

# Numerical Analysis of Bessel Function Roots and Applications in Physical Systems

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## ABSTRACT

We numerically computed and visualized the first five positive roots of the Bessel functions of the first kind,  $J_0(x)$ ,  $J_1(x)$ , and  $J_2(x)$ , using Python. These functions solve a second-order differential equation that occur in problems with radial symmetry. We used the `scipy` library to evaluate  $J_n(x)$  and locate its roots via the `fsolve` method. We interpreted the computed roots in the context of three physical materials: the radial wavefunction in a quantum infinite square well, heat conduction in cylindrical geometries, and vibrational modes of a circular drumhead. Our results show how Bessel function roots reflect physical boundary conditions and highlight the usefulness of numerical methods in modeling such systems.

*Keywords:* Bessel functions — root-finding — radial symmetry — boundary conditions — numerical analysis

## 1. INTRODUCTION

The Bessel functions of the first kind,  $J_n(x)$ , are solutions to the second-order linear differential equation:

$$z^2 \frac{d^2 w}{dz^2} + z \frac{dw}{dz} + (z^2 - n^2)w = 0, \quad (1)$$

where integer  $n = 0, 1, 2, \dots$  represents the order of the function. These functions appear in the separation of variables when solving partial differential equations in cylindrical or spherical coordinates, particularly in systems with radial symmetry [Abramowitz & Stegun \(1972\)](#). Bessel functions also have an integral

$$J_n(x) = \frac{1}{\pi} \int_0^\pi \cos(x \sin \theta - n\theta) d\theta \quad (2)$$

for integers  $n = 0, 1, 2, \dots$ , which is useful for understanding their oscillatory behavior.

Bessel functions were first introduced by German astronomer and mathematician Friedrich Wilhelm Bessel in the early 1800s during his study of planetary orbits. However, the equation itself had been investigated earlier by Bernoulli and Euler in problems involving vibrational membranes [Abramowitz & Stegun \(1972\)](#). Today, Bessel functions are foundational tools in mathematical physics and engineering. In quantum mechanics, they appear in radial solutions to the Schrödinger equation

for quantum wells [Hanson et al. \(2025\)](#), in the vibration modes of circular membranes like drumheads [Tamrin & Zakaria \(2022\)](#), and in heat conduction problems with cylindrical geometries [Neils \(2009\)](#).

In this project, we numerically compute the first five positive roots of the Bessel functions  $J_0(x)$ ,  $J_1(x)$ , and  $J_2(x)$ . These roots correspond to physically meaningful quantities such as resonance modes, cutoff frequencies, or quantized boundary values in systems with radial symmetry. We use Python to visualize each function, estimate the approximate locations of their roots, and apply the `fsolve` method from `scipy.optimize` to compute each root. Our goal is to obtain accurate root values for each order using numerical methods.

## 2. DATA AND OBSERVATIONS

This project focuses on computing and visualizing the first five roots of the Bessel functions of the first kind,  $J_n(x)$ , for orders  $n = 0, 1, 2$ . These functions solve the second-order linear differential equation:

$$z^2 \frac{d^2 w}{dz^2} + z \frac{dw}{dz} + (z^2 - n^2)w = 0, \quad (3)$$

which appear in physical systems with radial symmetry [Abramowitz & Stegun \(1972\)](#). The positive roots of  $J_n(x)$  represent physically meaningful quantities such as

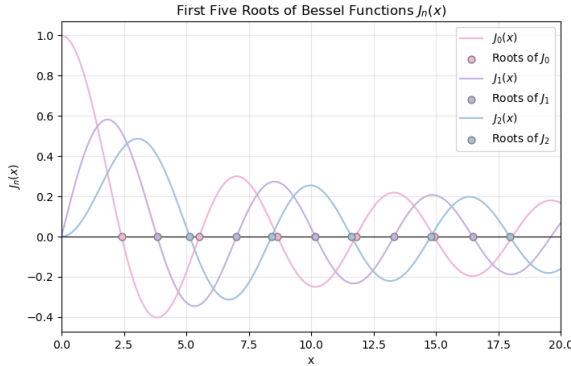
resonant frequencies or quantized energy levels in such systems.

We computed the Bessel functions using Python's `scipy.special.jv` method, which evaluates  $J_n(x)$  for an arbitrary order and argument. To find the roots numerically, we used the `fsolve` method from `scipy.optimize`, which refines an initial guess until it finds a point where the function crosses zero. We chose a set of initial guesses based on where the function appeared to cross the x-axis in the plot and passed those values into `fsolve` to calculate each root more precisely.

We used the following initial guesses:

- $J_0(x)$ : [2, 6.1, 8.6, 11.7, 15]
- $J_1(x)$ : [3.9, 7, 10.15, 13.1, 16.4]
- $J_2(x)$ : [5.1, 8.3, 11.8, 14.9, 18]

This gave us exactly five positive roots for each Bessel function. The resulting plots of  $J_0(x)$ ,  $J_1(x)$ , and  $J_2(x)$  over the domain  $x = 0$  to  $x = 20$  are shown in Figure 1, with the first five roots of each function marked as scatter points.



**Figure 1.** Bessel functions  $J_0(x)$ ,  $J_1(x)$ , and  $J_2(x)$  plotted from  $x = 0$  to  $x = 20$ , with their first five positive roots represented as circular markers.

### 3. RESULTS

The plots of  $J_0(x)$ ,  $J_1(x)$ , and  $J_2(x)$  from  $x = 0$  to  $x = 20$  show that Bessel functions of the first kind exhibit oscillatory behavior with a gradually decaying amplitude. As expected,  $J_0(x)$  begins at 1, while higher-order functions satisfy  $J_n(x) = 0$ . The zero crossings become slightly less frequent as  $x$  increases.

The computed roots correspond to the first five positive values of  $x$  for which  $J_n(x) = 0$  and are marked as circular points in Figure 1. These roots are important in radial problems where boundary conditions require

the function to vanish at a specific radius. For example, they determine the allowed energy levels in circular quantum wells, resonance frequencies in drumhead vibrations, and decay rates of thermal modes in cylindrical heat conduction.

Our numerical results followed the expected trend that root values increase with both the order  $n$  and the root index. The quantized structure highlights the physical significance of Bessel function solutions in bounded radial domains.

## 4. APPLICATIONS

Bessel functions appear in the solutions of physical problems that exhibit radial symmetry. In this section, we highlight three examples: the infinite square well in quantum mechanics, radial thermal diffusion, and wave propagation on a circular membrane. In each case, Bessel functions emerge from imposing boundary conditions on the radial part of a separable partial differential equation.

### 4.1. Quantum Mechanics: The Infinite Square Well

Bessel functions appear in the solution to the Schrödinger equation for a particle confined in a three-dimensional infinite spherical potential well. The potential is defined as:

$$V(x) = \begin{cases} 0, & \text{if } r < a \\ \infty, & \text{if } r \geq a \end{cases} \quad (4)$$

Inside the well ( $r < a$ ), the time-independent Schrödinger equation in spherical coordinates becomes:

$$\hat{H}\psi - E\psi = \hat{H}R_{n,l} - ER_{n,l} = 0 \quad (5)$$

$$\hat{H} = \frac{-\hbar^2}{2m} \left( \frac{\partial^2}{\partial r^2} + \frac{2}{r} \frac{\partial}{\partial r} - \frac{l(l+1)}{r^2} \right) + V(r) \quad (6)$$

where  $R_{n,l}(r)$  is the radial wavefunction,  $n$  is the principle quantum number,  $l$  is the angular momentum quantum number,  $m$  is the particle mass, and  $\hbar$  is the reduced Planck constant [Hanson et al. \(2025\)](#).

With  $V(r) = 0$  inside the well, the radial differential equation reduces to:

$$r^2 \frac{\partial^2 R_{n,l}}{\partial r^2} + 2r \frac{\partial R_{n,l}}{\partial r} + (k^2 r^2 - l(l+1)) R_{n,l} = 0 \quad (7)$$

This is the spherical Bessel differential equation where its solutions are spherical Bessel functions  $j_l(k_{n,l}r)$ . These relate to the ordinary Bessel functions by:

$$j_l(k_{n,l}r) = \sqrt{\frac{\pi}{2kr}} J_{l+\frac{1}{2}}(k_{n,l}r) \quad (8)$$

To satisfy the boundary condition that the wavefunction vanishes at  $r = a$ , we require that  $j_l(k_{n,l}a) = 0$ . This quantization conditions selects discrete values  $k_{n,l}$ , which in turn determine the allowed energy levels:

$$E_{n,l} = \frac{\hbar^2 k_{n,l}^2}{2ma^2} \quad (9)$$

The roots of the spherical Bessel functions determine the quantized energy spectrum of the particle in the well, which illustrates how the mathematical properties of Bessel functions have direct physical consequences in radially symmetric quantum systems [Weisstein \(2025\)](#).

#### 4.2. Thermal Diffusion

Bessel functions also appear in heat conduction problems with radial symmetry. In a two-dimensional circular region, the temperature distribution  $T(r, t)$  evolves according to the heat equation:

$$\frac{\partial T}{\partial t} = \alpha \nabla^2 T \quad (10)$$

$$= \alpha \left( \frac{1}{r} \frac{\partial T}{\partial r} + \frac{\partial^2 T}{\partial r^2} \right) \quad (11)$$

where  $\alpha = \frac{k}{\rho c_p}$  is the thermal diffusivity, with  $k$  as the thermal conductivity,  $\rho$  the material density, and  $c_p$  the specific heat capacity [Hahn & Özişik \(2021\)](#).

To solve this equation, we assume separable solutions of the form  $T(r, t) = X(r)\theta(t)$ , which leads us to two ordinary differential equations. The time-dependent part gives exponential decay:

$$\frac{d\theta}{dt} + \lambda^2 \alpha \theta = 0 \quad (12)$$

and the spatial part gives the Bessel differential equation of order zero:

$$\frac{d^2 X}{dr^2} + \frac{1}{r} \frac{dX}{dr} + \lambda^2 X = 0 \quad (13)$$

The general solution involves  $J_0(\lambda r)$ , and applying boundary conditions quantizes the solutions in terms of the roots  $\beta_n$  where  $J_0(\beta_n) = 0$  [Tsega et al. \(2017\)](#).

We impose the initial condition  $T(r, 0) = T_2$  and the boundary condition  $T(R, t) = 0$ , where  $R$  is the radius of the domain and  $T_1, T_2$  are constant values. The resulting dimensionless solution can be written as:

$$T^*(r, t) = \frac{T(r, t) - T(R, t)}{T(r, 0) - T(R, t)} = 2 \sum_{n=0}^{\infty} e^{-\beta_n^2 \frac{\alpha t}{R^2}} \frac{J_0(\beta_n \frac{r}{R})}{\beta_n J_1(\beta_n)} \quad (14)$$

which makes the full temperature solution:

$$T(r, t) = T^*(T(r, 0) - T(R, t)) + T(R, t) \quad (15)$$

The temperature function uses both the zeroth and first order Bessel function in the solution, and each mode decays exponentially at a rate set by  $\beta_n$ . Therefore, we see how the Bessel function determines the dynamics of thermal modes [Neils \(2009\)](#).

#### 4.3. Drum Wave Propagation

The vibration of a circular membrane, such as a drumhead, is another example of where Bessel functions appear. The vertical displacement of the membrane,  $z(r, \theta, t)$ , satisfies the two-dimensional wave equation:

$$\frac{\partial^2 z}{\partial t^2} = c^2 \nabla^2 z \quad (16)$$

where  $c = \sqrt{\frac{S}{\sigma}}$  is the wave speed,  $S$  is the surface tension, and  $\sigma$  is the surface mass density of the membrane.

Assuming symmetry about the axis (no  $\theta$ -dependence) and using separation of variables with  $z(r, \theta, t) = R(r)T(t)\Theta(\theta) = R(r)T(t)$ , we obtain two ordinary differential equations. The time-dependent part gives:

$$\frac{d^2 T}{dt^2} + \lambda^2 c^2 T = 0 \quad (17)$$

which describes simple harmonic motion in time, and the radial part becomes the Bessel differential equation of order  $n$ :

$$r^2 \frac{\partial^2 R}{\partial r^2} + r \frac{\partial R}{\partial r} + (\lambda^2 r^2 - n^2) R = 0. \quad (18)$$

with solutions  $R(r) = J_m(\lambda r)$ .  $J_m$  is the Bessel function of the first kind.

When we impose the boundary condition that the membrane is fixed at its edge ( $z(R_f, t) = 0$ ), we can quantize the allowable modes.  $\lambda$  must be one of the roots  $\lambda_{m,k}$  of the Bessel function  $J_m(\lambda r)$ , which satisfies:

$$J_m(\lambda_{m,k} R_f) = 0. \quad (19)$$

We can write the general solution as a double sum over all radial and angular modes:

$$z(r, t) = \sum_{m=0}^{\infty} \sum_{k=1}^{\infty} J_m(\lambda_{m,k} r) e^{-c \lambda_{m,n} t} \quad (20)$$

These quantized roots  $\lambda_{m,k}$  determine the resonant frequencies of the membrane. Each pair of integers  $(m, k)$  corresponds to a distinct vibration mode, where  $m$  is

the order of the Bessel function and  $k$  is the wavenumber. For  $m = 0$ , the Bessel function vanishes at the boundary. The final waveform  $z(r, t)$  describes the motion of the drumhead as a superposition of these modes evolving in time Tamrin & Zakaria (2022).

## 5. SUMMARY AND CONCLUSION

We explored the roots of Bessel functions of the first kind,  $J_n(x)$ , for orders  $n = 0, 1, 2$  using Python's `scipy` library. Our numerical approach combined the built-in evaluation of  $J_n(x)$  with the `fsolve` method to accurately identify the first five positive roots for each order. Our results aligned well with the expected theoretical behavior, including the trends in oscillation decay and root spacing.

The numerical methods we used were sufficient for our goals. Plotting the Bessel functions alongside their roots helped us visualize how the zero crossings correspond to boundary conditions in physical systems with radial symmetry. While our method was effective, it required manually selected initial guesses for each root. Future improvements include using analytical root approximations or symbolic solvers such as `SymPy` for better automation.

Overall, this project demonstrated how Bessel functions can be approached computationally, how their oscillatory behavior varies with order, and why their roots are important physical systems. Further work could involve exploring other families of Bessel functions or extending the root-finding method to more advanced applications involving partial differential equations where these functions appear due to symmetry. These additions would build a stronger understanding of the the mathematics and practical significance of Bessel functions across physics and applied fields.

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