.C. Berg, F.M. Harold, C. van der Drift, A proton motive force drives bacterial flagella, *Proc. Natl. Acad. Sci. U. S. A.* 74 (1977) 3060–3064.
- [123] D.W. Ribbons, E.A. Dawes, Environmental and growth conditions affecting the endogenous metabolism of bacteria, *Ann. N. Y. Acad. Sci.* 102 (1963) 564–586.
- [124] E.A. Dawes, D.W. Ribbons, Studies on the endogenous metabolism of *Escherichia coli*, *Biochem. J.* 95 (1965) 332–343.
- [125] C.S. Harwood, K. Fosnaugh, M. Dispensa, Flagellation of *Pseudomonas putida* and analysis of its motile behavior, *J. Bacteriol.* 171 (1989) 4063–4066.
- [126] P. Lewus, R.M. Ford, Temperature-sensitive motility of *Sulfolobus acidocaldarius* influences population distribution in extreme environments, *J. Bacteriol.* 181 (1999) 4020–4025.
- [127] F. Rivera-Chavez, S.E. Winter, C.A. Lopez, M.N. Xavier, M.G. Winter, S.-P. Nuccio, J.M. Russell, R.C. Laughlin, S.D. Lawhon, T. Sterzenbach, C.L. Bevins, R.M. Tsois, R. Harshey, L.G. Adams, A.J. Baumler, *Salmonella* uses energy taxis to benefit from intestinal inflammation, *PLOS Pathog.* 9 (2013) e1003267.
- [128] J.H. Carpenter, New measurement of oxygen solubility in pure and natural water, *Limnol. Oceanogr.* 11 (1966) 264–277.
- [129] H.N. Rasmussen, U.F. Rasmussen, Oxygen solubilities of media used in electrochemical respiration measurements, *Anal. Biochem.* 319 (2003) 105–113.
- [130] A. Boehm, M. Kaiser, H. Li, C. Spangler, C.A. Kasper, M. Ackermann, V. Kaever, V. Sourjik, V. Roth, U. Jenal, Second messenger-mediated adjustment of bacterial swimming velocity, *Cell* 141 (2010) 107–118.
- [131] B. Görke, J. Stülke, Carbon catabolite repression in bacteria: many ways to make the most out of nutrients, *Nat. Rev. Microbiol.* 6 (2008) 613–624.
- [132] H.C. Lai, J.C. Shu, S. Ang, M.J. Lai, B. Fruta, S. Lin, K.Y. Lu, S.W. Ho, Effect of glucose concentration on swimming motility in enterobacteria, *Biochem. Biophys. Res. Commun.* 231 (1997) 692–695.
- [133] H.C. Gerritsen, R. Sanders, A. Draaijer, C. Ince, Y.K. Levine, Fluorescence lifetime imaging of oxygen in living cells, *J. Fluoresc.* 7 (1997) 11–15.
- [134] K.J. Morris, M.S. Roach, W. Xu, J.N. Demas, B.A. DeGaff, Luminescence lifetime standards for the nanosecond to microsecond range and oxygen quenching of ruthenium(II) complexes, *Anal. Chem.* 79 (2007) 9310–9314.
- [135] C. Douarche, A. Buguin, H. Salman, A. Libchaber, *E. coli* and oxygen: a motility transition, *Phys. Rev. Lett.* 102 (2009) 198101.
- [136] T. Pilizota, J. Shaevitz, Plasmolysis and cell shape depend on solute outer-membrane permeability during hyperosmotic shock in *E. coli*, *Biophys. J.* 104 (2013) 2733–2742.
- [137] H. De Reuse, A. Danchin, The *ptsH*, *ptsI*, and *crr* genes of the *Escherichia coli* phosphoenolpyruvate-dependent phosphotransferase system: a complex operon with several modes of transcription, *J. Bacteriol.* 170 (1988) 3827–3837.
- [138] P.W. Postma, J.W. Lengeler, G.R. Jacobson, Phosphoenolpyruvate:carbohydrate phosphotransferase systems, in: F.C. Neidhardt (Ed.), *Escherichia coli* and *Salmonella*: Cellular and Molecular Biology, ASM Press, Washington, D.C., 1996, pp. 1149–1174.
- [139] Q.H. Tran, G. Unden, Changes in the proton potential and the cellular energetics of *Escherichia coli* during growth by aerobic and anaerobic respiration or by fermentation, *Eur. J. Biochem.* 251 (1998) 538–543.
- [140] S. Chattopadhyay, R. Moldovan, C. Yeung, X.L. Wu, Swimming efficiency of bacterium *Escherichia coli*, *Proc. Natl. Acad. Sci. U. S. A.* 103 (2006) 13712–13717.
- [141] J.I. Shioi, S. Matsuura, Y. Imae, Quantitative measurements of proton motive force and motility in *Bacillus subtilis*, *J. Bacteriol.* 144 (1980) 891–897.
- [142] O. Bèjà, L. Aravind, E.V. Koonin, M.T. Suzuki, A. Hadd, L.P. Nguyen, S.B. Jovanovich, C.M. Gates, R.A. Feldman, J.L. Spudich, E.N. Spudich, E.F. DeLong, Bacterial rhodopsin: evidence for a new type of phototrophy in the sea., *Science* 289 (2000) 1902–1906.
- [143] J.M. Walter, D. Greenfield, C. Bustamante, J. Liphardt, Light-powering *Escherichia coli* with proteorhodopsin., *Proc. Natl. Acad. Sci. U. S. A.* 104 (2007) 2408–2412.
- [144] M.J. Tipping, B.C. Steel, N.J. Delalez, R.M. Berry, J.P. Armitage, Quantification of flagellar motor stator dynamics through in vivo proton-motive force control, *Mol. Microbiol.* 87 (2013) 338–347.
- [145] I. Buttinoni, G. Volpe, F. Kümmel, G. Volpe, C. Bechinger, Active Brownian motion tunable by light, *J. Phys.: Condens. Matter* 24 (2012) 284129.
- [146] J.W. Dale, S.F. Park, *Molecular Genetics of Bacteria*, 5th ed., Wiley, Chichester, 2010.
- [147] T. Ara, M. Hasegawa, Y. Takai, Y. Okumura, M. Baba, K.A. Datsenko, M. Tomita, B.L. Wanner, H. Mori, Construction of *Escherichia coli* K-12 in-frame, single-gene knockout mutants: the Keio collection, *Mol. Syst. Biol.* 2 (2006), 2006.0008.
- [148] M. Kitagawa, T. Ara, M. Arifuzzaman, T. Ioka-Nakamichi, E. Inamoto, H. Toyonaga, H. Mori, Complete set of ORF clones of *Escherichia coli* ASKA library (a complete set of *E. coli* K-12 ORF archive): unique resources for biological research, *DNA Res.* 12 (2006) 291–299.
- [149] American Type Culture Collection, <http://www.atcc.org> (last accessed on 17.07.15).
- [150] *E. coli* Genetic Stock Centre at Yale, <http://cgsc.biology.yale.edu> (last accessed on 17.07.15).
- [151] National Collections of Industrial, Marine and Food Bacteria, <http://www.ncimb.com> (last accessed on July 17.07.15).
- [152] Deutsche Sammlung von Mikroorganismen und Zellkulturen, <http://www.dsmz.de> (last accessed on 17.07.15).
- [153] China General Microbiological Culture Collection, <http://www.cgmcc.net/english/Deposit.html> (last accessed on 17.07.15).
- [154] C.S. Barker, B.M. Prüß, P. Matsumura, Increased motility of *Escherichia coli* by insertion sequence element integration into the regulatory region of the *flhD* operon, *J. Bacteriol.* 186 (2004) 7529–7537.
- [155] B.J. Bachmann, Pedigrees of some mutant strains of *Escherichia coli* K-12, *Bacteriol. Rev.* 36 (1972) 525–557.
- [156] J.S. Parkinson, S.E. Houts, Isolation and behavior of *Escherichia coli* deletion mutants lacking chemotaxis functions, *J. Bacteriol.* 151 (1982) 106–113.
- [157] A.J. Wolfe, M.P. Conley, T.J. Kramer, H.C. Berg, Reconstitution of signaling in bacterial chemotaxis, *J. Bacteriol.* 169 (1987) 1878–1885.
- [158] H.-P. Cheng, G.C. Walker, Succinoglycan is required for initiation and elongation of infection threads during nodulation of Alfalfa by *Rhizobium meliloti*, *J. Bacteriol.* 180 (1996) 5183–5191.
- [159] M. Demerec, E.A. Adelberg, A.J. Clark, P.E. Hartman, A proposal for uniform nomenclature in bacterial genetics, *Genetics* 54 (1966) 61–76.
- [160] S.R. Maloy, K.T. Hughes, Strain collections and genetic nomenclature, *Methods Enzymol.* 421 (2007) 3–8.
- [161] EcoCyc, <http://ecocyc.org> (last accessed on 17.07.15).
- [162] H.C. Berg, L. Turner, *Biophys. J.* 58 (1990) 919–930.
- [163] J. Schwarz-Linek, J. Arlt, A. Jepson, A. Dawson, T. Vissers, D. Miroli, T. Pilizota, V.A. Martinez, W.C.K. Poon, *Escherichia coli* as a model active colloid: a practical introduction [dataset], University of Edinburgh, 2015, <http://dx.doi.org/10.7488/ds/290>.