# The De Giorgi Method Applications to Degenerate PDE

Logan F. Stokols

Department of Mathematics
The University of Texas at Austin

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#### **Outline**

- 1 Overview of the De Giorgi Method
- 2 Superquadratic Hamilton-Jacobi Equations
- 3 Hypoelliptic Fokker-Planck Equation
- 4 L<sup>2</sup>-stability of Shocks
- 5 SQG on Bounded Domains



#### Overview of De Giorgi Method

Consider the toy problem

$$\partial_t u - \operatorname{div}(A \nabla u) = 0$$

Given  $\lambda I \leq A \leq \Lambda I$  (in sense of positive matrices), parabolic, expect regularity



#### Overview of De Giorgi Method

Consider the toy problem

$$\partial_t u - \operatorname{div}(A \nabla u) = 0$$

Given  $\lambda I \leq A \leq \Lambda I$  (in sense of positive matrices), parabolic, expect regularity

In fact, 
$$\exists \alpha \in (0,1)$$
 s.t.  $\forall \varepsilon > 0$ ,  $\exists C > 0$ 

$$||u||_{C^{\alpha}([\varepsilon,\infty)\times\mathbb{R}^n)} \le C ||u(0,\cdot)||_{L^2(\mathbb{R}^n)}$$

c.f. De Giorgi ['57]



#### **Energy Inequality**

Let  $A \subseteq \mathbb{R}^n$ , [a, b] an interval,  $\varepsilon > 0$ 

Multiply by test function  $\phi(t,x)(u-k)_+$ , obtain

$$\sup_{[a,b]} \int_{A} (u-k)_{+}^{2} + \int_{a}^{b} \int_{A} |\nabla (u-k)_{+}|^{2} \lesssim \int_{a-\varepsilon}^{b} \int_{B_{\varepsilon}(A)} (u-k)_{+}^{2}$$



#### First De Giorgi Lemma

- $L^2$ -to- $L^{\infty}$  regularization
- global and local version
- proof by truncation, recursion



## First De Giorgi Lemma

- $L^2$ -to- $L^{\infty}$  regularization
- global and local version
- proof by truncation, recursion

Consider  $k \in \mathbb{N}$ 

$$Q_0 := [-2,0] \times B_2 \supseteq Q_1 \supseteq \cdots \supseteq Q_k \supseteq \cdots \supseteq [-1,0] \times B_1$$

and truncations

$$u_0 := (u - 0)_+ \ge u_1 \ge \cdots \ge u_k \ge \cdots \ge (u - 1)_+$$



#### First De Giorgi Lemma

#### Lemma

Let u solve parabolic equation, there exist  $\delta_0$  small so

$$\iint\limits_{Q_0}u_0^2\leq \delta_0$$

implies

$$\iint\limits_{\mathbb{R}^n} (u-1)_+ = 0 \qquad \equiv \qquad u \le 1 \text{ on } [-1,0] \times B_1$$



Energy inequality says

$$\sup_{Q_k} \int u_k^2 + \iint_{Q_k} |\nabla u_k|^2 \lesssim \iint_{Q_{k-1}} u_k^2$$



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Sobolev embedding says,  $\exists q > 2$ 

$$\int\int\limits_{Q_k}u_k^q\lesssim \left(\sup\limits_{Q_k}\int u_k^2+\int\int\limits_{Q_k}|
abla u_k|^2
ight)^{q/2}$$



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$$\iint\limits_{Q_k} u_k^q \lesssim \left(\sup\limits_{Q_k} \int u_k^2 + \iint\limits_{Q_k} |\nabla u_k|^2\right)^{q/2}$$

Basic Arithmetic says

$$\iint\limits_{O_k} u_k^2 \lesssim \iint\limits_{O_k} u_{k-1}^q$$



Put together,

$$\iint\limits_{Q_{k+1}}u_{k+1}^2\leq C_k\left(\iint\limits_{Q_k}u_k^2\right)^{1+\varepsilon}$$

if  $u_0$  is sufficiently small on  $Q_0$ , then  $(u-1)_+$  vanishes on  $[-1,0] \times B_1$ 



#### Second De Giorgi Lemma

- also called Isoperimetric Inequality
- quantitative version of "solutions to parabolic eqn have no jump discontinuities"
- can be non-constructive (compactness)
- proof varies significantly between applications



#### Second De Giorgi Lemma

#### Lemma

 $\exists \mu_0 > 0$  s.t., u solving parabolic equation, if

$$u \le 2$$
 on  $[-1, 4] \times B_3$ ,

and 
$$|\{u \ge 1\} \cap [2,4] \times B_2| \ge \delta_0,$$
  
and  $|\{u \le 0\} \cap [0,4] \times B_2| \ge \frac{1}{2} |[0,4] \times B_2|$   
then  $|\{0 < u < 1\} \cap [0,4] \times B_2| \ge \mu_0.$ 



#### Proof of Second De Giorgi Lemma

Assume false for all  $\mu_0$ , take sequence  $u_k$  of counterexamples  $\mu_0 = 1/k$ 

By first condition and energy inequality,  $u_k$  compact, has  $L_t^2(H_x^1)$  limit  $u_{\infty}$ 

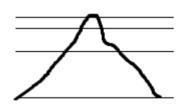
a.e. t fixed,  $\forall x$  either  $u_{\infty} \geq 1$  or  $u_{\infty} \leq 0$ 

 $u_{\infty}$  is  $\leq 0$  on most of [0,4] but  $\geq 1$  on some of [2,4], so energy increases suddenly in time, impossible



#### Hölder Continuity

Proof of Hölder continuity uses recursion apply two De Giorgi lemmas to rescalings of solution







## Superquadratic Hamilton-Jacobi Equation



## Superquadratic Hamilton-Jacobi Equation

e.g. 
$$\partial_t u + |\nabla u|^p - \varepsilon \Delta u = 0, \qquad \varepsilon \in \{0, 1, -1\}$$

- First considered by Lasry and Lions ['89]; Schwab ['13] (homogenization)
- Strongest known results Cardaliaguet ['09], Cannarsa and Cardaliaguet ['10], Cardaliaguet and Silvestre ['12]
- $\varepsilon = 0$  singularities can form, but solutions always continuous
- For p>2, continuous even for  $\varepsilon=-1$  [first order drives regularization]
- Chan and Vasseur ['17] use De Giorgi for  $\varepsilon=0$



#### Superquadratic Hamilton-Jacobi Equation

$$\partial_t u = H\left(t, x, u, \nabla u, D^2 u\right),$$

$$\Lambda^{-1} |\nabla u|^p - \operatorname{div}(A\nabla u) - f \le H\left(t, x, u, \nabla u, D^2 u\right) \le \Lambda |\nabla u|^p - \Lambda m^-(D^2 u) + \Lambda$$

p>2, A bounded unsigned matrix,  $f\in L^q,$   $m^-$  returns lowest negative eigenvalue

#### Theorem (S., Vasseur [CMS, '18])

Solutions (in appropriate weak sense) regularize from  $L^{\infty}(\mathbb{R}^+ \times \mathbb{R}^n)$  into  $C^{\alpha}([\varepsilon, \infty) \times \mathbb{R}^n)$ 



- De Giorgi method
- Adapted technique of Chan, Vasseur, overcome second-order term
- Combine divergence-form and non-divergence-form techniques
- Allow discontinuous A





Consider

$$\partial_t u + |\nabla u|^p + \Delta u = 0.$$



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$$\partial_t u + |\nabla u|^p + \Delta u = 0.$$

Using  $\varphi(t,x)(u-k)_+$  as test function, obtain energy inequality

$$\sup_{[-1,0]} \int_{B_1} (u-k)_+^2 + \int_{-1}^0 \int_{B_1} (u-k)_+ |\nabla (u-k)_+|^p$$

$$\lesssim \int_{-2}^{0} \int_{B_2} (u - k)_+^2 + \int_{-2}^{0} \int_{B_2} |\nabla (u - k)_+|^2$$



Consider

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$$\lesssim \int_{-2}^{0} \int_{B_2} (u - k)_+^2 + \int_{-2}^{0} \int_{B_2} |\nabla (u - k)_+|^2$$

Want to show, on  $Q_1 = [-1, 0] \times B_1$ , for some q > 2

$$\|(u-1)_+\|_{L^q(Q_1)} \lesssim \|(u)_+\|_{L^\infty_t(L^2_x)(Q_1)} + \|(u)_+|\nabla(u)_+|^p\|_{L^1(Q_1)}.$$



Coercivity when u large



Coercivity when u large

Strategy: consider two regions, *u* small and *u* big

- $u \text{ small} \Rightarrow L^q \text{ norm small}$
- u big  $\Rightarrow$  coercivity  $\Rightarrow L^q$  norm small



#### Coercivity when u large

Strategy: consider two regions, u small and u big

- $u \text{ small} \Rightarrow L^q \text{ norm small}$
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#### Implementation:

- $\|\nabla(u-1)_+\|_{L^p}^p \le \|u_+|\nabla u_+|^p\|_{L^1}$
- $\|(u-1)_+\|_{L^q} \lesssim \|(u-1)_+\|_{L^\infty(L^2)} + \|\nabla(u-1)_+\|_{L^p}$



# Hypoelliptic Fokker-Planck



## Hypoelliptic Fokker-Planck

e.g. 
$$\left[\partial_t + v \cdot \nabla_x\right] f + \left(-\Delta_v\right)^s f = 0$$

- Rarefied gas, neutral particles in plasma
- Imbert and Silvestre ['16]; Golse and Imbert and Mouhot and Vasseur ['16]
- Hypoelliptic: non-elliptic regularization, mixed elliptic/hyperbolic type
- Averaging Lemma (Golse et al '88): H<sup>s</sup> theory of hypoellipticity, regularity of averages for kinetic equation



#### Hypoelliptic Fokker-Planck

$$[\partial_t + v \cdot \nabla_x] f = \int K[f(w) - f(v)] dw + \sigma$$

 $K \approx |v-w|^{-(n+2s)}$ ,  $s \in (0,1)$ , K symmetric in  $(v,w) \mapsto (w,v)$  and in  $(v,v+y) \mapsto (v,v-y)$ 

#### Theorem (S. [SIMA, '19])

For f solution,  $f \in L^{\infty} \cap L^2_{t,x}(H^s_{\nu})$ ,  $\sigma \in L^2 \cap L^r$  for r >> 1, there exists  $\alpha \in (0,1)$  depending on kernel, C > 0 depending on domain and kernel s.t.

$$||f||_{C^{\alpha}([\varepsilon,\infty)\times\mathbb{R}^n\times B_1)}\leq C\left(||f||_{L^{\infty}}+||\sigma||_{L^r}\right).$$





Consider 
$$\Lambda = (-\Delta_{\nu})^{1/2}$$
,  $s \in (0, 1)$ ,

$$\partial_t f + v \cdot \nabla_x f + \Lambda^{2s} f = 0.$$

Note diffusion in v but not x! Energy inequality will have

$$||(f-\psi)_+||_{L^{\infty}_t(L^2_{x,v})} + ||\Lambda^s(f-\psi)_+||_{L^2_{t,x,v}}$$

on LHS



#### **Averaging Lemma**

Lemma (Bézard, '94): for  $\alpha = 1/(2(1+m))$ ,  $\Omega \subseteq \bar{\Omega} \subseteq \mathbb{R}^+ \times \mathbb{R}^n$ , and  $f, g \in L^2(\bar{\Omega} \times \mathbb{R}^n)$ , f compactly supported, we have

$$[\partial_t + v \cdot \nabla_x]f = g$$

implies

$$\left\| \int f \, dv \right\|_{H^{\alpha}(\Omega)} \lesssim \|f\|_{L^{2}(\bar{\Omega} \times \mathbb{R}^{n})} + \left\| (1 - \Delta_{\nu})^{-m/2} g \right\|_{L^{2}(\bar{\Omega} \times \mathbb{R}^{n})}.$$



Unfortunately: Can't apply lemma to  $(f - \psi)_+$  due to truncation (nonlocal)

Barrier function

$$0 \le \varphi(t, x)(f - \psi)_{+} \le F,$$

$$\|F\|_{L^{2}} + \left\| (1 - \Delta_{\nu})^{-m/2} \left[ \partial_{t} + \nu \cdot \nabla_{x} \right] F \right\|_{L^{2}} \le C \|\varphi(f - \psi)_{+}\|_{L^{2}}.$$

Now:

$$\left\| \int F \, dv \right\|_{H^{\alpha}} \le \left\| \varphi(f - \psi)_{+} \right\|_{L^{2}},$$

No regularity on f!



#### From averaging lemma:

$$\|(f-\psi)_+\|_{L^{2+\varepsilon}_{t,x}(L^1_{\nu})} \lesssim \|(f-\psi)_+\|_{L^2_{t,x,\nu}}$$



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$$||(f-\psi)_+||_{L^{2+\varepsilon}_{t,x}(L^1_{\nu})} \lesssim ||(f-\psi)_+||_{L^2_{t,x,\nu}}$$

From energy inequality:

$$\|(f-\psi)_+\|_{L^{\infty}_t(L^2_{x,v})} + \|(f-\psi)_+\|_{L^2_{t,x}(L^{2+\varepsilon}_v)} \lesssim \|(f-\psi)_+\|_{L^2_{t,x,v}} + \|(f-\psi)_+\|_{L^1_{t,x,v}}$$



## Hypoelliptic Fokker-Planck: Proof excerpt

From averaging lemma:

$$||(f-\psi)_+||_{L^{2+\varepsilon}_{t,x}(L^1_{\nu})} \lesssim ||(f-\psi)_+||_{L^2_{t,x,\nu}}$$

From energy inequality:

$$\|(f-\psi)_+\|_{L^{\infty}_{t}(L^2_{x,\nu})} + \|(f-\psi)_+\|_{L^2_{t,x}(L^{2+\varepsilon}_{\nu})} \lesssim \|(f-\psi)_+\|_{L^2_{t,x,\nu}} + \|(f-\psi)_+\|_{L^1_{t,x,\nu}}$$

Improvement in all three variables, for some q > 2

$$\|(f-\psi)\|_{L^q_{t,x,y}} \lesssim \|(f-\psi)_+\|_{L^2_{t,x,y}}$$



## Hypoelliptic Fokker-Planck: Proof excerpt

From averaging lemma:

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From energy inequality:

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Improvement in all three variables, for some q > 2

$$\|(f-\psi)\|_{L^q_{t,x,y}} \lesssim \|(f-\psi)_+\|_{L^2_{t,x,y}}$$

Control of regularity is means to an end, control of integrability is the end



# $L^2$ stability of Shocks



## $L^2$ stability of Shocks

Consider 1D scalar viscous conservation law

$$\partial_t u + \partial_x [Q(u)] = \varepsilon \partial_{xx} \eta'(u)$$

where  $\eta$ , Q uniformly convex and  $\varepsilon > 0$  arbitrary.

### Theorem (S. [Submitted])

For u a solution and s a sufficiently small shock solution, there exists Lipschitz  $\gamma(t)$  such that

$$\|u(\cdot,t)-s(\cdot-\gamma(t))\|_2$$

stable in time, up to a constant factor. Result is independent of  $\varepsilon$ .

c.f. Kang ['19], Kang & Vasseur ['19]



### **SQG** in Bounded Domains



#### **SQG** in Bounded Domains

e.g. 
$$\partial_t \theta + \left( \nabla^{\perp} (-\Delta)^{-1/2} \theta \right) \cdot \nabla \theta + \nu (-\Delta)^s \theta = 0$$

- atmospheric or ocean currents, used in weather modelling
- R<sup>2</sup>: Constantin, Majda, Tabak ['93]; Kiselev, Nazarov, Volberg ['08]; Caffarelli, Vasseur ['10]; Constantin, Vicol ['12]
- Bounded domain: Kriventsov ['15]; Novack, Vasseur ['18,'19]
- Best studied model by Constantin, Ignatova ['16]; Constantin, Ignatova, Nguyen [various]
- Boundary issues: Laplacian & gradient don't commute, Caffarelli-Stinga
   ['16] kernel representation degenerates

#### **SQG** in Bounded Domains

$$\begin{cases} \partial_t \theta + u \cdot \nabla \theta + \Lambda \theta = 0, \\ u = \nabla^{\perp} \Lambda^{-1} \theta. \end{cases}$$

 $\Omega \subseteq \mathbb{R}^2$  smooth bounded open,  $\Lambda := \sqrt{-\Delta_D}$  (defined spectrally),  $\Delta_D$  the Dirichlet Laplacian on  $\Omega$ 

### Theorem (S., Vasseur [ARMA, '20])

Let  $\Omega \subseteq \mathbb{R}^2$  a bounded set, initial data  $\theta_0 \in L^2(\Omega)$ There exists a global-in-time solution  $\theta$  to SQG such that: For any  $\varepsilon > 0$ , there exists  $\alpha \in (0,1)$  and a constant C so

$$\|\theta\|_{C^{\alpha}([\varepsilon,\infty)\times\Omega)} \leq C \|\theta_0\|_{L^2(\Omega)}$$
.

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## SQG in Bounded Domains: Proof excerpt



### SQG in Bounded Domains: Proof excerpt

Velocity u, energy inequality with cutoff  $\Psi$  has drift term on RHS

$$\int_{\Omega} u(\theta - \Psi)_{+} \cdot \mathrm{d}\Psi$$

recall u is Riesz transform of  $\theta \in L^{\infty}$ 

u bounded in:	$L^{\infty}$	BMO	$B^0_{\infty,\infty}$
$\theta \in L^{\infty} \Rightarrow u \in \underline{\hspace{1cm}}$	×	works on $\mathbb{R}^2$	complicated
$\int u \theta_+$ bounded	$\leq \int \theta_+$	John-Nirenberg	complicated
scaling invariant	<b>√</b>	✓	<b>√</b>



### Control on u: Littlewood Paley theory

Littlewood paley operators  $P_j = P_j(\Lambda)$ , functional calculus for  $\Lambda$ ,

Bernstein Inequalities

$$\|\Lambda^{s} P_{j} f\|_{p} \approx 2^{sj} \|f\|_{p},$$
  
$$\|\nabla \Lambda^{s} P_{j} f\|_{p} \approx 2^{(1+s)j} \|f\|_{p}.$$

Commutation Relation

$$||P_i \nabla P_j f||_p \lesssim \min(2^j, 2^i) ||f||_p$$
.

Bernstein: Iwabuchi, Matsuyama, Taniguchi ("Bilinear estimates in Besov spaces generated by the Dirichlet Laplacian" 2017)

### Control on u: high and low frequencies

Instead of Besov norm  $\sup_{j} \|P_{j} \nabla \Lambda^{-1} \theta\|_{\infty}$  consider

$$\nabla \Lambda^{-1} P_j \theta$$
.

To bound  $\int u\theta_+$ , decompose as

$$u_{\text{low}} = \sum_{j=-\infty}^{0} \nabla \Lambda^{-1} P_j \theta$$

which is Lipschitz,

$$u_{\text{high}} = \sum_{j=0}^{\infty} \nabla \Lambda^{-1} P_j \theta$$

which is in  $W^{-\varepsilon,\infty}$ .



# Thank you

