

we can write

$$T_{\mu_1 \dots \mu_r \lambda_1 \dots \lambda_s} = \frac{1}{i^{r+s}} \frac{\partial}{\partial a_1^{\mu_1}} \dots \frac{\partial}{\partial a_1^{\mu_r}} \frac{\partial}{\partial a_2^{\lambda_1}} \dots \frac{\partial}{\partial a_2^{\lambda_s}} \exp [i (a_1 \cdot k_1 + a_2 \cdot k_2)] \Big|_{a_1=a_2=0}. \quad (2.46)$$

Now, consider the integral

$$G^{(d)}(q^2) = \int \frac{d^d k_1}{(2\pi)^d} \int \frac{d^d k_2}{(2\pi)^d} \frac{\exp [i (a_1 \cdot k_1 + a_2 \cdot k_2)]}{c_1^{n_1} c_2^{n_2} c_3^{n_3} c_4^{n_4} c_5^{n_5}}. \quad (2.47)$$

To evaluate this integral we will use the following identity [110, 118]

$$\frac{1}{(k^2 - m^2)^n} = \frac{1}{i^n \Gamma(n)} \int_0^\infty d\alpha \alpha^{n-1} \exp [i\alpha (k^2 - m^2)]. \quad (2.48)$$

The integration variable α is called an alpha parameter and serves a similar purpose to the Feynman parameters introduced earlier. The resulting k_1 and k_2 loop integrals can be evaluated using the integral [118]

$$\int d^d k \exp [i (Ak^2 + 2q \cdot k)] = i \left[\frac{\pi}{iA} \right]^{\frac{d}{2}} \exp \left[-\frac{iq^2}{A} \right]. \quad (2.49)$$

Doing this, the result is

$$G^{(d)}(q^2) = \frac{i^{2-d}}{(4\pi)^d} \prod_{i=1}^5 \frac{1}{i^{n_i} \Gamma(n_i)} \int_0^\infty \frac{d\alpha_i}{[D(\alpha)]^{\frac{d}{2}}} \alpha_i^{n_i-1} \exp \left[i \left(\frac{Q(\alpha_i, a_1, a_2)}{D(\alpha)} - \sum_{j=1}^5 \alpha_j m_j^2 \right) \right]. \quad (2.50)$$

The functions $D(\alpha)$ and $Q(\alpha, a_1, a_2)$ are

$$D(\alpha) = \alpha_5 (\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4) + (\alpha_1 + \alpha_3) (\alpha_2 + \alpha_4), \quad (2.51)$$

$$\begin{aligned} Q(\alpha_i, a_1, a_2) &= [(\alpha_1 + \alpha_2) (\alpha_3 + \alpha_4) \alpha_5 + \alpha_1 \alpha_2 (\alpha_3 + \alpha_4) + \alpha_3 \alpha_4 (\alpha_1 + \alpha_2)] q^2 \\ &+ (a_1 \cdot q) Q_1 + (a_2 \cdot q) Q_2 + a_1^2 Q_{11}^2 + a_2^2 Q_{22}^2 + (a_1 \cdot a_2) Q_{12}, \end{aligned} \quad (2.52)$$