These equations are given by

$$i\frac{\partial\psi_1}{\partial t} = -\frac{\partial^2\psi_1}{\partial x^2} + v_1(x,t)\psi_1 + \sum_{k=1}^2 (g_{1k}(x,t)|\psi_k|^2)\psi_1,\tag{1}$$

$$i\frac{\partial\psi_2}{\partial t} = -\frac{\partial^2\psi_2}{\partial x^2} + v_2(x,t)\psi_2 + \sum_{k=1}^2 (g_{2k}(x,t)|\psi_k|^2)\psi_2,\tag{2}$$

where  $\psi_k = \psi_k(x,t)$ , the functions,  $v_k(x,t)$ , are the trapping potentials, and the  $g_{ij}$ , (i,j=1,2) describe the strength of the cubic nonlinearities. Note that we are using standard notation, with both the fields and coordinates dimensionless. Note also that we are supposing that our model describe 1D BEC, and then it should engender cigar-shaped traps to induce tight confinement in two transverse directions, leaving the condensate almost free along the x axis. In the superchemistry model of [7, 8], the above equations are amended by a term which goes as  $\sqrt{2}\alpha\psi_1\psi_2$  in Eq. (1) and with a term of  $\frac{\alpha}{\sqrt{2}}\psi_1^2$  in Eq. (2). These, number-nonconserving, terms arise from treating the Feshbach resonance as a coupling term between atoms and molecules.

It is interesting to comment here that the elimination of, say,  $\psi_2$  in favor of an effective equation for  $\psi_1$ , would necessarily end up with a GP equation with depletion effect. This could appear in the form of complex nonlinearity. We leave such study for a future work.

In the following we seek to reduce the above two equations into the following pair of equations

$$\mu_1 A_1 = -\frac{\partial^2 A_1}{\partial \zeta^2} + \sum_{k=1}^2 \left( G_{1k} |A_k|^2 \right) A_1, \tag{3}$$

$$\mu_2 A_2 = -\frac{\partial^2 A_2}{\partial \zeta^2} + \sum_{k=1}^2 \left( G_{2k} |A_k|^2 \right) A_2, \tag{4}$$

with  $\mu_i$  and the  $G_{ij}$  being constants. We achieve this with the use of the two*Anzatse*:

$$\psi_1(x,t) = \rho(x,t)e^{i\eta(x,t)}A_1(\zeta(x,t)),$$
 (5)

$$\psi_2(x,t) = \rho(x,t)e^{i\eta(x,t)}A_2(\zeta(x,t)). \tag{6}$$

However, we now have to have

$$\rho \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x} \left( \rho^2 \frac{\partial \eta}{\partial x} \right) = 0, \tag{7}$$

$$\frac{\partial \zeta}{\partial t} + 2 \frac{\partial \eta}{\partial x} \frac{\partial \zeta}{\partial x} = 0, \tag{8}$$

$$\frac{\partial}{\partial x} \left( \rho^2 \frac{\partial \zeta}{\partial x} \right) = 0. \tag{9}$$