

Acknowledgments

We are indebted to Andrey Mironov and Niels Obers for useful discussions.

Appendix A: Measure for integrating over reparametrization

Introducing

$$v_i = s_i - s_{i-1} \quad s_N = s_f, \quad (\text{A1})$$

we rewrite the measure of [6] for the integration over reparametrizations as

$$\begin{aligned} \int_{s_0}^{s_f} \mathcal{D}_{\text{diff}} s &\equiv \lim_{N \rightarrow \infty} \prod_{i=2}^{N-1} \int_{s_0}^{s_{i+1}} \frac{ds_i}{(s_{i+1} - s_i)} \int_{s_0}^{s_2} \frac{ds_1}{(s_2 - s_1)(s_1 - s_0)} \\ &= \lim_{N \rightarrow \infty} \prod_{i=1}^N \int_0^\infty \frac{dv_i}{v_i} \delta^{(1)}\left(s_f - s_0 - \sum_{j=1}^N v_j\right). \end{aligned} \quad (\text{A2})$$

The integration over v_i 's in Eq. (A2) can be represented through the integration over a scalar field as follows. Writing

$$v_i = e^{\psi_i/2}, \quad (\text{A3})$$

we have

$$\int_0^\infty \frac{dv_i}{v_i} \dots = \frac{1}{2} \int_{-\infty}^{+\infty} d\psi_i \quad (\text{A4})$$

and

$$\int_{s_0}^{s_f} \mathcal{D}_{\text{diff}} s = \lim_{N \rightarrow \infty} \frac{1}{2^N} \prod_{i=1}^N \int_{-\infty}^{+\infty} d\psi_i \delta^{(1)}\left(s_f - s_0 - \sum_{j=1}^N e^{-\psi_j/2}\right). \quad (\text{A5})$$

This represents the continuous measure as

$$\int_{s_0}^{s_f} \mathcal{D}_{\text{diff}} s = \int \mathcal{D}\psi \delta^{(1)}\left(s_f - s_0 - \int_{s_0}^{s_f} dt e^{-\psi(t)/2}\right), \quad (\text{A6})$$

where t is a certain parametrization of the contour (e.g. through the proper time) and $\mathcal{D}\psi$ is the usual measure

$$\int \mathcal{D}\psi = \prod_{s=s_0}^{s_f} \int_{-\infty}^{+\infty} d\psi(s). \quad (\text{A7})$$

The scalar field ψ , that appears in Eqs. (A1), (A3), is in fact a discretization of the boundary value of the Liouville field

$$\varphi(\tau, s) \Big|_{\text{boundary}} = \psi(s), \quad (\text{A8})$$