

# Quantum Measurement

## Zeno Effect

Leon Oleschko  
30.01.2025

*Modeling Quantum Hardware: open dynamics and control*  
Universität Konstanz

No phenomenon is a real until it is observed.

– John Archibald Wheeler 1970

# Historical Note

1900 Plank & Einstein: Blackbody Radiation

1920 Bohr, Heisenberg: Copenhagen interpretation

Born: Probabilistic interpretation  $P(m) = |\langle m | \psi \rangle|^2$

Schrödinger: Measurement Problem

1930 EPR Paradox

1932 von Neumann: *Mathematical Foundations of Quantum Mechanics*

1970 Decoherence Theory

Experimental Interest

# Projective (von Neumann) Measurement

Measurement Operator  $\hat{A} = \sum m|m\rangle$  on  $\psi$ :

$$p(m) = |\langle m|\psi\rangle|^2$$

$$\psi \xrightarrow{\text{Measuring } m} |m\rangle$$

More formally with projector  $\hat{M} \sim |m\rangle\langle m|$ :

$$\rho' \propto \hat{M}\rho\hat{M}$$

Neglegting Normalization and Degenercy: POVM Measurement

*Quantum Computation and Quantum Information*, M. A. Nielsen et al. 2010 (2.2.3)

*Quantum Measurement and Control*, Wiseman et al. 2010 (1)

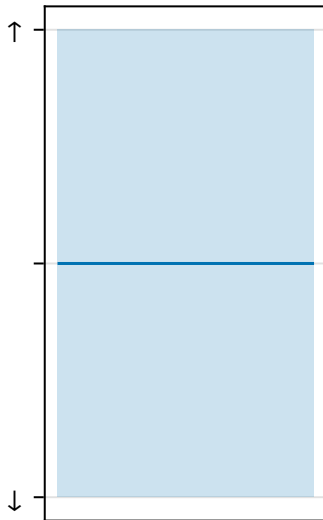
# Example: Superposition

$$H = \sigma_z$$

$$|\psi\rangle \propto |\uparrow\rangle + |\downarrow\rangle$$

$$i\partial_t \psi = H\psi$$

⇒ Superposition is stable



## Example: Superposition

$$H = \sigma_z$$

$$|\psi\rangle \propto |\uparrow\rangle + |\downarrow\rangle$$

$$i\partial_t \psi = H\psi$$

⇒ Superposition is stable

$$M = \sigma_z$$

$$p(\uparrow) = |\langle \uparrow | \psi \rangle|^2 = \frac{1}{2}$$



## Example: Decay

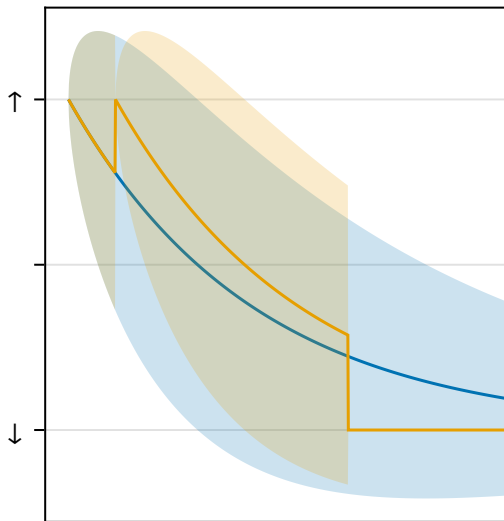
$$H = \sigma_z$$

$$|\psi\rangle \propto |\uparrow\rangle$$

$$M = \sigma_z$$

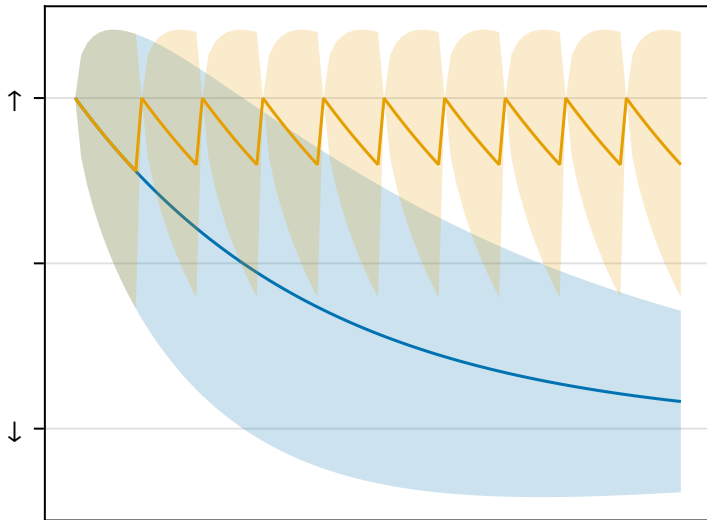
$$J = \kappa \sigma_-$$

$$\partial_t \rho = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$



**Strong Measurement**

## Example: Zeno



**Strong Measurement**



# Zeno Effect

- ▶ Zeno of Elea (460 BCE): Arrow paradox
- ▶ Misra and Sudarshan (1977):  
“The Zeno’s paradox in quantum theory”

*“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009*

# Zeno Effect

- ▶ Zeno of Elea (460 BCE): Arrow paradox
- ▶ Misra and Sudarshan (1977): “The Zeno’s paradox in quantum theory”
- ▶ Experimentally demonstrated (1990) with 5000  $^9\text{Be}^+$  ions at 250 mK

*“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009*

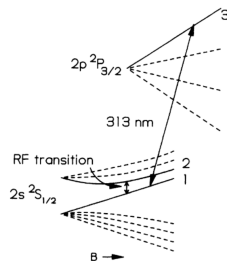


FIG. 2. Diagram of the energy levels of  $^9\text{Be}^+$  in a magnetic field  $B$ . The states labeled 1, 2, and 3 correspond to those in Fig. 1.

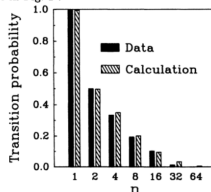


FIG. 3. Graph of the experimental and calculated  $1 \rightarrow 2$  transition probabilities as a function of the number of measurement pulses  $n$ . The decrease of the transition probabilities with increasing  $n$  demonstrates the quantum Zeno effect.

*“Quantum Zeno effect”, Itano et al. 1990*

# Zeno Effect

- ▶ Zeno of Elea (460 BCE): Arrow paradox
- ▶ Misra and Sudarshan (1977): “The Zeno’s paradox in quantum theory”
- ▶ Experimentally demonstrated (1990) with 5000  $^9\text{Be}^+$  ions at 250 mK
- ▶ Used in Magnetometers and possibly birds

*“Quantum Zeno effect explains magnetic-sensitive radical-ion-pair reactions”, Kominis 2009*

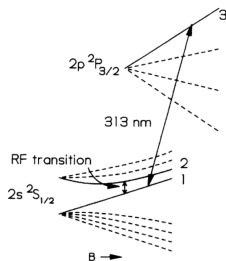


FIG. 2. Diagram of the energy levels of  $^9\text{Be}^+$  in a magnetic field  $B$ . The states labeled 1, 2, and 3 correspond to those in Fig. 1.

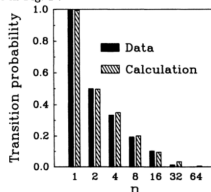


FIG. 3. Graph of the experimental and calculated  $1 \rightarrow 2$  transition probabilities as a function of the number of measurement pulses  $n$ . The decrease of the transition probabilities with increasing  $n$  demonstrates the quantum Zeno effect.

*“Quantum Zeno effect”, Itano et al. 1990*

# Practical Examples of strong measurements

## ► Trapped Ions

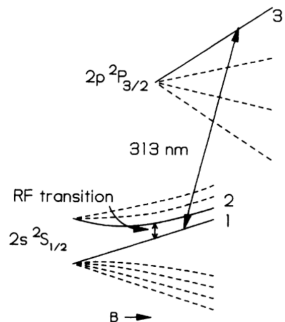
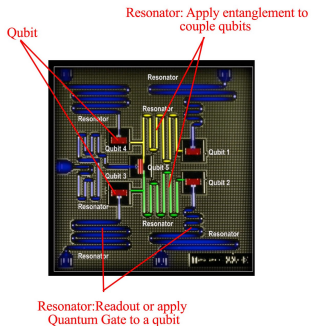


FIG. 2. Diagram of the energy levels of  ${}^9\text{Be}^+$  in a magnetic field  $B$ . The states labeled 1, 2, and 3 correspond to those in Fig. 1.

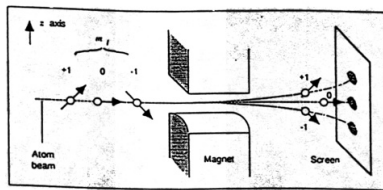
# Practical Examples of strong measurements

- ▶ Trapped Ions
- ▶ Superconducting Qubits



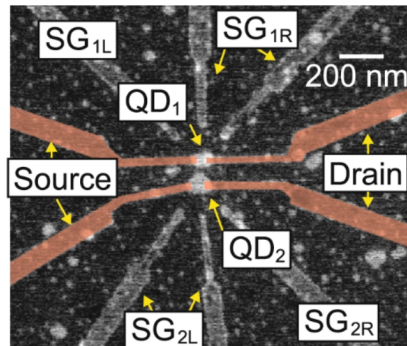
# Practical Examples of strong measurements

- ▶ Trapped Ions
- ▶ Superconducting Qubits
- ▶ Stern-Gerlach Experiment



# Practical Examples of strong measurements

- ▶ Trapped Ions
- ▶ Superconducting Qubits
- ▶ Stern-Gerlach Experiment
- ▶ Quantum Dot Charge Readout



# Practical Examples of strong measurements

- ▶ Trapped Ions
- ▶ Superconducting Qubits
- ▶ Stern-Gerlach Experiment
- ▶ Quantum Dot Charge Readout

⇒ Used when the quantum wave function is collapsed into a classical result

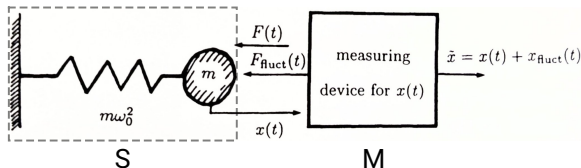
But breaks on:

- ▶ simultaneous measurement of conjugate variables
- ▶ optical measurements

*Quantum Measurement and Control*, Wiseman et al.  
2010



# Weak Measurement



Measured system  $S$  and Measurement Device  $S$  combined to:

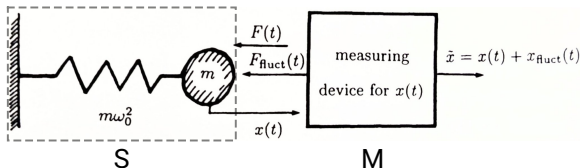
$$H = H_S \otimes 1 + 1 \otimes H_M + g(t) C_S \otimes P_M$$

Then a strong measurement  $X$  is performed on the  $M$   
this yields  $\langle C \rangle + \xi$  and partial collapse in  $S$

*Quantum measurement, Vladimir Braginsky et al. 1992 (Fig 8.4 modified)*

*"Introduction to Weak Measurements and Weak Values", Tamir et al. 2013*

# Weak Measurement



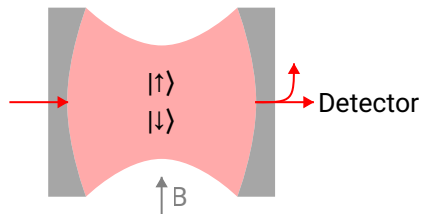
Continuous application leads to *SME* for  $S$ :

$$\partial_t \rho = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\} + (C\rho + \rho C^\dagger - \text{Tr}(C\rho + \rho C^\dagger)) \xi(t)$$

*Quantum measurement*, Vladimir Braginsky et al. 1992 (Fig 8.4 modified)

*"Introduction to Weak Measurements and Weak Values"*, Tamir et al. 2013

## Example: Weak measurement



$$H = \sigma_z + a^\dagger a + g \sigma_z a^\dagger a$$

$$J = \kappa a$$

$$C = \kappa \eta a$$

Atom

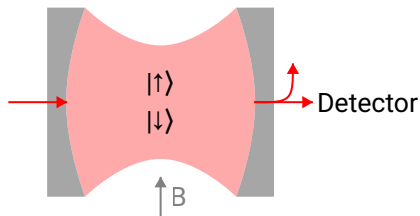
M

Dissipation

Measurement

*"Stochastic master equation for a probed system in a cavity", A. E. B. Nielsen et al. 2008*

# Rabi Oscillations Setup



$$H = g (a^\dagger a)(\sigma^+ \sigma^-)$$

$$+ g_s (\sigma^+ + \sigma^-)$$

$$- i\beta(a^\dagger - a)$$

$$J = \kappa a$$

$$C = \sqrt{\kappa\eta} a$$

Coupling

Magnetic

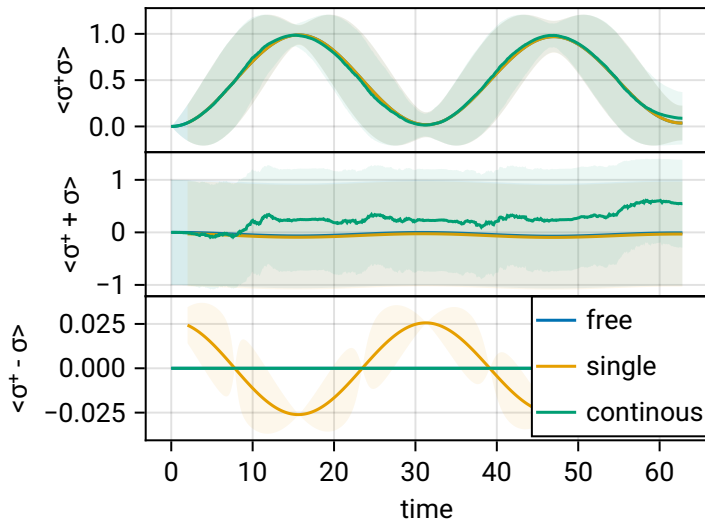
Optic

Dissipation

Measurement

*"Stochastic master equation for a probed system in a cavity", A. E. B. Nielsen et al. 2008*

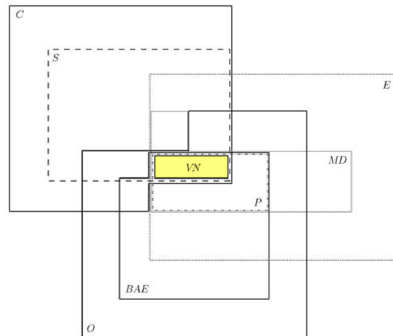
# Time evolution





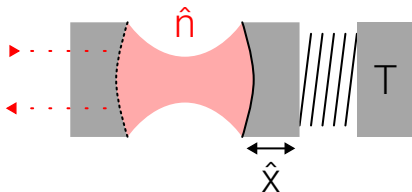
# Types of Measurements

Symbol	Name	Definition
$E$	Efficient	$\forall r, \exists \hat{M}_r, \mathcal{O}_r = \mathcal{J}[\hat{M}_r]$
$C$	Complete	$\forall \rho, \forall r, \mathcal{O}_r \rho \propto \mathcal{O}_r \hat{1}$
$S$	Sharp	$\forall r, \text{rank}(\hat{E}_r) = 1$
$O$	Of an observable $X$	$\forall r, \hat{E}_r = E_r(\hat{X})$
$BAE$	Back-action-evading	$O$ with $\forall \rho, \forall x \in \lambda(\hat{X}), \text{Tr}[\hat{\Pi}_x \rho] = \text{Tr}[\hat{\Pi}_x \mathcal{O} \rho]$
$MD$	Minimally disturbing	$E$ with $\forall r, \hat{M}_r = \hat{M}_r^\dagger$
$P$	Projective	$MD$ and $O$
$VN$	von Neumann	$P$ and $S$



Quantum Measurement and Control, Wiseman et al. 2010

# Problem Statement



*"Cavity optomechanics", Aspelmeyer et al. 2014*

*Quantum Optomechanics, Bowen et al. 2015*



# Hamiltonian

Optical Cavity  $\hat{a}$ ,  $\omega_o(\hat{x}_{\text{mech}}) = \omega_o + \frac{g}{\omega_o} \hat{x}_{\text{mech}}$ ; mechanical oscillations  $\hat{b}$ ,  $\omega_m$ ; coupling  $g$ ; Drive  $E$ ,  $\omega_L$

$$H = \underbrace{\omega_o a^\dagger a}_{\text{Cavity}} + \underbrace{\omega_m b^\dagger b}_{\text{Mechanical}} - \underbrace{g a^\dagger a (b + b^\dagger)}_{\text{Interaction}} + \underbrace{E(ae^{i\omega_L t} + a^\dagger e^{-i\omega_L t})}_{\text{Drive}}$$

*Rotating Wave Approximation* at  $\omega_L$  with  $\Delta = \omega_o - \omega_L$ ,  $a \rightarrow ae^{i\omega_L t}$ :

$$H_{\text{RWA}} = \Delta a^\dagger a + \omega_m b^\dagger b - g a^\dagger a (b^\dagger + b) + E(a + a^\dagger)$$

$\hbar = 1$

*Quantum Optomechanics*, Bowen et al. 2015 (2.3)

*QuantumOptics.jl*, Krämer et al. 2024 (Optomechanical Cavity)

# Hamiltonian Linearization (Currently not used)

$$H_{\text{RWA}} = \Delta a^\dagger a + \omega_m b^\dagger b - g a^\dagger a (b^\dagger + b) + E(a + a^\dagger)$$

Linearize  $a = \alpha + \delta a$ ,  $b = \beta + \delta b$ ; with  $\alpha, \beta$  steady state.

$$\begin{aligned} H_{\text{Interaction}} &= -g a^\dagger a (b^\dagger + b) \\ &\approx -\underbrace{g|\alpha|}_G (\delta a + \delta a^\dagger + \mathcal{O}(a^2 + \delta a \delta a^\dagger)) (\delta b + \delta b^\dagger + 2\beta) \\ a + a^\dagger &= |\alpha| + \delta a + \delta a^\dagger \sim \delta a + \delta a^\dagger \end{aligned}$$

Therefore for small  $G$ :

$$H_{\text{lin}} = \Delta \delta a^\dagger \delta a + \omega_m \delta b^\dagger \delta b - G(\delta a + \delta a^\dagger)(\delta b + \delta b^\dagger) + E(a + a^\dagger)$$

**Problem Statement**  $\sim \frac{\Delta}{2}(\hat{X}^2 + \hat{Y}^2) + \frac{\omega}{2}(\hat{Q}^2 + \hat{P}^2) - G\hat{X}\hat{Q} + E\hat{X}$

# Linearization in Quadratures

$$\begin{aligned}X &= \delta a + \delta a^\dagger & Q &= \delta b + \delta b^\dagger \\Y &= i(\delta a^\dagger - \delta a) & P &= i(\delta b^\dagger - \delta b) \\n &= \delta a^\dagger a & m &= \delta b^\dagger b\end{aligned}$$

$$H_{\text{RWA}} = \Delta n + \omega m - gnQ + EX$$

$$H_{\text{lin}} = \Delta n + \omega m - GXQ + EX$$

The drive  $EX$  is not getting lost in linearization.

There is no point in the simplification if solved numerically.

# Dissipation

Optical decay  $\kappa$ :

$$L = \sqrt{\kappa(n_T + 1)} \delta a + \sqrt{\kappa n_T} \delta a^\dagger$$

Mechanical resonator with  $\gamma$  and a thermal bath at the  $n$ -th thermal state:

$$+ \sqrt{\gamma(m_T + 1)} \delta b + \sqrt{\gamma m_T} \delta b^\dagger$$

*Quantum Optomechanics*, Bowen et al. 2015 (2.8)

# Implementation

truncated Fock Basis:  $F_{\text{optical}} \otimes F_{\text{mechanical}}$

definition of  $H, J$  with  $\delta a \otimes 1$

$$\psi(0) = |0\rangle \otimes |0\rangle$$

Time Evolution using the *Lindblad equation*:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$

# Continuous measurement

Lindblad Master Equation:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\}$$

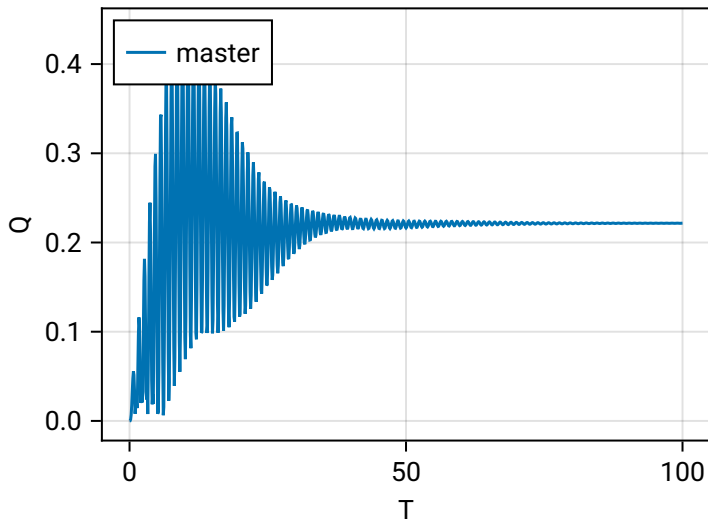
Stochastic Master Equation:

$$\dot{\rho} = -i[H, \rho] + J\rho J^\dagger - \frac{1}{2}\{J^\dagger J, \rho\} + (C\rho + \rho C^\dagger - \text{Tr}(C\rho + \rho C^\dagger)) \xi(t)$$

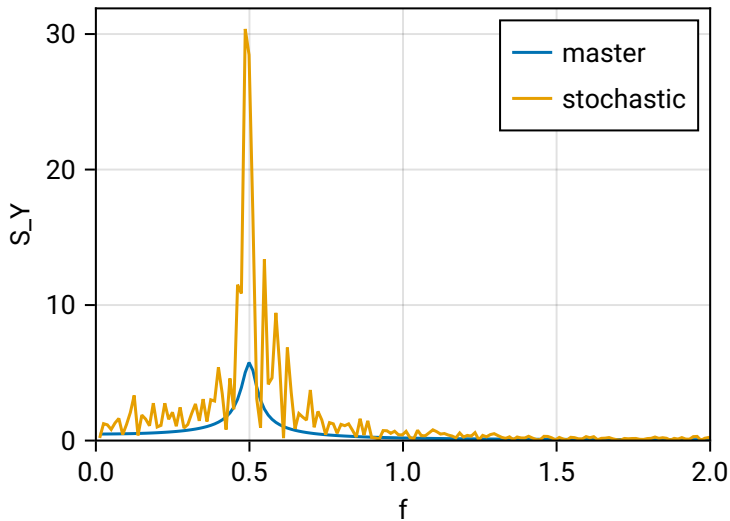
Let's look at the Quadrature  $C = \eta\sqrt{\kappa} (\delta a + \delta a^\dagger)$

*QuantumOptics.jl*, Krämer et al. 2024 (Stochastic Master equation, Quantum Zeno Effect)  
"A Straightforward Introduction to Continuous Quantum Measurement", Jacobs et al. 2006

## Time Evolution

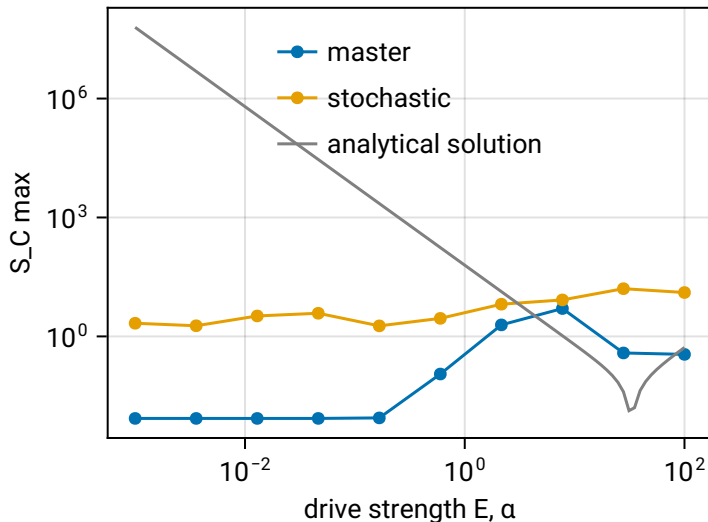


# Spectrum





## Power Dependence $\sim G$



Took  $\approx 1$  min of compute time. Why is the SME so much slower? True Random Values?  
Measurement

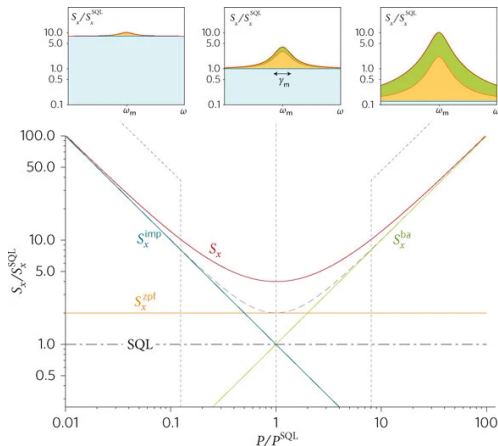
# Expectation

$$S_x = S_x^{\text{imp}} + S_x^{\text{ba}} + S_x^{\text{zpf}}$$

$S_x^{\text{imp}} \propto P^{-1} \omega^0$ : Imprecision /  
Shot noise

$S_x^{\text{ba}} \propto P$ : Back action

$S_x^{\text{zpf}}$ : Zero Point Fluctuation



*"Nanomechanical motion measured with an imprecision below that at the standard quantum limit", Teufel et al. 2009 (Fig. 1)*

# Understanding $S_{\text{det}}(\omega, P_{\text{in}})$

Where is the power dependence?

$$\bar{S}_{\text{det}}(\omega) = \frac{1}{8\eta\Gamma|C_{\text{eff}}|} + 2\Gamma|\chi(\omega)|^2|C_{\text{eff}}|$$

$$C_{\text{eff}}(\omega) = \frac{4g^2}{\kappa\Gamma(1-2i\omega/\kappa)^2}$$

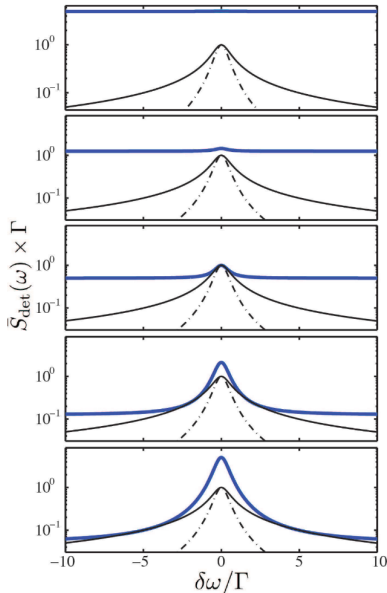
$$\chi(\omega) = \frac{\Omega}{\Omega^2 - \omega^2 - i\omega\gamma}$$

$\eta$ : Detection efficiency

$\Gamma = \gamma$ : Damping of oscillator

$$\bar{S}_{\text{det}}(\omega) = \frac{\kappa\Gamma|1-2i\omega/\kappa|}{8\eta\Gamma 4g^2} + 2\Gamma \frac{\Omega^2}{|\Omega^2 - \omega^2 - i\omega\gamma|^2} \frac{4g^2}{\kappa\Gamma(1-2i\omega/\kappa)^2}$$

Quantum Optomechanics, Bowen et al. 2015 (eq. 3.51)



$$P/P^{\text{SQL}} = \{0.1, 0.4, 1, 4, 10\}$$

Looking for a source that derives  $S_{\text{det}}(\omega, P)$ ?