

The final result for the real time noise correlator is then:

$$S^{\beta\beta'}(t, t') = \frac{(e^*)^2}{4\pi^2 a^2} e^{2\nu G^{\beta\beta'}(t-t')} (A(t)A^*(t') + A^*(t)A(t')) , \quad (36)$$

where we introduced the chiral green function of the bosonic fields:

$$G^{\beta\beta'}(t, t') = \langle T_K \{ \phi_r(t^\beta) \phi_r(t'^{\beta'}) \} \rangle - \frac{1}{2} \langle T_K \{ \phi_r(t^\beta)^2 \} \rangle - \frac{1}{2} \langle T_K \{ \phi_r(t'^{\beta'})^2 \} \rangle . \quad (37)$$

The double Fourier transform of this quantity, which will allow to relate it to the noise correlator, reads:

$$S^{\beta\beta'}(\Omega_1, \Omega_2) = \int \int dt dt' e^{i(\Omega_1 t + \Omega_2 t')} S(t, t') . \quad (38)$$

We now specify the periodic voltage modulation, which allows to write the tunneling amplitude in terms of a series of Bessel functions  $J_n$ :

$$A(t) = \Gamma_0 \sum_{n=-\infty}^{+\infty} e^{i(\omega_0 + n\omega_{AC})t} J_n \left( \frac{e^* V_1}{\hbar \omega_{AC}} \right) , \quad (39)$$

which gives the Fourier transform of non-symmetrised noise:

$$\begin{aligned} S^{\beta\beta'}(\Omega_1, \Omega_2) &= \frac{(e^*)^2 \Gamma_0^2}{2\pi^2 a^2} \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} J_n \left( \frac{e^* V_1}{\omega_{AC}} \right) J_m \left( \frac{e^* V_1}{\omega_{AC}} \right) \\ &\times \int \int dt dt' e^{i(\Omega_1 t + \Omega_2 t')} e^{2\nu G^{\beta\beta'}(t, t')} \cos(\omega_0(t - t') + \omega_{AC}(nt - mt')) . \end{aligned} \quad (40)$$

Next, it is convenient to perform a change of variable  $\tau = t - t'$  and  $\tau' = t + t'$ :

$$\begin{aligned} S^{\beta\beta'}(\Omega_1, \Omega_2) &= 2 \frac{(e^*)^2 \Gamma_0^2}{2\pi^2 a^2} \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} J_n \left( \frac{e^* V_1}{\omega_{AC}} \right) J_m \left( \frac{e^* V_1}{\omega_{AC}} \right) \\ &\times \int \int d\tau d\tau' e^{i(\Omega_1 - \Omega_2)\tau/2} e^{i(\Omega_1 + \Omega_2)\tau'/2} e^{2\nu G^{\beta\beta'}(\tau)} \cos \left( \left( \omega_0 + \frac{n+m}{2} \omega_{AC} \right) \tau + \frac{n-m}{2} \omega_{AC} \tau' \right) . \end{aligned}$$

Using standard trigonometric identities, one can write this expression as a product of separate integrals over  $\tau$  and  $\tau'$ . Integrals over  $\tau$  contain the (zero temperature) Green's function of the chiral fields and can be expressed in terms of Gamma function. The result has the form:

$$\begin{aligned} S^{\beta\beta'}(\Omega_1, \Omega_2) &= 2 \frac{(e^*)^2 \Gamma_0^2}{2\pi^2 a^2} \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} J_n \left( \frac{e^* V_1}{\omega_{AC}} \right) J_m \left( \frac{e^* V_1}{\omega_{AC}} \right) \\ &\left[ I_1(\Omega_1 + \Omega_2, \omega) I_2^{\beta\beta'}(\Omega_1 - \Omega_2, \omega_0, \omega) - I_3(\Omega_1 + \Omega_2, \omega) I_4^{\beta\beta'}(\Omega_1 - \Omega_2, \omega_0, \omega) \right] . \end{aligned} \quad (41)$$

The integrals  $I_1, I_2^{\beta\beta'}, I_3, I_4^{\beta\beta'}$  are defined and computed in the Appendix. The final result for the 4 Keldysh matrix elements of the noise correlator is:

$$\begin{aligned} S^{\beta-\beta}(\Omega_1, \Omega_2) &= 2 \frac{(e^*)^2 \Gamma_0^2}{4\pi^2 a^2} \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} J_n \left( \frac{e^* V_1}{\omega_{AC}} \right) J_m \left( \frac{e^* V_1}{\omega_{AC}} \right) \frac{\pi}{\Gamma(2\nu)} \left( \frac{a}{\nu_F} \right)^{2\nu} \times \\ &\left[ (1 - \beta \text{sgn}(\Omega_1 + \omega_0 + n\omega_{AC})) |\Omega_1 + \omega_0 + n\omega_{AC}|^{2\nu-1} \delta(\Omega_1 + \Omega_2 + (n-m)\omega_{AC}) \right. \\ &\left. + (1 - \beta \text{sgn}(\Omega_1 - \omega_0 - n\omega_{AC})) |\Omega_1 - \omega_0 - n\omega_{AC}|^{2\nu-1} \delta(\Omega_1 + \Omega_2 - (n-m)\omega_{AC}) \right] , \end{aligned} \quad (42)$$

$$\begin{aligned} S^{\beta\beta}(\Omega_1, \Omega_2) &= 2 \frac{(e^*)^2 \Gamma_0^2}{4\pi^2 a^2} \sum_{n=-\infty}^{+\infty} \sum_{m=-\infty}^{+\infty} J_n \left( \frac{e^* V_1}{\omega_{AC}} \right) J_m \left( \frac{e^* V_1}{\omega_{AC}} \right) \frac{\pi}{\Gamma(2\nu)} \left( \frac{a}{\nu_F} \right)^{2\nu} \frac{e^{-\beta i \pi \nu}}{\cos(\pi \nu)} \times \\ &\times \left[ |\Omega_1 + \omega_0 + n\omega_{AC}|^{2\nu-1} \delta(\Omega_1 + \Omega_2 + (n-m)\omega_{AC}) + |\Omega_1 - \omega_0 - n\omega_{AC}|^{2\nu-1} \delta(\Omega_1 + \Omega_2 - (n-m)\omega_{AC}) \right] . \end{aligned} \quad (43)$$