spherical elevator

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[1]: import numpy as np import matplotlib.pyplot as plt from IPython.display import HTML, Image from scipy.optimize import fsolve g=9.81
```

We consider a point mass m moving on a sphere with variable radius l(t) and centerpoint $\vec{R}(t)$ in a static gravitational potential $V(\vec{r}) \equiv mgr_z(t)$. We parametrize \vec{r} by some independent parameter t, i.e.

$$\vec{r}(t) = \vec{R}(t) + l(t) \begin{pmatrix} \cos(\phi(t))\sin(\theta(t)) \\ \sin(\phi(t))\sin(\theta(t)) \\ \cos(\theta(t)) \end{pmatrix}.$$

The Lagrangian of the system is given by

$$L[\theta(t), \dot{\theta}(t), \phi(t), \dot{\phi}(t)] = \frac{1}{2}m\dot{\vec{r}}(t)^2 - mgr_z(t)$$

and the Euler-Lagrange equations give

$$\ddot{\theta} = -\frac{1}{l} \left(+g\sin(\theta) - \cos(\theta)l\sin(\theta)\dot{\phi}^2 + 2l\dot{\theta} + \cos(\phi)\cos(\theta)\ddot{R}_x + \cos\theta\sin\phi\ddot{R}_y + \sin\theta\ddot{R}_z \right)$$

$$\ddot{\phi} = \frac{1}{l} \left(-2l\dot{\phi} - 2\cot\theta l\dot{\phi}\dot{\theta} + \csc\theta\sin\phi\ddot{R}_x - \cos\phi\csc\theta\ddot{R}_y \right),$$

which already implies that the physics of the point mass does not depend upon its mass. Hence, an elevators movement is the same, regardless of an additional player being on the elevator.

We defne

$$\vec{z}(t) = \begin{pmatrix} z_1(t) \\ z_2(t) \\ z_3(t) \\ z_4(t) \end{pmatrix} \equiv \begin{pmatrix} \theta(t) \\ \phi(t) \\ \dot{\theta}(t) \\ \dot{\phi}(t) \end{pmatrix},$$

such that

$$\dot{\vec{z}}(t) = \begin{pmatrix}
z_3 \\
z_4 \\
-\frac{1}{l} \left(+g\sin(z_1) - \cos(z_1)l\sin(z_1)\dot{z}_2^2 + 2\dot{l}\dot{z}_1 + \cos(z_2)\cos(z_1)\ddot{R}_x + \cos z_1\sin z_2\ddot{R}_y + \sin z_1\ddot{R}_z \right) \\
\frac{1}{l} \left(-2\dot{l}\dot{z}_2 - 2\cot z_1l\dot{z}_2\dot{z}_1 + \csc z_1\sin z_2\ddot{R}_x - \cos z_2\csc z_1\ddot{R}_y \right)$$

$$\equiv \text{RHS}(\vec{z}, l, \dot{l}, \ddot{\vec{R}}).$$

This means that, in order to track the time evolution of the point mass $\vec{r}(t)$, we need the length of the rope l(t), the rate of change in rope's length: $\dot{l}(t)$, and the acceleration of the sphere: $\ddot{\vec{R}}(t)$.

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[3]: def RHS(z, 1, 1d, Rdd):

"""

Returns the right-hand side (RHS) of the equation
\dot{z}(t) = RHS(z, l(t), \cot{l}(t), \dot{\vec{R}}(t))

Parameters:

z: np.ndarray of size 4
	z(t):=np.array([theta(t), phi(t), \dot{\theta}(t), \dot{phi}(t)])

l: floating point number
	Length of the rope at time t, i.e. l(t)

ld: floating point number
	Temporal derivative of the rope's length, i.e. \dot{l}

Rdd: np.ndarray of size 3
	Second derivative of sphere's position with respect to time, i.e. □

→\ddot{\vec{R}}(t)

"""
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z1 = z[0] # theta(t)
   z2 = z[1] # phi(t)
   z3 = z[2] \# \det\{theta\}(t)
   z4 = z[3] \# \det\{phi\}(t)
   z1d = z3 \# dot{z}_1 = z_3
   z2d = z4 \# \det\{z\}_2 = z_4
   z3d = -1./1*(
         g*np.sin(z1)
       - np.cos(z1)*l*np.sin(z1)*z2d**2.
       + 2.*ld*z1d
       + np.cos(z2)*np.cos(z1)*Rdd[0] + np.cos(z1)*np.sin(z2)*Rdd[1]
       + np.sin(z1)*Rdd[2]
   # EL eqs. give something like sin(theta)(... + phi'') = 0, i.e. if_{\sqcup}
\rightarrow theta==0, we can't divide by sin(theta).
   # This is a mathematical artifact of the spherical coordinates as (x, y, z):
→ R^2 -> R^3 is
   # not injective, i.e. if theta mod pi == 0, all points get mapped either to_{\sqcup}
\rightarrow the top or
   # the bottom of the sphere.
   # However, it should be more or less intuitive to notice that, if theta==0,,,
⇔seting phi''=0
   # comes with no physically relevant implications.
   if z1 % np.pi==0:
       z4d = 0
   else:
       z4d = 1./1*(
           - 2.*ld*z2d
           -2.*(1./np.tan(z1))*l*z1d*z2d
           + (1./np.sin(z1))*np.sin(z2)*Rdd[0]
           - np.cos(z2)/np.sin(z1)*Rdd[1]
       )
   return np.array([z1d, z2d, z3d, z4d])
   for k in range(N-1):
```

```
[4]: def explicit_euler(z, r, 1, 1d, R, Rdd, dt, N):
    for k in range(N-1):
        z[k+1] = explicit_euler_step(z[k], 1[k], 1d[k], Rdd[k], dt)
        r[k+1] = euclidean_coord(R[k+1], 1[k+1], z[k+1][0], z[k+1][1])

def implicit_euler(z, r, 1, 1d, R, Rdd, dt, N):
    for k in range(N-1):
        F = lambda x: x - z[k] - dt*RHS(x, 1[k+1], 1d[k+1], Rdd[k+1])
        z[k+1] = fsolve(F, z[k]+dt*RHS(z[k], 1[k], 1d[k], Rdd[k]))
        r[k+1] = euclidean_coord(R[k+1], 1[k+1], z[k+1][0], z[k+1][1])
```

```
[11]: # Example setup.
      T = 10.
      N = 10000
      dt = float( float(T)/float(N-1) )
      times = np.linspace(0., T, N)
      # Midpoint setup
      f1 = 0.5
      omega1 = 2.*np.pi*f1
      RO = 1.
      R = R0 * np.array([[np.cos(omega1*t), 0, 0] for t in times])
      Rdd = -(omega1**2) * R
      # Rope setup
      f2 = 0.1
      omega2 = 2.*np.pi*f2
      10 = 2.
      1 = np.array([10 for t in times])
      ld = np.array([0. for t in times])
      z0 = np.array([0, 0, 0, 0])
      r0 = euclidean_coord(R[0], 1[0], z0[0], z0[1])
      z = np.zeros((N, 4))
      r = np.zeros((N, 3))
      r[0] = r0
      z[0] = z0
      implicit_midpoint(z, r, 1, ld, R, Rdd, dt, N)
```

```
[13]: fig, ax = plt.subplots()
step = int(len(R[:, 0]) / 1000)
ax.scatter(r[::step, 0], r[::step, 2], s=1)
ax.scatter(R[::step, 0], R[::step, 2], s=1)
```

[13]: <matplotlib.collections.PathCollection at 0x7f6d30746c10>

