

# SOLITARY-WAVE FAMILIES OF THE MICRO-STRUCTURE PDE

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**ABSTRACT.** The Micro-Structure PDE is a model for ... Since solitary-wave solutions often play a central role in the long-time evolution of an initial disturbance, we consider such solutions here (via normal form approach) within the framework of reversible systems theory. On isolated curves in the relevant parameter region, the delocalized waves reduce to genuine embedded solitons. Directions for future work are also described.

## 1. INTRODUCTION

The nonlinear dispersive equation [2]

$$(1.1) \quad v_{tt} - bv_{xx} - \frac{\mu}{2} (v^2)_{xx} - \delta (\beta v_{tt} - \gamma v_{xx})_{xx} = 0$$

is known as the Micro-Structure PDE.

We use the theory of reversible systems and the method of normal forms to categorize the possible solitary waves of (1.1). As we shall see, several families of solitary waves exist in various regions of parameter space. Our main focus here will be on delineating the possible occurrence and multiplicity of solitary waves in different parameter regimes. Certain delicate questions relating to specific waves or wave families will form the basis of future work.

The remainder of this paper is organized as follows. In Section 2, we delineate the possible families of solitary waves in various parameter domains and on certain important curves using the theory of reversible systems. In Section 3 and 4, we next focus on the various transition curves and derive normal forms in their vicinity to confirm the existence of families of regular or delocalized solitary-wave solutions in their vicinity.

## 2. SOLITARY WAVES; LOCAL BIFURCATIONS

Solitary waves of (1.1) of the form  $v(x, t) = \phi(x - ct) = \phi(z)$  satisfy the fourth-order travelling wave ODE

$$(2.1) \quad \phi_{zzzz} - q\phi_{zz} + p\phi = \mathcal{N}[\phi]$$

where

$$(2.2) \quad \mathcal{N}[\phi] = -\Delta_1 \phi_z^2 - b\Delta_1 \phi \phi_{zz}$$

$$(2.3a) \quad z \equiv x - ct$$

$$(2.3b) \quad p \equiv 0$$

$$(2.3c) \quad q \equiv \frac{c^2 - b}{\delta(\beta c^2 - \gamma)}$$

$$(2.3d) \quad \Delta_1 \equiv \frac{\mu}{\delta(\beta c^2 - \gamma)}$$

Equation (2.1) is invariant under the transformation  $z \mapsto -z$  and is thus a reversible system. In this section we shall use the theory of reversible systems to characterize the homoclinic orbits to the fixed point of (2.1), which correspond to pulses or solitary waves of the Micro-Structure PDE in various regions of the  $(p, q)$  plane.

The linearized system corresponding to (2.1)

$$(2.4) \quad \phi_{zzzz} - q\phi_{zz} + p\phi = 0$$

has a fixed point

$$(2.5) \quad \phi = \phi_z = \phi_{zz} = \phi_{zzz} = 0$$

Solutions  $\phi = ke^{\lambda x}$  satisfy the characteristic equation  $\lambda^4 - q\lambda^2 + p = 0$  from which one may deduce that the structure of the eigenvalues is distinct in two regions of  $(p, q)$ -space. Since  $p = 0$  we have only two possible regions of eigenvalues. We denote  $C_0$  as the positive  $q$  axis and  $C_1$  the negative  $q$ -axis. First we shall consider the bounding curves  $C_0$  and  $C_1$  and their neighborhoods, then we shall discuss the possible occurrence and multiplicities of homoclinic orbits to (2.5), corresponding to pulse solitary waves of (1.1), in each region:

**Near  $C_0$ :** The eigenvalues have the structure  $\lambda_{1-4} = 0, 0, \pm\lambda$ , ( $\lambda \in \mathbb{R}$ ) and the fixed point (2.5) is a saddle-focus.

**Near  $C_1$ :** Here the eigenvalues have the structure  $\lambda_{1-4} = 0, 0, \pm i\omega$ , ( $\omega \in \mathbb{R}$ ). We will show by analysis of a four-dimensional normal form in Section 4 that there exists a *sech*<sup>2</sup> homoclinic orbit near  $C_1$ .

Having outlined the possible families of orbits homoclinic to the fixed point (2.5) of (2.4), corresponding to pulse solitary waves of (1.1), we now derive normal forms near the transition curves  $C_0$  and  $C_1$  to confirm the existence of regular or delocalized solitary waves in the corresponding regions of  $(p, q)$  parameter space.

### 3. NORMAL FORM NEAR $C_0$ : SOLITARY-WAVE SOLUTIONS

Using (2.4), the curve  $C_0$ , corresponding to  $\lambda = 0, 0, \pm\tilde{\lambda}$ , is given by

$$(3.1) \quad C_0 : p = 0, q > 0$$

Using (2.3c) implies

$$(3.2) \quad \frac{c^2 - b}{\delta(\beta c^2 - \gamma)} > 0$$

Denoting  $\phi$  by  $y_1$ , (2.1) may be written as the system

$$(3.3a) \quad \frac{dy_1}{dz} = y_2$$

$$(3.3b) \quad \frac{dy_2}{dz} = y_3$$

$$(3.3c) \quad \frac{dy_3}{dz} = y_4$$

$$(3.3d) \quad \frac{dy_4}{dz} = qy_3 - py_1 - (\Delta_1 y_2^2 + b\Delta_1 y_1 y_3)$$

We wish to rewrite this as a first order reversible system in order to invoke the relevant theory [1]. To that end, defining  $Y = \langle y_1, y_2, y_3, y_4 \rangle^T$  (3.3) may be written

$$(3.4) \quad \frac{dY}{dz} = L_{pq}Y - F_2(Y, Y)$$

where

$$(3.5) \quad L_{pq} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ q/3 & 0 & 1 & 0 \\ 0 & q/3 & 0 & 1 \\ q^2 - p & 0 & q/3 & 0 \end{pmatrix}$$

Since  $p = 0$  for (1.1), we have

$$(3.6) \quad \frac{dY}{dz} = L_{0q}Y - F_2(Y, Y)$$

where

$$(3.7) \quad F_2(Y, Y) = \langle 0, 0, 0, \Delta_1 y_2^2 + b\Delta_1 y_1 y_3 \rangle^T$$

Next we calculate the normal form of (3.6) near  $C_0$ . The procedure is closely modeled on [1] and many intermediate steps may be found there.

3.1. **Near  $C_0$ .** Near  $C_0$  the dynamics reduce to a 2-D Center Manifold

$$(3.8) \quad Y = A\zeta_0 + B\zeta_1 + \Psi(\epsilon, A, B)$$

and the corresponding normal form is

$$(3.9a) \quad \frac{dA}{dz} = B$$

$$(3.9b) \quad \frac{dB}{dz} = b\epsilon A + \tilde{c}A^2$$

Here,

$$(3.10) \quad \epsilon = \left( \frac{q^2}{9} - p \right) - \left( \frac{q}{3} \right)^2 = -p$$

measures the perturbation around  $C_0$ , and

$$(3.11a) \quad \zeta_0 = \langle 1, 0, -q/3, 0 \rangle^T$$

$$(3.11b) \quad \zeta_1 = \langle 0, 1, 0, -2q/3 \rangle^T$$

The linear eigenvalue of (3.9a) satisfies

$$(3.12) \quad \lambda^2 = b\epsilon$$

The characteristic equation of the linear part of (3.6) is

$$(3.13) \quad \lambda^4 - q\lambda^2 - \epsilon = 0$$

Hence, the eigenvalues near zero ( the Center Manifold ) satisfy  $\lambda^4 \ll \lambda^2$  and hence

$$(3.14) \quad \lambda^2 \sim -\frac{\epsilon}{q}$$

Matching (3.12) and (3.14)

$$(3.15) \quad b = -\frac{1}{q}$$

and only the nonlinear coefficient  $\tilde{c}$  remains to be determined in the normal form (3.9a).

In order to determine  $\tilde{c}$  (the coefficient of  $A^2$  in (3.9a) we calculate  $\frac{dY}{dz}$  in two ways and match the  $\mathcal{O}(A^2)$  terms.

To this end, using the standard 'suspension' trick of treating the perturbation parameter  $\epsilon$  as a variable, we expand the function  $\Psi$  in (3.8) as

$$(3.16) \quad \Psi(\epsilon, A, B) = \epsilon A \Psi_{10}^1 + \epsilon B \Psi_{01}^1 + A^2 \Psi_{20}^0 + AB \Psi_{11}^0 + B^2 \Psi_{02}^0 + \dots$$

where the subscripts denote powers of  $A$  and  $B$ , respectively, and the superscript denotes the power of  $\epsilon$ .

In the first way of computing  $dY/dz$ , we take the  $z$  derivative of (3.8) ( using (3.9a) and (4.10) ). The coefficient of  $A^2$  in the resulting expression is  $\tilde{c}\zeta_1$ . In the second way of computing  $dY/dz$ , use (3.8) and (4.10) in (3.4). The coefficient of  $A^2$  in the resulting expression is  $L_{0,q}\Psi_{20}^0 - F_2(\zeta_0, \zeta_0)$ . Hence

$$(3.17) \quad \tilde{c}\zeta_1 = L_{0,q}\Psi_{20}^0 - F_2(\zeta_0, \zeta_0)$$

Using (3.11a) and (3.7) and denoting  $\Psi_{20}^0 = \langle x_1, x_2, x_3, x_4 \rangle$  in (3.17) yields the equations

$$(3.18a) \quad 0 = x_2$$

$$(3.18b) \quad c = \frac{q}{3}x_1 + x_3$$

$$(3.18c) \quad 0 = \frac{q}{3}x_2 + x_4 \implies x_4 = 0 \text{ using (3.18b)}$$

and

$$(3.19) \quad -\frac{2q}{3}\tilde{c} = \frac{q}{3} \left( \frac{q}{3}x_1 + x_3 \right) + \frac{2q}{3} = \frac{q}{3}\tilde{c} + \frac{b\Delta_1}{3} \text{ using (3.18b)}$$

Hence we obtain

$$(3.20) \quad \tilde{c} = -\frac{b\Delta_1}{3}$$

Therefore, the normal form for (1.1) near  $C_0$  is

$$(3.21a) \quad \frac{dA}{dz} = B$$

$$(3.21b) \quad \frac{dB}{dz} = -\frac{\epsilon}{q}A - \frac{b\Delta_1}{3}A^2$$

#### 4. NORMAL FORM NEAR $C_1$ : POSSIBLE SOLITARY-WAVE SOLUTIONS

Using (2.4), the curve  $C_1$ , corresponding to  $\lambda = 0, 0 \pm i\omega$ , is given by

$$(4.1) \quad C_1 : p = 0, q < 0$$

Which implies

$$(4.2) \quad \frac{c^2 - b}{\delta(\beta c^2 - \gamma)} < 0$$

In order to investigate the possibility of  $\text{sech}^2$  homoclinic orbit in the neighborhood of  $C_1$  and delocalized solitary waves, we next compute the normal form near  $C_1$  following the procedure in [1].

Near  $C_1$  the dynamics reduce to a 4-D Center Manifold [1] Since all the eigenvalues are non-hyperbolic, the Center Manifold has the form (a nonlinear coordinate change [1])

$$(4.3) \quad Y = A\zeta_0 + B\zeta_0 + C\zeta_+ + \bar{C}\zeta_- + \Psi(\epsilon, A, B, C, \bar{C})$$

with a corresponding four-dimensional normal form

$$(4.4a) \quad \frac{dA}{dz} = B$$

$$(4.4b) \quad \frac{dB}{dz} = \bar{\nu}A + b_*A^2 + c_*|C|^2$$

$$(4.4c) \quad \frac{dC}{dz} = id_0C + i\bar{\nu}d_1C + id_2AC$$

Here  $C$  is complex,  $\bar{C}$  is the complex conjugate of  $C$ ,  $\epsilon, \zeta_0, \zeta_1$  are given previously and the two new complex eigenvectors co-spanning the Center Manifold are

$$(4.5) \quad \zeta_{\pm} = \left\langle 1, \lambda_{\pm}, 2q/3, \frac{\lambda_{\pm}}{3}q \right\rangle^T$$

Using (4.4c) and (3.9a)

$$(4.6) \quad \bar{\nu} = b\epsilon = -\frac{\epsilon}{q}$$

Also from the characteristic equation (3.13), the two non-zero (imaginary) roots are

$$(4.7) \quad \lambda^2 = \frac{q + \sqrt{q^2 + 4\epsilon}}{2} \approx q \text{ for } \epsilon \text{ small}$$

Hence

$$(4.8) \quad \lambda = \pm i\sqrt{-q}, q < 0$$

Matching this to the linear part of (4.4c) ( which corresponds to the imaginary eigenvalues),  $\lambda = id_0 = i\sqrt{-q}$  or

$$(4.9) \quad d_0 = \sqrt{-q}$$

If we do a dominant balance argument after the change of variable  $\tilde{\epsilon} = \sqrt{-3\alpha}$  on the characteristic equation as  $\lambda \rightarrow 0$  then we find  $d_1 = \frac{\sqrt{-3\alpha}}{18\alpha^2}$ . Using  $\alpha = q/3$  we find  $d_1 = \frac{\sqrt{-q}}{2q^2}$

The remaining undetermined coefficients in the normal form are the coefficients  $b_*$ ,  $c_*$  and  $d_2$  which correspond to the  $A^2$ ,  $|C|^2$  and  $AC$  terms respectively. In order to determine them, we follow the same procedure as in Section 3 and compute  $dY/dz$  in two distinct ways. We expand the function  $\Psi$  as

$$(4.10) \quad \Psi(\epsilon, A, B, C, \bar{C}) = \epsilon A \Psi_{1000}^1 + \epsilon B \Psi_{0100}^1 + A^2 \Psi_{2000}^0 + AB \Psi_{1100}^0 + AC \Psi_{1010}^0 + \epsilon C \Psi_{0010}^1 + \dots$$

with subscripts denoting powers of  $A$ ,  $B$ ,  $C$  and  $\bar{C}$ , respectively, and the superscript is the power of  $\epsilon$ . In the first way of computing  $dY/dz$  is computed by taking the  $z$  derivative of (4.3) ( using (4.4a) and (4.10)) and read off the coefficients of  $A^2$ ,  $\|C\|^2$ ,  $C\epsilon$  and  $AC$  terms.

In the second way,  $dY/dz$  is computed using (4.3) and (4.10) in (3.4) (with  $p = 0$  on  $C_1$  as given in (4.1)) and the coefficients of  $A$ ,  $B$ ,  $C$  and  $\bar{C}$  are once again read off.

Equating the coefficients of the corresponding terms in the two separate expressions for  $dY/dz$  yields the following equations:

$$(4.11a) \quad \mathcal{O}(A^2) : \quad b_* \zeta_1 = L_{0q} \Psi_{2000}^0 - F_2(\zeta_0, \zeta_0)$$

$$(4.11b) \quad \mathcal{O}(|C|^2) : \quad c_* \zeta_1 = L_{0q} \Psi_{0011}^0 - 2F_2(\zeta_+, \zeta_-)$$

$$(4.11c) \quad \mathcal{O}(\epsilon C) : \quad -\frac{i}{q} (d_1 \zeta_+ + d_0 \Psi_{0010}^1) = L_{0q} \Psi_{0010}^1 - F_2(\Psi_{0010}^1, \Psi_{0010}^1)$$

$$(4.11d) \quad \mathcal{O}(AC) : \quad i d_2 \zeta_+ + i d_0 \Psi_{1010}^0 = L_{0q} \Psi_{1010}^0 - 2F_2(\zeta_0, \zeta_+)$$

$$(4.11e)$$

where we have used the fact that  $F_2$  is a symmetric bilinear form. Equation (4.11b) is decoupled and yields  $c_* = 2\Delta_1 \left( \frac{1}{q} + \frac{2b}{3} \right)$ . The only coefficient left to determine is  $d_2$  which we shall compute now.

Using  $\Psi_{1010}^0 = \langle x_1, x_2, x_3, x_4 \rangle^T$  in (4.11e) implies

$$(4.12a) \quad i d_2 + i d_0 x_1 = x_2$$

$$(4.12b) \quad -d_0 d_2 + i d_0 x_2 = \frac{q}{3} x_1 + x_3$$

$$(4.12c) \quad \frac{2iq}{3} d_2 + i d_0 x_3 = \frac{q}{3} x_2 + x_4$$

$$(4.12d) \quad -\frac{q}{3} d_0 d_2 + i d_0 x_4 = \frac{q}{3} \left( \frac{q}{3} x_1 + x_3 \right) - \frac{2bq\Delta_1}{3}$$

Using (4.12a) in (4.12b), (4.12b) in (4.12d) and using these in (4.12c) yields  $d_2 = \frac{b\Delta_1}{3\sqrt{-q}}$ .

Therefore the normal form for (1.1) near  $C_1$  is

$$(4.13a) \quad \frac{dA}{dz} = B$$

$$(4.13b) \quad \frac{dB}{dz} = -\frac{\epsilon}{q} A - \frac{b\Delta_1}{3} A^2 + 2\Delta_1 \left( \frac{1}{q} + \frac{2b}{3} \right) |C|^2$$

$$(4.13c) \quad \frac{dC}{dz} = i\sqrt{-q}C - i\frac{\sqrt{-q}}{q^3} C\epsilon + i\frac{b\Delta_1}{3\sqrt{-q}} AC$$

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#### REFERENCES

- [1] IOOSS G., ADELMAYER, M., *Topics in Bifurcation Theory and Applications*, World Scientific, Singapore, 1998.
- [2] S. ROY CHOUDHURY, *Solitary-wave families of the Ostrovsky equation: An approach via reversible systems theory and normal forms*, Elsevier, (2007), pp. 1468–1479.