

# 1 MRI Notes

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### 1.1.1 Potential Form

The induction equation is given by

$$\partial_t \mathbf{b} = \nabla \times (\mathbf{u} \times \mathbf{b}) + \eta \nabla^2 \mathbf{b}$$

Using the following identity

$$\nabla \times \nabla \times \mathbf{f} = \nabla \nabla \cdot \mathbf{f} - \nabla^2 \mathbf{f}$$

and assumming  $\eta$  to be constant, we use  $\nabla \cdot \mathbf{b} = 0$ , giving

$$\partial_t \mathbf{b} = \nabla \times (\mathbf{u} \times \mathbf{b}) - \eta \nabla \times \nabla \times \mathbf{b}.$$

Then we define a vector potential  $\nabla \times \mathbf{A} \equiv \mathbf{b}$ , yielding

$$\begin{aligned} \partial_t \nabla \times \mathbf{A} &= \nabla \times (\mathbf{u} \times \mathbf{b}) - \nabla \times (\eta \nabla \times \mathbf{b}) \\ \partial_t \mathbf{A} &= \mathbf{u} \times \mathbf{b} - \eta \nabla \times \mathbf{b} + \nabla \phi \end{aligned}$$

where  $\phi$  is a scalar potential arising from “uncurling” the equation. We must then provide an additional constraint to fix  $\phi$ : the Coulomb gauge  $\nabla \cdot \mathbf{A} = 0$ . Therefore

$$-\nabla \times \mathbf{b} = -\nabla \times \nabla \times \mathbf{A} = \nabla^2 \mathbf{A}.$$

Next we decompose  $\mathbf{u} = \mathbf{u}_0 + \mathbf{u}'$  and  $\mathbf{b} = \mathbf{b}_0 + \mathbf{b}'$ . We assume the mean quantities  $\mathbf{u}_0$  and  $\mathbf{b}_0$  are themselves solutions to the original problem. If we consider only the 0th mode of  $\mathbf{b}$ , i.e.  $\mathbf{b} \cdot \hat{\mathbf{e}}_i \sim e^{i0}$  then clearly  $\nabla^2 \mathbf{b} = \mathbf{0}$  and therefore

$$\partial_t \mathbf{b}' = \nabla \times (\mathbf{u}_0 \times \mathbf{b}') + \nabla \times (\mathbf{u}' \times \mathbf{b}_0) + \nabla \times (\mathbf{u}' \times \mathbf{b}').$$

Using another identity

$$\nabla \times (\mathbf{A} \times \mathbf{B}) = \mathbf{A} \nabla \cdot \mathbf{B} - \mathbf{B} \nabla \cdot \mathbf{A} + \mathbf{B} \cdot \nabla \mathbf{A} - \mathbf{A} \cdot \nabla \mathbf{B}$$