

# latex-mimosi*s*

A minimal, modern L<sup>A</sup>T<sub>E</sub>X package for typesetting your thesis  
*by*  
Bastian Rieck

A document submitted in partial fulfillment of the requirements for the degree of  
*Technical Report*  
at  
MISKATONIC UNIVERSITY



# CONTENTS

1	INTRODUCTION	1
1.1	A note on units and notation . . . . .	1
2	DESCRIPTION OF THE SWISS CHEESE MODEL	3
2.1	Spacetime patches . . . . .	3
2.1.1	Friedmann-Robertson-Walker geometry . . . . .	3
2.1.2	Kottler geometry . . . . .	4
2.2	Matching conditions . . . . .	4



# 1 INTRODUCTION

There has been a debate over the past decade about whether the cosmological constant enters directly into the gravitational lensing equation.

## 1.1 A NOTE ON UNITS AND NOTATION

I use a comma to denote partial derivative. For example,  $x_{,t}$  refers to  $\frac{\partial x}{\partial t}$ . Throughout this work I use natural units such that  $c = G = 1$ .



# 2 DESCRIPTION OF THE SWISS CHEESE MODEL

## 2.1 SPACETIME PATCHES

Swiss-cheese (SC) models, also known as the Einstein-Straus models were introduced by Einstein and Strauss [??] to allow embedding of an inhomogeneous patch of spacetime within a homogeneous one (the "cheese"). Such a model is constructed by removing non-overlapping comoving spheres from the cheese and replacing them with a metric representing an appropriate condensed spherical mass distribution.

There are several reasons why this model was chosen. First, this is an exact solution of Einstein's equations, and hence it will allow us to properly investigate the higher order corrections that have been debated in literature. Previously some research (Simpson et al., 2010) has been done using a perturbative approach, but others (Ishak et al., 2010) have contested whether the approximations were valid. Secondly, by putting observers in the homogeneously expanding "cheese", it accounts for observers moving with the Hubble flow, which is a common objection[??] to Rindler and Ishak's [??] use of a static metric. Lastly, this model also takes the finite range of the mass into account by confining the influence of the central mass to the size of the hole.

This model has also previously been used to investigate the effect of local inhomogeneities on luminosity-redshift relations (Fleury et al., 2013). [??]

### 2.1.1 FRIEDMANN-ROBERTSON-WALKER GEOMETRY

Outside the hole, geometry is described by the Friedmann-Robertson-Walker (FRW) metric, given by

$$ds^2 = -dt^2 + a(t)^2 \left[ \frac{dr^2}{1 - kr^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2) \right] \quad (2.1)$$

where  $a$  is the scale factor and  $k$  represents the curvature. This is the line element for a homogeneous and expanding universe, with general spatial curvature. The scale factor  $a$  parametrizes the relative expansion of the universe, such that the relationship between physical distance and comoving distance between two points at a certain cosmic time  $t$  is given as

$$d_{\text{physical}} = a(t)d_{\text{comoving}}. \quad (2.2)$$

The scale factor also satisfies the Friedmann equation

$$H^2 \equiv \left( \frac{a,t}{a} \right) = \frac{8\pi G\rho}{3} + \frac{\Lambda}{3} - \frac{k}{a^2} \quad (2.3)$$

## 2 Description of the Swiss Cheese model

where  $\rho$  is the energy density of a pressureless fluid and  $H$  is the Hubble parameter.

It is common to introduce the cosmological parameters, where a subscript 0 refers to quantities evaluated today:

$$\Omega_m = \frac{8\pi G\rho_0}{3H_0^2}, \quad \Omega_\Lambda = \frac{\Lambda}{3H_0^2}, \quad \Omega_k = -\frac{k}{a_0^2 H_0^2} \quad (2.4)$$

and rewrite the Friedmann equation as

$$H^2 = H_0^2 \left[ \Omega_m \left( \frac{a_0}{a} \right)^3 + \Omega_k \left( \frac{a_0}{a} \right)^2 + \Omega_\Lambda \right]. \quad (2.5)$$

### 2.1.2 KOTTLER GEOMETRY

In this project we use a Kottler condensation in the Swiss-cheese for the central lensing mass. This is described by a Kottler metric (Kottler, 1918), which is the extension of the famous Schwarzschild metric to include a cosmological constant, given by

$$ds^2 = -f(R)dT^2 + \frac{dR^2}{f(R)} + R^2(d\theta^2 + \sin^2\theta d\phi^2) \quad (2.6)$$

with

$$f(R) = 1 - \frac{2M}{R} - \frac{\Lambda R^2}{3}, \quad (2.7)$$

where  $M$  is the mass of the central object. Unlike the FRW, this metric describes a static spacetime.

## 2.2 MATCHING CONDITIONS



## BIBLIOGRAPHY

- Fleury, Pierre, Hélène Dupuy, and Jean-Philippe Uzan (2013). “Interpretation of the Hubble diagram in a nonhomogeneous universe”. In: *Physical Review D* 87.12, p. 123526.
- Ishak, Mustapha, Wolfgang Rindler, and Jason Dossett (2010). “More on lensing by a cosmological constant”. In: *Monthly Notices of the Royal Astronomical Society* 403.4, pp. 2152–2156.
- Kottler, Friedrich (1918). “Über die physikalischen grundlagen der Einsteinschen gravitationstheorie”. In: *Annalen der Physik* 361.14, pp. 401–462.
- Simpson, Fergus, John A Peacock, and Alan F Heavens (2010). “On lensing by a cosmological constant”. In: *Monthly Notices of the Royal Astronomical Society* 402.3, pp. 2009–2016.