

Optical signatures of magnetic phase transitions in transition metal phosphorus sulfides

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Summary

Antiferromagnetic van der Waals layered semiconductors such as transition metal phosphorus sulfides are of great interest due to their unique combination of magnetic, optical, and structural properties. The layered structure allows for exfoliation and the study of property changes over the transition from three-dimensional to two-dimensional structures. The exotic thickness-dependent magnetic orders [1] are influenced by the varying magnetic and structural anisotropies for different materials [2, 3]. As the selected semiconductors have bandgaps in the optical and near-infrared, optical measurement techniques offer opportunities for fast in situ measurements of the magnetic structure on a limited sample volume.

In this work, polarization-resolved photoluminescence spectroscopy has been identified as the most promising optical measurement technique for studying the magnetic phase transitions in the selected materials. The magnetic splitting of the photoluminescence line of NiPS₃ has been demonstrated to be more complex than previously reported, with varying splitting, dependent on the linear polarization. In CrPS₄ the circular dichroism of the photoluminescence line was shown to be proportional to the magnetization, which allows the entire magnetic phase diagram to be studied with optical methods.

Zusammenfassung

Antiferromagnetische van-der-Waals Halbleiter wie Übergangsmetallphosphorsulfide sind aufgrund ihrer einzigartigen Kombination von magnetischen, optischen und strukturellen Eigenschaften von großem Interesse. Deren Schichtstruktur ermöglicht die Exfoliation und die Untersuchung von Eigenschaftsänderungen beim Übergang von dreidimensionalen zu zweidimensionalen Strukturen. Die exotischen dickenabhängigen magnetischen Ordnungen [1] werden durch die unterschiedlichen magnetischen und strukturellen Anisotropien für verschiedene Materialien beeinflusst [2, 3]. Da die ausgewählten Halbleiter, Bandlücken im optischen und nahen infrarot Bereich aufweisen, bieten optische Messtechniken die Möglichkeit, die magnetische Struktur in einem begrenzten Probenvolumen schnell in situ zu messen.

In dieser Arbeit wurde die polarisationsaufgelöste Photolumineszenzspektroskopie als die vielversprechendste optische Messtechnik zur Untersuchung der magnetischen Phasenübergänge in den ausgewählten Materialien identifiziert. Die magnetische Aufspaltung der Photolumineszenzlinie von NiPS₃ hat sich als komplexer erwiesen als bisher berichtet, mit einer variierenden Aufspaltung in Abhängigkeit von der linearen Polarisation. Bei CrPS₄ wurde gezeigt, dass der Zirkulardichroismus der Photolumineszenzlinie proportional zur Magnetisierung ist, wodurch das gesamte magnetische Phasendiagramm mit optischen Methoden untersucht werden kann.

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1 Introduction

Antiferromagnetic van der Waals materials such as transition metal phosphorus sulfides are of great interest due to their unique combination of magnetic, optical, and structural properties [5]. The structure of a specific material, NiPS₃, is illustrated in Figure 1.1 [4], demonstrating strong layers in the *a*-*b* plane, bonded by weak van der Waals interactions along the *c* axis. This allows for mechanical exfoliation of the flakes and enables the investigation of property changes over the transition from three-dimensional to two-dimensional structures [5]. The magnetic properties of these materials are dominated by the spin direction of the transition metal ions in the lattice [2]. They exhibit a wide range of thickness-dependent exotic magnetic orders [1], influenced by the varying anisotropy for different materials [2, 3].

However, determining these orders is challenging due to the antiferromagnetic structure, which results in a net magnetization of zero [2, 5, 6]. Especially fast measurements in the monolayer limit without disturbing the sample is difficult [1, 3], but important for studying phase changes and exploring applications in spintronics [1]. Optical measurements emerge as the most promising methods, as they can be used *in situ*, on a small sample volume, and offer spatial resolution. However, the commonly used Kerr effect faces limitations, as detecting small polarization rotations in thin layers, in the order of 5 μ rad has a low signal-to-noise ratio [1].

The goal of this project is to identify optical signatures of magnetic order or magnetic phase transitions in selected transition metal phosphorus sulfides, with the aim of developing a method suitable for fast *in situ* measurements of the magnetic structure in monolayers in the future. This project focuses on identifying and demonstrating a promising measurement technique, rather than on studying the magneto-dependent light-matter interaction process.

better wording?

Various optical measurement techniques were evaluated in the methodology section, with polarization-resolved photoluminescence spectroscopy identified as the most promising method. Subsequently, this method was applied to NiPS₃ and CrPS₄ in the results section.

While the phase transition in NiPS₃ was out of the experimental capabilities, the circular dichroism of the photoluminescence in CrPS₄ was used to observe phase transitions. The results

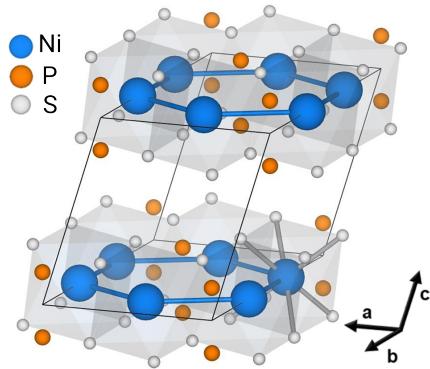


Figure 1.1

Crystal structure of NiPS₃ as an example for the studied materials. Weakly bonded layers in the *a*-*b* plane are held together by van der Waals forces along the *c* axis. The image is adapted from [4].

1.1 Properties of transition metal phosphorus sulfides

also demonstrate the limits of our current theoretical understanding of the photoluminescence process in these materials.

1.1 Properties of transition metal phosphorus sulfides

The utilized transition metal phosphorus sulfides or transition metal thiophosphates for this project are NiPS₃, CrPS₄, MnPS₃, and FePS₃. They were selected for their wide range of magnetic anisotropies [2, 3].

The magnetic properties of transition metal phosphorus sulfides are dominated by the spins of the transition metal ions in the lattice [2]. In the absence of an external magnetic field B and below the Néel temperature T_N , the magnetic order of the spins is antiferromagnetic. In this phase, the spins of the ions align in two or more sublattices, resulting in a net magnetization of zero [6, p.195].

This order can be altered by applying a sufficiently strong external magnetic field H to induce a magnetic phase transition, in which the spins align with the field [3]. Alternatively, the phase can be changed with increasing the temperature, as the Néel temperature is the temperature at which the magnetic order changes from antiferromagnetic to paramagnetic [2, 3].

All studied materials are semiconductors with a bandgap E_B in the visible to near-infrared range. The bandgap is the energy difference between the valence and conduction bands, which determines the energy required to excite an electron from the valence to the conduction band. For a photon to be absorbed by the semiconductor, the photon energy E_{ex} must exceed the bandgap E_B . The free electron in the conduction band can then recombine with a hole in the valence band, emitting a photoluminescence photon with energy E_B . The excess energy $E_{\text{ex}} - E_B$ is dissipated through processes such as phonon emission, which heats the sample, or intra-band photon emission.

Furthermore, more complex photoluminescence processes can occur, such as the formation of neutral excitons X , which are bound states of the electron and hole with a binding Energy E_X . When the exciton recombines a photoluminescence photon with energy $E_B - E_X$ is emitted below the bandgap. The photoluminescence process in the studied materials is reported to be dominated by the exciton recombination [7–9]. Although more complex processes are proposed [4]. The band edges change in response to the alignment of the spins of the metal ions, and therefore depend on the magnetic phase [8]. This change is in the order of 1 meV according to theoretical calculations [8]. The various photoluminescence processes can also influenced by the magnetic order [18].

This work evaluates whether these changes in optical properties can be used to study the magnetic phase transitions in the selected materials.

2 Methodology

2.1 Sample Preparation

The samples of all studied materials were prepared as a first step of the presented work. Initially, measurements were conducted on polycrystalline bulk crystals. However, to create a cleaner surface and increase the probability of obtaining monocrystalline flakes, the crystals were later exfoliated.

Prior to this work, bulk crystals of various transition metal phosphorus sulfides were grown using the vapor transport method, a conventional manufacturing technique for these materials [1, 5].

They are shown as they were at the beginning of this work in [Figure 2.1a](#), with approximate dimensions of $4 \times 2 \times 0.5$ mm. The thickness was quantified using thin-film interference, as detailed in [Section 2.2.2](#).

For the first measurements, the samples were mounted strain-free on a copper holder. They were fixed with a piece of black paper with a hole, that was glued to the copper holder at the corners, as shown in [Figure 2.1b](#). Upon examination under a microscope, the surface of the NiPS₃, CrPS₄, and MnPS₃ samples appeared smooth, as shown in [Figure 2.1c](#), while the FePS₃ sample exhibited a rough surface as depicted in [Figure 2.1d](#). The roughness appeared to result from contamination by an oxidized layer, as it could be removed by mechanical exfoliation.

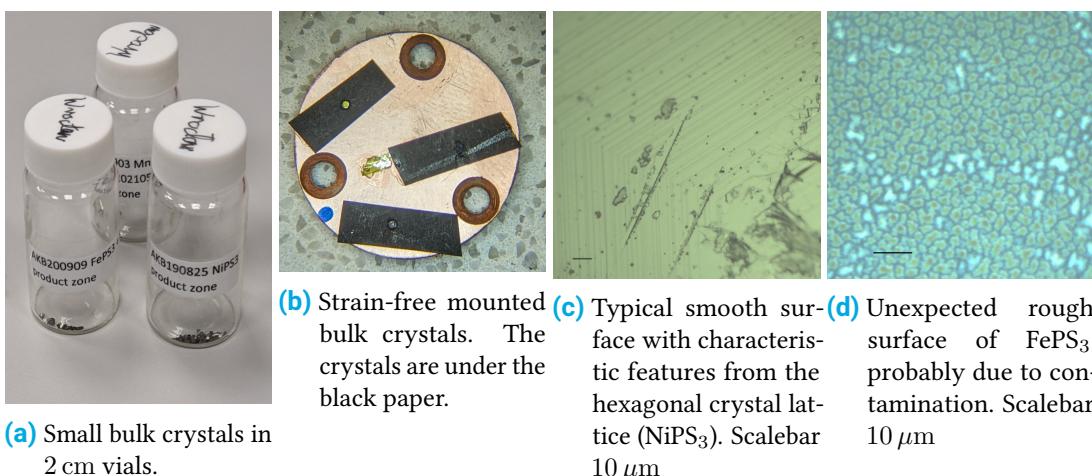


Figure 2.1 Photos of the bulk samples.

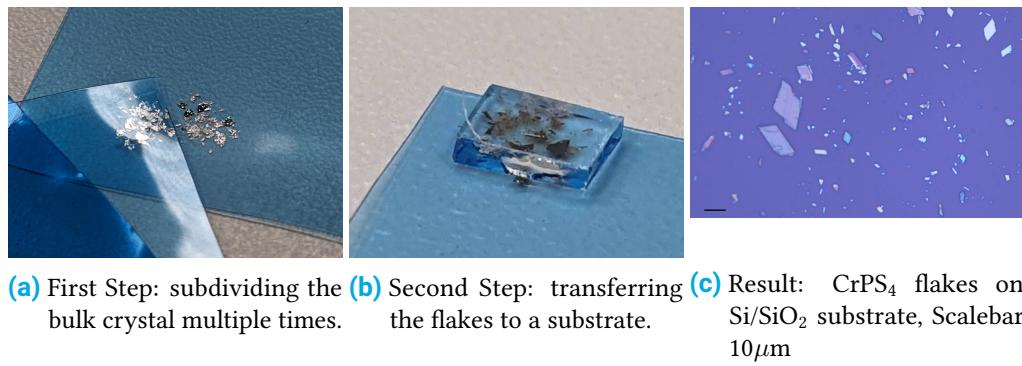


Figure 2.2 Photos at different stages in the exfoliation process.

2.1.1 Exfoliation

To achieve a cleaner surface, especially crucial for the FePS₃ crystals, and to increase the likelihood of obtaining monocrystalline flakes, the bulk crystals were mechanically exfoliated. The stages of this process are photographed in [Figure 2.2](#).

Initially, an as-grown bulk crystal with thickness d_0 is placed on adhesive tape. Subsequently, a second piece of tape is placed on top of the crystal and removed, causing the crystal to split into two flakes that adhere to the two pieces of tape. The flakes already on the tape are divided each time this is repeated into thinner flakes with the average thickness $\langle d_n \rangle = 1/2 \langle d_{n-1} \rangle$. For most samples, the number of repetitions n was 8. This step is depicted in [Figure 2.2a](#). The technique allows to produce monolayers with $\langle d_n \rangle = 0.8 \text{ nm}$ [10] from a $d_0 = 400 \text{ um}$ thick crystals, requiring just $n = \log_2 d_0 / \langle d_n \rangle \approx 19$ steps.

Finally, the flakes are transferred to a substrate. In [Figure 2.2b](#) the flakes are transferred to a glass substrate. In this step, only the thinner flakes adhere to the substrate when the tape is removed, as they are more flexible and can better deform to the substrate. The ratio of adhesion between the flakes and the substrate and the flakes and the tape determines the yield of flakes that remain on the substrate. The utilized adhesive tape, *Nitto Denko ELP BT-150E-CM*, was selected for its low adhesion. The adhesion between the flake and the substrate can be tuned by selecting different substrates. For instance, the yield on glass was 10 to 100 times lower than on a Si/Si₂ substrate. To get even more granularity the transfer was tried at elevated temperatures to reduce the adhesion. To increase the yield even more this transfer step was repeated multiple times for the glass substrate.

The results are flakes on a substrate, as seen under a microscope in [Figure 2.2c](#).

The resulting flakes exhibit a thickness of approximately 500 nm, as measured with atomic force microscopy in Section 3.2.1 ([Figure 3.10](#)). When compared to monolayers, these flakes can be considered bulk material.

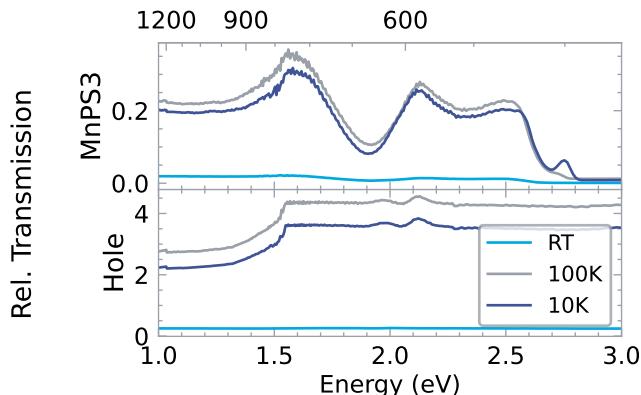


Figure 2.3 Unprocessed relative transmission spectra for NiPS₃ and a calibration Hole. Non-constant relative transmission of the hole with a jump at 1.55 eV.

2.2 Measurement Techniques

To identify a suitable method for investigating the optical properties, various optical spectroscopy techniques were evaluated without applying a magnetic field.

As the expected change in the bandgap due to magnetic field is in the order of 1 meV according to theoretical calculations [8], the techniques should be sensitive to small energy changes.

2.2.1 Transmission Spectroscopy

To begin the exploration of optical spectroscopy techniques, transmission spectroscopy was selected due to its independence from surface quality. Specifically, a commercial double-beam spectrophotometer, the *Cary 5000*, was employed. This spectrometer has two user-accessible beam paths, one for reference and the other for the sample. It uses a single light source and two detectors to calculate the relative transmissivity as $T = I_{\text{Sample}}/I_{\text{Reference}}$ between the two paths. For the initial measurements, MnPS₃ was selected because the 0.5 mm bulk crystals are already visually transparent. The relative transmissivity is shown in Figure 2.3 for a MnPS₃ crystal and a hole in the sample path. For both measurements the reference path was empty.

Surprisingly the relative transmission of the hole is not constant. The only difference between the paths are the Quartz windows of the cryostat in which the sample was mounted and the aperture of the sample holder. The non-uniform transmission spectrum observed can be attributed to a wrong calibration of the spectrometer around the transition from infrared to visible mode, occurring at 800 nm (corresponding to 1.55 eV).

Irrespective of the underlying cause, the measurement was reproducible and a measurement can be corrected by dividing it with a calibration spectrum of the hole and aligning the visible and infrared spectra in software. The corrected absorption spectrum ($\approx 1 - T$) is depicted in Figure 2.4 and reproduces the shape of the absorption spectrum in [11]. The bandgap is at 2.6 eV above which the spectrometer saturates due to insufficient signal through the sample.

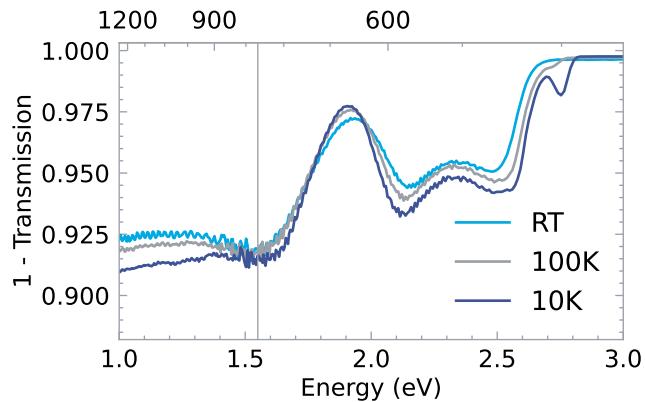
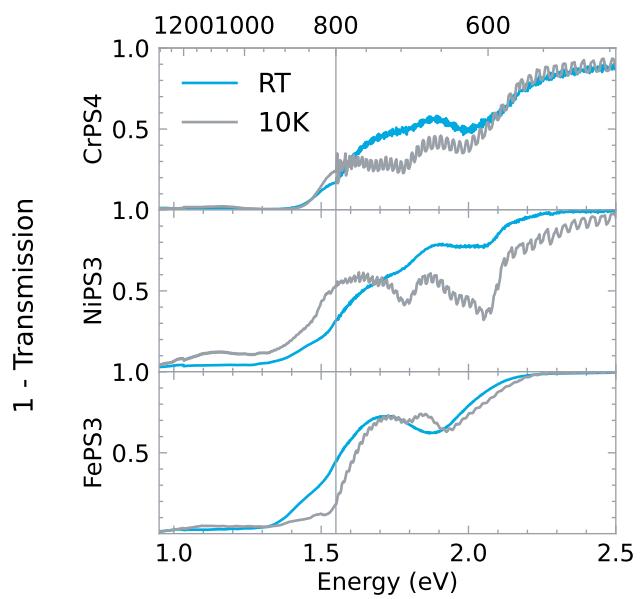


Figure 2.4 Absorption spectrum ($\approx 1 - T$) of MnPS₃ corrected with a calibration spectrum. The bandgap can be seen at 2.6 eV.

For the other samples the features depicted in Figure 2.5a correspond with those recorded in the literature [12–14]. The main problem with this setup is that the spectrometer saturates quickly in the sample or the reference path when the sample is too thin or too thick. It also has no way to individually display the intensity for the sample or reference path which makes adjustments difficult. In order to obtain any usable measurements, the samples were exfoliated to a specific optical transparency, as shown in Figure 2.5b. Additionally, a compensating aperture was placed in the reference beam. To hold the thick exfoliated flakes, they were not transferred to a substrate but mounted with the adhesive tape used for exfoliation. This method was successfully used down to 10 K without any issues.

However, this workflow is rather tedious and there is no way to use a magnet with the spectrometer away. There were plans to rebuild the transmission setup using a cryostat with a built-in super-conducting magnet within the presented project, but this was not realized due to time constraints and unexpected measurements discussed in Section 3.2.1.

Nevertheless, this technique remains promising as it allows for the independent measurement of the bandgap, circumventing luminescent phenomena.



(a) Corrected absorption spectra of different materials. Corresponding with features found in literature [12–14].



(b) Thick exfoliated samples, with specific transparency to not saturate the spectrometer, mounted with the tape used for exfoliation. From left to right: MnPS₃, FePS₃, NiPS₃, CrPS₄

Figure 2.5 Transmission measurements of multiple samples.

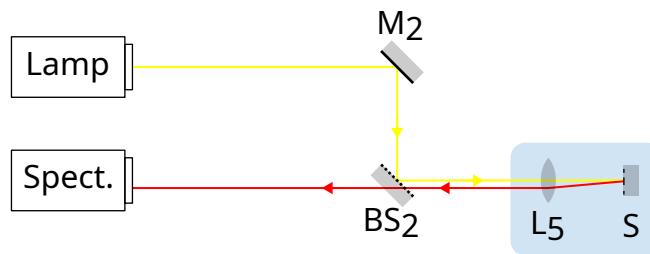


Figure 2.6 Setup used for reflection spectroscopy. The light from a halogen lamp is focused by L_5 on the sample S in the cryostat and the reflected light is collected through the beam-splitter BS_2 by a Spectrometer.

2.2.2 Reflection Spectroscopy

Subsequently, reflection spectroscopy using the setup shown in Figure 2.6 was evaluated. Here a halogen lamp was utilized to illuminate the sample. The light was focused by a lens (L_5) with a high numeric aperture on a small (in the order of 10 μm) spot. The reflected beam passed through a beam-splitter (BS_2) into a spectrometer. This technique only needs access to the sample from a single side, which simplifies the setup.

The unprocessed recorded spectra for bulk crystals of highly transparent MnPS_3 and the material with the most clearly visible bandgap edge (NiPS_3), are represented in Figure 2.7. Below the bandgap, the crystals are transparent, and a thin-film interference pattern is visible. Deconvolving the spectral shape of the lamp by dividing the signal from the sample with the spectrum of a mirror was attempted, but the objective lens L_5 in Figure 2.6 has a strong chromatic aberration, which alters the spectrum too much. In order to remove the lamp spectrum, a filter function f , for example, the *Empirical Mode Decomposition*, can be employed [15]. The result of the filtered signal is presented in Figure 2.8.

However, even with the filtering, the bandgap edge remains difficult to detect precisely, which makes the technique unsuitable for the detection of the expected small shifts in the bandgap.

Thickness Estimation

The thin film interference pattern can be utilized to estimate the thickness of the sample. In order for the total reflection to be strong, the top reflection and the bottom reflection have to be in phase. Specifically $2nd/\lambda = j$ with the count of whole cycle $j \in \mathbb{N}$, the refractive index n and 2 times the thickness d . Therefore the filtered reflected signal is similar to:

$$R(\lambda > \lambda_{\text{Bandgap}}) \approx \cos \left(2\pi \cdot \frac{2nd}{\lambda} \right) \quad (2.1)$$

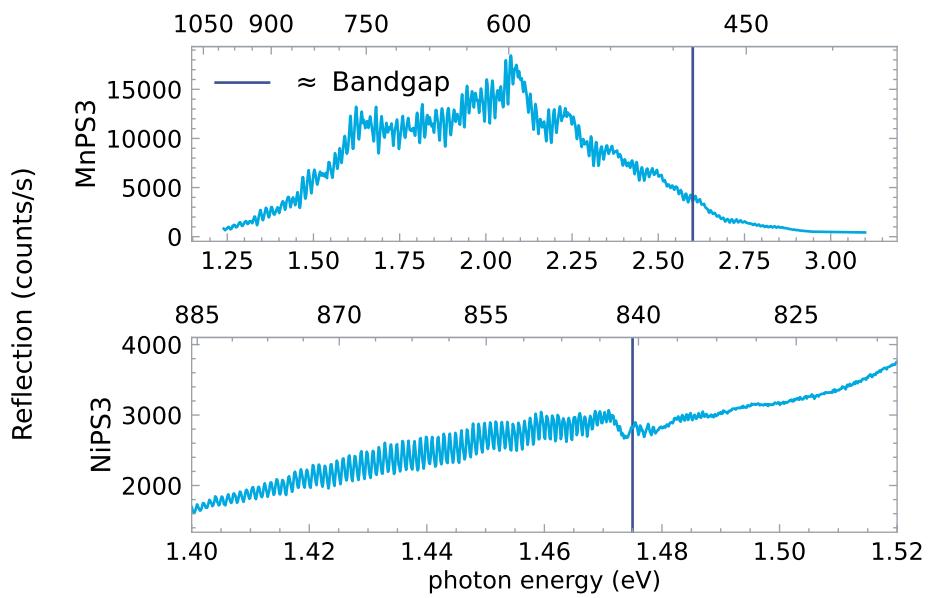


Figure 2.7 Reflection spectra of bulk crystals with broad spectrum from a halogen lamp used for illumination. Thin-film interference below the bandgap deduced from transmission measurements, shown as vertical lines.

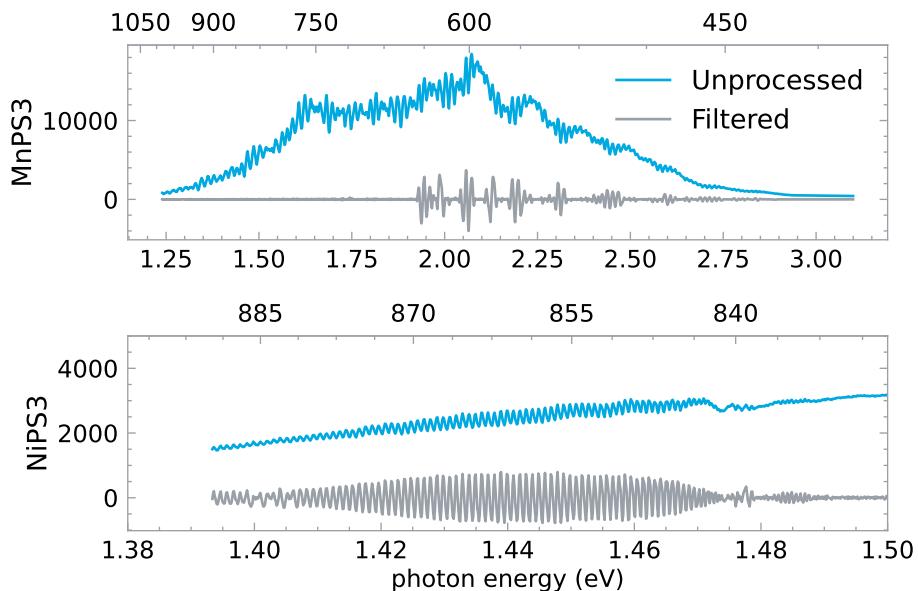


Figure 2.8 Reflection spectra filtered with empirical mode decomposition to eliminate the influence of the light source.

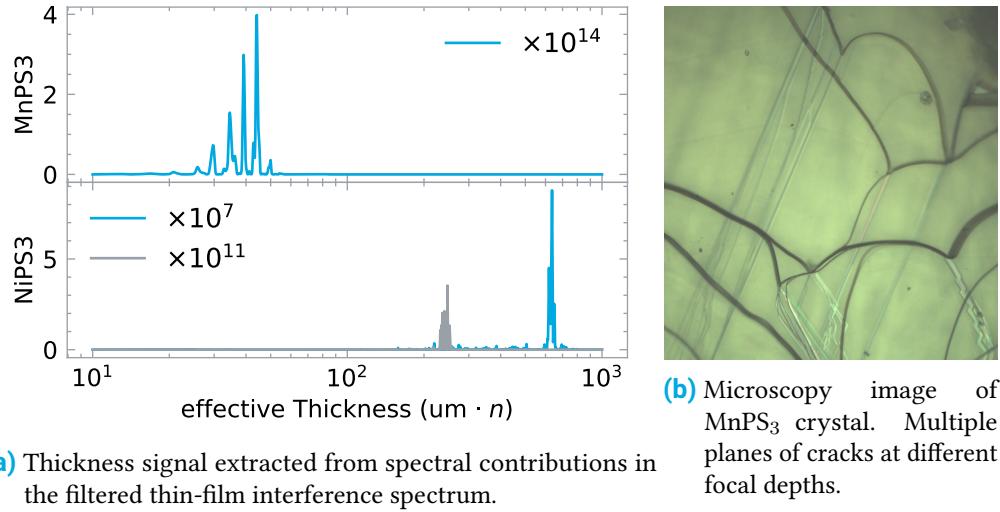


Figure 2.9 Thickness Estimation on thick crystal samples, where multiple breaks are visible in the MnPS₃ crystal.

which can be decomposed into the components for different thicknesses:

$$S(nd) = \mathcal{F} \left[f \left(R \left(\frac{1}{\lambda} \right) \right) \right] (2nd) \quad (2.2)$$

Where f is the filter function used to remove the influence of the light source. The Fourier Transform \mathcal{F} is evaluated using the Lomb-Scargle method [16] because the data is not evenly spaced. Even for evenly spaced data this method is useful as it is less sensitive to intensity noise than the discrete Fourier transform [16].

The thickness signal for different samples is depicted in Figure 2.9a. The NiPS₃ samples each have one effective thickness, in the order of 400 um. In contrast, the MnPS₃ sample displays multiple peaks in the range of 40 um. When observed with an optical microscope (see Figure 2.9b), the MnPS₃ sample has multiple layers of cracks at different depths, indicated by the varying focus. This matches with the recorded thickness signal.

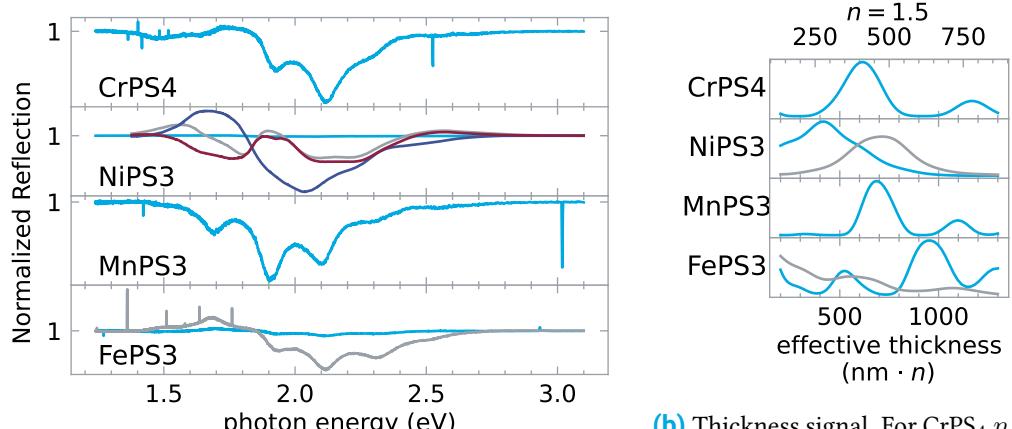
Exfoliated Flakes

The reflection spectra of exfoliated flakes were also measured. It is possible to normalize the spectrum from the flake to the reflection of the bare substrate next to the flake:

$$R_{\text{Normalized}} = \frac{R_{\text{Flake}} - Bkg}{R_{\text{Substrate}} - Bkg} \quad (2.3)$$

Where Bkg is a signal measured without illumination, it consists of an electronic dark noise from the spectrometer and stray light present in the room.

The results are displayed in Figure 2.10a. It is not possible to extract the position of the bandgap



(a) Relative reflectance of exfoliated flakes normalized to substrate. The complex spectrum is due to the thin-film interference of the sample layer and the SiO₂ layer of the substrate.

(b) Thickness signal. For CrPS₄ $n \approx 1.5$ from [17]. The thickness of the CrPS₄ flake is between 200 nm to 700 nm, as confirmed by atomic force microscopy.

Figure 2.10 Reflectance spectra from exfoliated flakes on Si/SiO₂ substrate.

edge in the signal. The hard-to-interpret shape is due to the thin film interference pattern of three interfaces. The first interface is the vacuum-flake, the next is the flake-SiO₂ of the substrate, and the last is the SiO₂-Si interface in the substrate.

This is also visible in the thickness signal in Figure 2.10b. For CrPS₄ the thickness of the flake is found to be between 200 nm to 700 nm by subsequent measurements with atomic force microscopy in Section 3.2.1, Figure 3.10b. With the refractive index of the flake $n \approx 1.5$ [17], the recorded thickness signal is consistent with the atomic force microscopy measurements. Additionally, a second peak at 1180 nm · n is observed, which may be attributed to the SiO₂ layer of the substrate. However, the substrate was derived from the same wafer for all samples, and there is no reason for the thickness of the SiO₂ layer to vary between the different samples.

Concluding, Reflection spectroscopy is not a suitable method for measuring the bandgap edge, even for the thick flakes used here and even less so for monolayers. Therefore, an alternative approach is required.

2.2.3 Photoluminescence Spectroscopy

A setup to measure polarization-sensitive photoluminescence is illustrated in Figure 2.11. The excitation laser at 532 nm or 647 nm is polarization-controlled using the polarizer P₁ and the retarder λ_1 . The mirror M₁ with the beam splitter BS₂ are aligned to minimize the reflection angle, to preserve the polarization as much as possible. The beam splitter BS₂ is a polka dot mirror because it minimizes the changes in polarization compared to a glass wedge. The detection path to the spectrometer is also polarization controlled with the retarder λ_2 and the polarizer P₂.

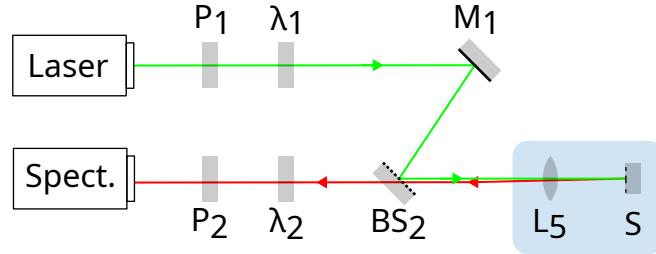


Figure 2.11 Used setup for photoluminescence measurements. Changed from the reflection setup by adding the polarizers $P_{1,2}$ and retarders $\lambda_{1,2}$ to control the polarisation of the excitation and the detection.

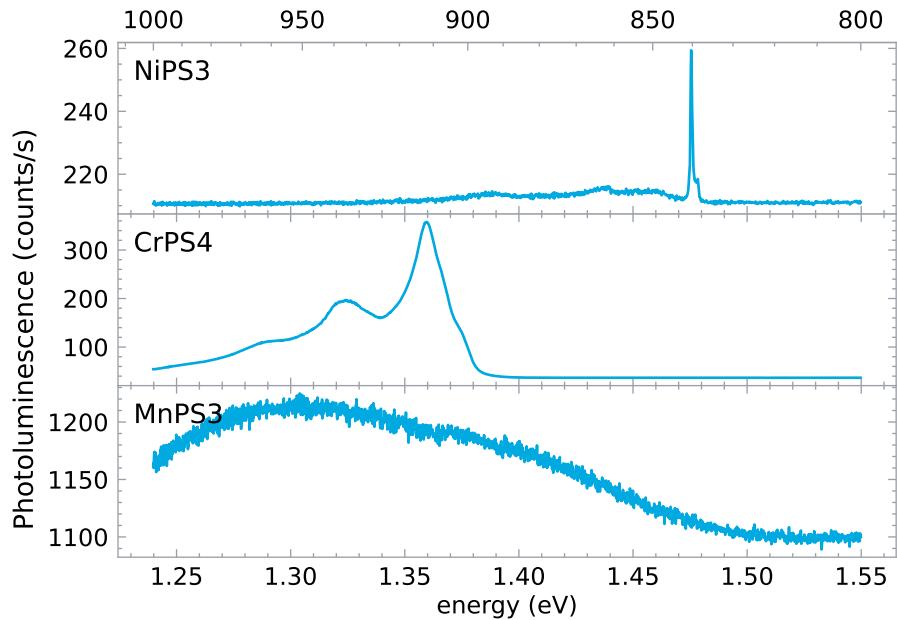


Figure 2.12 Measured Photoluminescence spectra at 10 K. The sharp main peak of NiPS₃ allows for the detection of small changes in the energy.

The retarders were configured to excite with circular polarization and detect with linear polarization or vice versa to prevent any bias in the excitation. While this may not be necessary given that the excitation polarization is lost in the complex luminescence process [18], it is a good practice to ensure that the measurements are not biased.

Figure 2.12 depicts the photoluminescence spectra of various samples at 10 K. A main photoluminescence peak is observed, with different widths for the different materials. This peak can be explained by the neutral exciton photoluminescence process in NiPS₃ and CrPS₄ [7, 19], as it is below the bandgap. Next to the main peak, multiple smaller peaks are visible. These are more

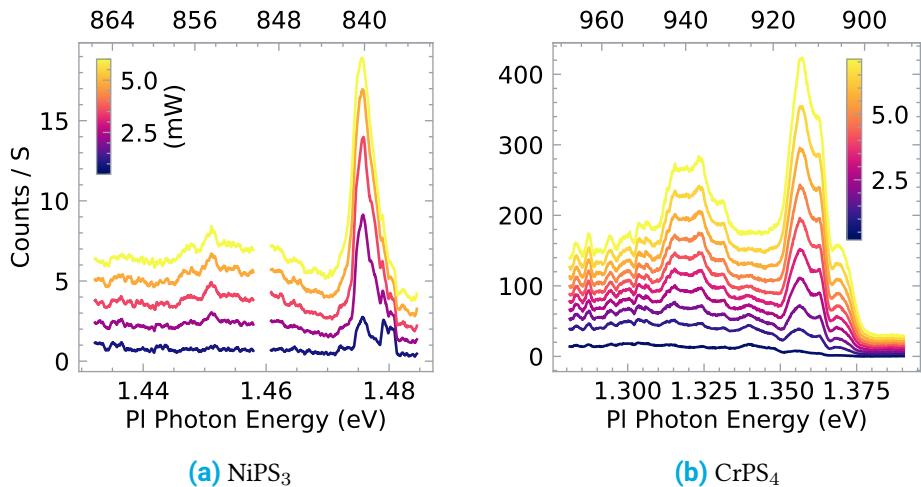


Figure 2.13 Changes of the photoluminescence spectrum for different excitation power at 10 K.

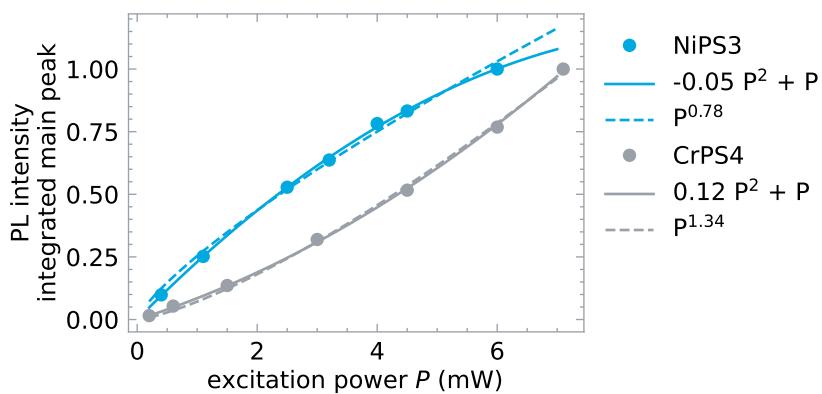


Figure 2.14 Excitation power dependence of the integrated main photoluminescence peak at 10 K with a quadratic and power law fit.

complex structures such as excitons in higher energy levels, charged excitons, or biexcitons [4, 7, 9, 19]. For this project, only the respective main peak is considered.

Despite exfoliation to remove the rough surface, no photoluminescence was detected for FePS₃ which is not consistent with the findings reported in [20].

It is standard practice to use an excitation power in the range of 5 uW to 5 mW to prevent damage to the sample [7, 9]. However, the damage threshold is dependent on the power density, and neither the beam size nor the damage threshold is reported. In order to get an understanding of the saturation of the sample the power scaling was measured. The resulting spectra for different excitation powers are shown in Figure 2.13. No unexpected drastic changes that would indicate damage to the sample are visible.

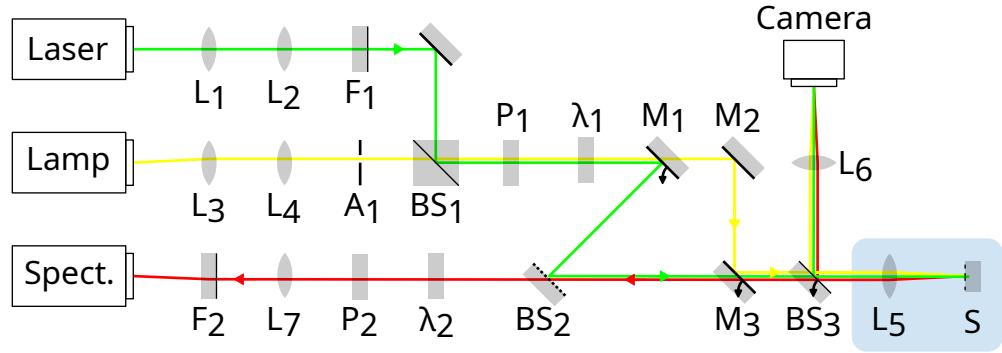


Figure 2.15 Detailed schematic of the entire measurement setup.

To gain a more comprehensive understanding, the intensity of the main peak can be integrated and plotted relative to the excitation power in [Figure 2.14](#). For NiPS₃, a quadratic function fitting indicates a linear scaling that begins to saturate at higher excitation powers. For CrPS₄, the scaling is faster than linear, which can be explained by a scaling probability of creating biexcitons additionally to the exciton photoluminescence process [21].

Since a power law fit is commonly used to describe the scaling the fit is also shown.

In order to achieve a good signal-to-noise ratio while avoiding saturating the sample, the excitation power was selected to be in the range of 1 to 4 mW. Even when illuminated with 20 mW, the sample did not exhibit any signs of damage.

However, following multiple measurements, the photoluminescence ceased in some NiPS₃ flakes. This is not due to the excitation, as it also occurred in flakes that were not excited. It is caused by oxidation during storage and manipulation of the samples. There were no visible changes to those samples on optical microscope images.

2.3 Measurement Setup

The optical table had a few additional components not depicted on the simplified schematics. The full schematic is presented in [Figure 2.15](#).

The sample *S* could be illuminated with a fiber-coupled array of diode lasers to excite the photoluminescence or a fiber-coupled halogen lamp to enable imaging of the sample with a camera, and subsequently a stronger white LED for the same purpose. The lenses *L*_{1–4} were used to expand the cross-section of the illuminating beam in order to utilize the entire aperture of *L*₅.

*F*₁ was an interference film filter used to block out the detection wavelengths from the excitation light. The aperture *A*₁ was mounted off-center to illuminate the sample in a dark-field configuration, thereby increasing the contrast in the camera image. These two illumination paths were combined with cube beam splitters BS₁. The polarizer *P*₁ and the retarder λ ₁ were

used to control the polarization of the light in the excitation path. Then the mirror M_1 and the beam splitter BS_2 were used to introduce the light into the detection path in a way that minimized the alteration of the polarization, by minimizing the reflection angle.

To get more light for the camera, M_1 could be flipped out and M_3 could be flipped in. This obstructs the spectrometer and compromises the polarization of the light, but increases the intensity by a factor of two. BS_3 was just used to image the surface of the sample with the camera. It was a film beam splitter with a large aperture that had to be removed for polarization-sensitive measurements, as it is strongly polarizing.

The objective lens L_5 was mounted on an x-y-z piezo stage in the cryostat to move and focus on different flakes of the samples.

The cryostat and the magnet are combined in a single unit since the 10 T superconducting magnet needs a helium reservoir. The sample was in a variable temperature insert with a controllable helium cooling flow and an electric heater.

The light from the sample then goes through the retarder λ_2 and the polarizer P_2 to control the analyzed polarization. The retarder λ_2 was mounted on a computer-controlled rotation stage. The light then was focused by the lens L_7 into the spectrometer. L_7 had a focal length of 20 cm arbitrarily chosen by the previous user. The interference filter F_2 was used to block out the excitation light from the spectrometer and is placed as close to the spectrometer as possible to block out all the stray excitation light.

The used spectrometer was a 500 nm Czerny-Turner imaging spectrometer (ANDOR SR-500i) with a piezo-cooled CCD sensor and a 1200 lines/mm grating.

3 Results

3.1 NiPS₃

First, NiPS₃ was investigated with polarization-resolved photoluminescence spectroscopy due to the sharp photoluminescence (PL) peak, which facilitates the detection of small changes in the energy. This was done with exfoliated flakes on a Si/SiO₂ substrate for easier handling.

When an external magnetic field H is applied to the sample in the natural a - b plane of the crystal, the main photoluminescence line observed from the plane normal splits into two lines. This is shown in Figure 3.1 for different detection polarizations P . The amount of splitting ΔE is different for different polarization directions. In the polarization direction where the external magnetic field H is aligned with the electric field of the photoluminescence signal $P \parallel H$, the splitting ΔE is larger than in the perpendicular direction $P \perp H$.

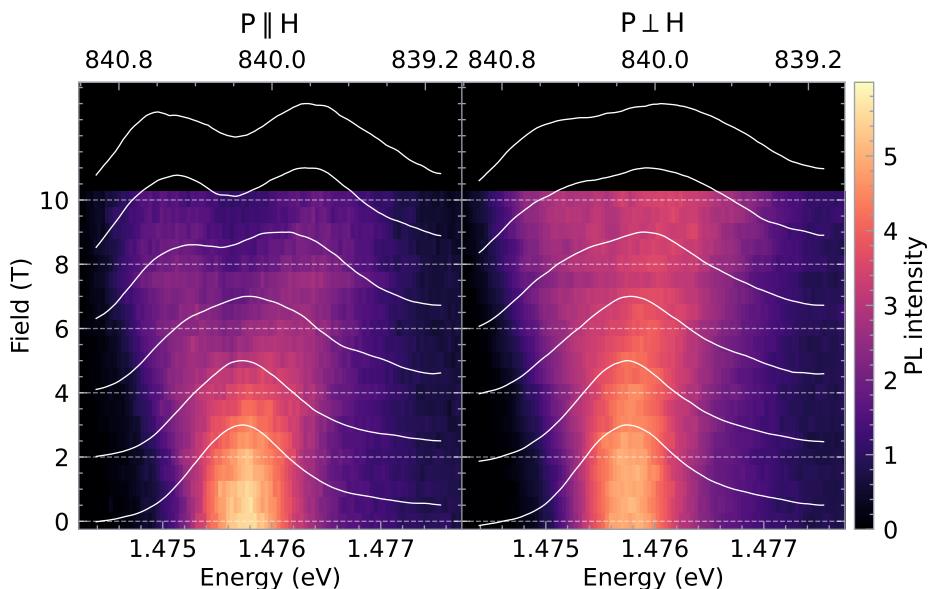


Figure 3.1 Photoluminescence intensity in NiPS₃ for different polarization directions P to applied external magnetic field H at 5 K.
For $P \parallel H$ the splitting ΔE is larger than for $P \perp H$.

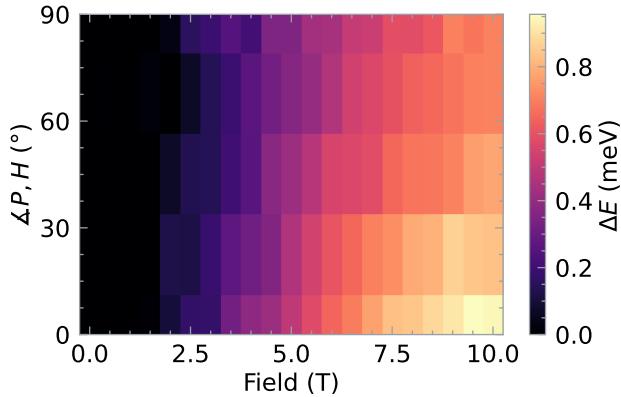


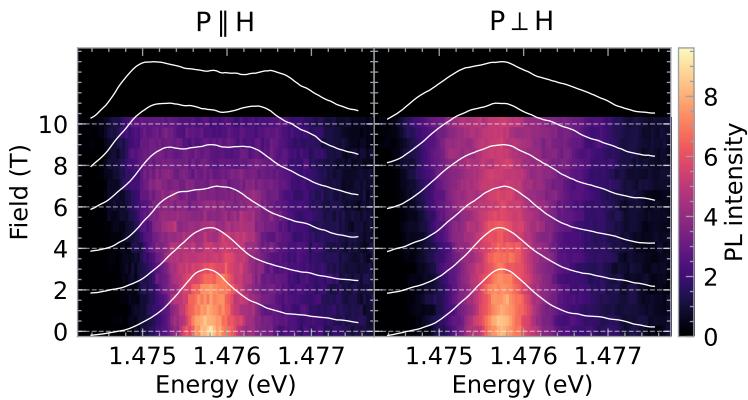
Figure 3.2 ΔE measured with bi-gaussian fit. The splitting is largest for $P \parallel H$ and high magnetic field and decreases with increasing angle and lower field.

To directly determine ΔE , the sum of two Gaussian peaks was fitted, with the peaks centered at $E \pm \Delta E/2$ with the same amplitudes I and widths σ . Given that the linewidth of the photoluminescence is significantly narrower [22] than the resolution of the spectrometer, σ was held constant. E and I were fitted independently for each measurement over field and polarization angle, as the turning retarder and the field-dependent moving sample holder altered the position of the focus on the spectrometer and the focus on the sample, thereby changing E and I respectively. The resulting splitting ΔE is shown over the field and polarisation angle in Figure 3.2.

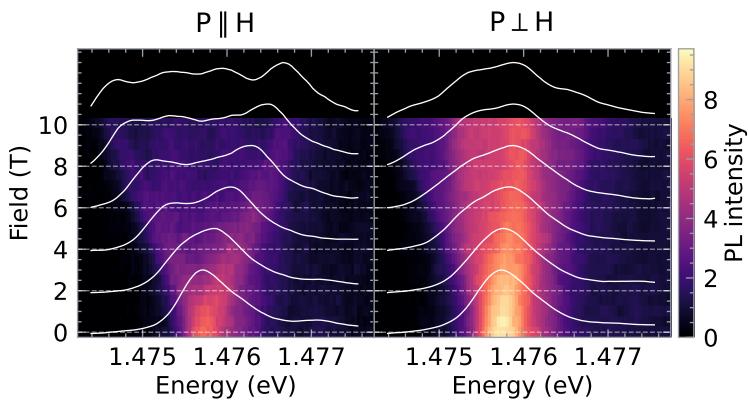
This measurement was repeated for different samples and at different temperatures. Representative results are shown in Figure 3.3. In most flakes, the splitting was not as clear as in Figure 3.1. It was either a broadening of the line as shown in Figure 3.3a or multiple overlapping splitting curves as shown in Figure 3.3b for 50 K.

At higher temperatures, the photoluminescence peak was broader and the splitting was not visible, as shown in Figure 3.3c. Above 50 K the photoluminescence could not be measured over the field, and thus the phase transition to paramagnetic at the Néel temperature of 155 K [2] can not be observed with this method.

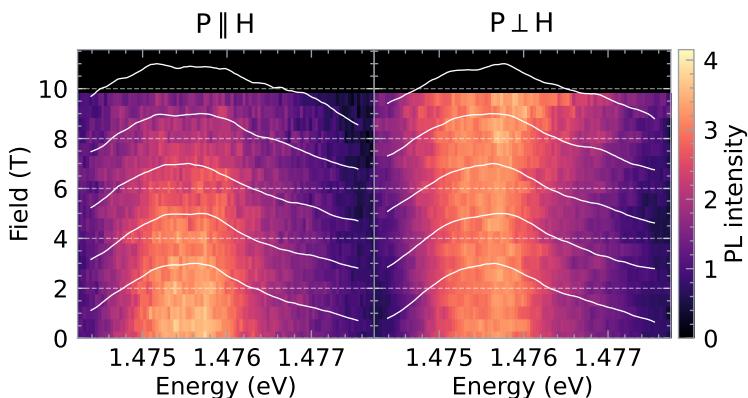
No hysteresis was observed over the magnetic field in any of the measurements and no change in circular polarization was detected.



(a) 10K: Many overlaid splitting curves caused by poly-crystalline flakes.

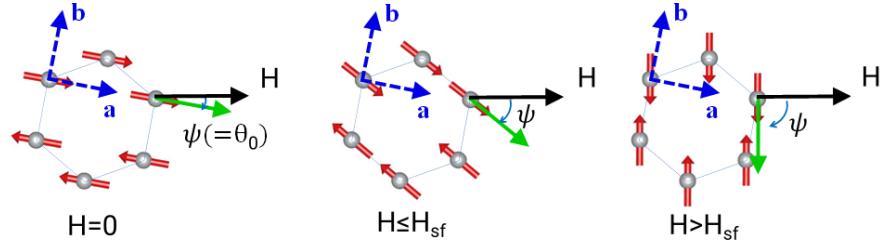


(b) 10K: Two overlapping splitting curves for 2 crystal flake.

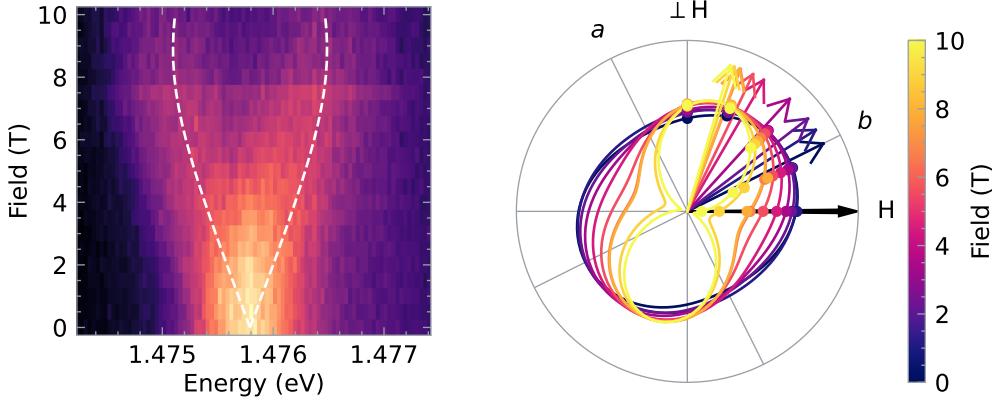


(c) 50K: The line is too broad to observe the splitting.

Figure 3.3 Splitting of the NiPS₃ photoluminescence line at different temperatures and different flakes.



(a) Model for the photoluminescence signal in NiPS₃, with the spins of the Ni²⁺-ions in the *a*-*b* crystal plane. Adapted from [22, Figure S2].



(b) Recorded data for $P \parallel H$ at 5 K, with a fit for ΔE with $\theta_0 = 27^\circ$ according to model Equation 3.2 from [22]. The model fits the data well if only one polarization is selected.

(c) Integrated photoluminescence signal of the main peak for different polarizations. With a \cos^2 fit, to show the polarization axis. The *a* and *b* crystal axis are inferred from the polarisation at $H = 0$ [18]. The polarisation rotates from perpendicular to the *a* crystal axis to perpendicular to the magnetic field.

Figure 3.4 Different models for the photoluminescence signal in NiPS₃.

To understand the observed splitting of the photoluminescence line, a model from [22] can be used, as the splitting was already observed there. A mean-field biaxial antiferromagnet model was proposed, that defines the angle ψ shown in Figure 3.4a between the spin direction of the Ni²⁺-ions and the external magnetic field \vec{H} . For $H = 0$ the spins align with the *a*-crystal axis, and thus, the initial angle $\theta_0 = \psi(H = 0) = \angle(a, H)$ can be defined. The coupling constants along the axis of the model result in an effective *g* factor and a spin-flop field H_{sf} . Consequently, the proposed model [22] yields the following expressions for the angle ψ and the splitting ΔE :

$$\tan 2\psi = \frac{\sin 2\theta_0}{\cos 2\theta_0 - \frac{H}{H_{sf}}^2} \quad (3.1)$$

$$\Delta E = \mu_B g H \cos \psi \quad (3.2)$$

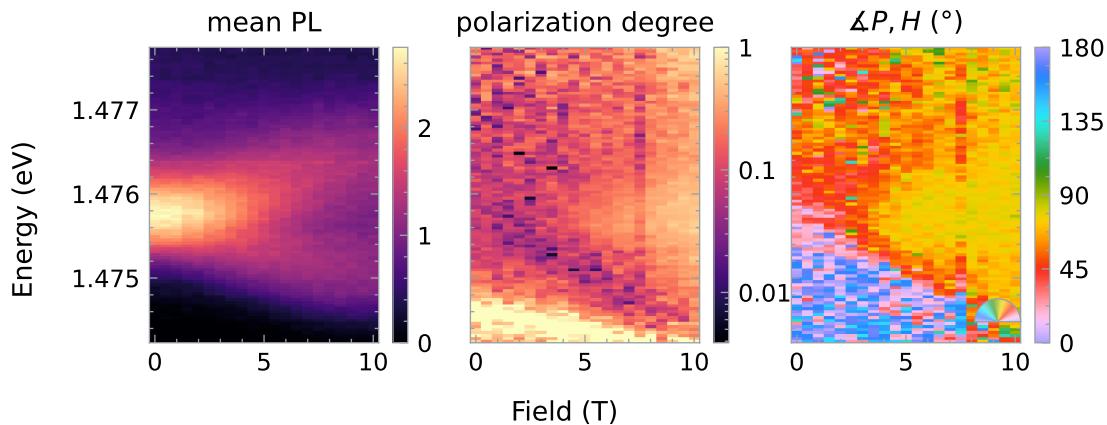


Figure 3.5 Results from the \cos^2 fit for the photoluminescence spectra of NiPS₃ at 5 K. The polarization of the photoluminescence peak does not behave uniformly.

This model does not describe polarization-resolved differences as ΔE is independent of the polarization angle. This is not consistent with the experimental observations, where the splitting is dependent on the polarization angle.

Nevertheless, by taking into account only one polarization, the model describes the data shown in Figure 3.4b. Even the same $H_{sf} = 10.55$ T and $g = 2.0$ from [22] are found to correctly fit the data.

Within this model, the overlaid splitting curves in Figure 3.3b can be understood as polycrystalline flakes with rotated orientations of the *a* and *b* crystal axis and therefore different splitting curves. This was confirmed by manufacturing a sample with more exfoliation steps to get thinner flakes, which yielded clearer splitting curves.

The non-zero linear polarisation degree of the observed photoluminescence is related to the nature of the underlying optical transition. The main photoluminescence peak is caused by the transition from the Zhang-Rice triplet to the Zhang-Rice singlet state [4]. The spin orientations of the free electrons in the Ni *d* orbital and the S *p* orbital is the same for the triplet state and the opposite for the singlet state [4], which results in a change in angular momentum during the transition. The difference in angular momentum ΔS is parallel to the spin alignment axis in the *a*-*b* plane. ΔS is carried away by a circularly polarized photon emitted in the *a*-*b* plane and is not detected [18] in the perpendicular direction. Viewed from the plane normal, the transition is a changing dipole that emits a linearly polarized photon with a polarization perpendicular to the spin alignment [18].

In previous studies [18], the linear polarization of the photoluminescence showed a magnetic-field-dependent rotation of the polarization from perpendicular to the *a* crystal axis to perpendicular to the magnetic field. This is reproduced in Figure 3.4c, which depicts the main peak of the photoluminescence signal measured with different polarisation angles and a \cos^2 fit. The peak intensity was integrated from 1.4732 eV to 1.4768 eV.

However, the peak does not behave uniformly. To examine the polarization direction, the

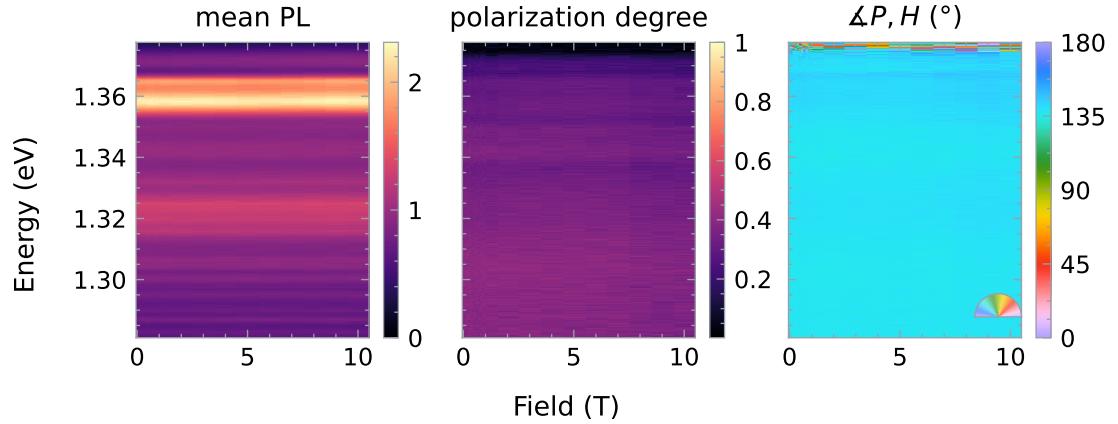


Figure 3.6 Linear polarization of CrPS₄ photoluminescence spectrum at 10 K over external magnetic field in plane. No change in the spectrum or the linear polarization is observed.

data can be displayed by fitting the \cos^2 function for every wavelength and field individually. The average intensity, polarization angle, and degree can then be extracted. This is shown in [Figure 3.5](#).

As previously stated the photoluminescence line splits less when $P \perp H$ than when $P \parallel H$. In the $\angle P, H$ -panel this is evident as a rotation of the polarization in the center of the line from 30 to 90°. The photoluminescence background below the main peak is polarized parallel to the magnetic field, resulting in an angle of 0 or 180°. In between there is a visible line where the angle is 45 degree. This simplifies the presentation of the data, as two images in [Figure 3.1](#) are combined into one.

In order to quantify the degree of polarization, it is necessary to define a reference level. The background level was selected as the photoluminescence signal next to the peak, evaluated as the mean between 1.442 eV to 1.467 eV.

The much lower polarization degree in [Figure 3.5](#) compared to near unity in [9] can be again explained by the polycrystalline nature of the flakes with rotated orientations of the crystal axis within the flakes and therefore different polarization directions.

Altogether, the linear polarization-resolved photoluminescence measurement shows that the luminescence process is more complex than previously understood [4, 22] and more theoretical work is needed to explain it fully.

3.2 CrPS₄

Following studies of NiPS₃, CrPS₄ was examined next, as the photoluminescence behavior of the material under a magnetic field was not yet documented. The photoluminescence signal in CrPS₄ exhibited a notably stronger and broader profile compared to NiPS3. No observable

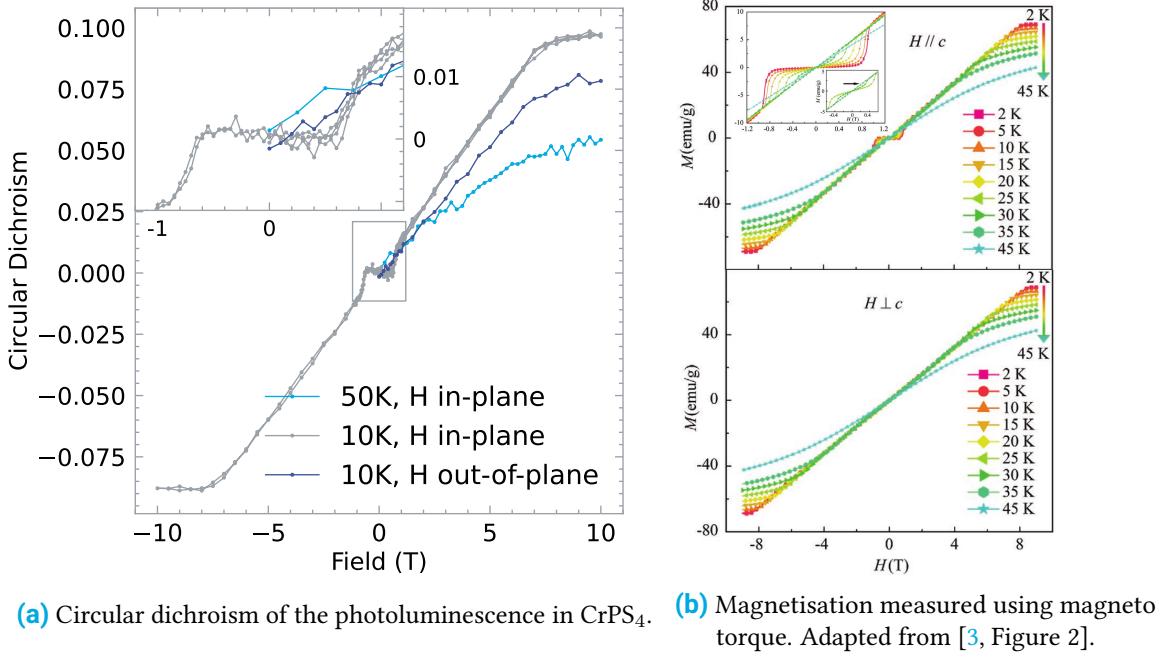


Figure 3.7 Comparison of circular dichroism of the photoluminescence line and the magnetization in CrPS₄. The circular dichroism is proportional to the magnetization.

splitting or displacement of the peak in the photoluminescence signal was detected in response to changes in the magnetic field. Additionally, even though the photoluminescence signal in CrPS₄ was strongly linear polarization (see Figure 3.6), no rotation over the external magnetic field was observed.

The circular dichroism of the photoluminescence in CrPS₄ in Figure 3.7a, had a clear non-hysterical change over the magnetic field.

It is evaluated by integrating the intensity over the main photoluminescence peaks from 1.28 eV to 1.39 eV. This reduces noise and is necessary to measure such small differences with the spectrometer. However, there is no significant difference in the circular dichroism over the wavelength, so the spectrometer is not essential for this measurement. The selection of the background level for calculating the dichroism is straightforward, as there is no significant difference if it is taken sufficiently far away from the photoluminescence peak or without the excitation.

However, the measured circular dichroism signal correlates perfectly with magneto torque measurements of the magnetization (see Figure 3.7b from [3]) in all measured magnetic field configurations and temperatures. However, the corresponding authors have thus far not provided their data for quantitative comparison. But at least from qualitative comparison, the circular dichroism of the photoluminescence line is proportional to the magnetization in CrPS₄. Additional theoretical work is necessary to explain this effect. The main photoluminescence peak is attributed to a simple transition between doublet states in the Cr³⁺-ions in [19] and

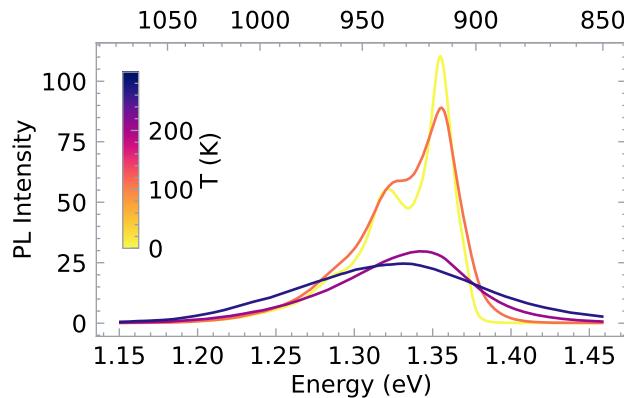


Figure 3.8 Photoluminescence of CrPS₄ for different temperatures. The photoluminescence signal gets broader but is still visible up to room temperature.

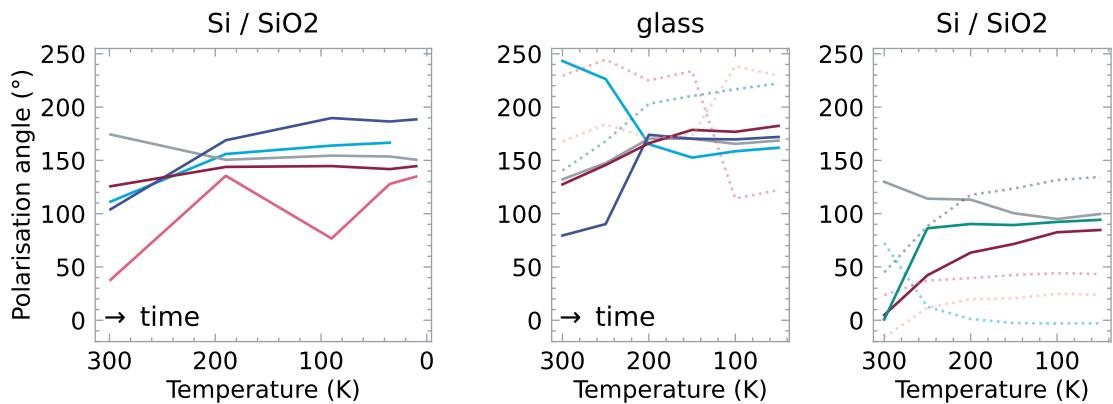
therefore the origin of the circular dichroism of the resulting emission is not clear.

Nevertheless, the observed phenomenon can already be used to optically measure the change in magnetization of the CrPS₄ subject to an external magnetic field. The magnetic phase transitions are visible as distinct changes in the magnetization curve in Figure 3.7a. The spin-flop transition at 0.8 T and the spin-flip phase transition at 7.5 T can be observed, matching the magneto torque measurements [3].

The method is not yet suitable for absolute magnetization measurement, as the circular dichroism to magnetization ratio differs between flakes. This might depend on the thickness of the flakes.

At higher temperatures, the photoluminescence signal broadens and the intensity decreases, as shown in Figure 3.8. Differently to NiPS₃ the photoluminescence signal is visible up to room temperature. This is promising for studying the phase transitions to paramagnetic at the Néel temperature of 38 K [3], which allows for this method to be used to study the entire magnetic phase diagram of CrPS₄.

In conclusion, the circular dichroism of the photoluminescence signal in CrPS₄ allows for fast and localized measurements of the magnetization, that were not possible with previous methods.



(a) First successful measurement.

(b) Partially reproduced measurement on different substrates.

Figure 3.9 Rotation of the linear polarization of the photoluminescence signal in CrPS₄ during cooling. The photoluminescence signal aligns in the same direction on some observed flakes.

3.2.1 Aligning Flakes

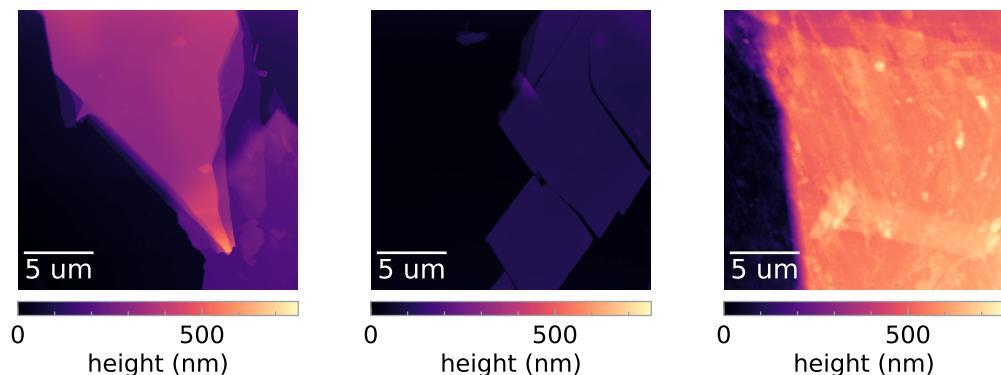
After several days of measurements on the first CrPS₄ sample, the linear polarization of the photoluminescence signal from all the studied flakes was aligned in the same direction. This is unexpected for out-of-plane magnetic field measurements, as there was assumed to be no inherent axis in the sample plane and therefore symmetry for the in-plane rotational degree of freedom.

This was measured more systematically by creating new samples without any alignment in the exfoliation process. For this, the adhesive tapes were rotated to arbitrary angles at each subdivision step and the transfer step was performed multiple times at different angles.

Then the linear polarization of several flakes was recorded during the cooling process, with particular care taken to ensure that no bias for one axis was introduced in the measurement. Images were taken at each temperature in order to determine if the entire flake rotated or just the photoluminescence signal.

The initial measurement presented in [Figure 3.9a](#) confirmed that the photoluminescence of all recorded flakes, which were randomly oriented at room temperature, aligned close to a single orientation at temperatures around 200 K. No rotation was observed in the images of the flakes, indicating that just the photoluminescence signal was rotating.

It was suspected that the Si lattice of the Si/SiO₂ substrate introduces an alignment direction. Therefore, a glass substrate was used, which should not have a preferred axis. The results are presented in [Figure 3.9b](#), along with a new Si/SiO₂ sample to reproduce the aforementioned result. However, in this experiment not all flakes aligned, which indicates a significant variance in the sample preparation. Nevertheless, the alignment directions are different, indicating that the directionality is not originating from the experimental setup. The images of the flakes show no correlation between the alignment and the shape or size of the flakes. Furthermore, the



(a) Normal flake, with a smooth surface. (b) Thin, torn flake, with cracks in the surface. (c) Flake with a rough surface, possibly glue residue.

Figure 3.10 Atomic force microscopy images of CrPS₄ flakes.

alignment does not correlate with the distance to neighboring flakes or the substrate material.

Following the measurements, atomic force microscopy images of the samples were recorded. The flakes had a thickness ranging from 200 nm to 700 nm, as shown in [Figure 3.10](#).

In the case of the sample where all flakes aligned, all flakes had a smooth surface like in [Figure 3.10a](#) and [Figure 3.10b](#).

On the samples where not all flakes aligned, some flakes had a rough surface, as shown in [Figure 3.10c](#). This rough patch extended over the flakes' boundary to the bare substrate. This is likely residue from the adhesive used in the exfoliation process.

Additionally, thinner flakes on the Si/SiO₂ substrate had cracks in the surface as shown in [Figure 3.10b](#). The observed cracks appear to be mechanical stress fractures. The origin of these cracks, whether due to unequal thermal expansion or from the mechanical exfoliation process, remains unclear.

The source of the alignment remains unknown. Further research is needed because the phenomenon can possibly have an unknown influence on similar measurements.

4 Conclusion and Outlook

In this work, several materials belonging to the family of transition metal phosphorus sulfides were studied. First, the samples were prepared from bulk crystals by mechanical exfoliation. This was followed by a thorough examination with several optical measurement techniques in order to identify the most promising method for studying the magnetic properties of these materials. Polarization-resolved photoluminescence spectroscopy was identified as the most promising method.

The findings reveal a higher complexity in the photoluminescence process of the studied materials than was previously reported.

Specifically, the linearly polarized photoluminescence of NiPS₃ exhibits a distinct splitting for different polarization directions, challenging current models [4, 18, 22]. Consequently, the objective now is to find a model for the photoluminescence process that can account for these measurements.

The investigation into CrPS₄ uncovered a correlation between the circular dichroism of the photoluminescence signal and the magnetization. The observed proportionality offers a novel method for measuring the magnetization that is faster and more localized than previous methods. With this method phase transitions have been observed, indicating its potential for studying the entire magnetic phase diagram of CrPS₄. However, this challenges our current understanding of the photoluminescence process in this material, necessitating further research into the underlying mechanisms. As of the time of writing, additional data is being recorded by the group in Warsaw to enhance the reliability of the results.

During these measurements, an unexpected alignment of the linear polarization of the photoluminescence signal was observed in CrPS₄ during cooling. Despite efforts, the cause of this alignment could not be identified within the scope of this work. Understanding the mechanism behind this alignment is crucial, as it may have implications for similar experiments. The proposed next steps in further investigation of this phenomenon include:

- Capturing atomic force microscopy images post-exfoliation to determine whether flakes are already torn.
- Modifying the temperature during the transfer step in exfoliation to investigate potential correlations with the rough surface, thereby confirming the presence of glue residue.
- Washing away the glue to confirm its stabilizing effect on the flakes.
- Experimenting with substrates with diverse thermal expansion coefficients to assess potential influence on the alignment.

However, all these procedures necessitate a statistical argument, demanding numerous samples. Nonetheless, it presents an opportunity for a future student to continue.

Bibliography

- [1] Sharidya Rahman, Juan F. Torres, Ahmed Raza Khan, and Yuerui Lu. "Recent developments in van der waals antiferromagnetic 2D materials: Synthesis, characterization, and device implementation". In: *ACS Nano* 15 (Nov. 2021), pp. 17175–17213. doi: [10.1021/acsnano.1c06864](https://doi.org/10.1021/acsnano.1c06864).
- [2] Pattayil A. Joy and Sridhar Vasudevan. "Magnetism in the layered transition-metal thiophosphates MPS_3 ($M = Mn, Fe$ and Ni)". In: *Physical Review B* 46 (Sept. 1992), pp. 5425–5433. doi: [10.1103/physrevb.46.5425](https://doi.org/10.1103/physrevb.46.5425).
- [3] Yuxuan Peng, Shilei Ding, Man Cheng, Qifeng Hu, Jie Yang, Fanggui Wang, Mingzhu Xue, Zhou Liu, Zhongchong Lin, Maxim Avdeev, and et al. "Magnetic structure and metamagnetic transitions in the van der waals antiferromagnet $CrPS_4$ ". In: *Advanced Materials* 32 (June 2020), 2001200:1–8. doi: [10.1002/adma.202001200](https://doi.org/10.1002/adma.202001200).
- [4] Soonmin Kang, Kangwon Kim, Beom Hyun Kim, Jonghyeon Kim, Kyung Ik Sim, Jae-Ung Lee, Sungmin Lee, Kisoo Park, Seokhwan Yun, Taehun Kim, Abhishek Nag, Andrew Walters, Mirian Garcia-Fernandez, Jiemin Li, Laurent Chapon, Ke-Jin Zhou, Young-Woo Son, Jae Hoon Kim, Hyeonsik Cheong, and Je-Geun Park. "Coherent many-body exciton in van der waals antiferromagnet $NIPS_3$ ". In: *Nature* 583 (July 2020), pp. 785–789. doi: [10.1038/s41586-020-2520-5](https://doi.org/10.1038/s41586-020-2520-5).
- [5] Fengmei Wang, Tofik A. Shifa, Peng Yu, Peng He, Yang Liu, Feng Wang, Zhenxing Wang, Xueying Zhan, Xiaoding Lou, Fan Xia, and Jun He. "New frontiers on van der waals layered metal phosphorous trichalcogenides". In: *Advanced Functional Materials* 28 (July 2018), 1802151:1–24. doi: [10.1002/adfm.201802151](https://doi.org/10.1002/adfm.201802151).
- [6] John M. D. Coey. "Antiferromagnetism and other magnetic order". In: *Magnetism and Magnetic Materials*. Cambridge University Press, 2010, pp. 195–230.
- [7] Ching-Hwa Ho, Tien-Yao Hsu, and Luthviyah Choirotul Muhimmah. "The band-edge excitons observed in few-layer $NIPS_3$ ". In: *npj 2D Materials and Applications* 5 (Jan. 2021), 8:1–10. doi: [10.1038/s41699-020-00188-8](https://doi.org/10.1038/s41699-020-00188-8).
- [8] Miłosz Rybak, Paulo E. Faria Junior, Tomasz Woźniak, Paweł Scharoch, Jaroslav Fabian, and Magdalena Birowska. "Magneto-optical anisotropies of two-dimensional antiferromagnetic MPX_3 from first principles". In: *Physical Review B* 109 (Feb. 2024), 054426:1–10. doi: [10.1103/physrevb.109.054426](https://doi.org/10.1103/physrevb.109.054426).

- [9] Kyle Hwangbo, Qi Zhang, Qianni Jiang, Yong Wang, Jordan Fonseca, Chong Wang, Geoffrey M. Diederich, Daniel R. Gamelin, Di Xiao, Jiun-Haw Chu, and et al. "Highly anisotropic excitons and multiple phonon bound states in a van der waals antiferromagnetic insulator". In: *Nature Nanotechnology* 16 (Mar. 2021), pp. 655–660. doi: [10.1038/s41565-021-00873-9](https://doi.org/10.1038/s41565-021-00873-9).
- [10] Cheng-Tai Kuo, Michael Neumann, Karuppannan Balamurugan, Hyun Ju Park, Soonmin Kang, Hung Wei Shiu, Jin Hyoun Kang, Byung Hee Hong, Moonsup Han, Tae Won Noh, and et al. "Exfoliation and Raman spectroscopic fingerprint of few-layer NIPS₃ van der waals crystals". In: *Scientific Reports* 6 (Feb. 2016), 20904:1–10. doi: [10.1038/srep20904](https://doi.org/10.1038/srep20904).
- [11] Vincenzo Grasso, Fortunato Neri, Pedro Perillo, Letteria Silipigni, and Mauro Piacentini. "Optical-absorption spectra of crystal-field transitions in MnPS₃ at low temperature". In: *Physical Review B* 44 (Nov. 1991), pp. 11060–11066. doi: [10.1103/physrevb.44.11060](https://doi.org/10.1103/physrevb.44.11060).
- [12] Resta A. Susilo, Bo Gyu Jang, Jiajia Feng, Qianheng Du, Zhipeng Yan, Hongliang Dong, Mingzhi Yuan, Cedomir Petrovic, Ji Hoon Shim, Duck Young Kim, and Bin Chen. "Band gap crossover and insulator-metal transition in the compressed layered CrPS₄". In: *npj Quantum Materials* 5 (Aug. 2020), 58:1–8. doi: [10.1038/s41535-020-00261-x](https://doi.org/10.1038/s41535-020-00261-x).
- [13] Eldad J Banda. "Optical absorption of NiPS₃ in the near-infrared, visible and near-ultraviolet regions". In: *Journal of Physics C: Solid State Physics* 19 (Dec. 1986), pp. 7329–7335. doi: [10.1088/0022-3719/19/36/023](https://doi.org/10.1088/0022-3719/19/36/023).
- [14] Adam K. Budniak, Szymon J. Zelewski, Magdalena Birowska, Tomasz Woźniak, Tatyana Bendikov, Yaron Kauffmann, Yaron Amouyal, Robert Kudrawiec, and Efrat Lifshitz. "Spectroscopy and structural investigation of iron phosphorus trisulfide—FePS₃". In: *Advanced Optical Materials* 10 (Feb. 2022), 2102489:1–10. doi: [10.1002/adom.202102489](https://doi.org/10.1002/adom.202102489).
- [15] Jiaxing Sun, Xinyu Li, Haojie Zhang, Jinlong Song, and Zhisong Li. "A method for measuring and calibrating the thickness of thin films based on infrared interference technology". In: *Results in Physics* 51 (Aug. 2023), 106727:1–12. doi: [10.1016/j.rinp.2023.106727](https://doi.org/10.1016/j.rinp.2023.106727).
- [16] Jeffrey D. Scargle. "Studies in astronomical time series analysis. II - statistical aspects of spectral analysis of unevenly spaced data". In: *The Astrophysical Journal* 263 (Dec. 1982), pp. 835–853. doi: [10.1086/160554](https://doi.org/10.1086/160554).
- [17] Alexandria R. Alcantara, Christopher Lane, Jason T. Haraldsen, and Roxanne M. Tutchton Tutchton. "Parameter-free treatment of a layered correlated van der waals magnet: CrPS₄". In: *Physical Review B* 108 (Oct. 2023), 155133:1–8. doi: [10.1103/physrevb.108.155133](https://doi.org/10.1103/physrevb.108.155133).
- [18] Xingzhi Wang, Jun Cao, Zhengguang Lu, Arielle Cohen, Hikari Kitadai, Tianshu Li, Qishuo Tan, Matthew Wilson, Chun Hung Lui, Dmitry Smirnov, Sahar Sharifzadeh, and Xi Ling. "Spin-induced linear polarization of photoluminescence in antiferromagnetic van der waals crystals". In: *Nature Materials* 20 (Apr. 2021), pp. 964–970. doi: [10.1038/s41563-021-00968-7](https://doi.org/10.1038/s41563-021-00968-7).

- [19] Suhyeon Kim, Sangho Yoon, Hyobin Ahn, Gangtae Jin, Hyesun Kim, Moon-Ho Jo, Changgu Lee, Jonghwan Kim, and Sunmin Ryu. “Photoluminescence path bifurcations by spin flip in two-dimensional CrPS₄”. In: *ACS Nano* 16 (Sept. 2022), pp. 16385–16393. doi: [10.1021/acsnano.2c05600](https://doi.org/10.1021/acsnano.2c05600).
- [20] Ellenor Geraffy, Shahar Zuri, Milosz Rybak, Faris Horani, Adam Budniak, Yaron Amouyal, Magdalena Birowska, and Efrat Lifshitz. “Crystal anisotropy implications on the intrinsic magnetic and optical properties in van der waals FePS₃”. In: *Research Square* (Sept. 2022). preprint (version 1), pp. 1–20. doi: [10.21203/rs.3.rs-2014052/v1](https://doi.org/10.21203/rs.3.rs-2014052/v1).
- [21] Yumeng You, Xiao-Xiao Zhang, Timothy C. Berkelbach, Mark S. Hybertsen, David R. Reichman, and Tony F. Heinz. “Observation of biexcitons in Monolayer WSe₂”. In: *Nature Physics* 11 (May 2015), pp. 477–481. doi: [10.1038/nphys3324](https://doi.org/10.1038/nphys3324).
- [22] Dipankar Jana, Piotr Kapuscinski, Ivan Mohelsky, Diana Vaclavkova, Ivan Breslavetz, Milan Orlita, Clément Faugeras, and Marek Potemski. “Magnon Gap excitations and spin-entangled optical transition in the van der waals antiferromagnet NiPS₃”. In: *Physical Review B* 108 (Sept. 2023), 115149:1–8. doi: [10.1103/physrevb.108.115149](https://doi.org/10.1103/physrevb.108.115149).
- [23] Yongheng Zhou, Kaiyue He, Huamin Hu, Gang Ouyang, Chao Zhu, Wei Wang, Sichen Qin, Ye Tao, Runfeng Chen, Le Zhang, Run Shi, Chun Cheng, Han Wang, Yanjun Liu, Zheng Liu, Taihong Wang, Wei Huang, Lin Wang, and Xiaolong Chen. “Strong neel ordering and luminescence correlation in a two-dimensional antiferromagnet”. In: *Laser & Photonics Reviews* 16 (June 2022), 2100431:1–10. doi: [10.1002/lpor.202100431](https://doi.org/10.1002/lpor.202100431).

Data Availability

All the recorded data and software used for the analysis and documentation are available publicly at https://github.com/leole100/bachelorarbeit_public.

Eigenständigkeitserklärung

Ich versichere hiermit, dass ich die anliegende Abschlussarbeit mit dem Thema: *Optical signatures of magnetic phase transitions in transition metal phosphorus sulfides* selbständig verfasst und keine anderen Hilfsmittel und Quellen als die angegebenen benutzt habe.

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Weiterhin versichere ich hiermit, dass die o.g. Arbeit noch nicht anderweitig als Abschlussarbeit einer Diplom-, Bachelor- bzw. Masterprüfung eingereicht wurde. Mir ist ferner bekannt, dass ich bis zum Abschluss des Prüfungsverfahrens die Materialien verfügbar zu halten habe, welche die eigenständige Abfassung der Arbeit belegen können.

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Als Urheber der anliegenden Arbeit stimme ich dem obigen Veröffentlichungsverfahren zu.

Konstanz, den 2. Mai 2024