

## Renormalization and Scale Setting

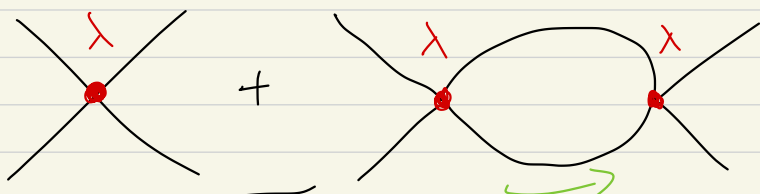
Renormalization/RG: bare parameters  $\rightarrow$  physics!

$N$  couplings  $\rightarrow N$  experiments to fix theory.

"Scale setting": choice of which inputs to use.

Continuum ex.:  $\phi^4$  theory, one-loop.

Scattering:  $\phi\phi \rightarrow \phi\phi$

$$i\mathcal{M} = \text{tree} + \text{one-loop}$$

$$\sim \int^{\Lambda} d^4k \frac{1}{k^2} \cdot \frac{1}{(p+k)^2}$$

Log divergent:  $-\mathcal{M} = \lambda + \frac{\lambda^2}{32\pi^2} \log \frac{s}{\Lambda^2}$

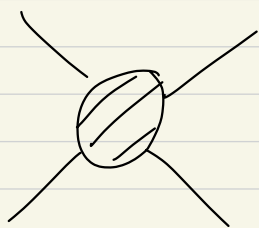
$\Lambda \rightarrow \infty$ ;  $\infty$  amplitude! Is the theory 'sick'?

( $s$ : Mandelstam  $E_{cm}^2$ )

Theory is okay! Only cross-section differences.

$$\mathcal{M}(s_1) - \mathcal{M}(s_2) \sim \left( \log \frac{s_1}{\Lambda^2} - \log \frac{s_2}{\Lambda^2} \right) = \log \frac{s_1}{s_2}$$

Define renormalized coupling at ref. energy  $s_0$ .



$$= -i\lambda_R \Rightarrow \lambda_R = \lambda + \frac{\lambda^2}{32\pi^2} \log \frac{s_0}{\Lambda^2} + \dots$$

Scattering in terms of  $\lambda_R$  is finite:

$$\mathcal{M}(s) = -\lambda_R - \frac{\lambda_R^2}{32\pi^2} \log \frac{s}{s_0} + \dots$$

Definition of  $\lambda_R$  is our renormalization condition.

NO divergences!\*

\* (except as  $s_0 \rightarrow \infty$ . No continuum limit on lattice!)

If we change  $s_0 \rightarrow s_0'$ ,  $\lambda_R$  changes too.

calculate dependence  $\rightarrow$  renormalization group.

$$\frac{d\lambda_R}{ds_0} = \frac{1}{s_0} \frac{\lambda^2}{32\pi^2}$$

Define the beta function:

$$\beta(\lambda_R) \equiv s_0 \frac{d\lambda_R}{ds_0} = +(\dots) \lambda_R^2$$



Arrows show UV  
flow (to  $s_0 \rightarrow \infty$ .)

Opposite is IR  
flow (to  $s_0 \rightarrow 0$ .)

( $\lambda_R \rightarrow 0$ .)

"Continuum limit"  $s_0 \rightarrow \infty$  flows into  $\lambda_R \rightarrow \infty$   
(perturbative control lost.)

Alternatively, hold  $\lambda_R(s_0)$  fixed, look at  
 $\lambda_R(s_1)$  w/  $s_1' < s_0$ . then  $\lambda_R(s_1) \rightarrow 0$  | "triviality"

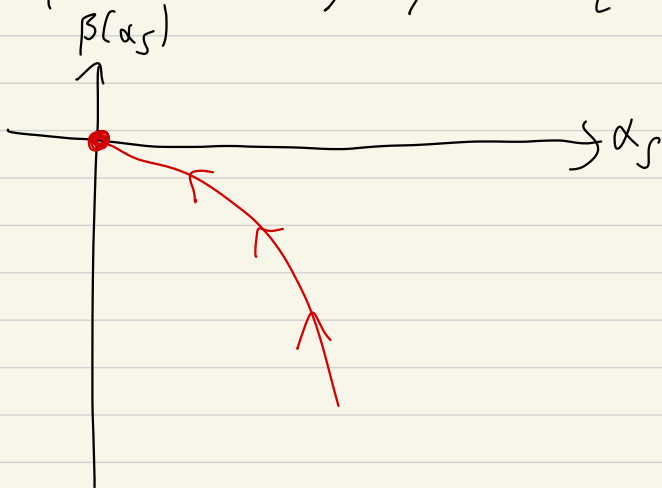
gauge coupling

For QCD first term in  $\beta(\alpha_s)$  ( $\alpha_s \equiv g_s^2/4\pi$ )  
has crucial minus sign.

(Gross, Wilczek '73; Politzer '73)

$$\Lambda \frac{d\alpha_s}{d\Lambda} \equiv \beta(\alpha_s) = -\frac{1}{2\pi} \left(11 - \frac{2}{3} N_f\right) \alpha_s^2 + \dots$$

( $N_f = \#$  of light, active quarks.)



(UV flow:  
 $\Lambda \rightarrow \infty$ )

asymptotic freedom

( $\Lambda \rightarrow \infty$  is  $\alpha_s \rightarrow 0$ .)

In QCD,  $\alpha_s$  is inconvenient to renormalize  
the theory (we can't scatter quarks/gluons.)

Instead, rely on dimensional transmutation

$$\alpha_s(\Lambda_0) = \alpha_{\text{ref}} \longleftrightarrow M_H = c_{\text{ref}} \Lambda_0$$

(roughly.)

(R. Sommer, 1401.3270)

Scale setting - focus on QCD.

Gauge coupling  $\alpha_s$  + quark masses  $m_u, m_d, m_s, m_c, m_b$ .

(Top quark is too heavy/unstable;  $\Theta = 0$ .)

Six experimental inputs (w/dependence on parameters;  
proton mass  $\nrightarrow m_b$ .)

can be any six processes ( $\pi\pi \rightarrow \pi\pi$  scattering,?)

We want highest precision  $\rightarrow$  hadron masses  
 $\rightarrow$  decay constants

Quark masses  $\leftrightarrow$  direct meson ( $\bar{q}q$ ) masses.

Heavy quarks:  $\begin{cases} M_{D_s} \sim m_c & (\bar{c}s) \\ M_{B_s} \sim m_b & (\bar{b}s) \end{cases}$

Light quarks: chiral perturbation theory  
(chiral symmetry restored at  $m_q = 0 \dots$ )

Implies the Gell-Mann-Oakes-Renner relation:

(GMOR):  $M_{q_1 q_2}^2 = B(m_{q_1} + m_{q_2})$   
(pseudoscalar)

works well for light, strange. Pions, kaons!

$$\begin{cases} \pi^+ \sim \bar{d}u & K^+ \sim \bar{s}u \\ \pi^0 \sim (\bar{u}u + \bar{d}d) & K^0 \sim \bar{s}d \end{cases}$$

(Flavor Lattice Averaging Group)

FLAG 2024  $\rightarrow$  "Edinburgh consensus" (no QED)

$$M_{\pi^+} = 135.0 \text{ MeV}$$

$$M_{K^+} = 491.6 \text{ MeV}$$

$$M_{K^0} = 497.6 \text{ MeV}$$

$$M_{D_s^+} = 1967 \text{ MeV}$$

$$M_{B_s^0} = 5367 \text{ MeV}$$

One more input to fix overall scale.

Simulation produces  $aM_{\pi^+}$ ; need to set  $a$  to fix  $M_{\pi^+} = 135 \text{ MeV}$ !

Three classes of scale setting:

A) Hadron masses;

B) Decay constants;

C) Theory scales.

Precision is essential!  $S$  to set scale.

$(aM_H)$  from calculation, precision  $\delta(aM_H)$ .

Then!  $M_H = \frac{aM_H}{aS} \cdot S_{\text{exp}}$  Frac error!

$$\frac{\delta M_H}{M_H} = \sqrt{\underbrace{\left(\frac{\delta(aM_H)}{aM_H}\right)^2}_{\text{lattice } aM_H \text{ error}} + \underbrace{\left(\frac{\delta(aS)}{aS}\right)^2}_{\text{lattice } aS \text{ error}} + \underbrace{\left(\frac{\delta S_{\text{exp}}}{S_{\text{exp}}}\right)^2}_{\text{experimental.}}}$$



(<https://pdg.lbl.gov/>)

$\delta S_{\text{exp}}/S_{\text{exp}}$  provides ultimate limit on error!

$d(aS)/aS$  also important.

A) Hadron masses - what are good options?

- Proton:  $\frac{\delta m_p^{\text{exp}}}{m_p} \sim 3 \times 10^{-8}$

But: exponential signal-to-noise degradation  
(see exercise)

-  $\pi, K$ : very precise, but we used them!

-  $n$ : also very precise, hard on lattice  
(disconnected diagrams...)

$\Omega$  baryon:  $\frac{\delta m_{\Omega}^{\text{exp}}}{m_{\Omega}^{\text{exp}}} \sim 0.02\%$   
(555)

Also exp. S/N problem; better than proton.

Lattice calcs  $\sim$  sub-percent  $\Omega M_{\Omega}$ .

B) Decay constants - more details next time.

$$\Gamma(\pi \rightarrow \mu \nu) \sim f_\pi^2; \quad \frac{\delta \Gamma^{\text{exp}}}{\Gamma^{\text{exp}}} \sim 10^{-2}$$

But disentangle EW physics. Systematics?  
( $f_\pi$  still commonly used.)

C) Theory scales:  $S_{\text{exp}}$  vs.  $\alpha S$  precision;

(systematic effects,  $N/S$ , or expensive.)

Theory scale  $1/T$  can be a proxy.

Easy to compute  $\alpha/T$  precisely; relate to  
some physical  $S_{\text{exp}}$  once.

$$\lim_{\alpha \rightarrow 0} \left( \frac{\alpha S}{\alpha/T} \right) = \boxed{TS}; \quad S_{\text{exp}} \rightarrow T_{\text{ref}}.$$

Future calculations use  $TM_H$  vs.  $\left( \frac{\alpha M_H}{\alpha S} \right) S_{\text{exp}}$ .

Static potential :  $r_0, r_1$

Gradient flow :  $\sqrt{t_0}, w_0$