

Identical particles

$$\mathcal{H} = L^2(\mathbb{R}^3, \mathbb{C}^2) := \left\{ \psi(\vec{r}, \sigma) \mid \sum_{\sigma \in \{\uparrow\downarrow, \downarrow\uparrow\}} \int_{\mathbb{R}^3} |\psi(\vec{r}, \sigma)|^2 d\vec{r} < \infty \right\}$$

$$\int d\vec{x} := \sum_{\sigma} \int_{\mathbb{R}^3} d\vec{r} \quad , \quad \vec{x} = (\vec{r}, \sigma)$$

$$\mathcal{H} = \left\{ \psi(\vec{x}) \mid \int d\vec{x} |\psi(\vec{x})|^2 < \infty \right\}.$$

2 particles  $\mathcal{H} \otimes \mathcal{H}$ .  $\psi(\vec{x}_1, \vec{x}_2)$

Permutation operator

$$P_{12} \psi(\vec{x}_1, \vec{x}_2) = \psi(\vec{x}_2, \vec{x}_1)$$

$$P_{12} \psi = a \psi. \quad P_{12}^2 = I. \Rightarrow a^2 = 1 \Rightarrow a = \pm 1.$$

$$H = -\frac{1}{2}(\Delta_{\vec{r}_1} + \Delta_{\vec{r}_2}) + V(\vec{r}_1, \vec{r}_2).$$

$$\begin{aligned} [H, P_{12}] \psi &= [V, P_{12}] \psi = V(\vec{r}_1, \vec{r}_2) \psi(\vec{x}_2, \vec{x}_1) \\ &\quad - V(\vec{r}_2, \vec{r}_1) \psi(\vec{x}_2, \vec{x}_1) \end{aligned}$$

$$V(\vec{r}_1, \vec{r}_2) = V(\vec{r}_2, \vec{r}_1) \Leftrightarrow [H, P_{12}] = 0.$$

$$H \psi = E \psi$$

$$\textcircled{1} \quad \psi(\vec{x}_1, \vec{x}_2) = \psi(\vec{x}_2, \vec{x}_1) \quad \text{Boson}$$

$$\textcircled{2} \quad \psi(\vec{x}_1, \vec{x}_2) = -\psi(\vec{x}_2, \vec{x}_1) \quad \text{Fermion}$$

$$\text{Ex. } H = -\frac{1}{2} \Delta_{\vec{r}_1} - \frac{1}{2} \Delta_{\vec{r}_2} - \frac{2}{|\vec{r}_1|} - \frac{2}{|\vec{r}_2|} + \frac{1}{|\vec{r}_1 - \vec{r}_2|}$$

$$A_2 = \bigwedge^2 L^2(\mathbb{R}^3; \mathbb{C}^2) \subset [L^2(\mathbb{R}^3, \mathbb{C}^2) \otimes L^2(\mathbb{R}^3, \mathbb{C}^2)]$$

$$[\vec{S}_{\text{tot}}^{+}, H] = [S_z^{+}, H] = 0.$$

$$\Psi(\vec{x}_1, \vec{x}_2) = \varphi(\vec{r}_1, \vec{r}_2) \chi(\sigma_1, \sigma_2)$$

Singlet :  $\chi$  anti-sym.  $\varphi$  sym.

triplet :  $\chi$  sym.  $\varphi$  anti-sym.

Simplest approximation "rank 1"

Singlet  $\varphi(\vec{r}_1, \vec{r}_2) = \phi(\vec{r}_1) \phi(\vec{r}_2)$

triplet  $\varphi(\vec{r}_1, \vec{r}_2) = \frac{1}{\sqrt{2!}} \begin{vmatrix} \phi_1(\vec{r}_1) & \phi_1(\vec{r}_2) \\ \phi_2(\vec{r}_1) & \phi_2(\vec{r}_2) \end{vmatrix}$

Slater determinant.

Proj. 1D Helium. Singlet / triplet energy.

$$L^2(\mathbb{R}, \mathbb{C}^2) \approx \text{span} \{ \tilde{\phi}_1(x), \tilde{\phi}_2(x) \} =: U$$

$$-\frac{1}{2} \frac{d^2}{dx^2} \tilde{\phi}_i - 2 V_c(x) \tilde{\phi}_i = \epsilon_i \tilde{\phi}_i, i=1,2.$$

① Project  $H$  to  $U \otimes U$ . Eigen-decomposition

and examine symmetry

② Find the best rank 1 approximation for singlet / triplet, and compare w. ①.

$N$ -particle system.

$$|\Psi\rangle \in \otimes^N L^2(\mathbb{R}^3; \mathbb{C}^2)$$

$$P_{ij} |\Psi(\vec{x}_1, \dots, \vec{x}_N)\rangle = |\Psi(\vec{x}_1, \dots, \overset{\uparrow}{\vec{x}_j}, \dots, \overset{\uparrow}{\vec{x}_i}, \dots, \vec{x}_N)\rangle$$

$\vdots \quad \quad \quad \vdots$

$$P_{ij}^2 = 1.$$

Postulation:  $|\Psi\rangle$  simultaneously eigenstate

of all  $P_{ij}$ . w. the same eigenvalue

$$(+1 \quad 0 \quad -1)$$

$$S_N := \text{Sym}^N L^2(\mathbb{R}^3; \mathbb{C}^2)$$

$$A_N := \Lambda^N L^2(\mathbb{R}^3; \mathbb{C}^2) \leftarrow \text{electrons}.$$

$N$  single particle functions

$$\{\psi_i(\vec{x})\}_{i=1}^N$$

$$\Psi_B(\vec{x}_1, \dots, \vec{x}_N) = C_B \sum_{\pi \in S_N} \psi_{\pi_1}(\vec{x}_1) \cdots \psi_{\pi_N}(\vec{x}_N)$$

$$\Psi_F(\vec{x}_1, \dots, \vec{x}_N) = C_F \sum_{\pi \in S_N} (-1)^\pi \psi_{\pi_1}(\vec{x}_1) \cdots \psi_{\pi_N}(\vec{x}_N)$$

$$= C_F \begin{vmatrix} \psi_1(\vec{x}_1) & \cdots & \psi_1(\vec{x}_N) \\ \vdots & \ddots & \vdots \\ \psi_N(\vec{x}_1) & \cdots & \psi_N(\vec{x}_N) \end{vmatrix}$$

$$\int \psi_i^*(\vec{x}) \psi_j(\vec{x}) dx = \delta_{ij}$$

$$\Rightarrow C_F = \frac{1}{\sqrt{N!}} \quad (\text{exer}).$$

exer: Write "rank 1" approx. to Helium atom using spin orbitals.

$$\psi(\vec{x}_1, \dots, \vec{x}_N) = 0. \quad \vec{x}_i = \vec{x}_j$$

Pauli exclusion principle.

$\phi(\vec{r})$  spin orbital.

$\phi(\vec{r}) \chi(\sigma)$  spatial orbital.

Quantum many body Hamiltonian.

$$\hat{H} = \sum_{i=1}^N -\frac{1}{2} \Delta \vec{r}_i + \sum_{i=1}^N V_{\text{ext}}(\vec{r}_i; \{R_I\})$$

T Ven

$$+ \sum_{i < j}^N \frac{1}{|\vec{r}_i - \vec{r}_j|} + \sum_{I < J}^M \frac{z_I z_J}{|\vec{R}_I - \vec{R}_J|}$$

Vee E<sub>II</sub>

# Variational principle

$$E_0 = \inf_{|\Psi\rangle \in \mathcal{A}_N, \langle \Psi | \Psi \rangle = 1} \langle \Psi | H | \Psi \rangle$$

Finite dimensional approx.

$$L^2(\mathbb{R}^3; \mathbb{C}^2) \approx U_d = \text{span}\{\phi_1(\vec{x}), \dots, \phi_d(\vec{x})\}$$

$$\mathcal{A}_{N,d} = \bigwedge^N U_d$$

$d > N$ . Pick  $d$  basis. antisymmetrize

$$\Rightarrow \bar{\Phi}_{[i_1 \dots i_N]} (\vec{x}_1, \dots, \vec{x}_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \phi_{i_1}(\vec{x}_1) & \dots & \phi_{i_1}(\vec{x}_N) \\ \vdots & \ddots & \vdots \\ \phi_{i_N}(\vec{x}_1) & \dots & \phi_{i_N}(\vec{x}_N) \end{vmatrix}$$

$\in A_{N,d} \subset A_N$ .

Thm.  $\{ \bar{\Phi}_{[i_1 \dots i_N]} \}$  is an orthonormal basis of  $A_{N,d}$

Pf: ① All det orthonormal.

$$\textcircled{2} \quad \text{span} \{ \varphi_{i_1}(\vec{x}_1) \dots \varphi_{i_N}(\vec{x}_N) \} = \bigodot^N U_d$$

anti symmetrize  $\Rightarrow A_{N,d}$

$$\text{Cor. } \dim A_{N,d} = \binom{d}{N}$$

Full configuration interaction

$$E_{0,N} = \inf \langle \Psi | H | \Psi \rangle$$

$$\Psi \in A_{N,d}$$

$$\langle \Psi | \Psi \rangle = 1.$$

Hartree - Fock:

best, single Slater determinant

$$\Psi(\vec{x}_1, \dots, \vec{x}_N) = \frac{1}{\sqrt{N!}} \begin{vmatrix} \psi_1(\vec{x}_1) & \dots & \psi_1(\vec{x}_N) \\ \vdots & \ddots & \vdots \\ \psi_N(\vec{x}_1) & \dots & \psi_N(\vec{x}_N) \end{vmatrix}$$

$$\psi_i(\vec{x}) = \sum_{\mu=1}^d \phi_{\mu i}(\vec{x}) c_{\mu i}$$

$\phi_{\mu i}(\vec{x}) = " \delta(\vec{r} - \vec{r}_{\mu}) \chi_{\mu}(\sigma)$  real space representation.

$A_N^{\circ}$ : all Slater det of  $N$  elec in  $A_N$ .

$A_{N,d}^{\circ}$ : all Slater det of  $N$  elec in  $A_{N,d}$ .

$$E_{HF} = \inf_{\Psi \in A_N^0} \langle \bar{\Psi} | H | \bar{\Psi} \rangle.$$

$$\langle \Psi | \bar{\Psi} \rangle = 1$$

$$E_c = E_0 - E_{HF} : \text{correlation energy.} \leq 0.$$

$$\bar{\Psi}(\vec{x}_1, \dots, \vec{x}_N) = \frac{1}{\sqrt{N!}} \sum_{\pi \in \text{Sym}(N)} (-1)^\pi \psi_{\pi(1)}(\vec{x}_1) \cdots \psi_{\pi(N)}(\vec{x}_N)$$

$$\langle \psi_i | \psi_j \rangle = \delta_{ij}$$

$$\langle \bar{\Psi} | \sum_{i=1}^N \left( \frac{1}{2} \Delta \vec{r}_i + V_{\text{ext}}(\vec{r}_i) \right) | \bar{\Psi} \rangle$$

$$= N \langle \Psi | -\frac{1}{2} \Delta_{\vec{r}_i} + V_{\text{ext}}(\vec{r}_i) | \Psi \rangle$$

$$= \frac{N!}{N!} \sum_{\pi, \pi' \in \text{Sym}(n)} (-1)^{\pi} (-1)^{\pi'} \langle \psi_{\pi(1)} | -\frac{1}{2} \Delta_{\vec{r}} + V_{\text{ext}} | \psi_{\pi'(1)} \rangle$$

$$\langle \psi_{\pi(2)} | \psi_{\pi'(2)} \rangle \cdots \langle \psi_{\pi(n)} | \psi_{\pi'(n)} \rangle$$

$$= \sum_{i=1}^N \langle \psi_i | -\frac{1}{2} \Delta_{\vec{r}} + V_{\text{ext}} | \psi_i \rangle$$

$$\langle \Psi | \sum_{i < j} \frac{1}{|\vec{r}_i - \vec{r}_j|} | \Psi \rangle$$

$$= \binom{N}{2} \langle \Psi | \frac{1}{|\vec{r}_i - \vec{r}_j|} | \Psi \rangle$$

$$= \frac{1}{2(N-2)!} \sum_{\substack{\pi, \pi' \\ \in \text{Sym}(N)}} (-1)^{\pi + \pi'} \int \psi_{\pi(1)}^*(\vec{x}_1) \psi_{\pi(2)}^*(\vec{x}_2) \frac{1}{|\vec{r}_1 - \vec{r}_2|} \psi_{\pi'(1)}(\vec{x}_1) \psi_{\pi'(2)}(\vec{x}_2) d\vec{x}_1 d\vec{x}_2$$

$$\langle \psi_{\pi(3)} | \psi_{\pi'(3)} \rangle \dots \langle \psi_{\pi(N)} | \psi_{\pi'(N)} \rangle$$

$$= \frac{1}{2} \sum_{\substack{i,j=1 \\ i \neq j}}^N \int \psi_i^*(\vec{x}_1) \psi_j^*(\vec{x}_2) \frac{1}{|\vec{r}_1 - \vec{r}_2|} \psi_i(\vec{x}_1) \psi_j(\vec{x}_2) d\vec{x}_1 d\vec{x}_2$$

$$- \sum_{\substack{i,j=1 \\ i \neq j}}^N \int \psi_i^*(\vec{x}_1) \psi_j^*(\vec{x}_2) \frac{1}{|\vec{r}_1 - \vec{r}_2|} \psi_j(\vec{x}_1) \psi_i(\vec{x}_2) d\vec{x}_1 d\vec{x}_2$$

$$\rho(\vec{x}) = N \int |\Psi(\vec{x}, \vec{x}_1, \dots, \vec{x}_N)|^2 d\vec{x}_1 \cdots d\vec{x}_N$$

$$\rho(\vec{r}) = \sum_{\sigma} \rho(\vec{r}, \sigma)$$

(exer)  $\rho(\vec{x}) = \sum_{i=1}^N |\psi_i(\vec{x})|^2.$

$$P(\vec{x}, \vec{x}') = \int \Psi(\vec{x}, \vec{x}_1, \dots, \vec{x}_N) \Psi^*(\vec{x}', \vec{x}_1, \dots, \vec{x}_N) d\vec{x}_1 \cdots d\vec{x}_N$$

$$\rho(\vec{x}) = P(\vec{x}, \vec{x}')$$

$$\mathcal{E}^{\text{ff}}(\{\psi_i\}_{i=1}^N) = \sum_{i=1}^N \langle \psi_i | -\frac{1}{2} \Delta_{\vec{r}} | \psi_i \rangle + \int V_{\text{ext}}(\vec{r}) \rho(\vec{r}) d\vec{r}$$

$$+ \frac{1}{2} \int \frac{\rho(\vec{r}) \rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' \\ - \frac{1}{2} \int \frac{|P(\vec{x}, \vec{x}')|^2}{|\vec{r} - \vec{r}'|} d\vec{x} d\vec{x}' + E_{\text{II}}$$

$G \vdash F$ .

UHF:

$$\psi_i(\vec{x}) = \varphi_i(\vec{r}) \chi_i(\sigma) \quad , \quad \chi_i(\sigma) = \langle \sigma | \uparrow \rangle \approx \langle \sigma | \downarrow \rangle.$$

RHF.  $N = 2N_{occ}$

$$\psi_i(\vec{x}) = \varphi_i(\vec{r}) \langle \sigma | \uparrow \rangle \quad , \quad i = 1, \dots, N_{occ}$$

$$\psi_{i+N_{occ}}(\vec{x}) = \varphi_i(\vec{r}) \langle \sigma | \downarrow \rangle.$$

$$E_0 \leq E^{GHF} \leq E^{UHF} \leq E^{RHF}.$$

(exer) Energy for UHF, RHF.

Sol fr HF.

RHF

$$\begin{aligned} \mathcal{E}^{\text{RHF}}(\{\varphi_i\}_{i=1}^{N_{\text{occ}}}) &= \sum_{i=1}^{N_{\text{occ}}} \int |\nabla \varphi_i(\vec{r})|^2 + \int V_{\text{ext}}(\vec{r}) \rho(\vec{r}) d\vec{r} \\ &+ \frac{1}{2} \iint \frac{\rho(\vec{r}) \rho(\vec{r}')}{|\vec{r}-\vec{r}'|} d\vec{r} d\vec{r}' - \int \frac{|P(\vec{r}, \vec{r}')|^2}{|\vec{r}-\vec{r}'|} d\vec{r} d\vec{r}'. \end{aligned}$$

$$\rho(\vec{r}) = 2 \sum_{i=1}^{N_{\text{occ}}} |\varphi_i(\vec{r})|^2$$

$$P(\vec{r}, \vec{r}') = \sum_{i=1}^{N_{\text{occ}}} \varphi_i(\vec{r}) \varphi_i^*(\vec{r}')$$

$$\delta \rho(\vec{r}) = 2 \sum_{i=1}^{N_{\text{occ}}} \delta \varphi_i^*(\vec{r}) \varphi_i(\vec{r}) + \varphi_i^*(\vec{r}) \delta \varphi_i(\vec{r})$$

$$\mathcal{L} = \epsilon^{\text{RHF}} - 2 \sum_{ij} \left( \langle \varphi_i | \varphi_j \rangle - \delta_{ij} \right) \lambda_{ji}$$

$$\frac{1}{2} \frac{\delta \mathcal{L}}{\delta \varphi_i^*(\vec{r})} = \left[ -\frac{1}{2} \Delta + V_{\text{ext}}(\vec{r}) + \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r}' \right] \varphi_i(\vec{r}) - \int \frac{P(\vec{r}, \vec{r}') \varphi_i(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r}' = \sum_{\alpha} \varphi_{j\alpha}(\vec{r}) \lambda_{ji}$$

$$H^{\text{RHF}} \varphi_i = \sum_j \varphi_j \lambda_{ji}$$

$$\lambda_{ji} = \langle \varphi_j | H^{\text{RHF}} | \varphi_i \rangle = \lambda_{ji}^*$$

$$\Lambda = U \Sigma U^* , \quad \Sigma = \begin{pmatrix} \varepsilon_1 & & \\ & \ddots & \\ & & \varepsilon_{N_{\text{occ}}} \end{pmatrix}$$

$$\psi_i(\vec{r}) = \sum_j \varphi_j(\vec{r}) U_{ji} \quad \hat{\Psi} = \hat{\Phi} U.$$

$$\rho(\vec{r}) = 2 \sum_{i=1}^{N_{occ}} |\psi_i(\vec{r})|^2. \quad P(\vec{r}, \vec{r}') = \sum_{i=1}^{N_{occ}} \psi_i(\vec{r}) \psi_i^*(\vec{r}')$$

$$H^{RHF}[4]\psi_i = \varepsilon_i \psi_i$$

① Solve RHF in a basis.

$$U_d = \{ \phi_1(\vec{r}), \dots, \phi_d(\vec{r}) \}.$$

$$\psi_i(\vec{r}) = \sum_{\mu=1}^d \phi_{\mu}(\vec{r}) C_{\mu i}$$

$$F_{\mu\nu} = \langle \phi_{\mu} | H^{\text{RHF}} | \phi_{\nu} \rangle = \underbrace{\langle \phi_{\mu} | -\frac{1}{2}\Delta + V_{\text{ext}} | \phi_{\nu} \rangle}_{h_{\mu\nu}}$$

$$+ 2 \sum_{\lambda\sigma} \int \frac{\phi_{\mu}^*(\vec{r}) \phi_{\nu}(\vec{r}) \phi_{\lambda}(\vec{r}') \phi_{\sigma}^*(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' P_{\lambda\sigma}$$

$$- \sum_{\lambda\sigma} \int \frac{\phi_{\mu}^*(\vec{r}) \phi_{\nu}(\vec{r}') \phi_{\lambda}(\vec{r}) \phi_{\sigma}^*(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' P_{\lambda\sigma}$$

$$P_{\lambda\sigma} = \sum_{i=1}^{N_{occ}} C_{\mu\lambda} C_{\nu\sigma}^* \quad \text{DM (in abasis)}$$

$$V_{pqrs} = \int \frac{\phi_p^*(\vec{r}) \phi_q^*(\vec{r}') \phi_r(\vec{r}) \phi_s(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' \equiv (pq|rs)$$

$$F_{\mu\nu} = h_{\mu\nu} + \sum_{\lambda\sigma} [2(\mu\sigma|\nu\lambda) - (\mu\sigma|\lambda\nu)] P_{\lambda\sigma}.$$

$O(d^4)$  cost.

$$S_{\mu\nu} = \langle \phi_\mu | \phi_\nu \rangle$$

$$FC = SC \circledast.$$

(exer) UHF HF eq.

Large basis . L16. QZ paper 18 .

density fitting .

Total energy computation (GHF)

$$E_0^{\text{HF}} = \sum_{i=1}^N \epsilon_i + \bar{E}_{\text{II}} ?$$

(exer) Correct after neglecting Hartree - Fock terms.

$$\begin{aligned} \sum_{i=1}^N \epsilon_i &= \frac{1}{2} \sum_{i=1}^N \int |\nabla_{\vec{r}} \psi_i(\vec{x})|^2 d\vec{x} + \int \rho V_{\text{ext}} \\ &+ \int \frac{\rho(\vec{r}) \rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' - \int \frac{|\rho(\vec{x}, \vec{x}')|^2}{|\vec{r} - \vec{r}'|} d\vec{x} d\vec{x}' \end{aligned}$$

$$E_o^{\text{HF}} = \sum_{i=1}^N \varepsilon_i + E_{II}$$

$$- \frac{1}{2} \int \frac{\rho(\vec{r}) \rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d\vec{r} d\vec{r}' + \frac{1}{2} \int \frac{|P(\vec{x}, \vec{x}')|^2}{|\vec{r} - \vec{r}'|} d\vec{x} d\vec{x}',$$

Double counting

(exb) Derive UHF eq. How many occupied states per spin channel?