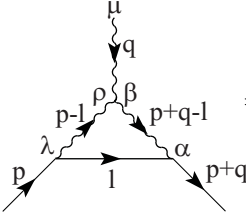


2.  = $-ie\Lambda_b^\mu(p, p+q)e^{ip\wedge q}$ (81)

$$= \int \frac{d^4 l}{(2\pi)^4} \frac{-ig_{\alpha\beta}}{(p+q-l)^2} \frac{-ig_{\lambda\rho}}{(p-l)^2} (-ie\gamma^\alpha) iS_0(l) (-ie\gamma^\lambda) \\ \times (2e)\gamma^{\mu\beta\rho}(q, -p-q+l, p-l) \frac{1}{2i} \left(1 - e^{2il\wedge q} e^{-2ip\wedge q}\right) e^{ip\wedge q},$$

whose planar part logarithmically diverges. In fact its pole part (PP) is given by

$$\text{PP}[-ie\Lambda_b^\mu(p, p+q)] = \frac{3}{16\pi^2} \frac{1}{\epsilon} \gamma^\mu \quad (82)$$

so that, in the minimal dimensional regularization scheme,

$$Z_{1F} = 1 - \frac{3}{16\pi^2} \frac{1}{\epsilon}, \quad (83)$$

which agrees with the result of previous calculation [12]. Contracting q_μ in the expression (81), we get

$$q_\mu \Lambda_b^\mu(p, p+q) = \int \frac{d^4 l}{(2\pi)^4} \frac{-ig_{\alpha\beta}}{(p+q-l)^2} \frac{-ig_{\lambda\rho}}{(p-l)^2} (-ie\gamma^\alpha) iS_0(l) (-ie\gamma^\lambda) \left(1 - e^{2il\wedge q} e^{-2ip\wedge q}\right) \\ \times \left\{ (p+q-l)^\rho q^\beta + (p-l)^\beta q^\rho - [(p+q-l) \cdot q + (p-l) \cdot q] g^{\beta\rho} \right\}, \quad (84)$$

which may be further simplified using

$$(p+q-l)^\rho q^\beta + (p-l)^\beta q^\rho = (p+q-l)^\rho (p+q-l)^\beta - (p-l)^\rho (p-l)^\beta, \text{ and} \\ (p+q-l) \cdot q + (p-l) \cdot q = (p+q-l)^2 - (p-l)^2, \quad (85)$$

to yield

$$q_\mu \Lambda_b^\mu(p, p+q) = \Sigma(p) - \Sigma(p+q) - \Sigma_{\text{np}}(p) + \Sigma_{\text{np}}(p+q) \quad (86) \\ + \int \frac{d^4 l}{(2\pi)^4} \frac{i}{(p+q-l)^2} \frac{i}{(p-l)^2} (ie)(\not{p} + \not{q} - \not{l}) iS_0(l) (ie)(\not{p} + \not{q} - \not{l}) \left(1 - e^{2il\wedge q} e^{-2ip\wedge q}\right) \\ - \int \frac{d^4 l}{(2\pi)^4} \frac{i}{(p+q-l)^2} \frac{i}{(p-l)^2} (ie)(\not{p} - \not{l}) iS_0(l) (ie)(\not{p} - \not{l}) \left(1 - e^{2il\wedge q} e^{-2ip\wedge q}\right),$$

where Σ is similar to Σ_{np} of Eq. (80), but without the phase factor.

For the last term in left-hand side of Eq. (76), we get

 = $iq^2 b(q^2)$ (87)