

## RESEARCH ARTICLE

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## Key Points:

- A much higher occurrence rate of weak  $B_z$  events is found than that of stronger  $B_z$
- Many of the weak  $B_z$  events originate from regions other than ICMEs and high-speed streams
- Even weak  $B_z$  events are often geoeffective and these include substorms and disturbances of the outer radiation belt

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Statistical properties and geoeffectiveness of southward interplanetary magnetic field with emphasis on weakly southward  $B_z$  events

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**Abstract** In this work we analyze the statistical properties and geoeffectiveness of southward interplanetary magnetic field (IMF)  $B_z$  ( $B_z$ ) events measured near the Earth from 1996 to 2015. We classify the  $B_z$  events into six contrasting classes according to the duration and intensity of  $B_z$ : two subclasses of long-lasting major intensity, two subclasses of long-lasting weak intensity, and two subclasses of short-lasting weak intensity. We find the following results. (1) The occurrence rate of  $B_z$  events is lowest for the long-lasting major-intensity  $B_z$  classes and increases toward shorter-lasting and weaker-intensity  $B_z$  classes. (2) Many of the  $B_z$  events in nearly all classes are associated with regions other than interplanetary coronal mass ejections and coronal hole high-speed streams. (3) Alfvén waves account for a significant fraction of the  $B_z$  events (e.g., ~60% for the short-lasting and weak-intensity  $B_z$  classes). (4) There are clear differences in solar cycle dependence of  $B_z$  events among the different  $B_z$  classes. (5) Even either the long-lasting or short-lasting weak-intensity class  $B_z$  events can often be geoeffective by triggering substorms of weak-to-medium intensity and by enhancing geosynchronous relativistic electron fluxes. In conclusion, we emphasize that even the weakly (in terms of either intensity or duration) southward IMF  $B_z$  should be considered significant from the viewpoint of their high occurrence rate and their geoeffectiveness compared to those of long-lasting major intensity.

## 1. Introduction

It is well known that geomagnetic storms are caused by a strong southward interplanetary magnetic field (IMF)  $B_z$ . There have been many studies, indicating that the intensity of a geomagnetic storm is closely related to the magnitude and duration of southward IMF  $B_z$ . For example, Gonzalez *et al.* [1994] found that long southward  $B_z$  (>3 h) almost always causes a strong geomagnetic storm (minimum  $Dst$  < −100 nT). Wu and Lepping [2002] examined 135 magnetic clouds and found that the  $Dst$  index correlates well with both the IMF  $B_z$  component and the solar wind electric field,  $V_{sw}B_z$ . They further found that the correlation coefficient for  $Dst$  versus IMF  $B_z$  increases dramatically when the solar wind speed exceeds 600 km/s. Gopalswamy *et al.* [2008] also reported that two thirds of their 99 magnetic clouds are geoeffective; the  $Dst$  index was highly correlated with the solar wind speed and the magnetic field in the magnetic clouds as well as their product.

Substorm occurrence is also related to IMF conditions. However, the relation is rather complex [Kamide, 2001]. While mostly it is the southward IMF  $B_z$  that is required for preconditioning prior to substorm occurrence, other situations can be also related to substorm occurrence such as northward turning of IMF  $B_z$  [e.g., Lyons *et al.*, 1997] and even continuously northward IMF  $B_z$  conditions [Lee *et al.*, 2010]. In fact, it is known that Alfvénic IMF fluctuations during coronal hole high-speed streams [Tsurutani and Gonzalez, 1987], involving continuous turning of IMF  $B_z$  between northward and southward directions, lead to repetitive substorms [Lee *et al.*, 2006; H.-J. Kim *et al.*, 2008]. In addition, storms usually accompany substorms, namely, the so-called “storm time substorms” [Kamide, 2001], as a southward IMF  $B_z$  is a common factor to occurrence of both storms and substorms. A short review of recent advances in substorm research can be found in Sergeev *et al.* [2012, and reference therein].

It is well known that relativistic electron fluxes in the outer radiation belt often change significantly in response to the major solar activity such as interplanetary coronal mass ejections (ICMEs) and coronal hole high-speed streams [e.g., Reeves *et al.*, 2003; Lyons *et al.*, 2005], which accompany southward IMF  $B_z$ .

conditions. The radiation belt dynamics can be affected in a sensitive way by storms and substorms, which produce energetic particles and plasma waves of various kinds (see a review by *Friedel et al.* [2002]). From a practical viewpoint, understanding the relativistic electron dynamics is critical as they can cause fatal failures of spacecraft operations through “deep dielectric charging” [*Gussenhoven et al.*, 1991; *Baker et al.*, 1994]. *Baker* [2000] reported that two spacecraft failures in 1998 occurred at the time when the  $>2$  MeV geosynchronous electron flux level reached  $\sim 2 \times 10^3/\text{cm}^2/\text{s}/\text{sr}$ .

Thus, it is of critical importance to understand the characteristics of southward IMF  $B_z$  conditions with different origins and the relative importance of their geoeffectiveness.

ICMEs and coronal hole high-speed streams are two types of major solar activity, which provides southward IMF  $B_z$  conditions [*Klein and Burlaga*, 1982; *Wilson*, 1987; *Lindsay et al.*, 1995; *Tsurutani et al.*, 2006; *Gopalswamy*, 2016]. Therefore, many previous studies of the geoeffectiveness of southward IMF  $B_z$  have concentrated on them. However, it is not unreasonable to consider the possibility that some southward IMF  $B_z$  conditions are not related to ICMEs and coronal hole high-speed streams. In fact, *Zhang and Moldwin* [2014] suggested that a quarter of the strongly southward IMF  $B_z$  events ( $t > 6$  h,  $B_z < -10$  nT) was not related to any well-defined solar wind structure. Also, *Zhang et al.* [2014] analyzed southward IMF  $B_z$  events defined by  $t > 1$  h and  $B_z < -5$  nT that were unrelated with any well-defined solar wind structure at 1 AU. They found that one third of the southward IMF events show the features of Alfvén waves and more than half of them originate from the slow solar wind. Motivated by these recent works [*Zhang and Moldwin*, 2014; *Zhang et al.*, 2014], we revisit the similar problem of the statistical features and geoeffectiveness of southward IMF  $B_z$  events measured near the Earth. In contrast to the previous works, we pay a special attention to the extent to which weakly southward IMF  $B_z$  events are significant in terms of occurrence rate and geoeffectiveness compared to strongly southward IMF  $B_z$  events.

For this work, we define six classes of southward IMF  $B_z$  events according to their intensity and duration. We derive the general statistical features of the six classes. In addition, our analysis includes determination of the extent to which the identified southward  $B_z$  events of each class are associated with ICMEs, coronal hole high-speed streams, and Alfvén waves. Finally, we check the geoeffectiveness of each of the different classes of southward  $B_z$  events.

The paper is organized as follows. In section 2, we describe the data used for the analysis and the methodology employed in this work. In section 3, we present the analysis results. Lastly, a summary and discussion are given in section 4.

## 2. Data and Methodology

To identify a southward IMF  $B_z$  event,  $B_s$ , we used the IMF data for 20 years from 1996 to 2015 which are available on NASA's OMNI website, <http://omniweb.gsfc.nasa.gov/>. We used the  $B_z$  data in the Geocentric Solar Ecliptic (GSE) coordinate system and at 5 min time resolution. It is  $B_z$  defined in the geocentric solar magnetospheric (GSM) coordinate system that can describe more suitably the interaction between IMF and the major magnetospheric phenomena such as storms and substorms. However, to avoid additional issues such as the diurnal variation of  $B_z$  in GSM, we use  $B_z$  in GSE for this study.

For this work, we determine a “ $B_s$  event” by requiring that IMF  $B_z$  is negative continuously for a certain time interval based on the 5 min time resolution data. We identified six groups of  $B_s$  events according to different durations and intensity of southward  $B_z$ : Specifically, they are “Long-Strong  $B_s$ ,” “Long-Medium  $B_s$ ,” “Long-Weak  $B_s \geq 3$  h,” “Long-Weak  $B_s \geq 1$  h,” “Short-Weak  $B_s 1$ ,” and “Short-Weak  $B_s 2$ .” The naming of each  $B_s$  event reflects duration and intensity of southward  $B_z$ : In this aspect, the six classes are well distinguished. The specific conditions for definition of each class are summarized in the second and third columns of Table 1. Long-Strong  $B_s$  class is defined by the condition that  $B_z$  remains southward for 3 h or longer during which its minimum value is  $-10$  nT or less. Long-Medium  $B_s$  class is defined similarly except that the minimum  $B_z$  is relaxed to be  $-5$  nT or less. Note that the Long-Medium  $B_s$  class automatically includes the Long-Strong  $B_s$  class. We compare these two classes with an intention to see if the stronger  $B_s$  condition can make any major difference. The next two classes, Long-Weak  $B_s \geq 3$  h and Long-Weak  $B_s \geq 1$  h, reflect southward  $B_z$  events with weaker intensity having  $B_z$  minimum between 0 and  $-5$  nT, but with different durations of southward

**Table 1.** The Criteria and Basic Statistics of  $B_z$  Event Classes<sup>a</sup>

$B_z$ Classes	Criteria of $B_z$		Statistics	
	Duration	Minimum	Event Numbers	Summed Duration <sup>a</sup>
Long-Strong $B_z$	$\geq 3$ h	$\leq -10$ nT	397	139 days (2%)
Long-Medium $B_z$	$\geq 3$ h	$\leq -5$ nT	2,240	616 days (8%)
Long-Weak $B_z \geq 3$ h	$\geq 3$ h	0 to $-5$ nT	3,680	780 days (11%)
Long-Weak $B_z \geq 1$ h	$\geq 1$ h	0 to $-5$ nT	15,663	1626 days (22%)
Short-Weak $B_z$ 1	30 min–60 min	0 to $-5$ nT	15,800	513 days (7%)
Short-Weak $B_z$ 2	15 min–45 min	0 to $-5$ nT	28,671	613 days (8%)

<sup>a</sup>The value in parenthesis is the percentage duration of each class during the 20 year measurement period.

interval,  $\geq 3$  h and  $\geq 1$  h, respectively. To be practical with these long lasting durations, i.e.,  $\geq 3$  h and  $\geq 1$  h, we allowed a very brief deviation of  $B_z$  from the southward direction to northward direction as long as this occurs only at a single data point of 5 min time resolution. Note that the Long-Weak  $B_z \geq 1$  h class automatically includes the Long-Weak  $B_z \geq 3$  h class. We test these classes with an intention to see if the longer-lasting  $B_z$  condition can make any major difference since we see no definite reason for having to select a specific duration when defining a “long-lasting” event. The last two classes, Short-Weak  $B_z$  1 and Short-Weak  $B_z$  2, represent southward  $B_z$  events with weak intensity ( $B_z$  minimum between 0 and  $-5$  nT) and short durations, 30 min–60 min and 15 min–45 min, respectively. The durations of these two Short-Weak classes overlap to some extent, and we have deliberately chosen these two durations with an intention to test any sensitivity of the statistical results to small but different durations. Figure 1 shows some examples of  $B_z$  event classes identified for an interval in 2015. The purpose with the six-class classification is to compare  $B_z$  events of different intensity and duration from perspectives of their origin and geoeffectiveness.

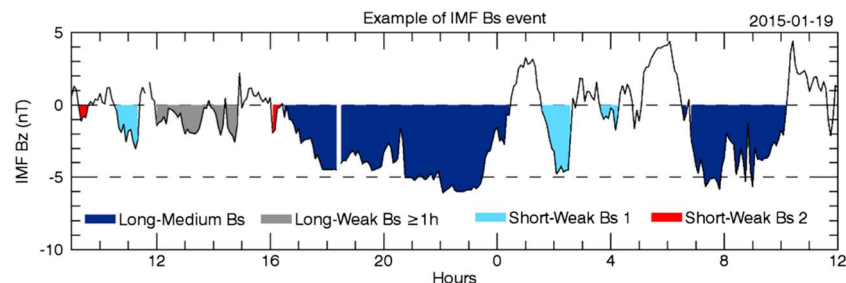
Our examination of the origin of the six classes of  $B_z$  checked occurrences of the major solar activity, namely, interplanetary coronal mass ejections and coronal hole high-speed streams, as well as for the existence of Alfvén waves in the solar wind. Also, we checked the dependence of each class on the sunspot cycle. In addition, we compared the six classes in terms of accompanying solar wind variables such as the solar wind speed, IMF magnitude, solar wind density, and temperature.

Our examination of the geoeffectiveness by  $B_z$  of the six classes focused on three major phenomena, namely, the occurrence of substorms (monitored by  $AL$  index), the occurrence of magnetic storms (based on  $SYM-H$  index), and the outer radiation belt disturbance (using  $>2$  MeV electron fluxes at geosynchronous orbit).

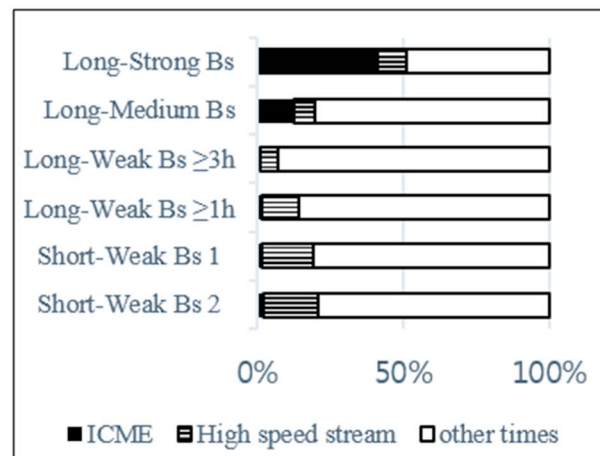
### 3. Statistical Results

#### 3.1. The Basic Statistics of $B_z$ Events

Table 1 shows a summary of the definition and statistics for the six classes of  $B_z$  events, which we obtained based on the methods described in section 2. The criteria for each class are summarized in the second and third columns in Table 1. The fourth column shows the number of independent events in each  $B_z$  class. The last column refers to the summed duration over all of the  $B_z$  events for each class, and the corresponding



**Figure 1.** An example of IMF  $B_z$  events. Four classes of  $B_z$  events out of the six classes in Table 1 are identified. The two short intervals of  $B_z < 0$  at ~05:00 UT and ~11:30 UT on 19 January 2015 are shorter than 15 min and do not satisfy the selection criteria in Table 1, and thus, they are not identified as an event here.



**Figure 2.** The solar events contribution to  $B_z$  classes.

events is smallest for Long-Strong  $B_z$  class and rapidly increases toward the classes of weaker intensity and shorter interval (from top to bottom rows). The statistics is somewhat different in terms of the summed durations. It shows that the summed durations are short for the Long-Strong and Long-Medium classes and increase for the Long-Weak  $B_z$  class  $\geq 1$  h being the longest duration among the six classes. The total durations of the two Short-Weak  $B_z$  classes are shorter mainly due to their definitions requiring short durations, although the occurrence percentage rates of their independent event classes are two of the highest.

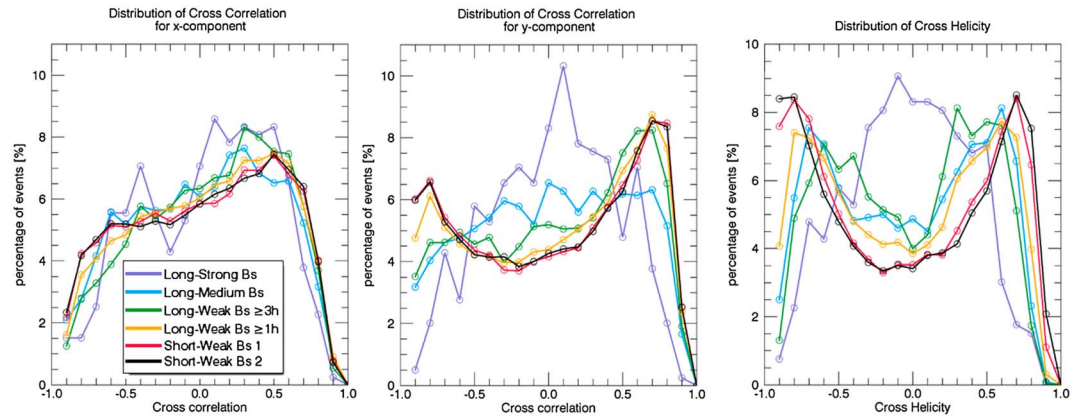
### 3.2. The Origin of Each Class of $B_z$ Events

We studied the origin of  $B_z$  events to determine the possible differences among the different  $B_z$  classes. First, for each  $B_z$  event, we checked their association with ICMEs and coronal hole high-speed streams. For determining the occurrence of ICME events, we relied mostly on the sharp changes in the IMF and the proton temperature and density as demonstrated in Kamide *et al.* [1998], Cane and Richardson [2003], and Richardson and Cane [2010] as well as on the website, <http://www.srl.caltech.edu/ACE/ASC/DATA/level3/icmetable2.htm>. For determining the occurrence of the coronal hole high-speed streams, we used a criterion of the solar wind speed being over 500 km/s continuously for 2 days or longer with a peak speed of 600 km/s or larger. The selected events were all confirmed by visual inspection for consistency with the generally known features of coronal hole high-speed streams including the existence of the compressed corotating interaction region followed by long-lasting Alfvénic fluctuating main body of the high-speed streams, and a repetitive occurrence every  $\sim 27$  days [e.g., Rosenberg and Coleman, 1980; Kamide *et al.*, 1998].

The results are presented in Figure 2, which shows that out of the Long-Strong  $B_z$  class events,  $\sim 41\%$  and  $\sim 10\%$  are associated with ICMEs and coronal hole high-speed streams, respectively. For the Long-Medium  $B_z$  class events, the association with ICMEs and coronal hole high-speed streams becomes less, being  $\sim 12\%$  and  $\sim 8\%$ , respectively. In contrast, the association with ICMEs becomes negligible for the other four  $B_z$  classes, which is not surprising because magnetic fields in ICMEs generally represent a significant enhancement over the background IMF. Instead, the association with coronal hole high-speed streams becomes greater for the Short-Weak  $B_z$  classes.

However, it should be emphasized that the majority of  $B_z$  events are associated with “other times,” ranging  $\sim 50\%$  for Long-Strong class to  $\sim 79\%$  for Short-Weak 2 class. This implies a possible origin of  $B_z$  due to reasons in the solar wind other than the periods of major solar activity. Specifically, the “other times” in Figure 2 must primarily be locations where the solar wind conditions have not well defined ICMEs and high-speed streams. However, they could partly include the interplanetary small-scale magnetic flux ropes (ISMFR) [Moldwin *et al.*, 2000; Feng *et al.*, 2008, 2010; Zhang *et al.*, 2013], and the conditions associated with shocks and discontinuities [Burlaga, 1968, 1970; Tsurutani and Ho, 1999]. In the present work, we have not identified ISMFRs, shocks, and discontinuities, which were, however, distinguished in the previous work by Zhang and Moldwin [2014]. In addition, we caution that in determining the association of  $B_z$  with ICMEs, we have not included the sheath region ahead of the ICME body. Thus, the “other times” includes the  $B_z$  events found in the sheath region

percentage rate (shown in parenthesis) out of the total 20 year time interval studied. Note that the summed durations in total correspond to  $<45\%$  of 20 years (It should be recalled that Long-Medium class automatically includes Long-Strong class, Long-Weak  $B_z \geq 1$  h class includes Long-Weak  $B_z \geq 3$  h class, and there is some overlap between the two Short-Weak classes). There is the remaining  $>55\%$  interval which consists of the southward  $B_z$  conditions that did not satisfy our selection criteria and of course the locations where the IMF  $B_z$  was northward. Not surprisingly, the number of  $B_z$



**Figure 3.** Distribution of cross correlations and cross helicity for the six classes of  $B_s$  events. (left and middle) Cross correlations refer to those between  $B_x$  and  $V_x$  and between  $B_y$  and  $V_y$ , respectively. (right) Cross helicity represents a degree of “Alfvénicity” based on the definition shown in the text.

ahead of ICMEs, but the number of such  $B_s$  events is small ( $<10\%$  of other times for the Long-Strong  $B_s$  class events to which the ICME contribution is largest).

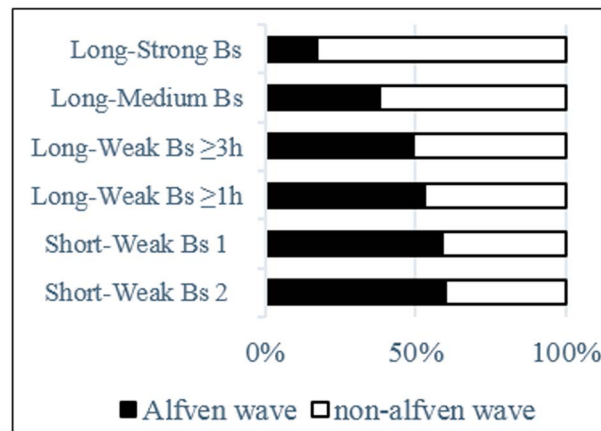
To examine the possible origins of  $B_s$  events from other than major solar activity, we checked for the existence of Alfvén waves in association with  $B_s$  events. We have done this by estimating the cross correlations between the magnetic field and solar wind velocity components [Belcher and Davis, 1971; Tsurutani and Ho, 1999]. Specifically, we have estimated them between  $B_x$  and  $V_x$  and between  $B_y$  and  $V_y$ . For further confirmation, we have also tested the cross helicity  $\sigma_c$  as in Riley *et al.* [1996], which is defined by  $\sigma_c = 2\delta V \cdot \delta V_A / (\delta V^2 + \delta V_A^2)$ , where  $\delta V$  and  $\delta V_A$  mean the solar wind velocity fluctuations and the Alfvén wave velocity fluctuations, respectively [Matthaeus and Goldstein, 1982]. When the magnitude of  $\sigma_c$  is close to unity, the fluctuations are considered to be Alfvénic. The computed results are shown in Figure 3. For Long-Strong  $B_s$  class events (violet lines), the distributions of the correlation coefficients and the cross helicity tend to peak around zero. In contrast, the distributions for the Long-Weak and Short-Weak  $B_s$  classes show the tendency for the peak correlations (in particular in y component) and helicity to occur at higher values. This tendency is most prominent for the Short-Weak  $B_s$  1 and 2 class events (black and red lines), the peaks occurring at  $\sim 0.5$  and  $\sim 0.7$  for the x and y component correlation coefficients, respectively, and  $\sim 0.7$  for the cross helicity. A similar tendency is also noted for the two Long-Weak classes (yellow and green lines). For the Long-Medium class, the distribution of cross helicity is similar to those of the Long-Weak and Short-Weak  $B_s$  classes, but those of cross correlations are less well distinguished.

If we choose to impose the specific requirement, simply to be practical, that the absolute values of cross correlation for either x or y component are  $>0.5$  to determine the existence of Alfvén waves, we obtain the results in Figure 4 showing the statistics of the contributions by the Alfvén waves to  $B_s$  events of each class. It is seen from Figure 4 that Alfvén wave contribution is larger when the  $B_s$  event is shorter in duration and/or weaker in intensity. In particular, for the four weak  $B_s$  classes, the Alfvén wave contribution is  $\sim 50\text{--}60\%$ . In contrast, for the Long-Strong  $B_s$  class, the contribution is  $\sim 17\%$ .

Here we add a short comment on the amplitude of Alfvén wave  $B_s$  in comparison with that of ICMEs. First, the average  $B_s$  values of ICMEs that were identified in the Long-Strong and Long-Medium classes are  $-9.7$  nT and  $-7.6$  nT, respectively. For the Alfvén waves of the same classes, the average  $B_s$  values are  $-9.3$  nT and  $-3.3$  nT, respectively. For the Alfvén waves of the other four weaker classes, the average  $B_s$  values range from  $-1.3$  nT to  $-1.7$  nT (Note that the definition of these four classes requires the  $B_s$  values range between 0 and  $-5$  nT).

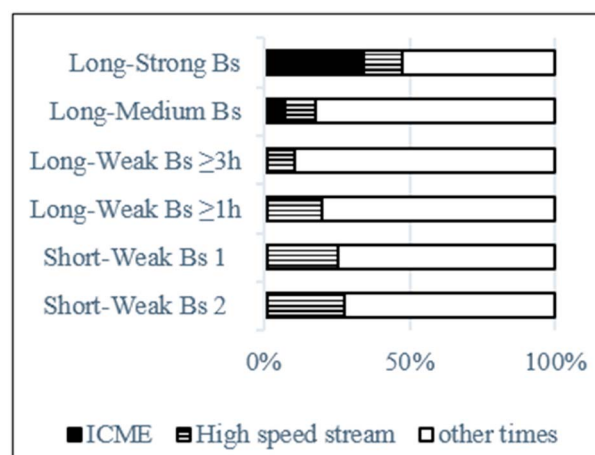
In addition, we investigated, though partially, the origin of Alfvén waves identified in Figure 4 in particular by determining the extent to which they are associated with ICMEs and coronal hole high-speed streams. The results are shown in Figure 5. The association of Alfvén waves with ICMEs and coronal hole high-speed streams is largest ( $\sim 48\%$  together) for Long-Strong  $B_s$  class, and it becomes much smaller for Long-Medium





**Figure 4.** Alfvén wave contribution to  $B_z$  events of the six classes.

column in Table 2 shows the fraction of Alfvén waves out of the  $B_z$  events that are associated with ICMEs/high-speed streams in Figure 2. Likewise, the third column in Table 2 refers to the fraction of Alfvén waves out of the  $B_z$  events that are associated with other times in Figure 2. With these, we intend to determine the extent to which Alfvén waves contribute to a group of  $B_z$  events with two different origins (ICMEs/high-speed streams versus other times) and compare the results among the six  $B_z$  classes. The second column in Table 2 clearly indicates that when  $B_z$  events are found in association with ICMEs or high-speed streams, the Alfvén wave contribution is quite large, ~70%–79%, for the four Long-Weak and Short-Weak classes while it is much lower for the Long-Strong and Long-Medium classes. Recall from Figure 2 that for the four Long-Weak and Short-Weak classes, the association of  $B_z$  events with high-speed streams is dominant over that with ICMEs and the situation is roughly opposite for the Long-Strong and Long-Medium classes. The third column in Table 2 shows that the Alfvén wave contribution is still substantial when  $B_z$  events are found at times other than during the well-defined solar activity, near 50% for the four Long-Weak and Short-Weak classes but less for the Long-Strong and Long-Medium classes. As mentioned in section 1, Zhang *et al.* [2014] reported that one third of southward IMF  $B_z$  events, which are defined by  $t > 1$  h and  $B_z < -5$  nT and by being unrelated with any well-defined solar wind structure at 1 AU, show the features of Alfvén waves. The  $B_z$  events in our classification that are closest to, but not precisely the same as, the  $B_z$  events in Zhang *et al.* [2014] are those of Long-Medium  $B_s$  class found in times other than ICMEs/high-speed streams. The fraction of Alfvén wave contribution to this  $B_s$  class is ~39% as shown in Table 2, which is not too different from the result of Zhang *et al.* [2014].



**Figure 5.** Origin of Alfvén waves identified in Figure 4.

class (~18%). For the four weak  $B_s$  classes, Alfvén wave events are more associated with coronal hole high-speed streams, which is well known [e.g., Tsurutani and Gonzalez, 1987; Lee *et al.*, 2006]. We emphasize that more than half of the identified Alfvén events for all of the six classes are not associated with either ICMEs or coronal hole high-speed streams, further study of which is left for future work.

In Table 2, we present the Alfvén wave results of Figures 4 and 5 in a different way in combination with Figure 2. Specifically, the second

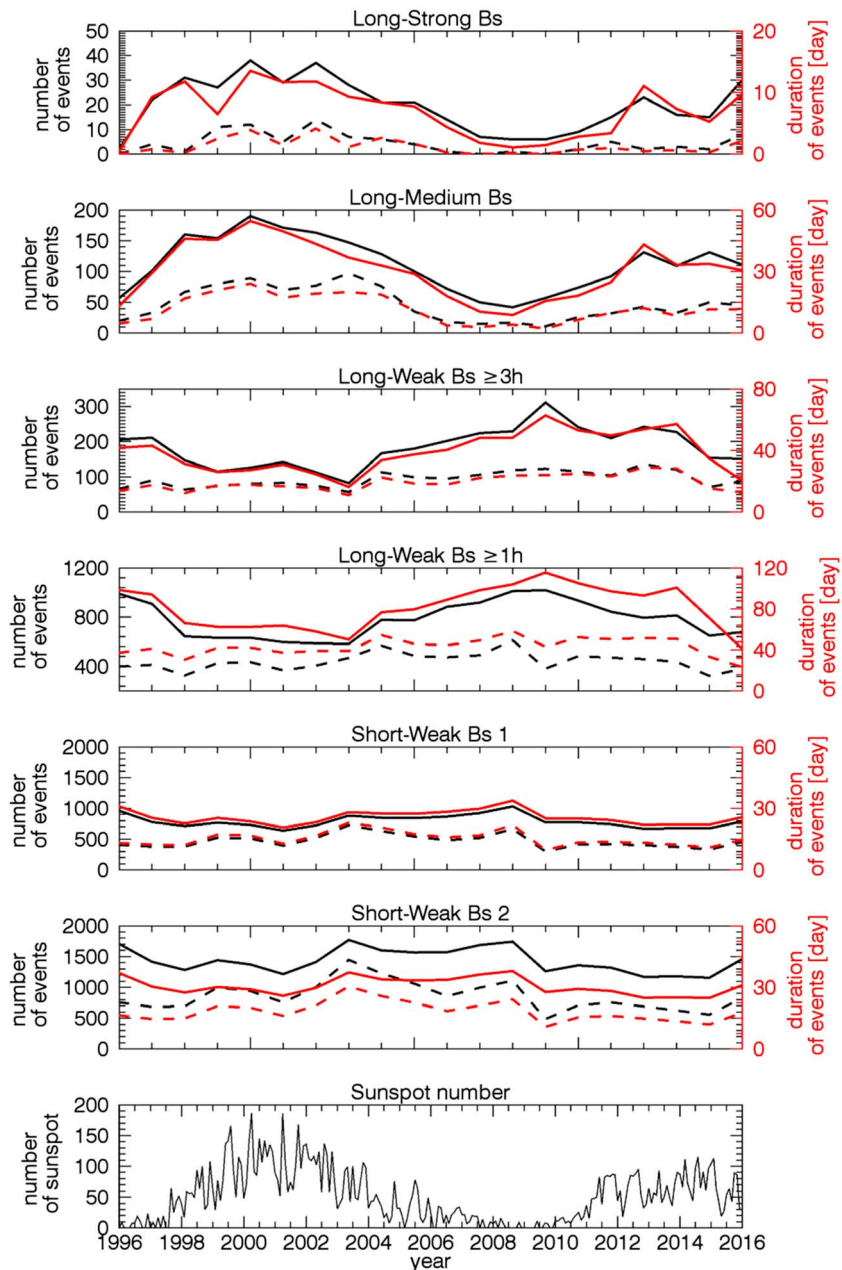
We checked the solar cycle dependence of each  $B_s$  class. Figure 6 shows the results where the number of the  $B_s$  events, and the corresponding durations are shown in black and red lines, respectively. In each panel, the dashed lines refer to the contribution by the Alfvén waves. The last panel is the sunspot number. First, there are clear differences in the solar cycle dependence among the six classes. The top two panels clearly show that the number of the events and the corresponding durations of the Long-Strong and Long-Medium  $B_s$

**Table 2.** Percentages of the Contributions by Alfvén Waves to  $B_s$  Events

$B_s$ Classes	ICME/High-Speed Stream <sup>a</sup> (%)	Other Times <sup>b</sup> (%)
Long-Strong $B_s$	15.9	18.1
Long-Medium $B_s$	33.6	39.1
Long-Weak $B_s \geq 3$ h	70.8	47.3
Long-Weak $B_s \geq 1$ h	74.6	49.5
Short-Weak $B_s$ 1	79.3	54.3
Short-Weak $B_s$ 2	79.3	54.9

<sup>a</sup>Refer to the percentage of Alfvén waves out of the  $B_s$  events which are associated with ICME/high-speed streams in Figure 2.

<sup>b</sup>Refer to the percentage of Alfvén waves out of the  $B_s$  events which are associated with other times in Figure 2.



**Figure 6.** Yearly distribution of the number of the events (black lines) and durations (red lines) for the six  $B_s$  class events, their associated Alfvén wave contributions (dashed lines), and the sunspot cycle (seventh panel) for the 20 year period.

classes have positive correlations with the sunspot cycle. This is reasonable for the Long-Strong  $B_s$  class to which the ICME contribution is largest (~41%) as shown in Figure 2, as the number of sunspots is proportional to the number of the CME events (and thus also of the ICME events). For the Long-Medium  $B_s$  class, the ICME contribution is substantially lower (~12%) as shown in Figure 2, and thus, the positive correlation with the sunspot cycle in Figure 6 was not necessarily expected. The Alfvén wave contributions also indicate positive correlations with the sunspot cycle. On the other hand, there is a strong asymmetry between the two solar cycles, indicating that the number and duration of the  $B_s$  events (and associated Alfvén waves) of the Long-Strong and Long-Medium  $B_s$  classes are larger and longer for solar cycle 23 than for solar cycle 24. The sunspot number declined by ~40% during solar cycle 24 compared to that during solar cycle 23. Gopalswamy *et al.* [2015a, 2015b] reported that the number of halo CMEs and magnetic clouds did not decline in cycle 24. However, it should be noted that our  $B_s$  events include both those associated with ICMEs and those not associated with ICMEs as shown in Figure 2. A further examination in this aspect is worthwhile to perform in the future.

Second, the third and fourth panels of Figure 6 show the tendency that the number of the events and durations for the two Long-Weak  $B_s$  classes are in inverse proportion to the number of sunspots. This is contrary to the Long-Strong and Long-Medium classes. However, a further examination on the intercycle dependence for the two Long-Weak  $B_s$  classes is desirable in the future. On the other hand, the Alfvén wave contributions to the two Long-Weak  $B_s$  classes do not show a significant dependence on solar cycle.

Finally, Figure 6 indicates that for the two Short-Weak  $B_s$  classes, the number of  $B_s$  events and their durations as well as the Alfvén wave contributions all indicate no obvious correlation with the solar cycle except that they appear slightly enhanced during the declining phase of solar cycle 23 and near 2009. Recall the results shown in Figure 2 that the contribution to these classes is partly due to the  $B_s$  events during coronal hole high-speed streams, which is usually dominant during the declining phase of solar cycle, and also due to the  $B_s$  events that are not associated with coronal hole high-speed streams. The overall flat distribution of  $B_s$  for these classes is likely due to the larger contribution by the latter type of  $B_s$  events, although a further study of this is also necessary in the future.

### 3.3. Characteristics of Solar Wind Parameters During $B_s$ Events of Each Class

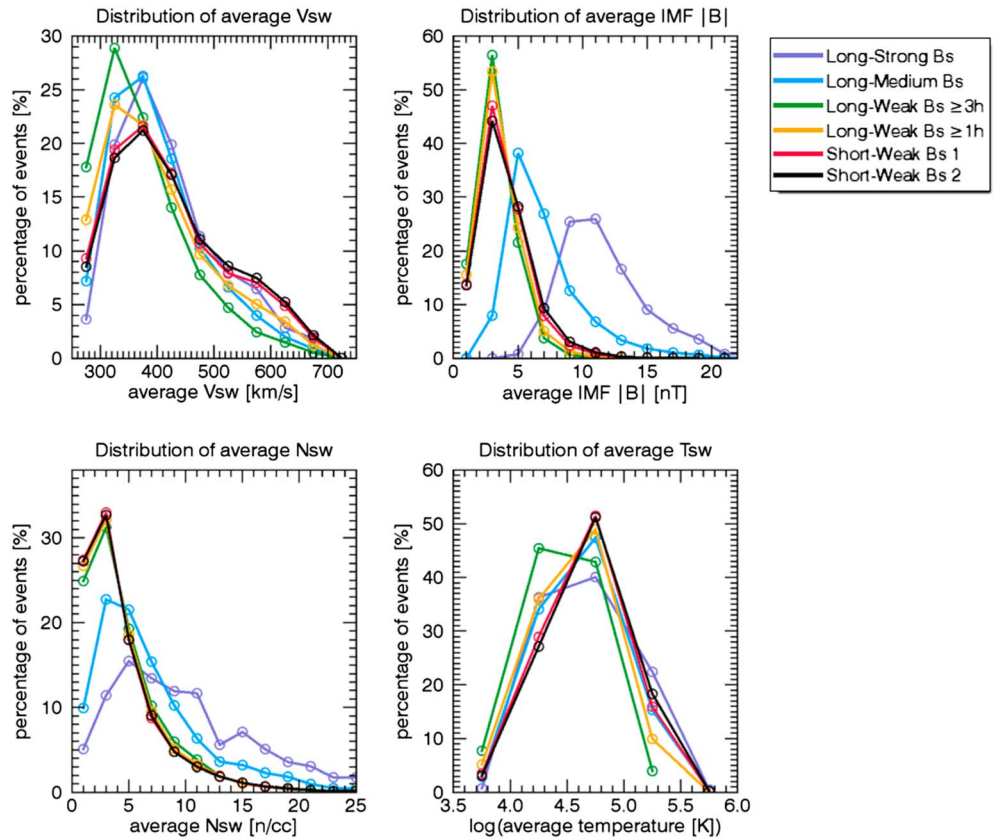
For  $B_s$  events of each class, we analyzed the solar wind conditions including IMF total  $B$ , ion number density, flow speed, and proton temperature. The results are summarized in Figure 7, which shows the distributions of the four solar wind parameters averaged over each  $B_s$  event interval with the six classes distinguished by different line colors.

First, the distribution of average solar wind speed indicates that the two Short-Weak  $B_s$  classes (red and black) have higher occurrence rates for the range of  $V_{sw}$  at ~500–700 km/s than the others do. This is mostly due to coronal hole high-speed streams, which contribute to the two Short-Weak classes to some extent (as shown in Figure 2). The distribution of the IMF total  $B$  indicates striking differences between the Long-Strong (violet) and Long-Medium (blue)  $B_s$  classes and the other four weak intensity classes. This is reasonable considering that a larger magnitude of  $B_z$  contributes more to total  $B$ . This feature is similar, but less prominent, for the solar wind ion density distribution. Finally, the only notable feature in the temperature distribution is that for Long-Weak  $B_s \geq 3$  h (green), the peak percentage occurs at a lower average temperature than those of the other classes. One would expect the temperature to be lower than the normal solar wind for the  $B_s$  events if they are related to ICMEs. This may be expected, but not seen in Figure 7, particularly for the Long-Strong  $B_s$  class to which the ICME contribution is largest. The discrepancy is due to the fact that the average in Figure 7 has been obtained over the entire number of  $B_s$  events, both the events found inside the ICME intervals and those found in other regions such as the higher-temperature sheath region ahead of ICMEs. We confirmed (not shown here) that when averaged over the  $B_s$  events that are related to ICMEs only (excluding the preceding sheath region), the temperature is indeed lower than that of the normal solar wind.

### 3.4. Geoeffectiveness in Terms of $AL$ , $SYM-H$ , and Geosynchronous Electron Fluxes

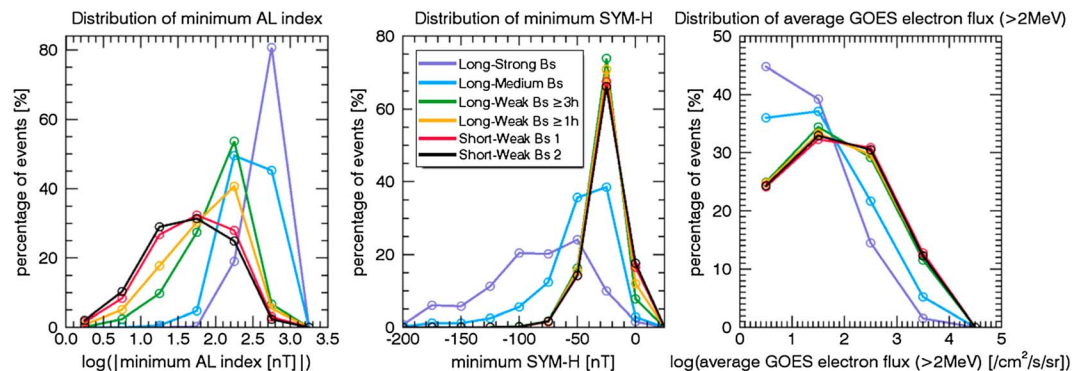
In order to assess the relative significance of  $B_s$  events from the viewpoint of their geoeffectiveness and compare this among the six  $B_s$  classes, we investigated the distributions of  $AL$ ,  $SYM-H$ , and geosynchronous  $>2$  MeV electron fluxes during  $B_s$  events of each class. The results are summarized in Figure 8, which





**Figure 7.** Distribution of solar wind parameters, solar wind speed  $V_{sw}$ , IMF  $B$ , ion number density  $N_{sw}$ , and proton temperature, as averaged over the duration of each of the  $B_z$  events.

shows clear differences among the six  $B_z$  classes. First, the left panel in Figure 8 shows the  $AL$  index distributions. Here we point out that  $AL$  is supposed to reflect the westward electrojet and it can serve as a substorm indicator when it decreases in association with a newly formed wedge current system due to substorm occurrence. At the same time, we stress that there are limitations to  $AL$  as a substorm indicator primarily due to the small number of magnetometer stations used (10–12 stations usually) and their uneven spatial distribution [e.g., *Gjerloev et al.*, 2004]. Nevertheless, since the substorms are characterized by multiple signatures in a complex way throughout the near-Earth magnetosphere, ionosphere, and on



**Figure 8.** Distribution of  $AL$ ,  $SYM-H$ , geosynchronous electron ( $> 2$  MeV) fluxes during  $B_z$  events of each class. For  $AL$  and  $SYM-H$ , the minimum values are taken during each  $B_z$  event interval, and for the geosynchronous electron fluxes measured on GOES spacecraft, the averages are taken over each  $B_z$  event interval.

**Table 3.** Percentages of the Geoeffectiveness by the  $B_s$  Events of Different Classes

$B_s$ Classes	Substorm <sup>a</sup> (%)	Storm <sup>b</sup> (%)	e Flux Disturbance <sup>c</sup> (%)
Long-Strong $B_s$	100	67.3	2.3
Long-Medium $B_s$	99.6	24.4	7.5
Long-Weak $B_s \geq 3h$	88.8	2.3	14.9
Long-Weak $B_s \geq 1h$	78.4	2.0	16.2
Short-Weak $B_s 1$	65.6	2.0	16.8
Short-Weak $B_s 2$	60.4	2.0	16.2

<sup>a</sup>Based on the criterion that  $AL_{\min} \leq -50$  nT [Allen and Kroehl, 1975].<sup>b</sup>Based on the criterion that  $SYM-H_{\min} \leq -36$  nT [Loewe and Prolss, 1997].<sup>c</sup>Based on the criterion that  $>2$  MeV electron flux level at geosynchronous orbit  $\geq 2 \times 10^3$  /cm<sup>2</sup>/s/sr [Baker, 2000; Kim et al., 2006].

the ground [e.g., Lui, 1996; Ohtani, 2004; H.-J. Kim et al., 2008; Lee et al., 2010; Akasofu et al., 2010; Sergeev et al., 2012], the use of  $AL$  is still practical for statistical studies. In the left panel of Figure 8, it is clear that for the Long-Strong  $B_s$  class, the distribution is greatly concentrated at the highest  $AL$  values compared to the other  $B_s$  classes. For the weaker and shorter  $B_s$  classes, the  $AL$  distribution is broader and its peak shifts toward lower  $AL$  values. In particular, for the two Short-Weak  $B_s$  classes, the peak occurs at  $\log |AL| \sim 1.7$  corresponding to  $AL \sim -50$  nT. We point out that there is no unique way to define a threshold value of  $AL$  beyond which a substorm is certain to occur [e.g., Newell and Gjerloev, 2011]. Nevertheless, if we follow the simple criterion by Allen and Kroehl [1975], an  $AL$  value below  $-50$  nT generally represents a beginning limit for substorm occurrence. Therefore, if we adopt this criterion, more than half of the Short-Weak  $B_s$  class events are associated with substorms having  $AL = -50$  nT or stronger. This implies that even a weak southward IMF  $B_z$  can often cause a substorm of weak-to-modest intensity. This finding is qualitatively consistent with the results shown in Kamide et al. [1977]. Table 3 summarizes the geoeffectiveness based on this criterion. Incidentally, it is well known that the onset of substorms is often preceded by a  $\sim 30$  min energy buildup growth phase in the magnetotail. Our results imply that the duration of the  $B_s$  events of the two Short-Weak classes may be sufficient for setting such a substorm growth phase.

The  $SYM-H$  distributions for the Long-Strong and Long-Medium classes show broad distributions, indicating some fraction of  $B_s$  events occur in association with major magnetic storms. However, the peaks are present in the bin between  $-50$  nT and  $-30$  nT. For the other four  $B_s$  classes, the  $SYM-H$  distributions are much narrower with the peaks at slightly above  $-30$  nT. Loewe and Prolss [1997] classified more than 1000 storms according to the intensity of  $Dst$  as weak ( $Dst_{\min} = -30$  nT to  $-50$  nT), moderate ( $Dst_{\min} = -50$  nT to  $-100$  nT), strong ( $Dst_{\min} = -100$  nT to  $-200$  nT), severe ( $Dst_{\min} = -200$  nT to  $-350$  nT), and great ( $Dst_{\min} < -350$  nT). The median values of  $Dst_{\min}$  are  $-36$  nT,  $-68$  nT,  $-131$  nT,  $-254$  nT, and  $-427$  nT, respectively, for the five groups. In a more recent work on geoeffectiveness of halo coronal mass ejections by Gopalswamy et al. [2007], they defined that a halo CME is geoeffective if  $Dst \leq -50$  nT. They found that the median  $Dst$  values for backside ( $-41$  nT),  $F$  limb ( $-58$  nT) and disk halos ( $-97$  nT) are close to the corresponding values for weak ( $-36$  nT), moderate ( $-68$  nT), and strong ( $-131$  nT) storms reported by Loewe and Prolss [1997]. We adopt the median  $Dst$  value for weak storms,  $-36$  nT, reported by Loewe and Prolss [1997] as the threshold for geoeffectiveness in the sense of storm occurrence. With this threshold, the percentage of the geoeffective  $B_s$  events is largest, 67.3%, for the Long-Strong  $B_s$ , becomes lower, 24.4%, for the Long-Medium  $B_s$  class, and is negligible for the other four weak classes. The statistics of the estimated geoeffectiveness is summarized in Table 3. Note that we use  $SYM-H$ , which is practically a high-resolution (1 min)  $Dst$  index and is essentially calculated by the same method as for  $Dst$  except for the use of different ground stations [Iyemori, 1990; Wanliss and Showalter, 2006].

Finally, the distributions of geosynchronous  $>2$  MeV electron fluxes also show a strong contrast between the Long-Strong and Long-Medium classes and the other four weaker classes. The contrast is characterized by the feature that the electron flux enhancements occur more often for the  $B_s$  events of the four weaker classes. Baker [2000] reported that two spacecraft failures in 1998 occurred at the times of the  $>2$  MeV electron flux level of  $\sim 2 \times 10^3$  /cm<sup>2</sup>/s/sr. Kim et al. [2006] used this as an approximate threshold value to define a geosynchronous relativistic electron event (GREE) in their statistical study. We are not aware of any obvious physical reasoning to define a threshold flux level, above which one can define GREEs. However, if we adopt the

criterion by *Kim et al.* [2006], ~15%–17% of the  $B_s$  events for the four weaker classes are GREEs. For the Long-Strong and Long-Medium classes, the association with the  $B_s$  events is more concentrated at low flux levels. If we continue to adopt the threshold given by *Kim et al.* [2006], only 2.3% of the  $B_s$  events of the Long-Strong class and 7.5% of the Long-Medium class correspond to GREEs, respectively (see Table 3.) This result implies that the dropout effect (loss) of the outer radiation belt electron fluxes is more important than acceleration effects (source) for the  $B_s$  conditions of the two classes, as will be discussed further in section 4. In a short summary, the occurrence rate of the geosynchronous electron disturbance due to the four weak  $B_s$  classes is larger by ~7 times than that due to the Long-Strong  $B_s$  class and by ~2 times than that due to the Long-Medium  $B_s$  class.

#### 4. Summary and Discussion

In this paper, we identified southward directed  $B_s$  events using IMF  $B_z$  values in GSE coordinates. We associated these with solar wind parameters over 20 years measured near the Earth, derived the statistical features of these events and determined their corresponding geoeffectiveness in terms of substorms, storms, and geosynchronous relativistic electron flux disturbances. For this study, we defined six classes of  $B_s$  events according to the intensity and duration of southward  $B_z$  conditions, which are Long-Strong, Long-Medium, Long-Weak  $\geq 3$  h, Long-Weak  $\geq 1$  h, Short-Weak 1, and Short-Weak 2. We aimed to compare the characteristics of  $B_s$  events among the six classes. The main results are summarized as follows.

1. The occurrence rate of independent  $B_s$  events is lower for longer lasting and more intense  $B_s$  classes and rapidly increases toward weaker and shorter  $B_s$  classes (Table 1).
2. Our Long-Strong  $B_s$  class has the highest association (~40%) with ICMEs. Coronal hole high-speed streams are more likely associated with the weaker and shorter  $B_s$  classes. However, many of the  $B_s$  events in nearly all classes are found in regions other than ICMEs and coronal hole high-speed streams (Figure 2).
3. Alfvén waves are present for ~17% (for Long-Strong class) and ~60% (for Short-Weak class 2) of the  $B_s$  events. Many of them originate during the times other than ICMEs and coronal hole high-speed streams (Figures 4 and 5).
4. There is a clear distinction of the solar cycle dependence of  $B_s$  among the different  $B_s$  classes: while the occurrence rate of the Long-Strong and Long-Medium  $B_s$  class events is in phase with the solar cycle, the Long-Weak  $B_s$  events show negative correlation between their occurrence frequency and the solar cycle. For the Short-Weak class  $B_s$  events, the correlation is much less obvious (Figure 6).
5. The Long-Strong and Long-Medium class  $B_s$  events are mostly associated with intense *AL* and *SYM-H* conditions. In contrast, the Short-Weak class  $B_s$  events can be geoeffective in terms of triggering substorms of weak-to-medium intensity although they are unlikely to be associated with major magnetic storms. The enhancement of geosynchronous relativistic electron fluxes is more often caused by the Long-Weak and Short-Weak class  $B_s$  events than the Long-Strong and Long-Medium  $B_s$  events (Figure 8 and Table 3).

As the most important conclusion, we emphasize the potential importance of the weakly (in intensity and/or duration) southward  $B_z$  conditions for triggering substorms of weak-to-moderate intensity and for resulting in an enhancement of geosynchronous relativistic electron fluxes. This effect should not be underestimated because the weak  $B_s$  events have a much higher occurrence rate, regardless of solar cycle phase, compared to the strongly southward IMF  $B_z$  conditions which can cause the major storms and substorms. This conclusion implies that there is a necessity that the models for IMF  $B_z$  prediction [e.g., *Jackson et al.*, 2015] should cover a wide range of southward  $B_z$  conditions.

Our methodology with six classes of  $B_s$  events misses other  $B_s$  conditions that do not satisfy our selection criteria and is different from that of previous work by others [e.g., *Zhang and Moldwin*, 2014] (see the last paragraph of this section for a more detailed comparison with the work by *Zhang and Moldwin* [2014]). Nevertheless, in our classification of  $B_s$  events, the six classes are in contrast in terms of duration and/or intensity of  $B_s$ . Additionally, as a main goal, the present work demonstrates the extent to which the relatively weaker or shorter  $B_s$  conditions are significant relative to the stronger and longer  $B_s$  events from the perspectives of their occurrence rates and geoeffectiveness.

We found that some fraction of  $B_s$  is due to Alfvén waves. It is well known that a continuous Alfvén wave train is the main feature of coronal hole high-speed streams which can last several days as observed near the Earth

and cause the high-intensity long-duration continuous AE activity (named HILDCAA) [e.g., *Tsurutani and Gonzalez, 1987*]. However, in our study, many of the Alfvén waves are present at regions other than coronal-hole high-speed streams, a further study of which is left for a future project. Incorporation of an Alfvén wave contribution to  $B_s$  in the models for IMF  $B_z$  prediction may be worthwhile to consider. Also, while ICMEs and coronal hole streams represent large-scale structures, Alfvén waves can be of a small scale. Our study compares geoeffectiveness mainly in terms of  $B_s$  without considering the effect related to the differing spatial scales. A future study on this issue is worthwhile to perform.

We found that geosynchronous relativistic electron flux enhancements more often occur during weakly southward  $B_z$  conditions than for the more strongly southward  $B_z$  conditions which usually lead to major storms. Statistically, the high levels of geosynchronous relativistic electron flux are caused mostly by Corotating interaction region (CIR)-driven storms rather than by CME-driven storms [*Borovsky and Denton, 2006*], and high-speed streams have long been known to be a critical factor to determine the flux level of geosynchronous relativistic electrons [*Paulikas and Blake, 1979*]. Thus, the weakly southward  $B_z$  embedded in high-speed streams should be a condition favorable for geosynchronous relativistic electron flux enhancement. However, details of the net flux change involve more complex processes than the simple solar wind speed. This can be explained as follows. First, the geosynchronous electron flux level as observed at satellites is the result of delicate balance between particle source and loss [e.g., *Reeves et al., 2003; Hwang et al., 2004; Kim et al., 2006; Lee et al., 2013*]. In particular, the geosynchronous relativistic electron fluxes vary during the evolution of the storms. During the storm main phase the geosynchronous electrons can suffer from major flux dropouts by drift loss through the magnetopause into the solar wind leading to shrinkage of the outer radiation belt [e.g., *K. C. Kim et al., 2008; Kim et al., 2010*]. In addition, atmospheric precipitation by plasma waves can contribute to flux dropout [e.g., *Miyoshi et al., 2008; Li et al., 2014*]. Therefore, a strong  $B_s$  event, if identified during the times overlapping the storm main phase, must be associated with decrease of geosynchronous relativistic electron fluxes. During the recovery phase of storms, the geosynchronous electron flux may or may not recover to the prestorm level depending on the significance of electron acceleration relative to loss effects. Consequently, for a strong  $B_s$  event which is identified during the times overlapping the storm recovery phase, the change of geosynchronous relativistic electron fluxes may either increase or decrease. Therefore, on average the geosynchronous relativistic electron flux associated with strong  $B_s$  conditions such as during the storm times is not necessarily high. In contrast, for the weakly southward  $B_z$  conditions which do not cause a major storm but can often lead to substorms, the trapping region of the outer radiation belt electrons can expand beyond geosynchronous orbit distances and thus have no loss due to electron drift into the magnetopause. This condition favors keeping the geosynchronous electron flux maintained at a nominal (or even enhanced) level unless another major loss process exists. The substorms that occur during the weak  $B_s$  events can supply the seed particles to geosynchronous altitude which can generate plasma waves responsible for particle acceleration. Our result in Figure 8 showing the frequent occurrence of the flux enhancement for the Long-Weak and Short-Weak  $B_s$  events implies that often during these times loss processes do not dominate and instead during substorms source effects contribute to the geosynchronous electron flux to a larger extent.

Recently, *Zhang and Moldwin [2014]* performed a similar work to ours. Compared to their work, our work is distinguished in the following way. First, while they distinguished several well-defined solar wind structures in association with their  $B_s$  events such as magnetic cloud, ICMEs without MC signatures, stream-interaction regions, interplanetary small-scale magnetic flux ropes, and shocks, our groups are ICMEs without distinguishing well-defined MC from the others, and high-speed streams without excluding the corotating interaction regions ahead of them. From the results in *Zhang and Moldwin [2014]*, we note that the fractional number of the interplanetary small-scale magnetic flux ropes and shocks is small and would be included in “other times” by us. Second, there is distinction between our study and that of *Zhang and Moldwin [2014]* in categorizing  $B_s$  events. While *Zhang and Moldwin [2014]* show nine classes of  $B_s$  events in their Figure 3 and Table 1, our work employs six classes of  $B_s$  events. There is little overlap between the two studies in the categorization of  $B_s$  events. While there is no reason for having to accept some specific set of definition of  $B_s$  events against others, it is clear from comparisons between Figure 3 of *Zhang and Moldwin [2014]* and Figure 6 of ours that the solar cycle dependence of  $B_s$  events of a specific class depends on specific definition of  $B_s$  events. Most notably, we found an anticorrelation between the  $B_s$  events of the Long-Weak  $B_s \geq 1$  h and



$\geq 3$  h classes and the solar cycle as shown in Figure 6 above, which was not explicitly shown in Zhang and Moldwin [2014]. Fourth, we have assessed the contribution by Alfvén waves for all of the six classes of  $B_z$  events, and this was not done in Zhang and Moldwin [2014]. Though studied in Zhang et al. [2014], this was done only for southward IMF  $B_z$  events defined by  $t > 1$  h and  $B_z < -5$  nT, unrelated to any well-defined solar wind structure. Lastly and most importantly, while Zhang and Moldwin [2014] examined the geoeffectiveness of various types of  $B_z$  events in relation mostly to the magnetic storms with only a limited study of substorms, we checked three aspects of geoeffectiveness, namely, magnetic storms, substorms, and geosynchronous electron flux disturbances. We further found an important contribution from weakly southward  $B_z$  events to substorms and geosynchronous electron flux disturbances. Overall, we find that our work is complementary, rather than contradictory, to the works by Zhang and Moldwin [2014] and Zhang et al. [2014], although a direct quantitative comparison cannot be made precisely due to the different methods used by these authors and us.

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