

Using orthogonality of the Legendre polynomials it can be expressed as

$$C_\alpha = \frac{2}{a_\perp \sqrt{\pi L}} \left[\sum_{n=0}^{\infty} \left(\frac{\cot \sqrt{\epsilon_\alpha - \lambda n}}{(\epsilon_\alpha - \lambda n)^{3/2}} + \frac{1}{(\epsilon_\alpha - \lambda n) \sin^2 \sqrt{\epsilon_\alpha - \lambda n}} \right) \right]^{-1/2}, \quad (10)$$

where the rescaled energy ϵ_α is given by $\epsilon_\alpha \equiv \lambda E_\alpha / (2\hbar\omega_\perp)$ and $\lambda \equiv (L/a_\perp)^2$ is the aspect ratio.

For the Fermi-Huang interaction the eigenenergy equation (4) attains the form

$$\frac{\partial}{\partial r} \left[r \sum_{n,l} \frac{\langle \rho | n \rangle \langle z | l \rangle \langle n | 0 \rangle \langle 0 | 0 \rangle}{E - E_{nl}} \right]_{r=0} = \frac{1}{V} \quad (11)$$

In the limit $r \rightarrow 0$, the eigenstate is spherically symmetric. This allows us to deal with the z -axis only (see [6]). Substitution of the noninteracting wavefunctions (6) and (7) and energies (8) with the subsequent summation over l leads to

$$\sqrt{\lambda} \frac{\partial}{\partial z} \left[z \sum_{n=0}^{\infty} \frac{\cos(2\sqrt{\epsilon_\alpha - \lambda n}(z/L - 1/2))}{\sqrt{\epsilon_\alpha - \lambda n} \sin \sqrt{\epsilon_\alpha - \lambda n}} \right]_{z=0} = \frac{a_s}{a_\perp}. \quad (12)$$

This sum contains both the regular part and the irregular one, the latter being proportional to z^{-1} . The regular part can be extracted using the identity

$$\lim_{z \rightarrow 0} \left[\sum_{n=0}^{\infty} (\lambda n - \epsilon)^{-1/2} \exp(-2\sqrt{\lambda n - \epsilon} z/L) - \frac{L}{\lambda z} \right]$$

$$= -\frac{1}{\sqrt{\lambda}} \zeta \left(\frac{1}{2}, -\frac{\epsilon}{\lambda} \right) \quad (13)$$

(see [6]) where $\zeta(\nu, \alpha)$ is the Hurwitz zeta function (see, e. g., [7, 8]). Finally we arrive at the transcendental equation for the eigenenergies (4) in [1]. The summands in the sums Eqs. (3) and (4) in [1] and in Eq. (10) decay exponentially with n , leading to the fast converging series. Note that the imaginary parts of the two terms in the left hand side of Eq. (4) in [1] cancel each other automatically.

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