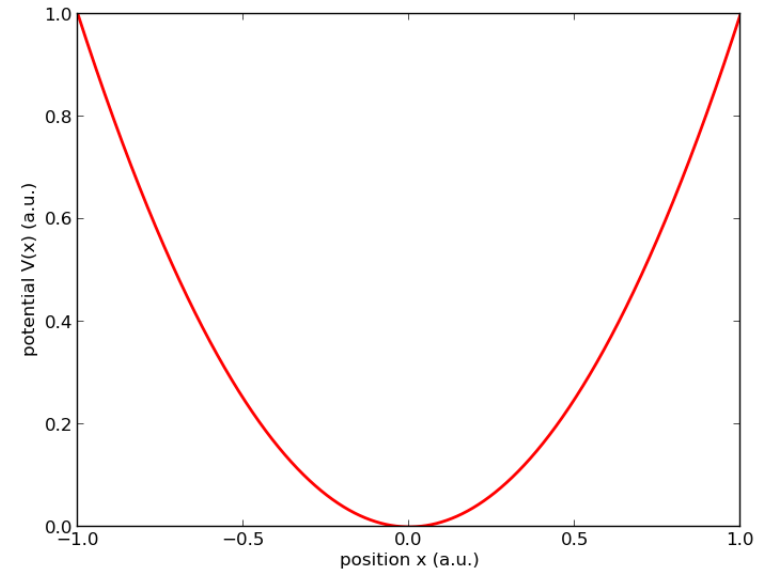
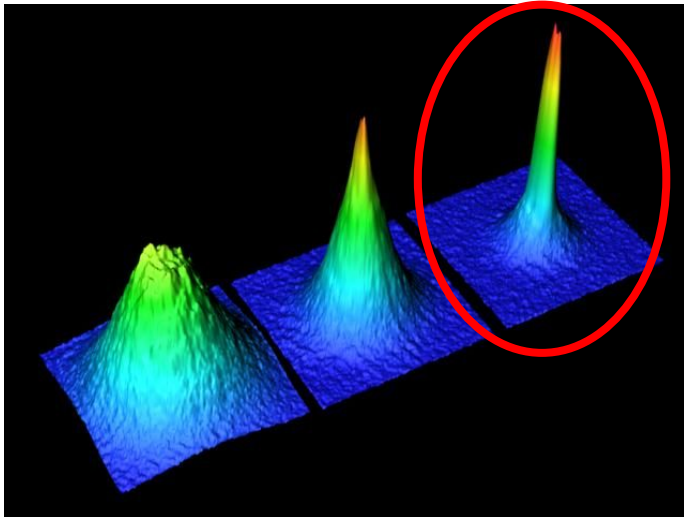
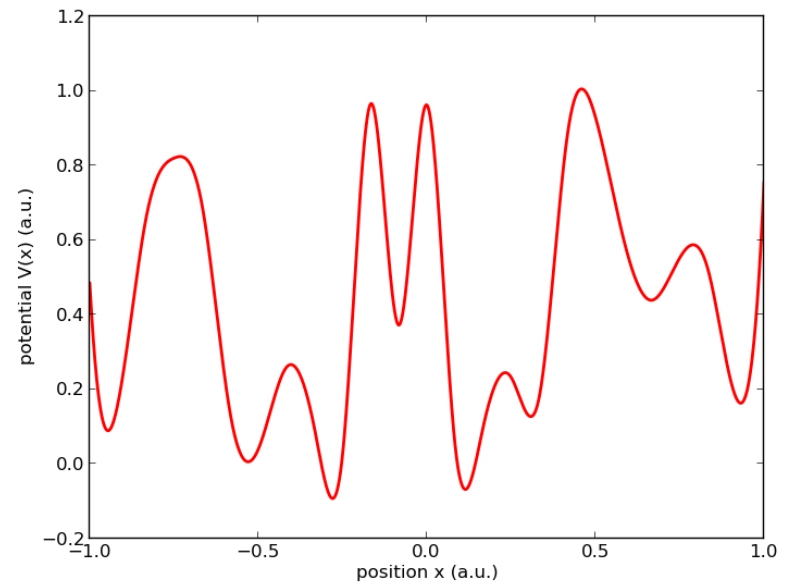
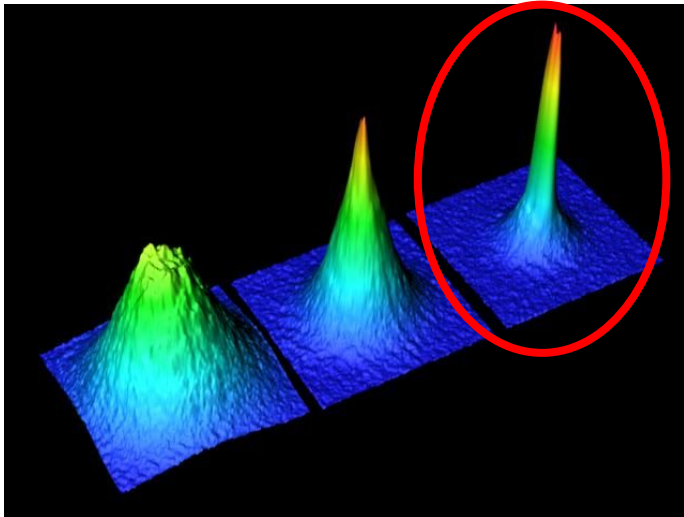

Numerical solving Gross-Pitaevskii-equation

Journal club 9-21-2012
Kris

Goal:



Goal:



Goal:

- Find ground state wave function of time-independent Schrödinger equation
 - 1D for simplicity
 - arbitrary potentials
 - easy to implement and use
 - include atom-atom s-wave interactions in mean field
 - > solve GPE
-

Harmonic oscillator: analytic

$$\hat{H} = \hat{T} + \hat{V} = \frac{\hat{p}^2}{2m} + V(x)$$

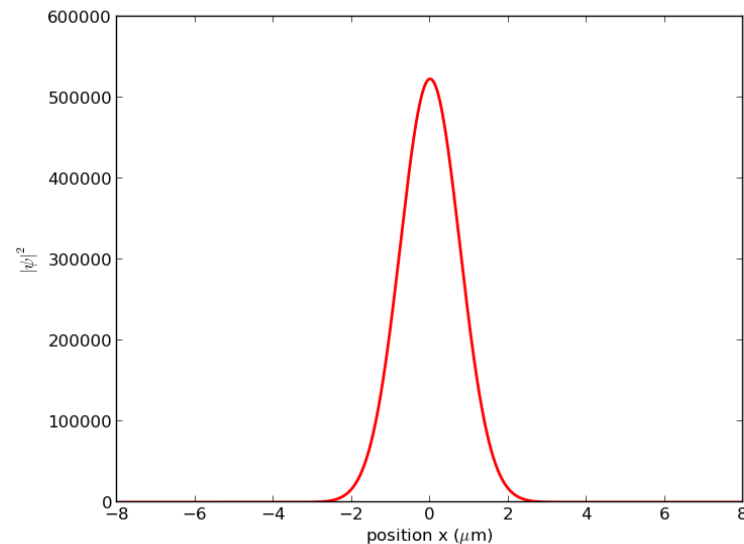
Harmonic oscillator: analytic

$$\hat{H} = \hat{T} + \hat{V} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$$

Harmonic oscillator: analytic

$$\hat{H} = \hat{T} + \hat{V} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2$$

$$\psi_n(x) = \frac{1}{\sqrt{2^n n!}} \cdot \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} \cdot e^{-\frac{m\omega x^2}{2\hbar}} \cdot H_n\left(\sqrt{\frac{m\omega}{\hbar}}x\right), \quad n = 0, 1, 2, \dots$$



Method I: Direct diagonalization

$$\hat{H}\Psi(x) = E\Psi(x)$$

- Find suitable basis $|n\rangle$
- Write down Hamiltonian in matrix form

$$H_{m,n} = \langle m|\hat{H}|n\rangle$$

- solve the eigenvalue problem
-

Method I: Direct diagonalization

$$H = T + V = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2$$

Method I: Direct diagonalization

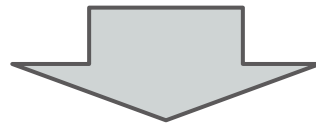
$$H = T + V = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2$$

$$V(x) = \sum_b c_b e^{ik_b x} \quad k_b = \frac{2\pi b}{L}$$

Method I: Direct diagonalization

$$H = T + V = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \frac{1}{2} m \omega^2 x^2$$

$$V(x) = \sum_b c_b e^{ik_b x} \quad k_b = \frac{2\pi b}{L}$$



$$H = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \sum_b c_b e^{ik_b x}$$

Method I: Direct diagonalization

$$H = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \sum_b c_b e^{ik_b x}$$

- Find suitable basis

$$|n\rangle = \frac{1}{\sqrt{L}} e^{ik_n x}$$

Method I: Direct diagonalization

$$H = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + \sum_b c_b e^{ik_b x} \quad |n\rangle = \frac{1}{\sqrt{L}} e^{ik_n x}$$

- Write down Hamiltonian in matrix form

$$H_{m,n} = \langle m | \hat{H} | n \rangle$$

$$T_{m,n} = \frac{\hbar^2 k_n^2}{2m} \delta_{m,n}$$

$$V_{m,n} = c_b \delta_{m,n+b}$$

Method I: Direct diagonalization

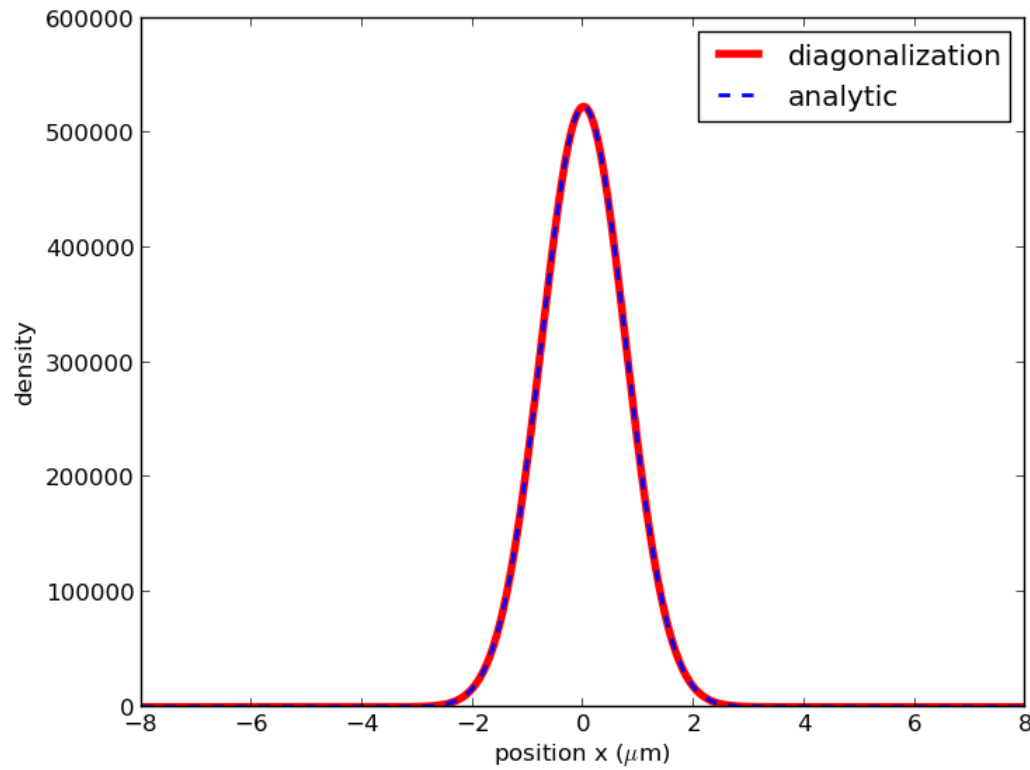
```
x,dx = linspace(-8e-6,8e-6,N,endpoint=False,retstep=True)
k = 2*pi*fftshift(fft.fftfreq(len(x), dx))

T = diag(k**2*hbar**2/(2.*m)+1j*0)

potential = .5*m*omega**2*x**2
potential_ft = concatenate((zeros(N/2),fftshift(fft(potential))/N,zeros(N/2)))
V = array([roll(potential_ft,i) for i in xrange(N)]][:,N-1:]

eva,eve= eig(T+V)
```

Method I: Direct diagonalization



Imaginary time propagation

- find approximation of the ground state of a Hamiltonian

Imaginary time propagation

- find approximation of the ground state of a Hamiltonian
- assume general wave function $|\psi\rangle = \sum_k c_k |k\rangle$
with $\hat{H}|k\rangle = \epsilon_k |k\rangle$ and $\epsilon_0 = 0$ for simplicity

Imaginary time propagation

- find approximation of the ground state of a Hamiltonian
- assume general wave function $|\psi\rangle = \sum_k c_k |k\rangle$
with $\hat{H}|k\rangle = \epsilon_k |k\rangle$ and $\epsilon_0 = 0$ for simplicity
- apply the time propagator $e^{i\Delta_t \hat{H}/\hbar}$ in imaginary time
i.e. $\Delta_t = -i\tau$

$$|\psi(\tau)\rangle = e^{-\hat{H}\tau/\hbar} \sum_k c_k |k\rangle = \sum_k e^{-\epsilon_k \tau} c_k |k\rangle$$

- not unitary -> norm not preserved
-

Split-Step Technique

- In general $\hat{H} = \hat{T} + \hat{V}$

Split-Step Technique

- In general $\hat{H} = \hat{T} + \hat{V}$ and $[\hat{T}, \hat{V}] \neq 0$

Split-Step Technique

- In general $\hat{H} = \hat{T} + \hat{V}$ and $[\hat{T}, \hat{V}] \neq 0$

$$e^{-i\Delta t \hat{H}/\hbar} = e^{-i\Delta t(\hat{T} + \hat{V})/\hbar} \neq e^{-i\Delta t \hat{T}/\hbar} e^{-i\Delta t \hat{V}/\hbar}.$$

Split-Step Technique

- In general $\hat{H} = \hat{T} + \hat{V}$ and $[\hat{T}, \hat{V}] \neq 0$

$$e^{-i\Delta t \hat{H}/\hbar} = e^{-i\Delta t(\hat{T}+\hat{V})/\hbar} \neq e^{-i\Delta t \hat{T}/\hbar} e^{-i\Delta t \hat{V}/\hbar}.$$

$$e^{-\frac{i}{\hbar} \hat{H} \Delta t} = e^{\tilde{T} + \tilde{V}} = e^{\tilde{T}} \cdot e^{\tilde{V}} \cdot e^{-\frac{1}{2}[\tilde{T}, \tilde{V}] + \mathcal{O}(\Delta t^3)},$$

Split-Step Technique

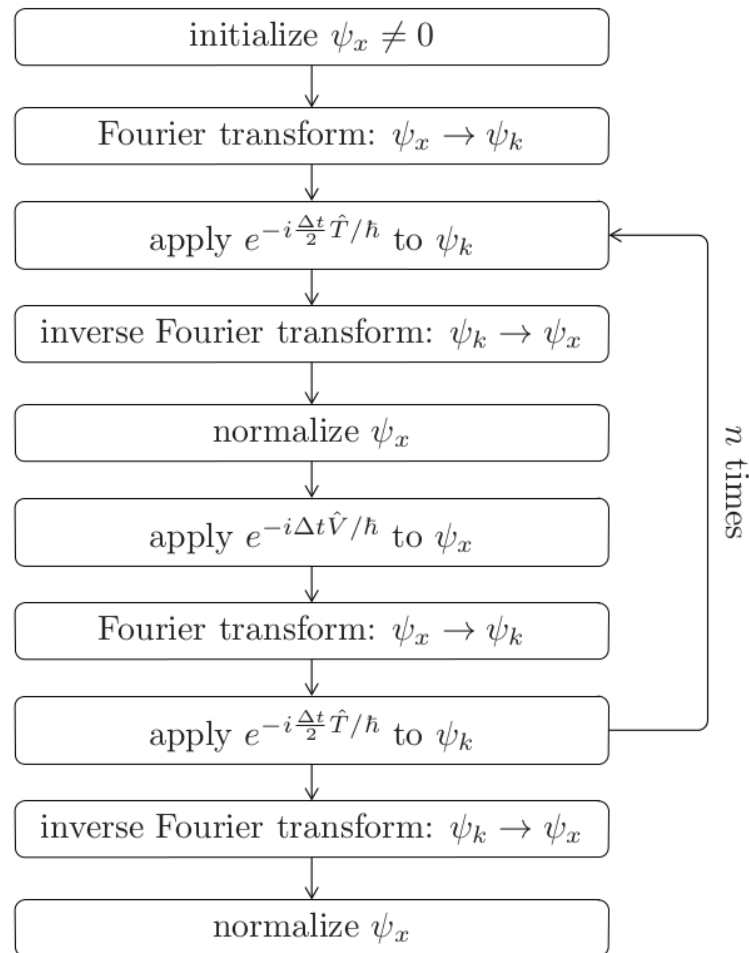
- In general $\hat{H} = \hat{T} + \hat{V}$ and $[\hat{T}, \hat{V}] \neq 0$

$$e^{-i\Delta t \hat{H}/\hbar} = e^{-i\Delta t(\hat{T}+\hat{V})/\hbar} \neq e^{-i\Delta t \hat{T}/\hbar} e^{-i\Delta t \hat{V}/\hbar}.$$

$$e^{-\frac{i}{\hbar} \hat{H} \Delta t} = e^{\tilde{T} + \tilde{V}} = e^{\tilde{T}} \cdot e^{\tilde{V}} \cdot e^{-\frac{1}{2}[\tilde{T}, \tilde{V}] + \mathcal{O}(\Delta t^3)},$$

$$\begin{aligned} e^{-\frac{i}{\hbar} H \Delta t} &= e^{\tilde{T}/2 + \tilde{V} + \tilde{T}/2} \\ &= e^{\tilde{T}/2} \cdot e^{\tilde{V} + \tilde{T}/2} \cdot e^{-\frac{1}{2}[\tilde{T}/2, \tilde{T}/2 + \tilde{V}] + \mathcal{O}(\Delta t^3)} = e^{\tilde{T}/2} \cdot e^{\tilde{V} + \tilde{T}/2} \cdot e^{-\frac{1}{2}[\tilde{T}/2, \tilde{V}] + \mathcal{O}(\Delta t^3)} \\ &= e^{\tilde{T}/2} \cdot e^{\tilde{V}} \cdot e^{\tilde{T}/2} \cdot e^{-\frac{1}{2}[\tilde{V}, \tilde{T}/2] + \mathcal{O}(\Delta t^3)} \cdot e^{-\frac{1}{2}[\tilde{T}/2, \tilde{V}] + \mathcal{O}(\Delta t^3)} \\ &= e^{\tilde{T}/2} \cdot e^{\tilde{V}} \cdot e^{\tilde{T}/2} \cdot e^{\mathcal{O}(\Delta t^3)} \end{aligned}$$

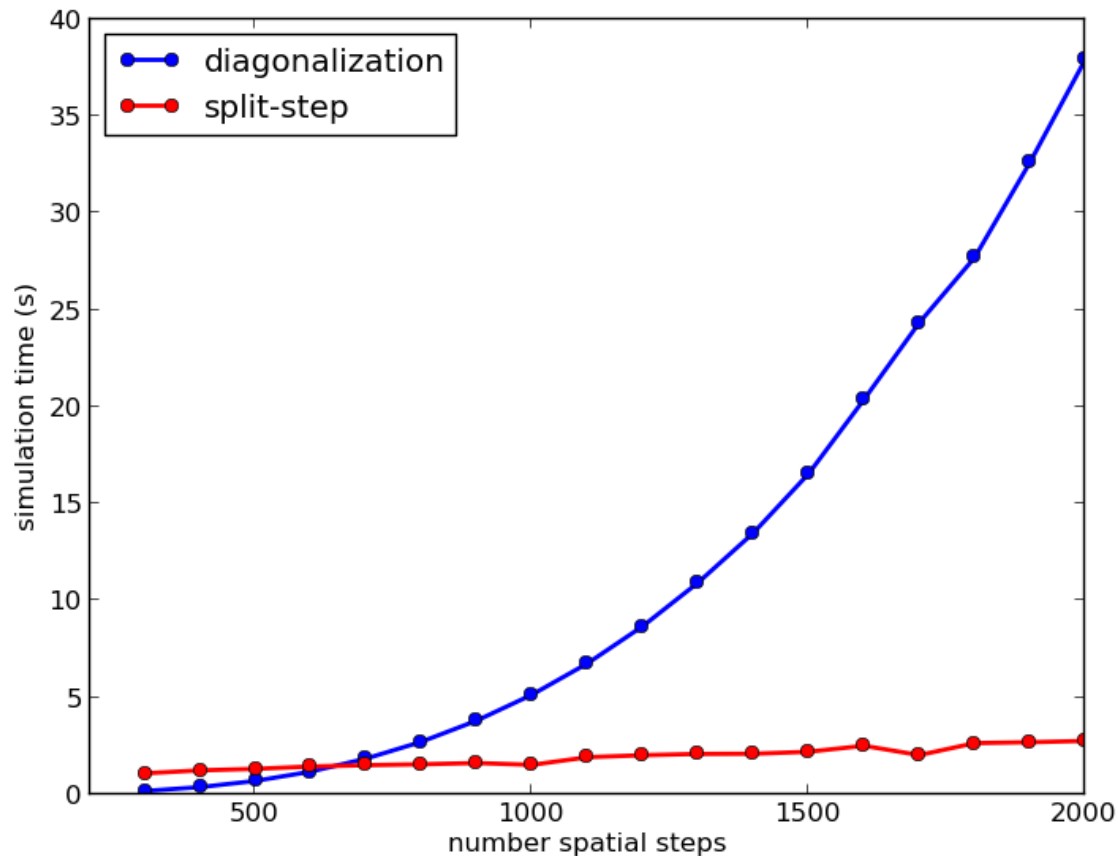
Split-Step Technique



Imaginary-time split-step Technique

```
class ssft():
    def __init__(self, lx, steps = 1001):
        self.p = parameters.tp_parameters()
        self.lx = lx
        self.steps = steps
        self.x, self.dx = scipy.linspace(-self.lx/2., self.lx/2., self.steps, endpoint=False, retstep=True)
        self.dk = 2*pi/self.lx
        self.k = 2*pi*scipy.fftpack.fftfreq(len(self.x), self.dx) ## k-grid
    def flat_psi(self):
        return self.Normalize(scipy.ones(self.x.shape))
    def rand_psi(self):
        return self.Normalize(scipy.rand(self.x.shape[0]))
    def Normalize(self, psi):
        return (1./((scipy.absolute(psi)**2).sum()*self.dx))**.5*psi
    def Normalize_k(self, psi_k):
        return (1./((scipy.absolute(psi_k)**2).sum()*self.dk))**.5*psi_k
    def ssft(self, potential, dt=-.1j*50e-6, max_its = 10000, start_psi = None, precision= 1e-8):
        self.dt = dt
        self.T = hbar**2*self.k**2/(2*self.p.m)
        self.exp_T = scipy.exp(-1j*self.dt/(2.*hbar)*self.T)
        self.exp_2T = scipy.exp(-1j*self.dt/(hbar)*self.T)
        if start_psi == None:
            psi = self.rand_psi()
        else:
            psi = self.Normalize(start_psi)
        psi_old = scipy.array(psi)
        exp_prefactor = -1j*self.dt/hbar
        psik = scipy.fftpack.fft(psi)
        psik *= self.exp_T
        psi = scipy.fftpack.ifft(psik)
        for i in scipy.arange(max_its):
            if scipy.iscomplex(self.dt):
                psi = self.Normalize(psi)
                psi *= scipy.exp(exp_prefactor*potential)
                psik = scipy.fftpack.fft(psi)
                psik *= self.exp_2T
                psi = scipy.fftpack.ifft(psik)
                if precision != None:
                    if ((scipy.absolute(psi-psi_old)/scipy.absolute(psi)).sum() < precision)*(i>50):
                        break
                else:
                    psi_old = psi
            if scipy.iscomplex(self.dt):
                psi = self.Normalize(psi)
                psi *= scipy.exp(-1j*self.dt/hbar*potential)
                psik = scipy.fftpack.fft(psi)
                psik *= self.exp_T
                psi = scipy.fftpack.ifft(psik)
        return self.Normalize(psi)
```

Split-step vs. diagonalization



Interactions

Gross-Pitaevskii equation:

$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) + \frac{4\pi\hbar^2 a}{m} |\psi(x)|^2 \right] \psi(x) = \mu \psi(x)$$

Interactions

Gross-Pitaevskii equation:

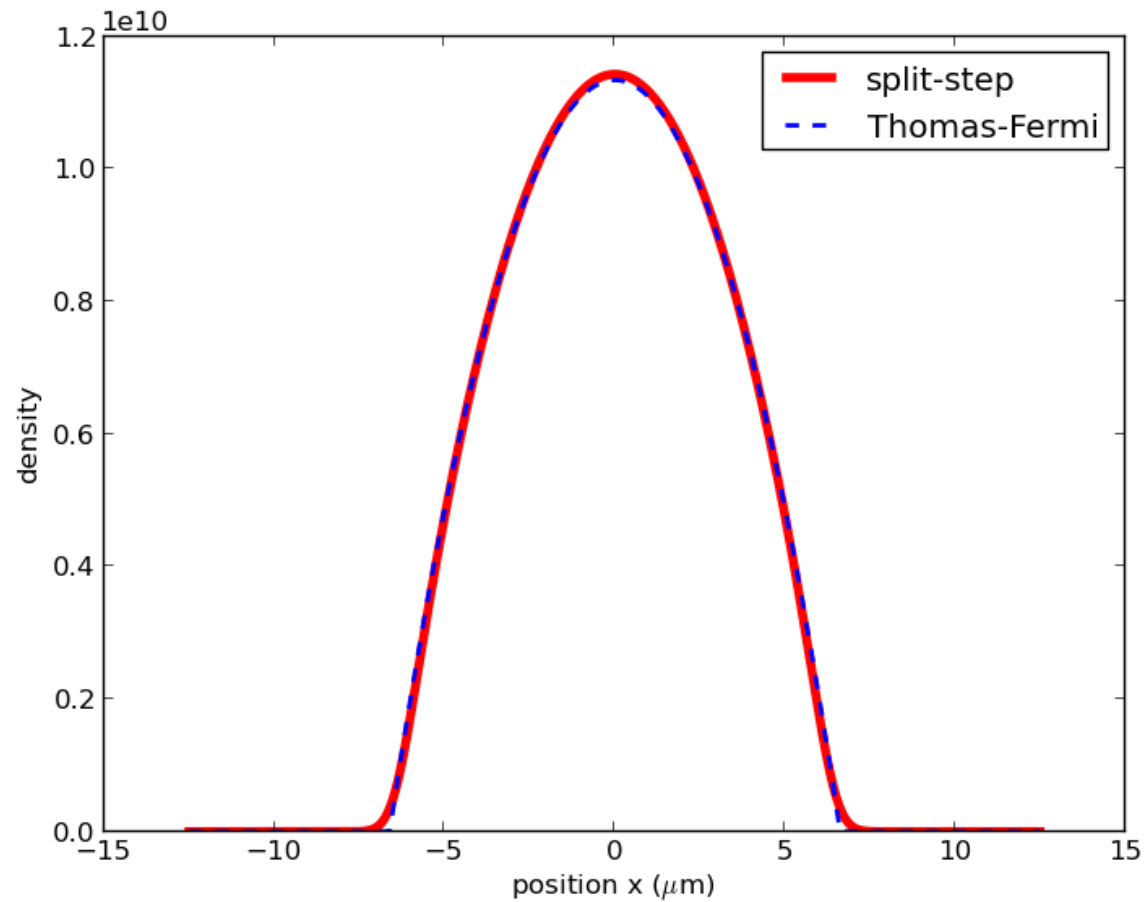
$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) + \frac{4\pi\hbar^2 a}{m} |\psi(x)|^2 \right] \psi(x) = \mu \psi(x)$$

Thomas-Fermi approximation:

$$[V(x) + U_0 |\psi(x)|^2] \psi(x) = \mu \psi(x)$$

$$n(x) = |\psi(x)|^2 = \frac{[\mu - V(x)]}{U_0}$$

Split-step



Problems, pitfalls, etc...

- Programming language

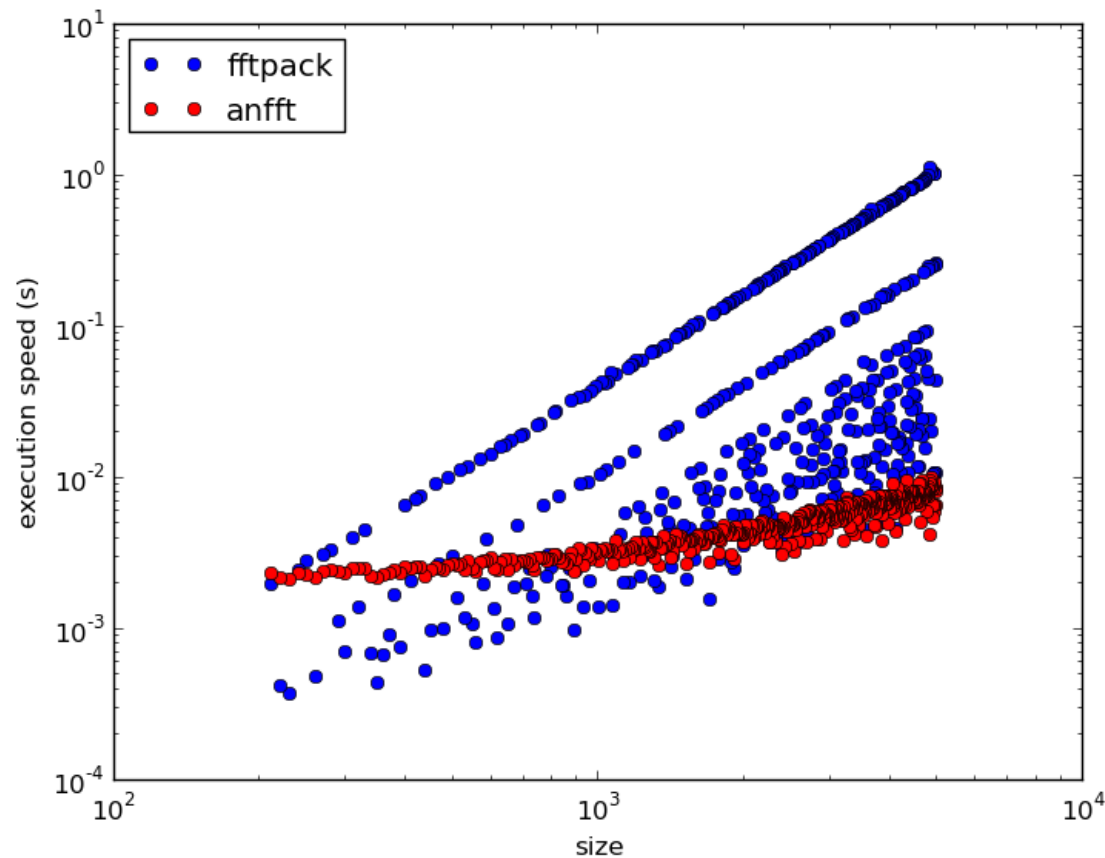
Problems, pitfalls, etc...

- Programming language
- FFT

Problems, pitfalls, etc...

- Programming language
 - FFT
 - speed
-

Comparison: FFTPACK vs. FFTW



Problems, pitfalls, etc...

- Programming language
 - FFT
 - speed
 - alignment
-

Problems, pitfalls, etc...

- Programming language
 - FFT
 - speed
 - alignment
 - scattering length
-

Scattering length in 1d & 2d

$$\frac{4\pi\hbar^2 a}{m} |\psi|^2$$

Scattering length in 1d & 2d

$$\frac{4\pi\hbar^2 a}{m} |\psi|^2$$

$$a_{1d} = \left(\frac{15Na}{\bar{a}} \right)^{2/5} \left(\frac{\hbar\bar{\omega}}{2} \right)^{3/2} \frac{\sqrt{2m}}{3\omega_x \hbar 2\pi N}$$

$$a_{2d} = \frac{5^{4/5} \left(\frac{aN}{\bar{a}} \right)^{2/5} \bar{\omega}^2}{2 \cdot 3^{1/5} N \pi \omega_x \omega_y},$$