

1 Question 1

To shift continuous data with a convolution, a shifted delta function kernel is used (where the delta is the kernel of the identity convolution, and a shift in the kernel produces a shift in the convolution result.) For the discrete case, consider these discrete convolutions:

$$\begin{aligned}\sum_{x'=0}^N f[x']g_0[x-x'] &= f[x]g_0[y] = f[x] \\ \sum_{x'=0}^N f[x']g_1[x-x'] &= f[x+1]g_1[y+1] = f[x+1] \\ &\vdots \\ \sum_{x'=0}^N f[x']g_n[x-x'] &= f[x+n]g_n[y+n] = f[x+n]\end{aligned}$$

where y is the only index of the array g_0 where its value is 1 and not 0. I.e. the convolution of f with g_0 at index x 'picks out' the value of f at index x . Therefore this g_0 vector is the discrete convolution identity kernel. Similarly, g_1 and g_n are the analogues of the shifted delta function kernels, where the non-zero indices are $y+1$ and $y+n$ respectively.

For odd N , the index where vector g_0 is non-zero is just the middle index. For even N , this index was found to be the greater of the two middle indices. Therefore it is easy to produce any g_n by taking g_0 and shifting its non-zero index by n .

Here is the code for taking the convolution via FFT and then for shifting an array using this FFT convolution with the g_n vector described above:

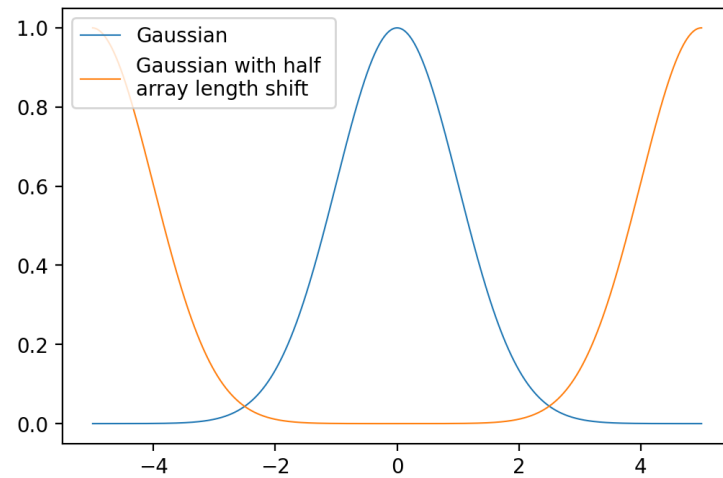
```
def convolve(f,g):
    f_ft=np.fft.fft(f)
    g_ft=np.fft.fft(g)
    conv=np.fft.ifft(f_ft*g_ft)
    return np.fft.fftshift(conv) # FFT shift for negative frequencies

def conv_shift(x,n):
    size=len(x)
    delta_shift=np.zeros(size)
    delta_shift[int(size/2)+n]=1 # int(size/2) index works for even or odd size
    return convolve(x,delta_shift)

N=2000
x=np.linspace(-5,5,N)
y=gaussian(x,0,1)
y_shift=conv_shift(y,N//2-1)
plt.plot(x,y)
plt.plot(x,y_shift)
```

Since 0 is the first index, for even N , `int(size/2)` gives the greater of the two middle indices. For odd N , `int(size/2)` rounds down, giving the middle index.

Here is the result of a shifted Gaussian (next page):



2 Question 2

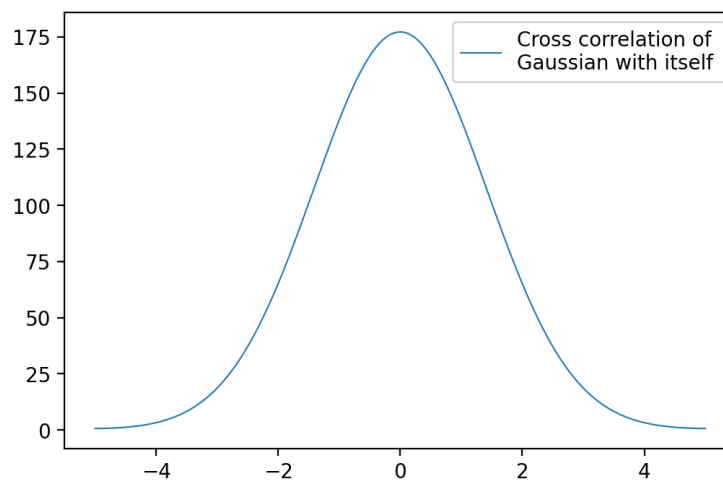
a)

To calculate the correlation of two arrays, the following function was written:

```
def correlation(f,g):
    f_ft=np.fft.fft(f)
    g_ft=np.fft.fft(g)
    corr=np.fft.ifft(f_ft*np.conj(g_ft))
    return np.fft.fftshift(corr)
```

The correlation of a Gaussian with itself was calculated and plotted:

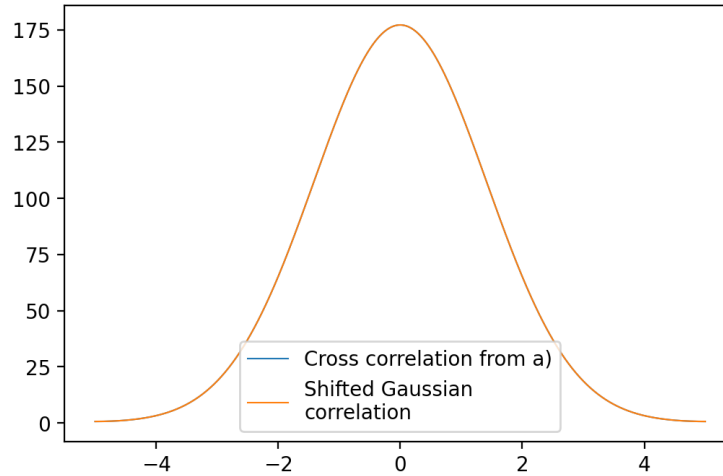
```
x=np.linspace(-5,5,1001)
y=gaussian(x,0,1)
corr=correlation(y,y)
plt.plot(x,corr)
plt.legend(['Cross correlation of Gaussian with itself'])\
```



b)

Now to calculate and plot the correlation of a shifted Gaussian with itself and compare to the result from a):

```
y_shift=conv_shift(y,100)
shift_corr=correlation(y_shift,y_shift)
plt.plot(x,corr)
plt.plot(x,shift_corr)
```



There is no difference if both arrays being correlated are shifted by the same amount. If we look at the convolution formula:

$$(f * g)(x) = \int_{-\infty}^{+\infty} f(y)g(x+y)dy$$

The convolution of the two functions $f'(x) = f(x + \delta)$ and $g'(x) = g(x + \delta)$ is:

$$(f' * g')(x) = \int_{-\infty}^{+\infty} f(y + \delta)g(x + y + \delta)dy$$

We can change the integration variable to $y' = y + \delta$, and since y ranges over all values, the integration bounds are unchanged, giving:

$$\begin{aligned}(f' * g')(x) &= \int_{-\infty}^{+\infty} f(y')g(x + y')dy' \\ &\Rightarrow f' * g' = f * g\end{aligned}$$

Therefore it is not a surprise that cross correlation between the shifted Gaussian and itself is the same as in part a).

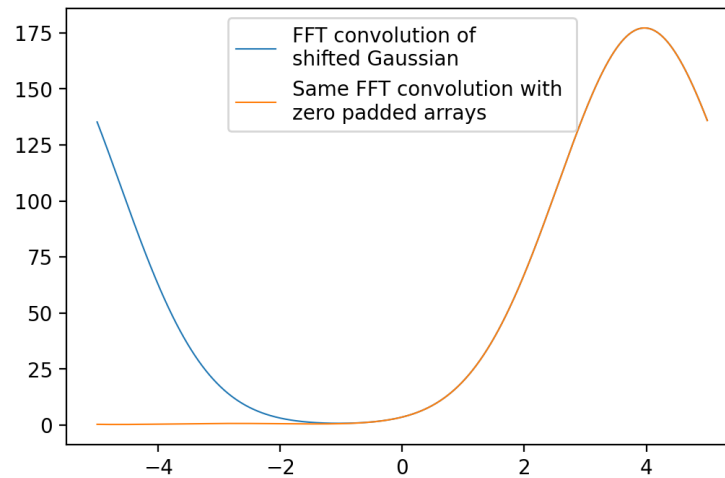
3 Question 3

To avoid the error at the edges of convolution due to the circulant nature of the DFT, the input arrays are first zero-padded on both ends. Since the arrays are now 0 at the end, the FFT produces less error at the endpoints since they are no longer forced to have the same value. After the padded arrays are convolved, the ends of the result array are removed in accordance with the number of zeros padding the input arrays. Here is the code:

```
def padded_conv(f,g,npad):
    pad=np.zeros(npad)
    fpad=np.append(pad,f)
    fpad=np.append(fpad,pad)
    gpad=np.append(pad,g)
    gpad=np.append(gpad,pad)
    conv_pad=convolve(fpad,gpad)
    return conv_pad[npad:-npad]
```

A shifted Gaussian was then convolved with itself using this function and compared to the non-padded result:

```
x=np.linspace(-5,5,1001)
y=gaussian(x,0,1)
yshift=conv_shift(y,200)
window=np.cos(np.pi*x/(x[-1]-x[0]))**2
conv=convolve(yshift,yshift)
conv_pad=padded_conv(yshift,yshift,200)
plt.plot(x,conv)
plt.plot(x,conv_pad)
```



The circulant nature of original FFT convolution is apparent on the left, while the zero-padded result does not have this error.

4 Question 4

a)

The partial sum for the geometric series is:

$$\sum_{i=0}^{N-1} \alpha^x = \frac{1 - \alpha^N}{1 - \alpha}$$

Using $\alpha = \exp(-2\pi ik/N)$:

$$\sum_{i=0}^{N-1} \exp(-2\pi ik/N)^x = \frac{1 - \exp(-2\pi ik/N)^N}{1 - \exp(-2\pi ik/N)} = \frac{1 - \exp(-2\pi ik)}{1 - \exp(-2\pi ik/N)}$$

b)

Taking the limit of the partial sum above:

$$\begin{aligned}\lim_{k \rightarrow 0} \frac{1 - \exp(-2\pi i k)}{1 - \exp(-2\pi i k/N)} &=_{LHR} \lim_{k \rightarrow 0} \frac{\frac{d}{dk}[1 - \exp(-2\pi i k)]}{\frac{d}{dk}[1 - \exp(-2\pi i k/N)]} = \lim_{k \rightarrow 0} \frac{2\pi i \exp(-2\pi i k)}{(2\pi i/N) \exp(-2\pi i k/N)} \\ &= \frac{2\pi i}{2\pi i/N} = N\end{aligned}$$

If k is an integer, then:

$$1 - \exp(-2\pi i k) = 0$$

Also, if k is not a multiple of N , k/N is not an integer and therefore:

$$\begin{aligned}1 - \exp(-2\pi i k/N) &\neq 0 \\ \Rightarrow \frac{1 - \exp(-2\pi i k)}{1 - \exp(-2\pi i k/N)} &= 0\end{aligned}$$

c)

Writing the analytic sum for the DFT of a sine wave with frequency k_0 :

$$F(k) = \sum_{x=0}^{N-1} f(x) \exp(-2\pi i k x) = \sum_{x=0}^{N-1} \exp(2\pi i k_0 x) \exp(-2\pi i k x) = \sum_{x=0}^{N-1} \exp[-2\pi i (k - k_0) x]$$

Using the result from a):

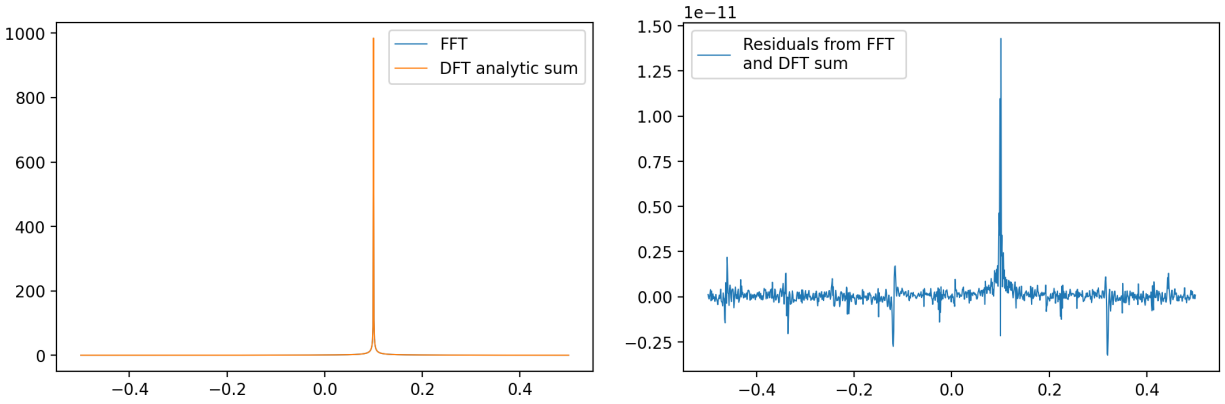
$$F(k) = \sum_{x=0}^{N-1} \exp[-2\pi i N(k - k_0)x/N] = \frac{1 - \exp[-2\pi i N(k - k_0)]}{1 - \exp[-2\pi i N(k - k_0)/N]} = \frac{1 - \exp[2\pi i N(k_0 - k)]}{1 - \exp[2\pi i (k_0 - k)]}$$

Using this code to compare the FFT vs. the DFT sum:

```
def dft_sum(k, k0, N):
    a = 1 - np.exp(-2J*np.pi*N*(k-k0))
    b = 1 - np.exp(-2J*np.pi*(k-k0))
    return a/b

k0 = 0.1
N = 1001
x = np.linspace(0, N-1, N)
y = np.exp(2J*np.pi*k0*x)
ft = np.fft.fft(y)
k = np.fft.fftfreq(N, np.abs(x[1]-x[0]))
plt.plot(np.fft.fftshift(k), np.fft.fftshift(np.abs(ft))) # fftshift to plot neg. freqs
plt.plot(np.fft.fftshift(k), np.fft.fftshift(np.abs(dft_sum(k, k0, N))))
```

The plots and the residuals are on the next page:



As seen in the plots above, the FFT and DFT sum are in agreement, with an error which is essentially just machine precision.

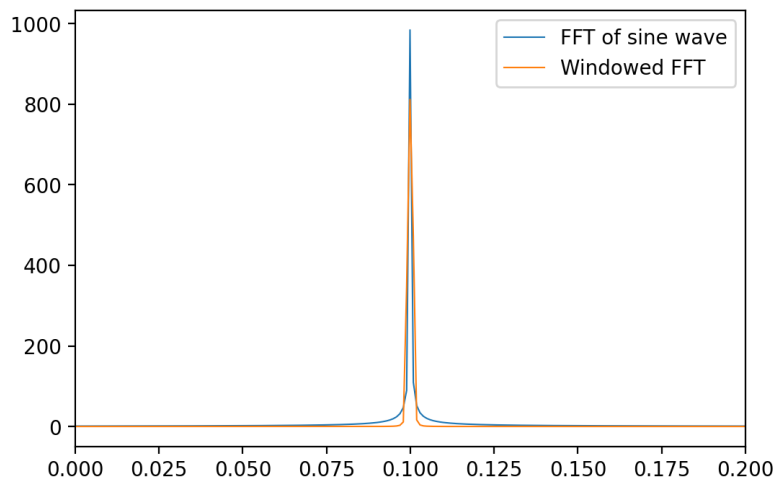
We can also see in the plot on the right the phenomenon of spectral leakage. We expect a delta function when Fourier transforming a pure frequency sine wave, however since we are taking a discrete Fourier transform, we expect a non-perfect delta spike with some finite width. In this case, the spike has a small width but is otherwise quite close to a delta function.

d)

Here is the code used for the windowed Fourier transform of a pure frequency sine wave and the resulting plot:

```
def window_fun(x,N):
    return 0.5-0.5*np.cos(2*np.pi*x/N)

ft_wndw=np.abs(np.fft.fft(y*window_fun(x,N)))
norm=np.sqrt(np.mean(window_fun(x,N)**2))
ft_wndw=ft_wndw/norm
ft_wndw_shift=np.fft.fftshift(ft_wndw)
plt.plot(kshift,ftshift)
plt.plot(kshift,ft_wndw_shift)
plt.xlim([0,0.2])
```



As seen in the plot above, the spectral leakage drops quite dramatically, with sharper corners at the bottom of the spike instead of the gradual decay that leaks into nearby frequencies.

e)

Printing the FFT of this window function (N=1001) in python gives:

$$[5.00500000e+02 -2.50250000e+02 1.45623712e-14 \dots 4.09901694e-15 1.45623712e-14 -2.50250000e+02]$$

$$= [N/2 -N/4 0 \dots 0 -N/4]$$

Where if we take the absolute value (for negative frequencies) we get:

$$[N/2 N/4 0 \dots 0 N/4]$$

Since the multiplication of the window in real space corresponds to a convolution in Fourier space, we will consider a discrete convolution:

$$(F * G)(k) = \int_{-\infty}^{\infty} F(k')G(k - k')dk' \rightarrow (F * G)[k] = \sum_{k'=-\infty}^{k'=+\infty} F[k']G[k - k']$$

If $G[k]$ is the window function $[N/2, -N/4, 0, \dots, 0, -N/4]$:

$$G[k - k'] \neq 0 \Rightarrow k - k' = -1, 0, 1 \Rightarrow k' = k - 1, k, k + 1$$

Therefore the sum reduces to three terms:

$$(F * G)[k] = F[k - 1]G[-1] + F[k]G[0] + F[k + 1]G[1]$$

$$= -\frac{N}{4}F[k - 1] + \frac{N}{2}F[k] - \frac{N}{4}F[k + 1]$$

Since the convolution will be normalized to avoid power loss from the window function, we can remove the factors of N now instead of dividing by a much larger normalization factor later:

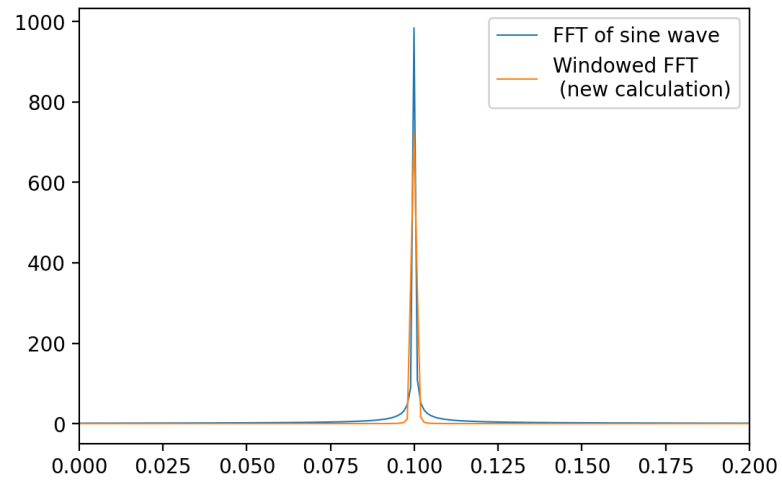
$$F[k] \rightarrow -F[k - 1]/4 + F[k]/2 - F[k + 1]/4$$

Implementing this into code and plotting:

```
ft_wndw_new=np.empty(N)
for i in range(N):
    if i==N-1: # For wrap around, use special case for last array element
        ft_wndw_new[i]=-ft[i-1]/4+ft[i]/2+-ft[0]/4
    else:
        ft_wndw_new[i]=-ft[i-1]/4+ft[i]/2+-ft[i+1]/4

plt.plot(kshift,ftshift)
norm_new=np.sqrt(np.mean(ft_wndw_new**2))
ft_wndw_new=np.abs(ft_wndw_new)/norm
plt.plot(kshift,np.fft.fftshift(ft_wndw_new))
plt.xlim([0,0.2])
```

Plot on next page:



As expected, this is the same as the previous plot.

5 Question 5

a)