

$$D_i D_j \left( \Psi - \Phi \right)^{(2)} + \left\{ (-\Delta + 2\partial_\eta^2 + 4\mathcal{H}\partial_\eta - 2K) \Psi^{(2)} + (2\mathcal{H}\partial_\eta + 2\partial_\eta \mathcal{H} + 4\mathcal{H}^2 + \Delta + 2K) \Phi^{(2)} \right\} \gamma_{ij} \\ - \frac{1}{a^2} \partial_\eta \left( a^2 D_{(i} \nu_{j)}^{(2)} \right) + \frac{1}{2} (\partial_\eta^2 + 2\mathcal{H}\partial_\eta + 2K - \Delta) \chi_{ij}^{(2)} - 8\pi G \left( \partial_\eta \varphi \partial_\eta \varphi_2 - a^2 \varphi_2 \frac{\partial V}{\partial \varphi}(\varphi) \right) \gamma_{ij} = \Gamma_{ij}, \quad (6.22)$$

where  $\Gamma_0, \Gamma_i, \Gamma_{ij}$  are the collection of the quadratic term of the first-order perturbations as follows:

$$\Gamma_0 := 4\pi G \left( (\partial_\eta \varphi_1)^2 + D_i \varphi_1 D^i \varphi_1 + a^2 (\varphi_1)^2 \frac{\partial^2 V}{\partial \varphi^2} \right) - 4\partial_\eta \mathcal{H} \left( \Phi^{(1)} \right)^2 - 2 \Phi^{(1)} \partial_\eta^2 \Phi^{(1)} - 3D_k \Phi^{(1)} D^k \Phi^{(1)} - 10 \Phi^{(1)} \Delta \Phi^{(1)} \\ - 3 \left( \partial_\eta \Phi^{(1)} \right)^2 - 16K \left( \Phi^{(1)} \right)^2 - 8\mathcal{H}^2 \left( \Phi^{(1)} \right)^2 + D_l D_k \Phi^{(1)} \chi^{lk} + \frac{1}{8} \partial_\eta \chi_{lk}^{(1)} \partial_\eta \chi^{lk} + \mathcal{H} \chi_{kl}^{(1)} \partial_\eta \chi^{lk} \\ - \frac{3}{8} D_k \chi_{lm}^{(1)} D^k \chi^{lm} + \frac{1}{4} D_k \chi_{lm}^{(1)} D^l \chi^{mk} - \frac{1}{2} \chi_{lm}^{(1)} \Delta \chi_{lm}^{(1)} + \frac{1}{2} K \chi_{lm}^{(1)} \chi^{lm}; \quad (6.23)$$

$$\Gamma_i := 16\pi G \partial_\eta \varphi_1 D_i \varphi_1 - 4\partial_\eta \Phi^{(1)} D_i \Phi^{(1)} + 8\mathcal{H} \Phi^{(1)} D_i \Phi^{(1)} - 8 \Phi^{(1)} \partial_\eta D_i \Phi^{(1)} + 2D^j \Phi^{(1)} \partial_\eta \chi_{ji}^{(1)} - 2\partial_\eta D^j \Phi^{(1)} \chi_{ij}^{(1)} \\ - \frac{1}{2} \partial_\eta \chi_{jk}^{(1)} D_i \chi^{kj} - \chi_{kl}^{(1)} \partial_\eta D_i \chi^{lk} + \chi^{kl} \partial_\eta D_k \chi_{il}^{(1)}; \quad (6.24)$$

$$\Gamma_{ij} := 16\pi G D_i \varphi_1 D_j \varphi_1 + 8\pi G \left\{ (\partial_\eta \varphi_1)^2 - D_l \varphi_1 D^l \varphi_1 - a^2 (\varphi_1)^2 \frac{\partial^2 V}{\partial \varphi^2} \right\} \gamma_{ij} - 4D_i \Phi^{(1)} D_j \Phi^{(1)} - 8 \Phi^{(1)} D_i D_j \Phi^{(1)} \\ + \left( 6D_k \Phi^{(1)} D^k \Phi^{(1)} + 4 \Phi^{(1)} \Delta \Phi^{(1)} + 2 \left( \partial_\eta \Phi^{(1)} \right)^2 + 8\partial_\eta \mathcal{H} \left( \Phi^{(1)} \right)^2 + 16\mathcal{H}^2 \left( \Phi^{(1)} \right)^2 + 16\mathcal{H} \Phi^{(1)} \partial_\eta \Phi^{(1)} - 4 \Phi^{(1)} \partial_\eta^2 \Phi^{(1)} \right) \gamma_{ij} \\ - 4\mathcal{H} \partial_\eta \Phi^{(1)} \chi_{ij}^{(1)} - 2\partial_\eta^2 \Phi^{(1)} \chi_{ij}^{(1)} - 4D^k \Phi^{(1)} D_{(i} \chi_{j)k}^{(1)} + 4D^k \Phi^{(1)} D_k \chi_{ij}^{(1)} - 8K \Phi^{(1)} \chi_{ij}^{(1)} + 4 \Phi^{(1)} \Delta \chi_{ij}^{(1)} - 4D^k D_{(i} \Phi^{(1)} \chi_{j)k}^{(1)} \\ + 2\Delta \Phi^{(1)} \chi_{ij}^{(1)} + 2D_l D_k \Phi^{(1)} \chi^{lk} \gamma_{ij} + \partial_\eta \chi_{ik}^{(1)} \partial_\eta \chi_j^{(1)k} - D^k \chi_{il}^{(1)} D_k \chi_j^{(1)l} + D^k \chi_{il}^{(1)} D^l \chi_{jk}^{(1)} - \frac{1}{2} D_i \chi^{lk} D_j \chi_{lk}^{(1)} \\ - \chi_{lm}^{(1)} D_i D_j \chi^{ml} + 2 \chi_{lm}^{(1)} D_l D_{(i} \chi_{j)m}^{(1)} - \chi_{lm}^{(1)} D_m D_l \chi_{ij}^{(1)} \\ - \frac{1}{4} \left( 3\partial_\eta \chi_{lk}^{(1)} \partial_\eta \chi^{kl} - 3D_k \chi_{lm}^{(1)} D^k \chi^{ml} + 2D_k \chi_{lm}^{(1)} D^l \chi^{mk} - 4K \chi_{lm}^{(1)} \chi^{lm} \right) \gamma_{ij}. \quad (6.25)$$

Here, we used Eqs. (4.8), (5.12), (5.14), (5.16) and (5.18).

The tensor part of Eq. (6.22) is given by

$$(\partial_\eta^2 + 2\mathcal{H}\partial_\eta + 2K - \Delta) \chi_{ij}^{(2)} = 2\Gamma_{ij} - \frac{2}{3} \gamma_{ij} \Gamma_k{}^k - 3 \left( D_i D_j - \frac{1}{3} \gamma_{ij} \Delta \right) (\Delta + 3K)^{-1} \left( \Delta^{-1} D^k D_l \Gamma_k{}^l - \frac{1}{3} \Gamma_k{}^k \right) \\ + 4 \left\{ D_{(i} (\Delta + 2K)^{-1} D_{j)} \Delta^{-1} D^l D_k \Gamma_l{}^k - D_{(i} (\Delta + 2K)^{-1} D^k \Gamma_{j)k} \right\}. \quad (6.26)$$

This tensor mode is also called the second-order gravitational waves.

Further, the vector part of Eqs. (6.21) and (6.22) yields the initial value constraint and the evolution equation of the vector mode  $\nu_j^{(2)}$ :

$$\nu_i^{(2)} = \frac{2}{\Delta + 2K} \{ D_i \Delta^{-1} D^k \Gamma_k - \Gamma_i \}, \quad \partial_\eta \left( a^2 \nu_i^{(2)} \right) = \frac{2a^2}{\Delta + 2K} \{ D_i \Delta^{-1} D^k D_l \Gamma_k{}^l - D_k \Gamma_i{}^k \}. \quad (6.27)$$