

Marginally-Stable Thermal Equilibria of Rayleigh-Bénard Convection

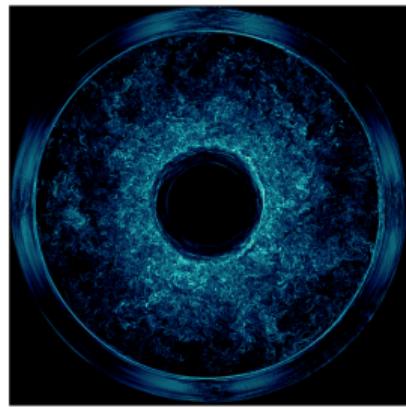
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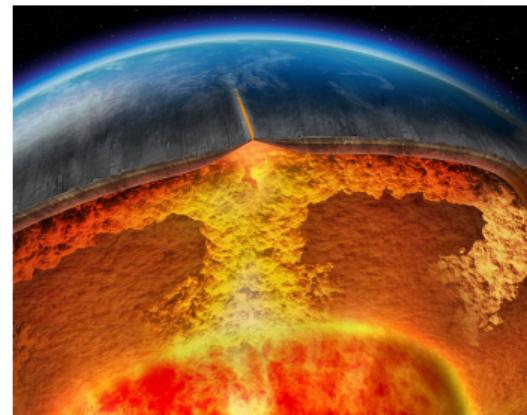


Northwestern
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Why are we Still Studying Convection? Applications



Astrophysics

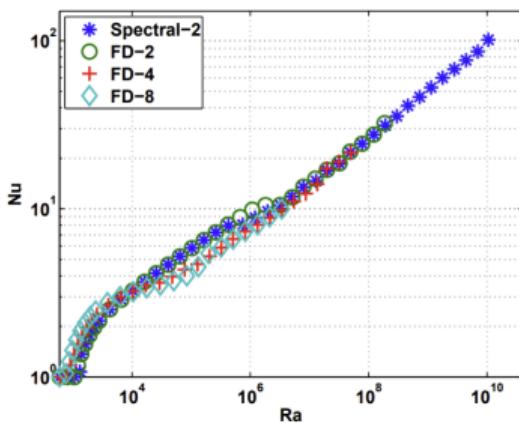


Geophysics

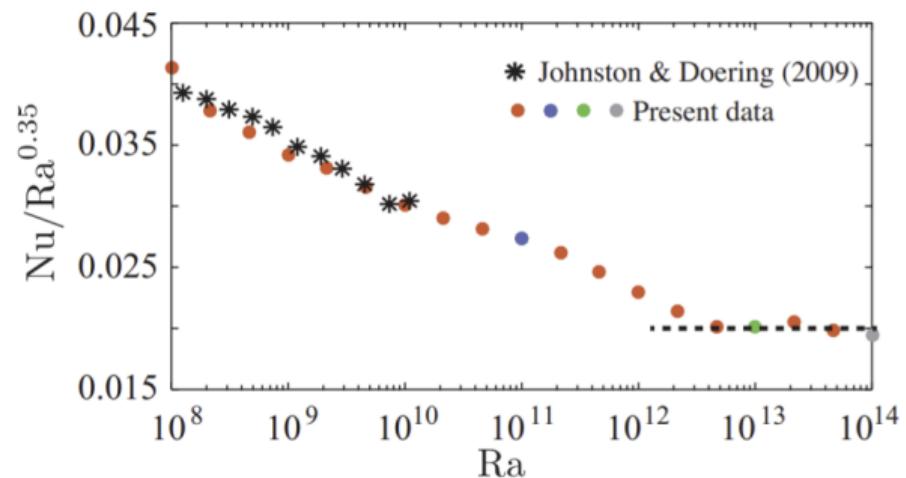


Engineering

Why are we Still Studying Convection? It's Elusive



(Johnson and Doering
2009)



(Zhu et al 2019)

Boussinesq Convection: The PDE

- Incompressible kinematics with buoyancy term

$$\nabla \cdot \mathbf{u} = 0 \quad (1)$$

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + T \hat{z} + \mathcal{R} \nabla^2 \mathbf{u} \quad (2)$$

$$\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T = \mathcal{P} \nabla^2 T \quad (3)$$

where \mathbf{u} , T , and p and velocity, temperature, and pressure respectively.

- Momentum diffusivity: $\mathcal{R} = \sqrt{\frac{\text{Pr}}{\text{Ra}}}$
- Thermal diffusivity: $\mathcal{P} = \frac{1}{\sqrt{\text{Pr Ra}}}$.

Boussinesq Convection: Miscellaneous Prescriptions

- Nondimensionalized on the freefall timescale
- Domain: $\mathcal{D} = \{(x, z) \mid x \in (0, 4), z \in (-1/2, 1/2)\}$
- Prescribed-temperature: $T|_{z=-1/2}(x, t) = 1/2, T_{z=1/2}(x, t) = -1/2$
and no-slip $\mathbf{u}|_{z=\pm 1/2}(x, t) = \mathbf{0}$ boundary conditions
- Fixed Prandtl number $\text{Pr} = 1$, varied Rayleigh number $10^6 \leq \text{Ra} \leq 10^9$

Rayleigh's Linear Stability Analysis

- Assume variables can be decomposed like

$$\mathbf{u}(x, z, t) = \mathbf{u}'(x, z, t) = u'(x, z, t)\hat{x} + w'(x, z, t)\hat{z} \quad (4)$$

$$T(x, z, t) = \bar{T}(z, t) + T'(x, z, t) \quad (5)$$

$$p(x, z, t) = \bar{p}(z, t) + p'(x, z, t). \quad (6)$$

- The Eigenvalue Problem (EVP) is given by

$$\nabla \cdot \mathbf{u}' = 0 \quad (7)$$

$$\frac{\partial \mathbf{u}'}{\partial t} = -\nabla p' + T' \hat{z} + \mathcal{R} \nabla^2 \mathbf{u}' \quad (8)$$

$$\frac{\partial T'}{\partial t} + \frac{\partial \bar{T}}{\partial z} w' = \mathcal{P} \nabla^2 T' \quad (9)$$

Solutions to the Linear Problem

Solutions are given by

$$w'(x, z, t) = A \Re \left[W(z) e^{i(k_x x - st)} \right] \quad (10)$$

$$u'(x, z, t) = A \Re \left[U(z) e^{i(k_x x - st)} \right] \quad (11)$$

$$T'(x, z, t) = A \Re \left[\theta(z) e^{i(k_x x - st)} \right] \quad (12)$$

$$p'(x, z, t) = A \Re \left[P(z) e^{i(k_x x - st)} \right] \quad (13)$$

where A is the (undetermined) mode amplitude, $s = \omega + i\sigma$ is the eigenvalue,

and the allowed wavenumbers $k_x \in \left\{ \frac{n\pi}{2} \mid n \in \mathbb{N} \right\}$.

The Quasilinear Model

- The 1D quasilinear initial value problem (IVP)

$$\frac{\partial \bar{T}}{\partial t} + \frac{\partial}{\partial z} \langle w' T' \rangle_x = \mathcal{P} \frac{\partial^2 \bar{T}}{\partial z^2} \quad (14)$$

- Evolution due to **Advection**: $\frac{\partial}{\partial z} \langle w' T' \rangle_x$ and **Diffusion**: $\mathcal{P} \frac{\partial^2 \bar{T}}{\partial z^2}$
- Derived by manipulating the nonlinear and linear equations
- Can be solved in conjunction with the EVP provided the amplitude A is known

→ **Marginal Stability**

Marginal Stability Constraint

- We assume the perturbations evolve on a shorter timescale than the background state
- We impose **Marginal Stability** at each timestep

$$\max_{k_x} \{\sigma\} = 0. \quad (15)$$

- The growth rate σ is obtained by solving the EVP for some $\bar{T}(z, t)$
- Does not allow us to solve for A directly → discrete root-finding methods

Solving for the Perturbation Amplitude A

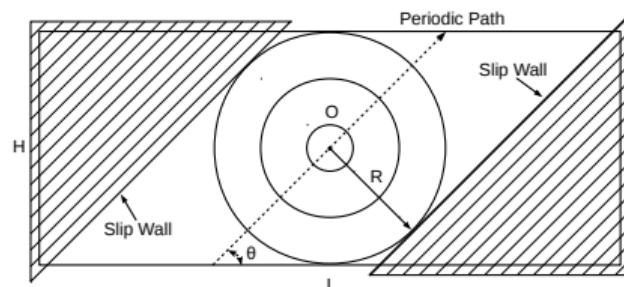
- 1 Given some guess for the amplitude A and a fixed timestep Δt
- 2 We evolve the initial (marginally stable) temperature profile $\bar{T}(z, t) \rightarrow \bar{T}(z, t + \Delta t)$
- 3 Solve the EVP using $\bar{T}(z, t + \Delta t)$
- 4 This renders $\sigma(A)$
- 5 Use Newton's method with finite differences to solve $\sigma(A) = 0$

The Amplitude Guess

- Consider advection and diffusion separately on $t \rightarrow t + \Delta t$
- Solve $\frac{\partial \bar{T}}{\partial t} + \langle w' T' \rangle_x = 0$, yielding $\bar{T}_{adv} \rightarrow \sigma_{adv}$
- Solve $\frac{\partial \bar{T}}{\partial t} = \mathcal{P} \frac{\partial^2 \bar{T}}{\partial z^2}$, yielding $\bar{T}_{adv} \rightarrow \sigma_{diff}$
- For small timesteps, we assume diffusion and advection act independently, i.e.

$$0 = \Delta\sigma \approx A^2 \sigma_{adv} + \sigma_{diff} \quad \rightarrow \quad A^2 \approx -\frac{\sigma_{diff}}{\sigma_{adv}} \quad (16)$$

Position animation + Make Table



To play the video, the compiled PDF should be moved out from the "Tmp" directory

$m_x \times m_y$	L_1 error	L_1 order	L_2 error	L_2 order	L_∞ error	L_∞ order
40×20	$3.536\text{e-}2$	—	$6.097\text{e-}2$	—	$4.105\text{e-}1$	—
80×40	$9.113\text{e-}3$	1.956	$2.497\text{e-}2$	1.288	$1.997\text{e-}1$	1.039
160×80	$2.034\text{e-}3$	2.163	$6.548\text{e-}3$	1.931	$5.236\text{e-}2$	1.931
320×160	$5.114\text{e-}4$	1.992	$1.640\text{e-}3$	1.997	$1.278\text{e-}2$	2.035
640×320	$1.287\text{e-}4$	1.990	$4.097\text{e-}4$	2.001	$3.119\text{e-}3$	2.034
1280×640	$3.233\text{e-}5$	1.993	$1.024\text{e-}4$	2.000	$7.818\text{e-}4$	1.996



Ordinary text

- **A 3D, high-resolution, parallelized, gas-solid flow solver**
 - Establishes a numerical framework for the direct simulation of gas-solid flows.
 - Solves coupled and interface-resolved fluid-fluid, fluid-solid, and solid-solid interactions.
 - Addresses shocked flow conditions, irregular and moving geometries, and multibody contact and collisions.
- **Advancement in understanding particle clustering and jetting**
 - Demonstrates a valid statistical dissipative property in solving explosively dispersed granular materials with respect to Gurney velocity.
 - Extends the time range of the velocity scaling law with regard to Gurney energy in the Gurney theory from the steady-state termination phase to the unsteady evolution phase.
 - Proposes an explanation for particle clustering and jetting instabilities to increase the understanding of experimental observations.

Thank you for your attention!



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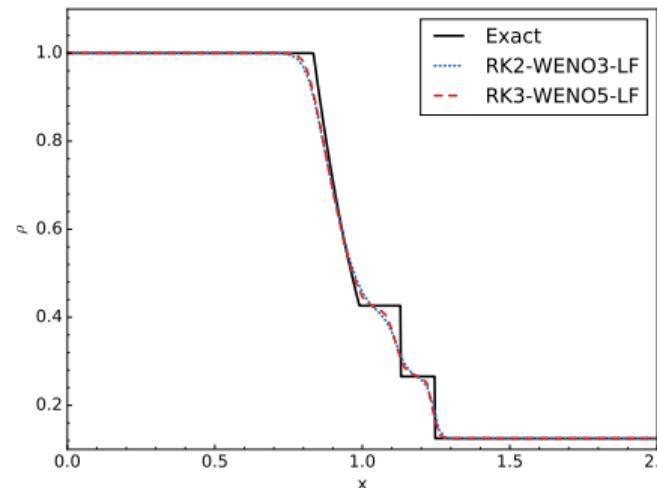
Part I

Appendix

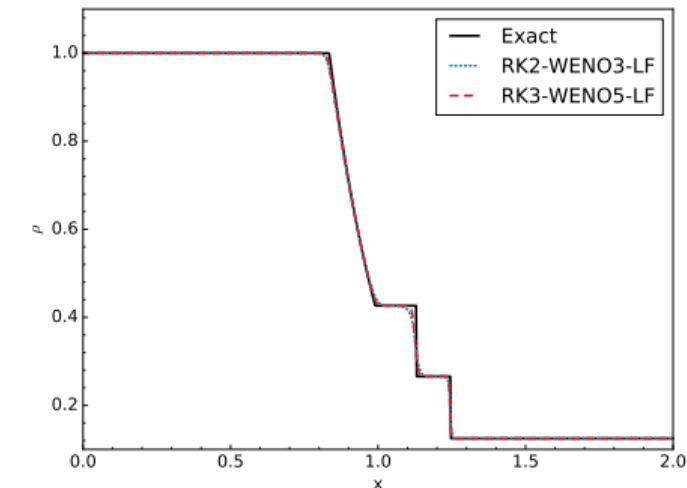
6 Appendix

Classic Beamer Style
References

Sod's problem (Sod 1978)



$n = 100$



$n = 500$

$$\rho = 1; \quad u = 0; \quad p = 1 \quad \text{if } 0 \leq x < 1$$

$$\rho = 0.125; \quad u = 0; \quad p = 0.1 \quad \text{if } 1 < x \leq 2$$

References I

Sod, G. A. (1978). "A survey of several finite difference methods for systems of nonlinear hyperbolic conservation laws". In: J. Comput. Phys. 27.1, pp. 1–31.