



FIG. 2: Left: Dark matter annihilation into the dark photon A', which decays into charged leptons such as electrons and/or muons, can explain the cosmic-ray electron and/or positron excesses seen by PAMELA, Fermi, ATIC, HESS, and other experiments. Right: Dark matter scattering into an excited state off nuclei through A' exchange in direct dark matter detection experiments can explain the annual modulation signal observed by DAMA/LIBRA, and the null results of other direct detection experiments.

ing is dominated by an inelastic process,

$$\chi N \to \chi^* N,$$
 (9)

in which the dark matter χ scatters off a nucleus N into an excited state χ^* with mass splitting $\delta \approx 100$ keV [67]. The kinematics of these reactions is also remarkably consistent with all the distinctive properties of the nuclear recoil spectrum reported by DAMA/LIBRA. In addition, the INTEGRAL telescope [71] has reported a 511keV photon signal near the galactic center, indicating a new source of $\sim 1\text{-}10$ MeV electrons and positrons. This excess could be explained by collisions of $\mathcal{O}(100$ GeV-1 TeV) mass dark matter into $\mathcal{O}(\text{MeV})$ excited states in the galaxy [72] — dark matter excited by scattering decays back to the ground state by emitting a soft e^+e^- pair. The 511keV excess then arises from the subsequent annihilation of the produced positrons.

The existence of an A' may also help explain various other particle physics anomalies [27] such as the anomalous magnetic moment of the muon $((g-2)_{\mu})$ [73] and the HyperCP anomaly [74].

While these experimental hints provide an urgent motivation to look for an A', it is important to emphasize the value of these searches in general. There has never been a systematic search for new GeV-scale force carriers that are weakly coupled to Standard Model particles. Nothing forbids their existence, and their discovery would have profound implications for our understanding of nature. A relatively simple experiment using the facilities available at, for example, Jefferson Laboratory and Mainz will probe a large and interesting range of A' masses and couplings.

C. Current Limits on Light U(1) Gauge Bosons

Constraints on new A''s that decay to e^+e^- and the search reach of an experiment using the spectrometers of Hall A at Jefferson Laboratory are summarized in Figure 1. Shown are constraints from electron and muon anomalous magnetic moment measurements, a_e and a_μ [27], the BaBar search for $\Upsilon(3S) \to \gamma A' \to \gamma \mu^+ \mu^-$, and three beam dump experiments, E137, E141, and E774 [3]. The constraints from a_μ and the BaBar search assume that the A' couples to muons — this is the case, for example, if it mixes with the photon. If it only couples to electrons, then the constraints on α'/α and $m_{A'}$ in the region to which the proposed experiment is sensitive are weaker than $\alpha'/\alpha \lesssim 10^{-4}$.

We refer the reader to [3, 27] for details on existing constraints. Here, we briefly review the constraint on $e^+e^- \rightarrow \gamma A' \rightarrow \gamma \mu^+\mu^-$ derived from the BaBar search [75]. If the A' couples to both electrons and muons, this is the most relevant constraint in the region probed by the proposed experiment. The analysis of [75] was in fact a search for $\Upsilon(3S)$ decays into a pseudoscalar $a, \Upsilon(3S) \to \gamma a \to \gamma \mu^+ \mu^-$, but can be interpreted as a limit on A' production because the final states are identical. Using $\mathcal{L}_{\rm int} \sim 30~{\rm fb^{-1}}$ of data containing $\sim 122 \times 10^6 \Upsilon(3S)$ events, a 90% C.L. upper limit of roughly $(1-4) \times 10^{-6}$ on the $\gamma \mu^+ \mu^-$ branching fraction was found for $m_{A'} \sim 2m_{\mu} - 1$ GeV. This search would thus be sensitive to about $\sim 100-500$ events with $e^+e^- \rightarrow \gamma A' \rightarrow \gamma \mu^+\mu^-$. Requiring that $\sigma(e^+e^- \to \gamma A') \times BR(A' \to \mu^+\mu^-) \times \mathcal{L}_{\rm int} \lesssim 500$, where $BR(A' \to \mu^+ \mu^-) = 1/(2 + R(m_{A'}))$ for $m_{A'} > 2m_{\mu}$ with $R = \frac{\sigma(e^+e^- \to \text{hadrons}; E = m_{A'})}{\sigma(e^+e^- \to \mu^+\mu^-; E = m_{A'})}, \text{ and rescaling the resulting}$