

Quantum Measurement

Zeno Effect

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30.01.2025

Modeling Quantum Hardware: open dynamics and control
Universität Konstanz

No phenomenon is a real until it is observed.

– John Archibald Wheeler 1970

Historical Note

1900 Plank & Einstein: Blackbody Radiation

1920 Bohr, Heisenberg: Copenhagen interpretation

Born: Probabilistic interpretation $P(m) = |\langle m | \psi \rangle|^2$

Schrödinger: Measurement Problem

1930 EPR Paradox

1932 von Neumann: *Mathematical Foundations of Quantum Mechanics*

1970 Decoherence Theory

Experimental Interest

Projective (von Neumann) Measurement

Measurement Operator $\hat{A} = \sum m|m\rangle$ on ψ :

$$p(m) = |\langle m|\psi\rangle|^2$$

$$\psi \xrightarrow{\text{Measuring } m} |m\rangle$$

More formally with projector $\hat{M} \sim |m\rangle\langle m|$:

$$\rho' \propto \hat{M}\rho\hat{M}$$

Neglegting Normalization and Degenercy: POVM Measurement

Quantum Computation and Quantum Information, M. A. Nielsen et al. 2010 (2.2.3)

Quantum Measurement and Control, Wiseman et al. 2010 (1)

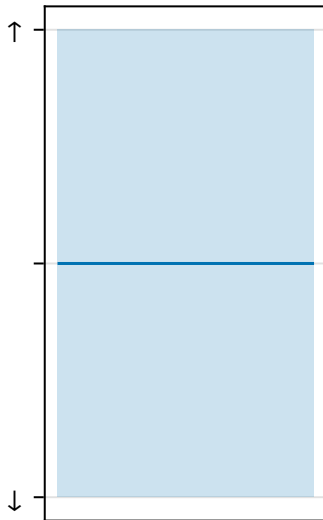
Example: Superposition

$$H = \sigma_z$$

$$|\psi\rangle \propto |\uparrow\rangle + |\downarrow\rangle$$

$$i\partial_t \psi = H\psi$$

⇒ Superposition is stable



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⇒ Superposition is stable

$$M = \sigma_z$$

$$p(\uparrow) = |\langle\uparrow|\psi\rangle|^2 = \frac{1}{2}$$



Example: Decay

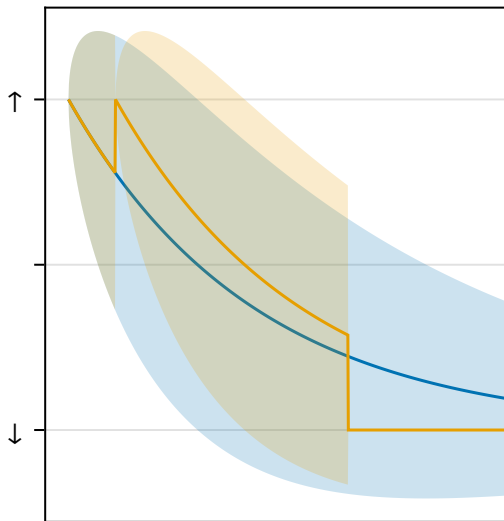
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$$M = \sigma_z$$

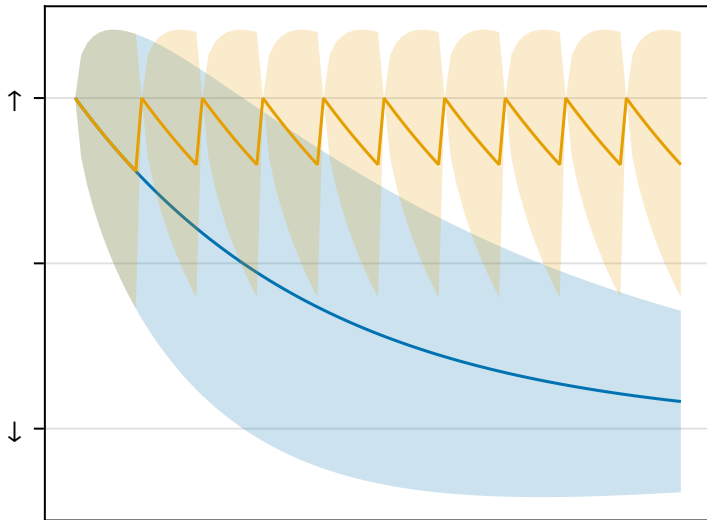
$$J = \kappa \sigma_-$$

$$\partial_t \rho = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$



Strong Measurement

Example: Zeno



Strong Measurement

Zeno Effect

- ▶ Zeno of Elea (460 BCE): Arrow paradox
- ▶ Misra and Sudarshan (1977):
“The Zeno’s paradox in quantum theory”

“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009

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- ▶ Experimentally demonstrated (1990) with 5000 $^9\text{Be}^+$ ions at 250 mK

“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009

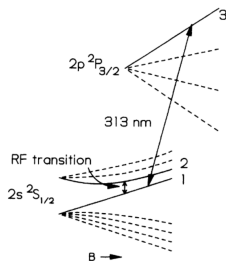


FIG. 2. Diagram of the energy levels of $^9\text{Be}^+$ in a magnetic field B . The states labeled 1, 2, and 3 correspond to those in Fig. 1.

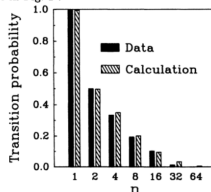


FIG. 3. Graph of the experimental and calculated $1 \rightarrow 2$ transition probabilities as a function of the number of measurement pulses n . The decrease of the transition probabilities with increasing n demonstrates the quantum Zeno effect.

“Quantum Zeno effect”, Itano et al. 1990

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- ▶ Misra and Sudarshan (1977): “The Zeno’s paradox in quantum theory”
- ▶ Experimentally demonstrated (1990) with 5000 $^9\text{Be}^+$ ions at 250 mK
- ▶ Used in Magnetometers and possibly birds

“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009

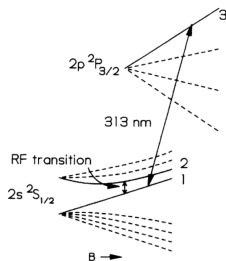


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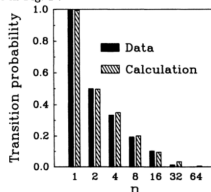


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Practical Examples of strong measurements

► Trapped Ions

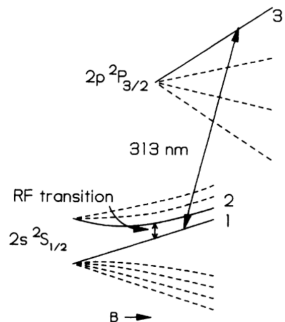
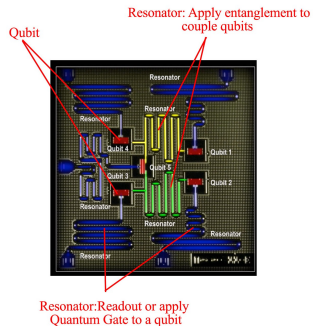


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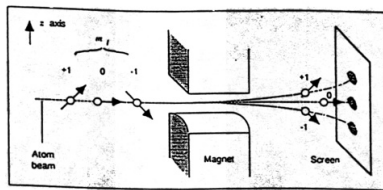
Practical Examples of strong measurements

- ▶ Trapped Ions
- ▶ Superconducting Qubits



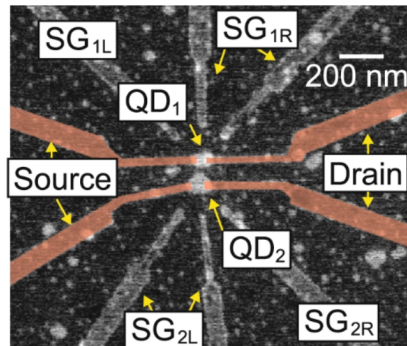
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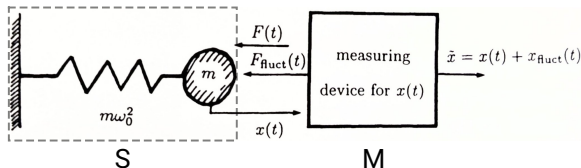
⇒ Used when the quantum wave function is collapsed into a classical result

But breaks on:

- ▶ simultaneous measurement of conjugate variables
- ▶ optical measurements

Quantum Measurement and Control, Wiseman et al.
2010

Weak Measurement



Measured system S and Measurement Device S combined to:

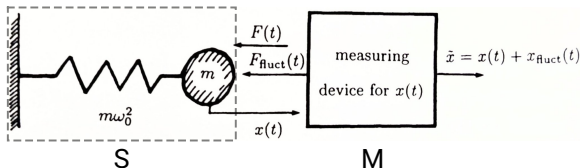
$$H = H_S \otimes 1 + 1 \otimes H_M + g(t) C_S \otimes P_M$$

Then a strong measurement X is performed on the M
this yields $\langle C \rangle + \xi$ and partial collapse in S

Quantum measurement, Vladimir Braginsky et al. 1992 (Fig 8.4 modified)

"Introduction to Weak Measurements and Weak Values", Tamir et al. 2013

Weak Measurement



Continuous application leads to *SME* for S :

$$\partial_t \rho = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\} + (C\rho + \rho C^\dagger - \text{Tr}(C\rho + \rho C^\dagger)) \xi(t)$$

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Example: Weak measurement

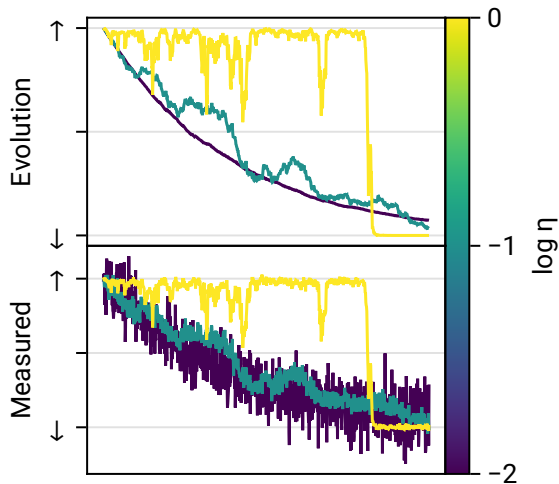
$$H = \sigma_z$$

$$\psi = \uparrow$$

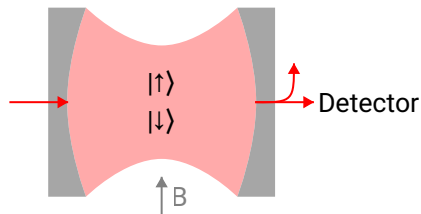
$$M = \sigma_z$$

$$J = \kappa \sigma_-$$

$$C = \eta \sigma_z$$



Example: Weak measurement



$$H = \sigma_z + a^\dagger a + g \sigma_z a^\dagger a$$

$$J = \kappa a$$

$$C = \kappa \eta a$$

Atom

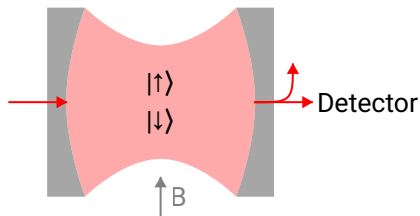
M

Dissipation

Measurement

"Stochastic master equation for a probed system in a cavity", A. E. B. Nielsen et al. 2008

Rabi Oscillations Setup



$$H = g (a^\dagger a)(\sigma^+ \sigma^-)$$

$$+ g_s (\sigma^+ + \sigma^-)$$

$$- i\beta(a^\dagger - a)$$

$$J = \kappa a$$

$$C = \sqrt{\kappa\eta} a$$

Coupling

Magnetic

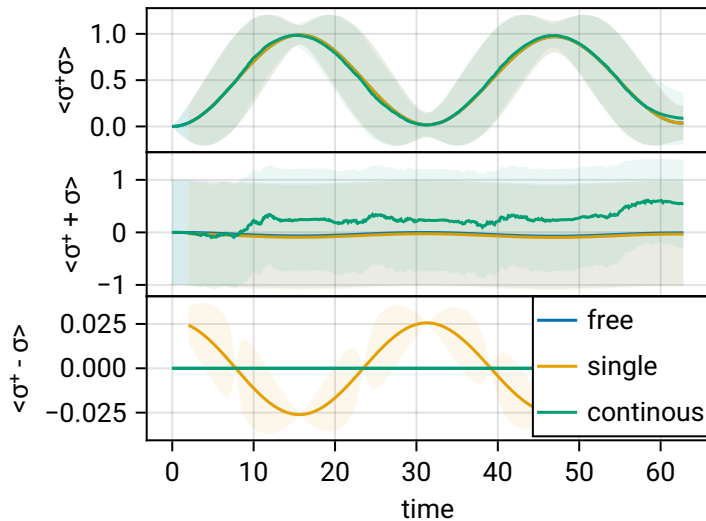
Optic

Dissipation

Measurement

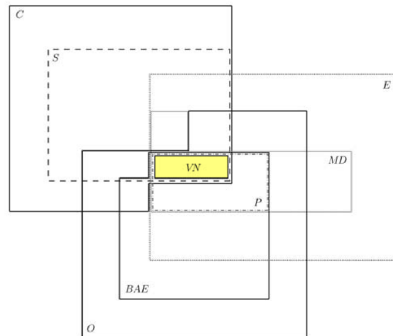
"Stochastic master equation for a probed system in a cavity", A. E. B. Nielsen et al. 2008

Time evolution



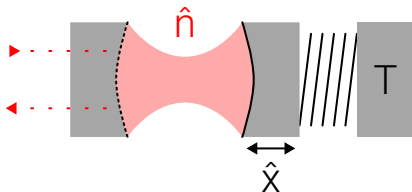
Types of Measurements

Symbol	Name	Definition
E	Efficient	$\forall r, \exists \hat{M}_r, \mathcal{O}_r = \mathcal{J}[\hat{M}_r]$
C	Complete	$\forall \rho, \forall r, \mathcal{O}_r \rho \propto \mathcal{O}_r \hat{1}$
S	Sharp	$\forall r, \text{rank}(\hat{E}_r) = 1$
O	Of an observable X	$\forall r, \hat{E}_r = E_r(\hat{X})$
BAE	Back-action-evading	O with $\forall \rho, \forall x \in \lambda(\hat{X}), \text{Tr}[\hat{\Pi}_x \rho] = \text{Tr}[\hat{\Pi}_x \mathcal{O} \rho]$
MD	Minimally disturbing	E with $\forall r, \hat{M}_r = \hat{M}_r^\dagger$
P	Projective	MD and O
VN	von Neumann	P and S



Quantum Measurement and Control, Wiseman et al. 2010

Problem Statement



"Cavity optomechanics", Aspelmeyer et al. 2014

Quantum Optomechanics, Bowen et al. 2015

Hamiltonian

Optical Cavity \hat{a} , $\omega_o(\hat{x}_{\text{mech}}) = \omega_o + \frac{g}{\omega_o} \hat{x}_{\text{mech}}$; mechanical oscillations \hat{b} , ω_m ; coupling g ; Drive E , ω_L

$$H = \underbrace{\omega_o a^\dagger a}_{\text{Cavity}} + \underbrace{\omega_m b^\dagger b}_{\text{Mechanical}} - \underbrace{g a^\dagger a (b + b^\dagger)}_{\text{Interaction}} + \underbrace{E(ae^{i\omega_L t} + a^\dagger e^{-i\omega_L t})}_{\text{Drive}}$$

Rotating Wave Approximation at ω_L with $\Delta = \omega_o - \omega_L$, $a \rightarrow ae^{i\omega_L t}$:

$$H_{\text{RWA}} = \Delta a^\dagger a + \omega_m b^\dagger b - g a^\dagger a (b^\dagger + b) + E(a + a^\dagger)$$

$\hbar = 1$

Quantum Optomechanics, Bowen et al. 2015 (2.3)

QuantumOptics.jl, Krämer et al. 2024 (Optomechanical Cavity)

Hamiltonian Linearization (Currently not used)

$$H_{\text{RWA}} = \Delta a^\dagger a + \omega_m b^\dagger b - g a^\dagger a (b^\dagger + b) + E(a + a^\dagger)$$

Linearize $a = \alpha + \delta a$, $b = \beta + \delta b$; with α, β steady state.

$$\begin{aligned} H_{\text{Interaction}} &= -g a^\dagger a (b^\dagger + b) \\ &\approx -\underbrace{g|\alpha|}_G (\delta a + \delta a^\dagger + \mathcal{O}(a^2 + \delta a \delta a^\dagger)) (\delta b + \delta b^\dagger + 2\beta) \\ a + a^\dagger &= |\alpha| + \delta a + \delta a^\dagger \sim \delta a + \delta a^\dagger \end{aligned}$$

Therefore for small G :

$$H_{\text{lin}} = \Delta \delta a^\dagger \delta a + \omega_m \delta b^\dagger \delta b - G(\delta a + \delta a^\dagger)(\delta b + \delta b^\dagger) + E(a + a^\dagger)$$

Problem Statement $\sim \frac{\Delta}{2}(\hat{X}^2 + \hat{Y}^2) + \frac{\omega}{2}(\hat{Q}^2 + \hat{P}^2) - G\hat{X}\hat{Q} + E\hat{X}$

Linearization in Quadratures

$$\begin{aligned}X &= \delta a + \delta a^\dagger & Q &= \delta b + \delta b^\dagger \\Y &= i(\delta a^\dagger - \delta a) & P &= i(\delta b^\dagger - \delta b) \\n &= \delta a^\dagger a & m &= \delta b^\dagger b\end{aligned}$$

$$H_{\text{RWA}} = \Delta n + \omega m - gnQ + EX$$

$$H_{\text{lin}} = \Delta n + \omega m - GXQ + EX$$

The drive EX is not getting lost in linearization.

There is no point in the simplification if solved numerically.

Dissipation

Optical decay κ :

$$L = \sqrt{\kappa(n_T + 1)} \delta a + \sqrt{\kappa n_T} \delta a^\dagger$$

Mechanical resonator with γ and a thermal bath at the n -th thermal state:

$$+ \sqrt{\gamma(m_T + 1)} \delta b + \sqrt{\gamma m_T} \delta b^\dagger$$

Quantum Optomechanics, Bowen et al. 2015 (2.8)

Implementation

truncated Fock Basis: $F_{\text{optical}} \otimes F_{\text{mechanical}}$

definition of H, J with $\delta a \otimes 1$

$$\psi(0) = |0\rangle \otimes |0\rangle$$

Time Evolution using the *Lindblad equation*:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$

Continuous measurement

Lindblad Master Equation:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$

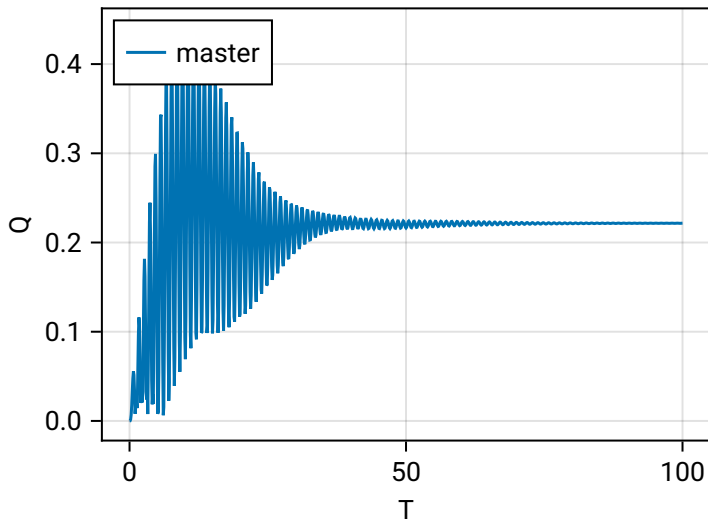
Stochastic Master Equation:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\} + (C\rho + \rho C^\dagger - \text{Tr}(C\rho + \rho C^\dagger)) \xi(t)$$

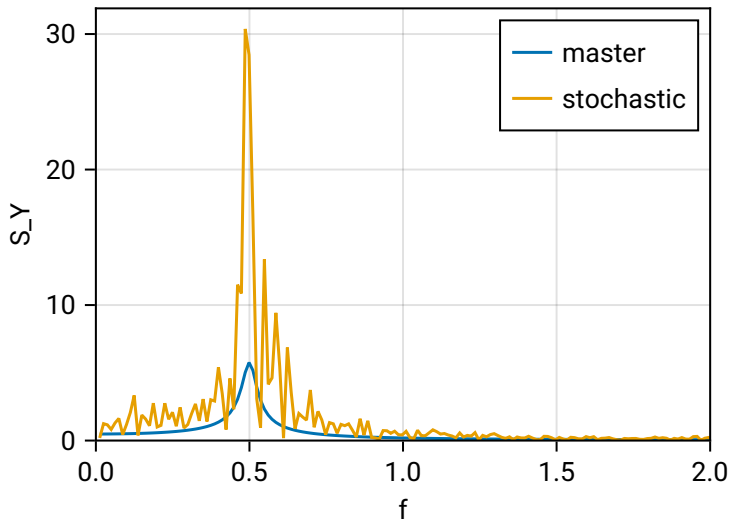
Let's look at the Quadrature $C = \eta\sqrt{\kappa} (\delta a + \delta a^\dagger)$

QuantumOptics.jl, Krämer et al. 2024 (Stochastic Master equation, Quantum Zeno Effect)
"A Straightforward Introduction to Continuous Quantum Measurement", Jacobs et al. 2006

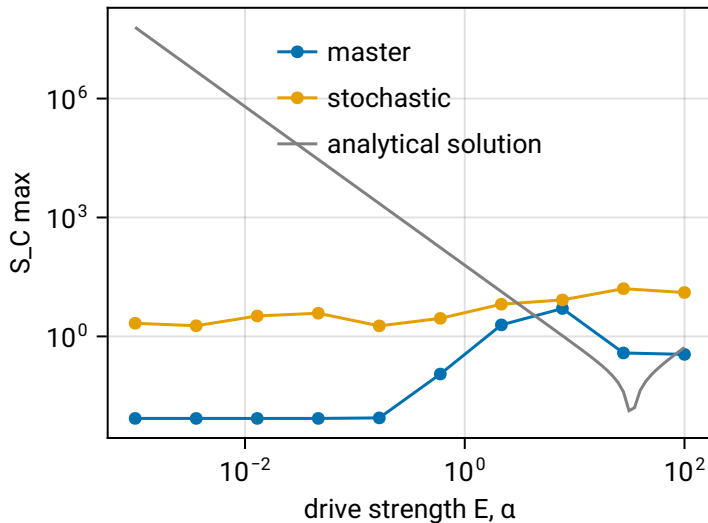
Time Evolution



Spectrum



Power Dependence $\sim G$



Took ≈ 1 min of compute time. Why is the SME so much slower? True Random Values?
Measurement

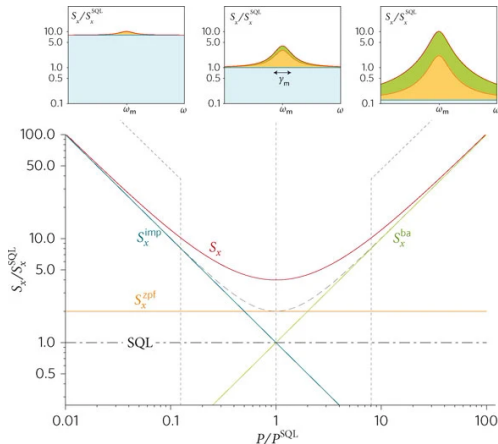
Expectation

$$S_x = S_x^{\text{imp}} + S_x^{\text{ba}} + S_x^{\text{zpf}}$$

$S_x^{\text{imp}} \propto P^{-1} \omega^0$: Imprecision /
Shot noise

$S_x^{\text{ba}} \propto P$: Back action

S_x^{zpf} : Zero Point Fluctuation



"Nanomechanical motion measured with an imprecision below that at the standard quantum limit", Teufel et al. 2009 (Fig. 1)

Understanding $S_{\text{det}}(\omega, P_{\text{in}})$

Where is the power dependence?

$$\bar{S}_{\text{det}}(\omega) = \frac{1}{8\eta\Gamma|C_{\text{eff}}|} + 2\Gamma|\chi(\omega)|^2|C_{\text{eff}}|$$

$$C_{\text{eff}}(\omega) = \frac{4g^2}{\kappa\Gamma(1-2i\omega/\kappa)^2}$$

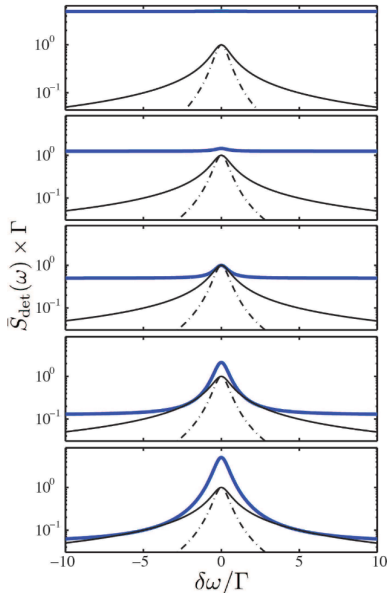
$$\chi(\omega) = \frac{\Omega}{\Omega^2 - \omega^2 - i\omega\gamma}$$

η : Detection efficiency

$\Gamma = \gamma$: Damping of oscillator

$$\bar{S}_{\text{det}}(\omega) = \frac{\kappa\Gamma|1-2i\omega/\kappa|}{8\eta\Gamma 4g^2} + 2\Gamma \frac{\Omega^2}{|\Omega^2 - \omega^2 - i\omega\gamma|^2} \frac{4g^2}{\kappa\Gamma(1-2i\omega/\kappa)^2}$$

Quantum Optomechanics, Bowen et al. 2015 (eq. 3.51)



$$P/P^{\text{SQL}} = \{0.1, 0.4, 1, 4, 10\}$$

Looking for a source that derives $S_{\text{det}}(\omega, P)$?