

### III. HIGHER-ORDER ELECTRON-CORRELATION CORRECTIONS TO THE TRANSITION ENERGIES

Electron-electron interaction within the basic principles of QED is described by exchange of virtual photons. The one-photon exchange leads to the operator

$$I(\omega) = e^2 \alpha_1^\mu \alpha_2^\nu D_{\mu\nu}(\omega, \mathbf{r}_{12}), \quad (8)$$

where  $D_{\mu\nu}$  is the photon propagator, which in the Coulomb gauge is written as

$$\begin{aligned} D_{00}(\omega, \mathbf{r}_{12}) &= \frac{1}{4\pi r_{12}}, \quad D_{i0} = D_{0i} = 0 \quad (i = 1, 2, 3), \\ D_{il}(\omega, \mathbf{r}_{12}) &= \int \frac{d\mathbf{k}}{(2\pi)^3} \frac{\exp(i\mathbf{k} \cdot \mathbf{r}_{12})}{\omega^2 - \mathbf{k}^2 + i0} \left( \delta_{il} - \frac{k_i k_l}{\mathbf{k}^2} \right) \quad (i, l = 1, 2, 3), \end{aligned} \quad (9)$$

$r_{12} = |\mathbf{r}_{12}| = |\mathbf{r}_1 - \mathbf{r}_2|$ ,  $\mathbf{r}_i$  is the position vector of the  $i$ th electron, and  $\alpha^\mu = (1, \boldsymbol{\alpha})$  are the Dirac matrices.

Expanding expression (9) in powers of the photon frequency one can derive a simplified form of the interaction. The low-frequency limit of this interaction consists of two parts, referred to as the Coulomb and the Breit interaction,

$$V(i, j) = V_C(i, j) + V_B(i, j) = \frac{\alpha}{r_{ij}} - \alpha \left[ \frac{\boldsymbol{\alpha}_i \cdot \boldsymbol{\alpha}_j}{2r_{ij}} + \frac{(\boldsymbol{\alpha}_i \cdot \mathbf{r}_{ij})(\boldsymbol{\alpha}_j \cdot \mathbf{r}_{ij})}{2r_{ij}^3} \right]. \quad (10)$$

The most traditional approach for the treatment of the electron-electron interaction in relativistic many-electron atoms consists in using so-called Breit approximation. In this approximation the total Hamiltonian can be represented as the sum of the one-electron Dirac Hamiltonians and the Coulomb and Breit electron-electron interactions, projected on the positive-energy Dirac's states. In this way one gets the Dirac-Coulomb-Breit equation. Traditional methods for solving the Dirac-Coulomb-Breit equation are the many-body perturbation theory (MBPT) [48, 49], the multi-configuration Dirac-Fock method [50], and the configuration-interaction (CI) method [3, 34]. All these methods treat the one-photon exchange exactly and the higher-order electron correlation is accounted for within the Breit approximation only.

The current level of experimental accuracy demands rigorous QED calculations of two-photon exchange contributions, which for  $n = 2$  states of Li-like ions were performed in Refs. [6, 8–11, 13, 51]. Meanwhile rigorous QED calculations of three- and more photon exchange contributions have not been performed up to now. For high- $Z$  few-electron ions