

Quantum Optics 2016 Spring – Homework Set #4

Date: 5/11/2016

Due: 5/18/2016

1. [Plenio and Knight, Rev. Mod. Phys. (1998)] This exercise is to implement the quantum trajectory (or called quantum jump) approach for the Lindblad master equation

$$\dot{\rho} = -\frac{i}{\hbar}[H_0, \rho] - \frac{\gamma}{2}\mathcal{L}(\rho),$$

where the Lindblad superoperator $\mathcal{L}^-(\rho) = c^\dagger c \rho + \rho c^\dagger c - 2c\rho c^\dagger$ corresponds to the decoherence. We call c the jumping operator that represents a “decay channel”. It can be done by stochastic numerical processes that deal with a pure state $|\psi(t)\rangle$ and the following steps:

- 1) For a small time interval δt , the probability of having a quantum jump (a decay event) is $\delta p \equiv \gamma \delta t \langle \psi | c^\dagger c | \psi \rangle$, where γ is the decay rate. To simulate this, generate a random number r , which is uniformly distributed over $[0, 1]$.
- 2) If $r > \delta p$, then there is no jump event within δt . Evolve the state $|\psi(t)\rangle \rightarrow |\psi(t + \delta t)\rangle_{\text{no jump}}$ according to an effective non-Hermitian Hamiltonian $H_{\text{eff}} = H - i\hbar\frac{\gamma}{2}c^\dagger c$ and properly normalized.
- 3) If $r < \delta p$, then there is a jump event within δt . The evolved state becomes $|\psi(t + \delta t)\rangle_{\text{jump}} = \frac{c|\psi(t)\rangle}{\sqrt{\langle \psi(t) | c^\dagger c | \psi(t) \rangle}}$.
- 4) Considering an initial state $\psi(t_i = 0)$, repeat the above procedures until the final time t_f . Then you get a single “trajectory”, i.e., $|\psi(t)\rangle$ from time $t_i \rightarrow t_f$, and therefore an evolution curve of the observable of interest.
- 5) Average the observable over many such trajectories.

Now consider the following cases:

- a) The density matrix from t to $t + \delta t$

$$|\psi(t)\rangle\langle\psi(t)| \rightarrow \delta p |\psi\rangle_{\text{jump}}\langle\psi| + (1 - \delta p) |\psi\rangle_{\text{no jump}}\langle\psi|.$$

Prove that by taking $\frac{\delta\rho}{\delta t}$ and setting $\delta t \rightarrow dt$, one can recover the Lindblad master equation.

- b) Consider a two-level system $|\psi(t)\rangle = a(t)|e\rangle + b(t)|g\rangle$ under no interaction, i.e. $H = 0$, with the jumping operator $c = \sigma^- = |g\rangle\langle e|$ and an initial condition $a(0) = 1$, $b(0) = 0$. Run a computer program that plots $a(t)$ for $\gamma t = 0 \rightarrow 3$ for a single trajectory, average over 10, and 1000 trajectories.
- c) Consider the same system as (b), but only reset the jumping operator $c = \sigma_z$ and a different initial condition $a(0) = \sqrt{2/3}$ and $b(0) = \sqrt{1/3}$. Use a computer program to determine $\rho(t \rightarrow \infty)$.
- d) Consider an equally spaced three-level cascade atom with the ground state $|g\rangle$, the excited state $|e\rangle$, and an intermediate state $|m\rangle$. Then the corresponding Hamiltonian is given by $H_0 = \hbar\omega(2|e\rangle\langle e| + |m\rangle\langle m|)$. If a single mode laser field with frequency ω and intensity χ is applied to the atom, in the interaction

picture, $V_I = \hbar\chi(|e\rangle\langle m| + |m\rangle\langle g| + h.c.) - \hbar\delta(2|e\rangle\langle e| + |m\rangle\langle m|)$, where $\delta = \omega - \omega_0$. Include the Lindblad processes:

$$-\frac{\gamma}{2}\mathcal{L}(\rho) = -\frac{\gamma}{2}\left(c^+c^-\rho + \rho c^+c^- - 2c^-\rho c^+\right),$$

where $c^- = |g\rangle\langle m| + |m\rangle\langle e|$ and $c^+ = (c^-)^\dagger$. Consider the case that $\chi = 0.5\gamma$ and all population is initially located in $|e\rangle$. Use the quantum trajectory technique to determine $\rho_{ee}(t)$ for all times, and the steady state ρ_{mm} and ρ_{mg} (for $t \rightarrow \infty$).