

1 MRI Notes

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1.1.1 Potential Form

The induction equation is given by

$$\partial_t \mathbf{b} = \nabla \times (\mathbf{u} \times \mathbf{b}) + \eta \nabla^2 \mathbf{b}$$

Using the following identity

$$\nabla \times \nabla \times \mathbf{f} = \nabla \nabla \cdot \mathbf{f} - \nabla^2 \mathbf{f}$$

and assumming η to be constant, we use $\nabla \cdot \mathbf{b} = 0$, giving

$$\partial_t \mathbf{b} = \nabla \times (\mathbf{u} \times \mathbf{b}) - \eta \nabla \times \nabla \times \mathbf{b}.$$

Then we define a vector potential $\nabla \times \mathbf{A} \equiv \mathbf{b}$, yielding

$$\begin{aligned} \partial_t \nabla \times \mathbf{A} &= \nabla \times (\mathbf{u} \times \mathbf{b}) - \nabla \times (\eta \nabla \times \mathbf{b}) \\ \partial_t \mathbf{A} &= \mathbf{u} \times \mathbf{b} - \eta \nabla \times \mathbf{b} + \nabla \phi \end{aligned}$$

where ϕ is a scalar potential arising from “uncurling” the equation. We must then provide an additional constraint to fix ϕ : the Coulomb gauge $\nabla \cdot \mathbf{A} = 0$. Therefore

$$-\nabla \times \mathbf{b} = -\nabla \times \nabla \times \mathbf{A} = \nabla^2 \mathbf{A}.$$

Next we decompose $\mathbf{u} = \mathbf{u}_0 + \mathbf{u}'$ and $\mathbf{b} = \mathbf{b}_0 + \mathbf{b}'$. We assume the mean quantities \mathbf{u}_0 and \mathbf{b}_0 are themselves solutions to the original problem. If we consider only the 0th mode of \mathbf{b} , i.e. $\mathbf{b} \cdot \hat{\mathbf{e}}_i \sim e^{i0}$ then clearly $\nabla^2 \mathbf{b} = \mathbf{0}$ and therefore

$$\partial_t \mathbf{b}' = \nabla \times (\mathbf{u}_0 \times \mathbf{b}') + \nabla \times (\mathbf{u}' \times \mathbf{b}_0) + \nabla \times (\mathbf{u}' \times \mathbf{b}').$$

Using another identity

$$\nabla \times (\mathbf{A} \times \mathbf{B}) = \mathbf{A} \nabla \cdot \mathbf{B} - \mathbf{B} \nabla \cdot \mathbf{A} + \mathbf{B} \cdot \nabla \mathbf{A} - \mathbf{A} \cdot \nabla \mathbf{B}$$

Momentum Equation

Verbatim from Jeff Oishi’s MRI paper:

$$\frac{D\mathbf{u}'}{Dt} + f \hat{\mathbf{z}} \times \mathbf{u}' + S u'_x \hat{\mathbf{y}} + \nabla p' + \nu \nabla \times \boldsymbol{\omega}' = B_0 \partial_z \mathbf{b}'$$

where f is the coriolis parameter, S is the background shearing rate, and $B_0\hat{\mathbf{z}}$ is a uniform background magnetic field. This equation is linearized wrt perturbations. Accordingly, the material derivative goes like

$$\begin{aligned}\frac{D}{Dt} &\equiv \partial_t + \mathbf{u} \cdot \nabla \\ &= \partial_t + Sx\partial_y\end{aligned}$$

due to the background velocity $\bar{\mathbf{u}} = Sx\hat{\mathbf{y}}$. In the nonlinear case we have

$$= \partial_t + (Sx\hat{\mathbf{y}} + \mathbf{u}') \cdot \nabla$$

From inspection and stuff, the irrotational momentum equation goes like

$$\frac{D\mathbf{u}}{Dt} + \nabla p + \nu \times \boldsymbol{\omega} = \mathbf{b} \cdot \nabla \mathbf{b}$$

Next we generalize $\mathbf{u} = \mathbf{u}' + Sx\hat{\mathbf{y}}$ and $\mathbf{b} = \mathbf{b}' + B_0\hat{\mathbf{z}}$, giving

$$\underline{\partial_t \mathbf{u}' + \mathbf{u}' \cdot \nabla \mathbf{u}' + Sx\partial_y \mathbf{u}' + Su'_x \hat{\mathbf{y}}} + \nabla p + \nu \nabla \times \boldsymbol{\omega} = B_0 \partial_z \mathbf{b}' + \mathbf{b}' \cdot \nabla \mathbf{b}'$$

where the material derivative $\frac{D\mathbf{u}}{Dt}$ consists of the underlined terms