Quantum States and Spectra of Gases

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Abstract

This experiment introduces and works with the emission spectrum of various gases, comparing known data with experimentally obtained data. When an electron transitions from different quantum states, it releases a photon, containing a packet of light and energy. Different gases found in discharge tubes were charged with a high voltage to emit photons to analyse their behaviors. The emission of the photon directly correlates to the spectral lines on the emission spectrum. Results show that the emission lines of hydrogen and the helium cation behave differently. The emission spectrum was also used to portray how humans perceive colors and how it differentiates the colors from white light. Colored solutions had white light shone at them, and the results show that the highest point of transmission corresponded to the color of the solution on the color spectrum.

I Introduction

When an electron transitions from a lower energy state to a higher energy state, it releases a photon, which is comprised by a spectrum of multiple waves of electromagnetic radiation. The energy of this photon is the energy difference between the higher and lower quantum states of the chemical compound. Since each element's emission spectrum is unique, it can be used to study the behavior of different elements and identify unknown elements.

In this experiment, a gas discharge tube is placed in a high voltage source that causes electrons to be heavily charged and travel rapidly toward the anode of the tube. When these electrons collide with the gas particles in the tube, the gas particles can become positively charged (ionized); consequently, the gas particle loses an electron, which will also travel towards the anode of the tube. Contrarily, the now positively charged gas particle (cation) will travel toward the cathode of the tube. The process of ionization causes the total energy of the neutral atom to increase; in effect, the electron will transition to a different energy level, leaving a vacancy in its previous electron shell. In the aforementioned process, energy is released in the form of light (photon), and the ion energy is decreased. This light can be represented as a spectrum of electromagnetic radiation across different wavelengths.

When the element found in the gas discharge tube is heated, it will begin to glow (glow discharge), releasing energy in the form of light. With the help of a spectrometer, the different wavelengths in the components of the light, emitted from the photon, can be separated and placed on the spectrum using external software, PASCO Spectrometry.

Quantum physicists Max Planck and Niels Bohr uncovered a relationship between the energy of the photon released and the transition of the particle between the quantum states. For a hydrogen atom or any particle similar to a hydrogen atom, such as He⁺, the energy of the photon emitted during the particle transition from quantum state m to n can be calculated as follows:

$$hf = R_{EH}Z^2(\frac{1}{m^2} - \frac{1}{n^2}) \tag{1}$$

The empirical formula for a visible spectrum of hydrogen was discovered by Johann Balmer, named the Balmer series:

$$hf = R_{EH}(\frac{1}{2^2} - \frac{1}{n^2}) \tag{2}$$

Finally, to calculate the energy of any stationary state of the hydrogen (or hydrogen-like) atom, Bohr's postulates are used to obtain the formula for E_n :

$$E_n = \frac{Z^2 k_e e^2}{2n^2 a_0} \approx -\frac{13.6Z^2}{n^2} \tag{3}$$

where Z is the number of protons and n is the quantum state of the atom.

To explain why the energy of a hydrogen atom at the n-th energy level is a negative value, it should be understood that the nucleus of the atom exerts an electrostatic force on the electron to hold the atom together. The energy of the electron consists of both potential and kinetic energy, so that when the electron is infinitely far away from the nucleus, the potential energy is zero. However, as the electron becomes infinitely close to the nucleus, the potential energy approaches negative infinity.

The above observations can be expressed as limits, where r is the distance between the nucleus and the electron:

$$\lim_{r \to \infty} PE = 0$$

$$\lim_{r \to 0} PE = -\infty$$

The reason for the value of the energy being negative is because the electron requires energy to be supplied to it to overcome the strong, attractive force of the nucleus if it wishes to escape the atom.

II Materials and Methods

II.I Materials:

For the experiments, the equipment used were as follows:

- Mercury (Hg), Hydrogen (H), Helium (He), and "unknown" gas discharge tubes
- Computer with USB port
- PASCO wireless spectrometer PS-2600
- Fiber optics cable PS-2601 (probe)
- High voltage power supply
- Software Spectrometry (resolution \pm 2-3 nm)
- Cuvettes of different water-diluted substances
- Clamp

II.II Methods

II.II.1 Gas Discharge Tube (Part 1):

- 1. Turn the spectrometer on using the ON/OFF Button and make sure all three LED bulbs are glowing.
- 2. Make sure the power supply for a discharge tube is unplugged, and carefully insert a Hg gas tube into the power supply to plug it in on both ends.
- 3. Place the rectangular end of the fibre optics cable in the cuvette opening of the spectrometer, making sure the arrow on top of the rectangular end showed the direction toward the built-in detector.
- 4. Fix the probe end of the cable in the clamp vertically above the tube and point the probe to the tube center at about 0.5 cm from the tube.
- 5. Plug the wire of the power supply into the net. Turn on the power supply with the red switch on its side.
- 6. Used the Spectrography software and begin to analyse the light of the gas tube. Start a recording of the test spectrum, and adjust parameters to get the optimal recording.
- 7. Record data for about 2 minutes, and then turn off power supply.

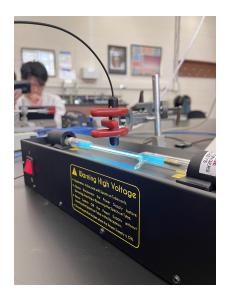


Figure 1: Mercury gas discharge tube with the high voltage source powered on; the clamp is holding the fiber optic probe that detects the emission from the tube.

II.II.2 Gas Discharge Tube (Part 2):

- 1. Repeat steps 1-5 of part 1 with the hydrogen tube instead of the mercury tube.
- 2. Calculate quantum state energies of hydrogen at n = 3, 4 and 5, and expect energies of the Balmer series spectral lines using Equation 2.
- 3. Convert energies of spectral lines to wavelengths, and compare to experimental values.

II.II.3 Gas Discharge Tube (Part 3):

- 1. Repeat steps 1-5 of part 1 with the helium tube instead of the mercury tube.
- 2. Determine wavelength of measured values, compared them to the reference wavelengths at each electron configuration, and calculate change in quantum number.

II.II.4 Gas Discharge Tube (Part 4):

- 1. Repeat steps 1-5 of part 1 with the unknown gas tube instead of the mercury tube.
- 2. Determine the unknown gas through calibration values, and the standard spectra stored on the software.

II.II.5 Dyed Substance Solutions (Part 5):

- $1.\,$ Remove probe, and calibrate the spectrograph to dark.
- 2. Insert the clear cuvette into the cuvette holder, with the glossy side facing the white light and calibrate the spectrograph to the clear cuvette.
- 3. Switch to the blue cuvette and measure levels of transmittance and absorbance with the blue cuvette under the white light.
- 4. Repeat step 3 with the red and green cuvette.

II.II.6 Dyed Substance Solutions (Part 6):

- 1. Insert yellow cuvette with the glossy side facing the violet and blue light. Find the absorption edge, which is the wavelength at which absorption jumps significantly.
- 2. Use the wavelength of the absorption edge to calculate the band gap energy of the dye.
- 3. Determine whether the yellow dye demonstrates semiconductor properties.

III Data and Analysis

III.I Part I:

When acquiring the data, there were several variables to take into account, such as the distance between the fiber optic probe and the gas discharge tube, the number of scans to average and the number of points to smooth the spectral curve.

Increasing the distance between the fiber optic probe and gas discharge tube decreased the intensity of spectral line measured, which is explained by the fact that the probe detects the light waves emitted by the gas discharge tube. As light travels through space, its waves will propagate in every direction, filling up the space it travels through; however, as the light travels away from the source, more area is filled, which means the light intensity must be distributed evenly throughout the area. Thus, the intensity of the light waves decrease proportionally with the increase of distance from the source.

If D is the distance from the source of light and I is the intensity of the light wave, the relationship can be modelled as:

$$D \propto \frac{1}{I^2}$$

Adjusting the number of scans to average did not appear to have a great effect on the spectral curve, but adjusting the number of points to smooth decreased the intensity of spectral curve. This is because there are not adequate data points being collected by the software, which forces to smooth by flattening.

Color	Violet	Violet	Blue	Green	Yellow	Yellow
Expected λ (nm)	404.6565	407.7837	435.8328	546.0735	576.9598	579.0663
Expected Energies (eV)	3.0663	3.0428	2.8469	2.2722	2.1506	2.1427
Experiment λ (nm)	403.467	407.178	437.708	543.389	574.123	576.350
Experiment Energy (eV)	3.0753	3.0473	2.8347	2.2834	2.1612	2.1528

Table 1: Experimental and calculated data from Experiment 1.

The slope and y-intercept of the line of best fit is (from Figure 2 in Appendix B):

• Slope: $m = 1.02 \pm 0.02$; Y-Intercept: $b = -6.59 \pm 7.20$

The goodness of fit criteria applied to this line of best fit are reduced chi-squared and residuals:

• Reduced chi-squared: $\chi^2 = 3.222$; Residuals: $\sum_{i=1}^{N} r_i = -11.147$

The slope and y-intercept of the line of best fit is (from Figure 3 in Appendix B):

• Slope: $m = 0.99 \pm 0.01; \text{ Y-Intercept: } b = 0.030 \pm 0.022$

The goodness of fit criteria applied to this line of best fit are reduced chi-squared and residuals:

• Reduced chi-squared: $\chi^2=0.016$; Residuals: $\sum_{i=1}^N r_i=-0.002$

III.II Part II:

Quantum State (n)	3	4	5
Energy at Quantum State n (eV)	-1.511	-0.850	-0.544
Balmer Series at Quantum State n (eV)	1.89	2.55	2.86
Theoretical λ (nm)	660.83	487.20	435
Experiment λ (nm)	653.05	485.303	_

Table 2: Calculated energies and wavelengths of different quantum states of hydrogen gas.

Relative intensity of the lines may vary for the same atoms from experiment to experiment because the spectrum reflects the movement from one energy level to another. Since there are several energy levels, the intensity between each combination may be different. For instance, Balmer predicted one set of intensities for hydrogen, based specifically off of an initial energy level of n = 2. The Balmer equation can predict other combinations as well, but may give different results.

III.III Part III:

For a helium ion transitioning from ground quantum state m = 2 to n = 3, the energies are calculated as follows:

•
$$E_2 = -\frac{13.6(2)^2}{2^2} = -13.6 \text{ eV}$$

•
$$E_3 = -\frac{13.6(2)^2}{3^2} = -6.044 \text{ eV}$$

which translates to the energy released by the photon to be ~ 7.55 eV. The wavelength corresponding to this energy released can be calculated as follows: $E = \frac{hc}{\lambda} \rightarrow \lambda = \frac{hc}{E} = \frac{(6.626 \times 10^{-34})(3 \times 10^8)}{7.55 \times 1.6 \times 10^{-19}} = 164 \text{ nm}$

The calculated and measured wavelengths are not equal because $\lambda = 164$ nm is not on the visible light spectrum and cannot be detected as a point of photon release. The most defined point at which the photon is released is $\lambda = 588$ nm, which also corresponds to the helium ion transitioning from quantum states n = 3 to m = 2.

Electron	Electron	Reference	Reference	Measured	Energy	Energy	Δn	Δl	ΔJ
configu-	configu-	Inten-	wave-	wavelength	of the	of the			
ration of	ration of	sity in	length	of the	quantum	quantum			
the initial	the final	arbitrary	of the	emission	state, E_n	state,			
state	state	units	emission	line, $\lambda_{\rm exp} \pm$	(eV)	E_m , (eV)			
			line, λ_r	$\Delta \lambda_{\rm exp} \ ({\rm nm})$					
			(nm)						
1s2p (1)	$1s^2(0)$	1000	58.43339		-3.34	-24.57	1	1	1
1s3s (1)	1s2p (2)	200	706.5190	702.098 ± 3	-1.84	-3.60	1	-1	-1
1s3p (1)	1s2s (1)	500	388.8648	391.919 ± 3	-1.56	-4.75	1	1	0
1s3d (3)	1s2p (2)	500	587.5621	584.151 ± 3	-1.49	-3.60	1	1	1
1s3d (2)	1s2p (1)	100	667.8151	664.594 ± 3	-1.492	-3.35	1	1	1
1s3p (1)	1s2s (0)	100	507.5678	500.148 ± 3	-1.51	-3.95	1	1	1
1s4d (1)	1s2p (2)	200	447.1480	448.362 ± 3	-0.83	-3.60	2	1	-1

Table 3: Experimental and calculated data of the emission lines of He^+ and its transition from quantum state n=3 to m=2.

Through experimentation, it can be concluded that 6 spectral lines contribute to the transition of the helium ion from n = 3 to m = 2. The hydrogen emission spectrum only has 2 spectral lines since its singular electron cannot travel between different quantum states as easily, the electron configuration of hydrogen being $1s^{1}$.

Equation 1 cannot be used to calculate the energy of a neutral helium atom since it can only be applied to hydrogen-like atoms. In other words, a neutral helium particle has two electrons orbiting the nucleus while a helium cation only has one electron orbiting the nucleus, qualifying it as a hydrogen-like atom.

III.IV Part IV:

Based on the spectrography of the unknown gas, we have come to the conclusion that the gas is likely krypton. As per Table 4, the actual, calculated, and measured values of krypton wavelengths were very similar. The notable difference between the actual and calculated is the fact that in our experiment, there was no peak at the wavelength of 769 mm. However, the other measured peaks were within a few nanometres of the actual measures. In addition, the ratio of intensity between the three peaks that were measured in this experiment was roughly 18:15:9, which is not perfectly accurate to the actual ratio, of roughly $1:\frac{1}{2}:\frac{1}{4}$, but is similar (ratio is $1:\frac{5}{6}:\frac{1}{2}$).

The values of the expected wavelength below are chosen from Appendix 2 for **Krypton** (**Kr**). The first 5 entries were ignored because the spectrometer did not detect any release in energy until $\lambda = 757.604$ nm.

Expected Wavelength, λ (nm)	Observed λ (nm)	Calibrated λ (nm)	Relative Intensity (to strongest)
759	757.604	766.166	0 (17.829)
769			
810	810.2	819.814	-2.9 (14.929)
826	826.803	836.749	-8.775 (9.054)

Table 4: Compiled and measured data used to determine the unknown gas.

III.V Part V:

When white light passes through a coloured solution, bands of the spectrum are absorbed by the solution, and the remaining wavelengths are transmitted through the solution. The transmitted wavelengths are what gives the coloured solution its particular colour. In the results of this part of the experiment, the spectrography of each coloured solution shows a dip in

absorbance and a peak in transmittance over the range of that colour in the spectrum of white light. For instance, for the blue solution, the range over which absorbance dips and transmission peaks is from 432 to 481 nm, as per Table 5. Additionally, the range for the red solution is from 635 nm onward, and the green solution has a range of from 500 to 523 nm. These ranges can all be seen on the spectrography for the respective solutions, in Figures 11 and 12 of the appendix.

Solution Color	Measured Transmittance (%)	Measured Absorption	Calculated Absorption	λ (Peak)
Blue	15.284	0.837	0.816	456.128
Red	99.5	≈ 0	0.0022	635
Green	9.662	1.063	1.015	508.726

Table 5: Absorption and transmittance properties of the colored solutions.

The wavelength values at which the peak transmission occurs also corresponds to its color on the spectrum. For instance, at $\lambda = 456.128$ nm, it occurs at the blue strip on the color spectrum, which is why the blue color of the solution can be seen. The same reasoning can be applied to the other two colored solutions as well.

III.VI Part VI:

With 405 nm excitation: absorption edge $\rightarrow \lambda \in [500, 540]$. With 500 nm excitation: absorption edge $\rightarrow \lambda \in [495, 520]$.

Band gap energy @ 405 nm and @ 500 nm, respectively:

$$E_g = \frac{1240}{\lambda} = [2.296, 2.48], [2.385, 2.505]$$

Band gap for 3 types of materials — metals: theoretically 0 eV, insulators: > 3 eV, semi-conductors: < 3 eV. The band gap for the yellow dye falls in the range of < 3 eV, thus, this substance is a semi-conductor.

Semiconductors are characterized by having a certain level of electrical conductivity, making them distinct from insulators, but not to the extent that traditional conductors such as certain metals do. Insulators have a large band gap, generally greater than 3 eV. True conductors have no band gap. Semiconductors generally have a much smaller band gap than insulators, generally around 1 eV. In the case of the yellow solution, the band gap ranges from 2.3 to 2.5 eV.

IV Conclusion and Errors:

Through analysis of various spectrographs, several conclusions were reached. Intensity of spectral emissions varied as movement between energy levels occurred, as observed through analysis of hydrogen. Through observation of the intensity at expected wavelengths, an unknown gas could be correctly identified. With a high level of certainty, the spectrum of transmittance of different solutions could be observed. Lastly, a solution could be correctly identified as a conductor, insulator, or semiconductor based on band gap

The main source of error in this lab was the missing peak of the krypton spectrograph for determining the unknown gas. This most likely occurred because the wavelength that the peak should have been at was very close to another wavelength - 769 nm for the missing and 759 nm for the nearby peak. On the spectrograph, the peak was quite wide, and therefore the peak that should have been at 769 may have been absorbed into the nearby more intense peak. However, the other three peaks of the krypton spectrograph were very close to the expected wavelength values, and were of the expected intensity ratio so that it could be confidently predicted that the unknown gas was krypton.

A Appendices

A.I Additional Calculations:

A.I.1 Step 1:

The expected and experimental energies from Table 1 were calculated using the energy of a photon equation, where $h = 6.626 \times 10^{-34} \text{ J} \cdot \text{s}$ and $c = 3.0 \times 10^8 \text{ m/s}$:

$$E = \frac{hc}{\lambda} \times 6.242 \times 10^{18} \text{ eV}$$

Expected Energy	Experimental Energy		
$\mathbf{E}_1 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{404.6565} \times 6.242 \times 10^{18} = 3.0663 \text{ eV}$	$\mathbf{E}_1 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{403.467} \times 6.242 \times 10^{18} = 3.0753 \text{ eV}$		
$\mathbf{E}_2 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{407.7837} \times 6.242 \times 10^{18} = 3.0428 \text{ eV}$	$\mathbf{E}_2 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{407.178} \times 6.242 \times 10^{18} = 3.0473 \text{ eV}$		
$\mathbf{E}_3 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{435.8328} \times 6.242 \times 10^{18} = 2.8469 \text{ eV}$	$\mathbf{E}_3 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{437.708} \times 6.242 \times 10^{18} = 2.8347 \text{ eV}$		
$\mathbf{E}_4 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{546.0735} \times 6.242 \times 10^{18} = 2.2722 \text{ eV}$	$\mathbf{E}_4 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{543.389} \times 6.242 \times 10^{18} = 2.2834 \text{ eV}$		
$\mathbf{E}_5 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{576.9598} \times 6.242 \times 10^{18} = 2.1506 \text{ eV}$	$\mathbf{E}_5 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{574.123} \times 6.242 \times 10^{18} = 2.1612 \text{ eV}$		
$\mathbf{E}_6 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{579.0663} \times 6.242 \times 10^{18} = 2.1427 \text{ eV}$	$\mathbf{E}_6 = \frac{6.626 \times 10^{-34} \cdot 3 \times 10^8}{576.350} \times 6.242 \times 10^{18} = 2.1528 \text{ eV}$		

The uncertainties of the slopes and y-intercepts of the linear fits were calculated using the set of standard deviation formulas for the slope of a line:

$$\begin{split} \Delta &= N \sum_i x_i^2 - (\sum_i x_i)^2 \\ s_{y,x}^2 &= \frac{1}{N-2} \sum_i [y_i - (b+mx)]^2 \\ s_m^2 &= N \frac{s_{y,x}^2}{\Delta} \\ s_b^2 &= \frac{s_{y,x}^2 \sum_i x_i^2}{\Delta} \end{split}$$

where the slopes of the line can be expressed as m \pm s_n and the y-intercept of the line can expressed as b \pm s_b.

Uncertainties for Measured vs. Experimental Wavelength	Uncertainties for Measured vs. Experimental Energy
$\Delta = 6 \times (1477235) - (2942.215)^2 = 206787$	$\Delta = 6 \times (41.14) - (15.52)^2 = 5.952$
$\mathbf{s}_{y,x}^2 = \frac{1}{6-2} \times 28.99 = 7.249$	$\mathbf{s}_{y,x}^2 = \frac{1}{6-2} \times 3.25 \times 10^{-4} = 8.13 \times 10^{-5}$
$\mathbf{s}_m^2 = 6 imes rac{7.249}{206787} = 0.0002, \mathbf{s}_m = 0.015$	$\mathbf{s}_m^2 = 6 imes rac{8.13 imes 10^{-5}}{5.952} = 8.2 imes 10^{-5}, \mathbf{s}_m = 0.009$
$\mathbf{s}_b^2 = rac{7.249 imes 2942.215}{206787} = 51.78, \mathbf{s}_b = 7.20$	$\mathbf{s}_b^2 = rac{8.13 imes 10^{-5} imes 41.14}{5.952} = 5.6 imes 10^{-4}, \mathbf{s}_b = 0.024$

The goodness of fits criteria used were reduced chi-squared and residuals. To calculate the aforementioned methods, the formulas are:

$$\chi^2 = \sum_{i=1}^{N} \frac{[y_i - f(x_i)]^2}{\sigma_{y_i}^2}$$
$$r_i = y - y_i$$

Goodness of Fit (Graph 1) Goodness of Fit (Graph 2)
$$\chi^{2} = \sum_{i=1}^{N} \frac{[y_{i} - f(x_{i})]^{2}}{\sigma_{y_{i}}^{2}} = \mathbf{3.22} \qquad \chi^{2} = \sum_{i=1}^{N} \frac{[y_{i} - f(x_{i})]^{2}}{\sigma_{y_{i}}^{2}} = \mathbf{0.0165}$$
$$r_{i} = y - y_{i} = -\mathbf{11.15} \qquad r_{i} = y - y_{i} = -\mathbf{0.0022}$$

A.I.2 Step 2:

Calculating the energies for the different quantum states uses equation 3:

$$E_n = \frac{Z^2 k_e e^2}{2n^2 a_0} \approx -\frac{13.6Z^2}{n^2}$$

•
$$E_3 = -\frac{13.6(1)^2}{3^2} = -1.511 \text{ eV}$$

•
$$E_4 = -\frac{13.6(1)^2}{4^2} = -0.85 \text{ eV}$$

•
$$E_5 = -\frac{13.6(1)^2}{5^2} = -0.544 \text{ eV}$$

Calculating the Balmer series for the different quantum states uses equation 2:

$$hf = R_{EH}(\frac{1}{2^2} - \frac{1}{n^2})$$

• hf (n=3) =
$$13.6(\frac{1}{4} - \frac{1}{3^2}) = 1.89 \text{ eV}$$

• hf (n=4) =
$$13.6(\frac{1}{4} - \frac{1}{4^2}) = 2.55 \text{ eV}$$

• hf (n=5) =
$$13.6(\frac{1}{4} - \frac{1}{5^2}) = 2.86$$
 eV

Calculating the wavelengths uses the energy of a photon equation:

•
$$\lambda_3 = \frac{(6.626 \times 10^{-34})(3 \times 10^8)}{1.88 \times 1.6 \times 10^{-19}} = 660.83 \text{ nm}$$

•
$$\lambda_4 = \frac{(6.626 \times 10^{-34})(3 \times 10^8)}{2.55 \times 1.6 \times 10^{-19}} = 487.20 \text{ nm}$$

•
$$\lambda_5 = \frac{(6.626 \times 10^{-34})(3 \times 10^8)}{2.856 \times 1.6 \times 10^{-19}} = 435 \text{ nm}$$

A.I.3 Step 3:

To calculate the energy of the quantum state, E_n , there is an intermediate step that involves determining the energy of the reference wavelength of the emission line, λ_r .

 E_n is calculated by:

$$E_n = E_m + \frac{hc}{\lambda_r}$$

Let E_i represent the intermediary energy defined by $\frac{hc}{\lambda}$

$$E_{i} = \frac{hc}{\lambda_{r}} \qquad E_{n}$$

$$E_{i1} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{1}}} = 21.23 \qquad E_{r1} = E_{m1} + E_{i1} = 21.23 - 24.57 = -3.34$$

$$E_{i2} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{2}}} = 1.76 \qquad E_{r2} = E_{m2} + E_{i2} = 1.76 - 3.60 = -1.84$$

$$E_{i3} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{3}}} = 3.19 \qquad E_{r3} = E_{m3} + E_{i3} = 3.19 - 4.75 = -1.56$$

$$E_{i4} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{4}}} = 2.11 \qquad E_{r4} = E_{m4} + E_{i4} = 2.11 - 3.60 = -1.49$$

$$E_{i5} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{5}}} = 1.858 \qquad E_{r5} = E_{m5} + E_{i5} = 1.858 - 3.35 = -1.492$$

$$E_{i6} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{6}}} = 2.44 \qquad E_{r6} = E_{m6} + E_{i6} = 2.44 - 3.95 = -1.51$$

$$E_{i7} = \frac{(6.626 \times 10^{-34})(3 \times 10^{8})}{\lambda_{r_{7}}} = 2.77 \qquad E_{r7} = E_{m7} + E_{i7} = 2.77 - -3.60 = -0.83$$

A.I.4 Step 4:

To obtain the calibrated wavelengths, the line of best fit for wavelength was used, where $\lambda_{\text{true}} = 1.02\lambda_{\text{measured}}$ - 6.59.

•
$$\lambda_1 = 1.02 \times (757.604) - 6.59 = 766.166 \text{ nm}$$

- $\lambda_2 = 1.02 \times (810.2)$ 6.59 = 819.814 nm
- $\lambda_1 = 1.02 \times (826.803)$ 6.59 = 836.749 nm

A.I.5 Step 5:

To calculate the absorbance given a transmittance (in percent), the following formula can be used:

$$A = -\log_{10}\left(\frac{T}{100}\right)$$

where A is the absorption and T is the transmittance percentage.

- $A_1 = -\log_{10}(\frac{15.284}{100}) = 0.816$
- $A_2 = -\log_{10}(\frac{99.5}{100}) = 0.0022$
- $A_3 = -\log_{10}(\frac{9.662}{100}) = 1.015$