$\chi_i(t<0)=0$ . The Fourier transform of Eq. 3 for a purely ac field  $E(t)=E\cos(\omega t)$  gives:

$$\frac{P(\omega')}{\epsilon_0} = \frac{E}{2} \left[ \chi_1(\omega) + \frac{3E^2}{4} \chi_3(-\omega, \omega, \omega) + \dots \right] \delta(\omega' - \omega) 
+ \frac{E}{2} \left[ \chi_1(-\omega) + \frac{3E^2}{4} \chi_3(\omega, -\omega, -\omega) + \dots \right] \delta(\omega' + \omega) 
+ \frac{E^3}{8} \chi_3(\omega, \omega, \omega) \delta(\omega' - 3\omega) 
+ \frac{E^3}{8} \chi_3(-\omega, -\omega, -\omega) \delta(\omega' + 3\omega) + \dots,$$
(4)

where the polarization P and the susceptibilities  $\chi_i$  are now taken in the frequency domain and the dots indicate again infinite sums involving higher order terms. The response P(t)to  $E(t) = E\cos(\omega t)$  can thus be written

$$P(t)/\epsilon_0 = Re \left[ (E\chi_1(\omega) + 3/4E^3\chi_{\bar{3}}(\omega) + ...)e^{-i\omega t} \right] + Re \left[ 1/4E^3\chi_3(\omega)e^{-i3\omega t} + ... \right] + ....$$
 (5)

To obtain Eq. (5), we have used the fact that because  $\chi_1$  and  $\chi_3$  are real in the time domain, their Fourier transform verify  $\chi_1^*(\omega) = \chi_1(-\omega)$  and  $\chi_3^*(\omega_1, \omega_2, \omega_3) = \chi_3(-\omega_1, -\omega_2, -\omega_3)$  (the star denotes the complex conjugate), and the invariance of  $\chi_3$  by permutation of its arguments. For simplicity, we write  $\chi_3(\omega) = \chi_3(\omega, \omega, \omega)$  and  $\chi_3(\omega) = \chi_3(-\omega, \omega, \omega)$ . Eq. (5) can be written

$$P(t)/\epsilon_0 = E(\chi_1' \cos \omega t + \chi_1'' \sin \omega t) + 3/4E^3(\chi_3' \cos \omega t + \chi_3'' \sin \omega t) + \dots + 1/4E^3(\chi_3' \cos 3\omega t + \chi_3'' \sin 3\omega t) + \dots,$$
(6)

where the susceptibilities  $\chi_i$  are given as a function of their real and imaginary parts  $\chi'_i$  and  $\chi''_i$ . For practical applications, the modulii and arguments  $|\chi_i|$  and  $\delta_i$  are rather used:

$$P(t)/\epsilon_0 = E |\chi_1| \cos(\omega t - \delta_1) + 3/4E^3 |\chi_{\bar{3}}| \cos(\omega t - \delta_{\bar{3}}) + \dots + 1/4E^3 |\chi_3| \cos(3\omega t - \delta_3) + \dots$$
(7)

We see in the rhs in Eqs 5-7 that the nonlinear susceptibility of interest, namely  $\chi_3$