c for repulsive fermions c 0.14 0.12 0.10 0.08 0.04 0.02

FIG. 20: The scaling function c for repulsive fermions as a function of μ/T .

-3

The free energy density then takes the simple form:

-5

-6

$$\mathcal{F} = -\frac{1}{\beta} \sum_{a} \int (d\mathbf{k}) \left[-s_a \log \left(1 - s_a e^{-\beta \varepsilon_a} \right) - \frac{1 - y_a^{-1}}{2(e^{\beta \varepsilon_a} - s_a)} \right]$$
(106)

-2

For the two-component fermion defined by the action (3), spin up particles scatter with spin down, and $G_{\uparrow\downarrow} = G_{\downarrow\uparrow}$. Thus, if the chemical potentials are equal, $\mu_{\uparrow} = \mu_{\downarrow} = \mu$, then $\varepsilon_{\uparrow} = \varepsilon_{\downarrow}$, and the thermodynamics is just two copies of the single component fermion described above. In particular, the density and c are doubled, but the energy per particle scaling functions are the same.

X. CONCLUSIONS

We have shown that the S-matrix based approach to quantum gases developed in [20] leads to a new treatment of the scale-invariant unitary limit, where all of the thermodynamic scaling functions can be computed as a function of μ/T . Though our