

Evanescent light scattering Optical Tweezers

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Abstract auf Englisch (10-15 Zeilen) Hello, here is some text without a meaning. This text should show what a printed text will look like at this place. If you read this text, you will get no information. Really? Is there no information? Is there a difference between this text and some nonsense like “Huardest gefburn”? Kjift – not at all! A blind text like this gives you information about the selected font, how the letters are written and an impression of the look. This text should contain all letters of the alphabet and it should be written in of the original language. There is no need for special content, but the length of words should match the language. Hello, here is some text without a meaning. This text should show what a printed text will look like at this place. If you read this text, you will get no information. Really? Is there no information? Is there a difference between this text and some nonsense like “Huardest gefburn”? Kjift – not at all! A blind text like this gives you information about the selected font, how the letters are written and an impression of the look. This text should contain all letters of the alphabet and it should be written in of the original language. There is no need for special content, but the length of words should match the language.

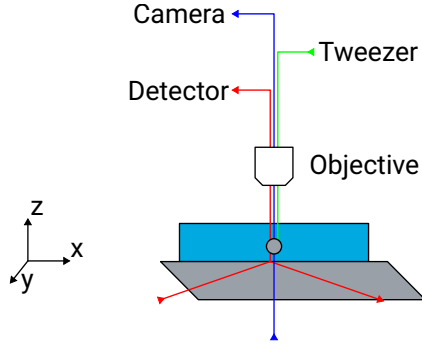


Figure 1 Schematic of the experimental setup. Optical tweezers in green, transmission light microscopy in blue and total internal reflection microscopy in red.

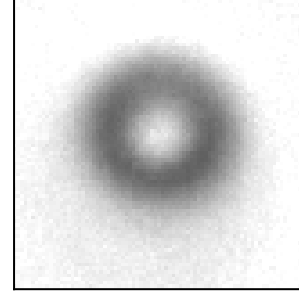


Figure 2 Transmission light microscopy of an observed particle. The radius of the particle is $14(2)$ px, equivalent to $1.86(27)$ μm .

1 Introduction

2 Methods

This experiment operates with light on small particles (in the range of $1 \mu\text{m}$ [1]) in an aqueous solution. This is shown in Figure 1 as the gray circle in the blue box.

Using an optical tweezers setup the particle can be trapped in the center of the experiment. This results in a 3 dimensional harmonic potential, with a spring constant that is roughly proportional to the trap strength.

To directly observe the particle transmission light microscopy is used. For this a blue led is used to illuminate the particle from below. This allows the particle position to be measured in x and y direction.

To measure the position in z direction, total internal reflection microscopy is used. For this a red laser is bounced off a prism to create an evanescent wave in the aqueous solution. This wave scatters off the particle into a sensitive photo detector.

3 Results

All recorded data and the analysis is available at www.github.com/leoolle100/fp2.

3.1 Transmission Light Microscopy

In the first section of the experiment, the particles were observed using a transmission light microscope setup. For this images with a resolution of 600×800 px were recorded with a frequency of

10 Hz for 10 min. The magnification of the microscope was assumed to be $0.13319672 \mu\text{m}/\text{px}$ [1], this is the main systematic error of this measurement procedure.

The images were normalized with a black (illumination off) and white (illumination on, particle not in frame) reference image, to remove the influence of dust in the imaging elements. The particle that was used for this experiment is shown in Figure 2, after the normalization.

To determine the trajectory of the particle, an effective center of mass was calculated for each frame:

$$\vec{r}(t) = \iint \vec{r} \cdot (1 - I(\vec{r}, t))^2 d\vec{r} \quad (1)$$

The density of the resulting trajectory is shown in Figure 3 for different optical trap stiffnesses. The approximate radius of the optical trap of $2.5 \mu\text{m}$ is drawn as a red circle. For the trap stiffness of 0, the particle is free to wander around, for the higher stiffnesses like 1.01 the particle is mostly confined to the trap.

For a weak trap like 0.75 the particle is still mostly confined to the trap, but can escape the linear trap region and randomly wander around. This happened multiple times during the shown measurement.

Mean Square Displacement

A method to describe the trajectory of a random walk is the mean square displacement (MSD) [2].

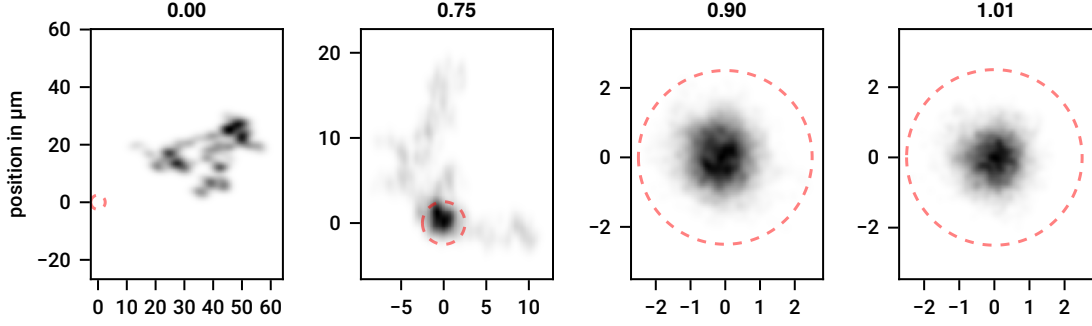


Figure 3 Density of recorded particle positions, grouped by optical trap stiffness. The red circle indicates the approximate radius of the optical trap $2.5 \mu\text{m}$.

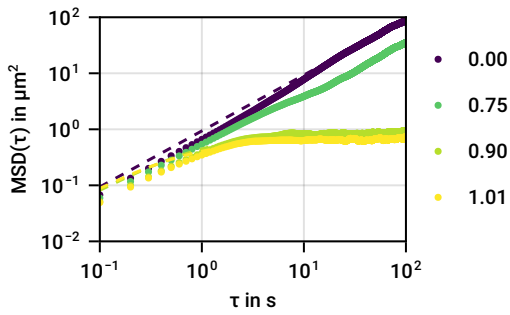


Figure 4 Mean Square Displacement for different optical trap stiffnesses, fit: Equation 3

This is defined as the average of the squared distance of the particle from the starting point:

$$\text{MSD}(t) = \langle \Delta r^2(t) \rangle = \frac{1}{N} \sum_i (x_i(t) - x_i(0))^2 \quad (2)$$

Here a different implementation using the autocorrelation of the velocity was used, to achieve a more stable result. This was implemented by [3]. The resulting MSD for different optical trap stiffnesses is shown in Figure 4.

The MSD can be described by the following model:

$$\text{MSD}(\tau) = \frac{1}{\frac{1}{D_0 \tau} + \frac{1}{\text{MSD}(\infty)}} \quad (3)$$

For a free particle the $\text{MSD}(\infty) = \infty$ and the MSD grows linearly with D_0 over time [1, 2].

The linear fit for the free particles, works as described with a $D_0 = 0.9357(25) \mu\text{m}^2/\text{s}$. This is roughly equivalent with the expected value of $D_0 = \frac{k_B T}{6\pi\eta r} = 1.22(18) \mu\text{m}^2/\text{s}$ [1], with $r = 1.86(27) \mu\text{m}$ and $\eta = 0.955 \text{ mPa} \cdot \text{s}$.

For a confined particle the MSD reaches a plateau at $\text{MSD}(\infty)$ [1]. For the higher spring stiffnesses (0.90, 1.01), the MSD reaches a plateau and the spring stiffness are shown in Figure 6.

For the lower measured spring stiffness (0.75), the MSD does not reach a plateau, as it partially escapes the trap and wanders around (see Figure 3). Therefore this procedure is not adequate for such low spring stiffnesses.

Potential

A more detailed analysis of the Potential can be done by looking at the distribution of the particle positions. For this the probability density function (PDF) of the particle positions has to be estimated. For this the kernel density estimation (KDE) with a Gaussian kernel was used [4]. The resulting 2 dimensional PDF is shown in Figure 3.

As the density is small for large parts of the explored 2d space, the data is reduced by aggregating into the image coordinates x and y . As the Potential should be rotationally symmetric, the data from the x and y axis should not significantly deviate.

Using the maxwell boltzmann relations [1], the potential at a position p can be calculated from the PDF:

$$V(p) = -\frac{\log \text{PDF}(p)}{k_B T} \quad (4)$$

The resulting potential is shown in Figure 5. The part of the measurements with a $\text{PDF}(x) > 0.05$ are used for a quadratic fit, to estimate the spring constant. The resulting spring constants are shown in Figure 6 and are in good agreement with the previous estimates.

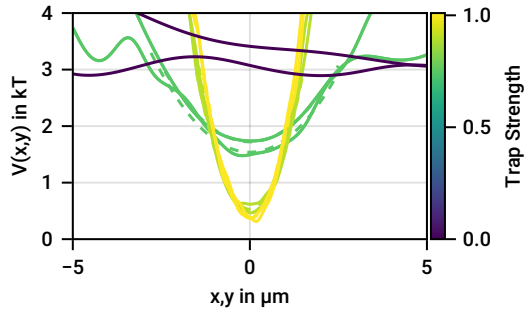


Figure 5 Measured Potential, with quadratic fit.

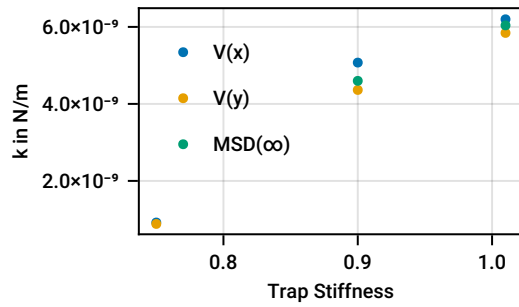


Figure 6 Differently measured spring constants.

For the lower spring constant of 0.75, the potential begins quadratic, but flattens after approximately $2.5 \mu\text{m}$. This is due to the limited size of the trap, but the spring constant can still be estimated.

This method allows for the measurement of a spring constant as low as $0.9011(45) \text{ nN/m}$, which means over the used length scales that forces in the order of 1 fN can be measured.

3.2 Total Internal Reflection Microscopy

4 Discussion

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References

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