Networks Project: Numerical analysis of the Barabási-Albert Model

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Abstract: We investigate the behaviour of the Barabási-Albert Model.

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1 Pure preferential attachment

1.1 Degree distribution

1.1.1 Theory

The master equation that describes the evolution of the BA model is given by

$$n(k, t+1) = n(k, t) + m\Pi(k-1, t)n(k-1, t) - m\Pi(k, t)n(k, t) + \delta_{k, m}$$
(1)

where k is the total degree of a vertex, n(k,t) is the number of nodes at time t with total degree k, and the probability Π for choosing the existing vertex depends on the model.

In the pure preferential attachment model, we choose an existing edge with $\Pi_{pa} \propto k$, which after normalizing gives $\Pi_{pa} = k/2E(t)$ where E(t) is the number of edges and 2E(t) is the normalization constant corresponding to the total degree of the network. Assuming E(0) = mN(0), the number of edges at a given time t is given by E(t) = mN(t), so we get $\Pi = k/2mnN(t)$. Since we are concerned with the degree distribution of the model at large t, we consider the long-time ansatz $n(k,t) \to N(t)p_{\infty}(k)$.

Substituting these terms into the master equation, we obtain

$$p_{\infty}(k) = \frac{1}{2}[(k-1)p_{\infty}(k-1) - kp_{infty}(k)] + \delta_{k,m}$$
 (2)

It is clear that $p_{\infty}(k < m) = 0$, since m edges are added at every stage. So there are 2 cases to consider when solving for the above equation: k = m and k > m.

We first consider the case when k > m. In this case, $\delta_{k,m} = 0$ and we can rearrange Equation 2 to get

$$\frac{p_{\infty}(k)}{p_{\infty}(k+1)} = \frac{k-1}{k+2}$$
 (3)

To solve this equation, we can substitute in a trial solution of the form

$$f(z) = A \frac{\Gamma(z+1+a)}{\Gamma(z+1+b)} \tag{4}$$

where $\Gamma(z)$ is the Gamma function, which is an extension of the factorial function, with its argument shifted by one, to all real and complex nnumbers except the non-positive integers. Its central property is that

$$\Gamma(z+1) = z\Gamma(z), \ \Gamma(1) = 1. \tag{5}$$

Substituting the trial solution in Equation 4 gives

$$\frac{A\Gamma(z+1+a)}{\Gamma(z+1+b)} \times \frac{\Gamma(z+b)}{A\Gamma(z+a)} \tag{6}$$

which indeed simplifies to give (z+a)/(z+b), using the property in Equation 5 that $\Gamma(z+a+1)/\Gamma(z+a)=z+a$.

Substituting a = -1 and b = 2, we get the solution for Equation 3 in terms of A and the Gamma function:

$$p_{\infty}(k) = A \frac{\Gamma(k)}{\Gamma(k+3)} \tag{7}$$

which simplifies to

$$p_{\infty}(k) = \frac{A}{k(k+1)(k+2)}.$$
 (8)

For the second case of k = m, Equation 2 becomes

$$p_{\infty}(m) = \frac{1}{2}[(m+1)p_{\infty}(m-1) - mp_{\infty}(m)] + 1.$$
 (9)

However, we already know that $p_{\infty}(k < m) = 0$, that is, $p_{\infty}(m-1) = 0$. Using this, and rearranging Equation 9, we get

$$p_{\infty}(m) = \frac{2}{m+2}.\tag{10}$$

Substituting k = m and Equation 10 into Equation 8, we get

$$\frac{A}{m(m+1)(m+2)} = \frac{2}{m+2},\tag{11}$$

giving us the constant A as

$$A = 2m(m+1). (12)$$

For this constant to be physically reasonable, we need to check that the probability satisfies normalization, that is, we need to prove

$$\sum_{k=m}^{\infty} p_{\infty}(k) = 2m(m+1) \sum_{k=m}^{\infty} \frac{1}{k(k+1)(k+2)} = 1.$$
 (13)

The term in the summation of Equation 13 can be expanded as a partial fraction:

$$\sum_{k=m}^{\infty} \frac{1}{k(k+1)(k+2)} = \sum_{k=m}^{\infty} \frac{1}{2k} - \sum_{k=m}^{\infty} \frac{1}{k+1} + \sum_{k=m}^{\infty} \frac{1}{2(k+2)}$$
 (14)

By writing out the first few terms of each summation, we can see that most terms cancel:

$$\frac{\frac{1}{2m} - \frac{1}{m+1} + \frac{1}{2(m+2)}}{+ \frac{1}{2(m+1)} - \frac{1}{m+2} + \frac{1}{m+3}} + \frac{1}{2(m+4)} + \frac{1}{2(m+4)} - \frac{1}{m+4} + \frac{1}{2(m+5)} + \dots$$
(15)

ining terms we get the relation in Equation 13

and from the remaining terms we get the relation in Equation 13

$$\sum_{k=m}^{\infty} p_{\infty}(k) = 2m(m+1) \left(\frac{1}{2m} - \frac{1}{m} + \frac{1}{2(m+1)} \right) = 2m(m+1) \frac{1}{2m(m+1)} = 1 \quad (16)$$

Hence, we can confirm that the complete exact solution for the probability distribution in the long time limit is

$$p_{\infty}(k) = \frac{2m(m+1)}{k(k+1)(k+2)}.$$
(17)

1.1.2 Numerical analysis

To leverage its speed, c++ code was used to generate graph data, while python was used for data analysis due to its wide range of data analysis tools, such as numpy, scipy, and pandas.

There were two main concerns when doing the numerical simulation:

- 1. What should the initial graph G_0 be?
- 2. Should self loops and multiple edges be allowed?
- 3. Of what order should m and N be respectively?

The initial graph was

Computational efficiency of our algorithm is important considering that bigger datasets give better and more reliable statistical results. The following algorithm was used In this algorithm, the array M holds the list of edges represented by pairs of vertices, for example, the vertices at M[0] and M[1] are connected, M[2] and M[3] are connected, and so on. In this list, the number of occurences of a vertex is equal to its degree, so it can be used as a sample pool to achieve preferential attachment. To choose m neighbours for each new vertex, we then sample from M. This is also equivalent to choosing an edge at random and then choosing a vertex at random from the edge.

Algorithm 1 Algorithm for preferential attachment

```
\triangleright N > m
Require: number of vertices N, minimum degree m
 1: Initialize our graph q
 2: Initialize M as an empty array of length 2Nm
 3: for v in [0, ..., n-1] do
        g.addVertex()
                                                                          \triangleright add new vertex to q
        for i in [0, ..., m-1] do
 5:
            M[2(vm+i)] \leftarrow v
 6:
           draw r uniformly at random
 7:
           from [0, ..., 2(vm + i)]
                                                              \triangleright Choose random vertex from M
 8:
           M[2(vm+i)+1] \leftarrow M[r]
                                                       \triangleright Add edge between vertex M[r] and v
 9:
10:
11: for i in [0, ..., nm-1] do
                                                   ▶ Add all edges stored in M into the graph
12:
        Add edge (M[2i], M[2i+1]) to graph g
```

Clearly this approach produces self loops and multiple edges, especially at the beginning before m vertices have been created. However, since we are concerned with the limit of $N \to \infty$, these effects are insignificant. Also, while unsatisfactory, multiple edges and self-loops do not affect our theoretical result. In the large N limit, the probability of getting multiple edges or self loops is small, and so for simplicity this algorithm was used without modification.

To check that the model was implemented correctly, the degree distribution generated by the model compared checked against the **networkx** implementation of the BA model. By using the same random seed, we could check that degree distribution between the two models were exactly the same. Graphs of fewer than 10 nodes were also generated to ensure that the points followed some basic constraints. This gives confidence that the algorithm is working as expected.

To investigate the degree distribution, the model was run for fixed N but varying m. Several numerical runs were performed for each set of values to improve statistics. It was found that generally above $N = 10^4$, finite size scaling effects are insignificant (see section 1.2.2). In order to reduce finite size effects, $N = 10^6$ was used.

Visually, as can be seen from ??, the numerical results seem to agree with the theoretical model after log-binning, until finite sized effects begin to kick in at large k. Alternatively, we can look at the complementary cumulative distribution function (ccdf) to observe the behaviour of the fat tail more clearly. This is shown in ??. We can see that near the fat tail, the numerical ccdf goes slightly higher than the theoretical, before falling off. This is consistent across different m values, and the first bump increases in size with increasing m.

A Kolmogorov-Smirnov (KS) test was used to quantify the goodness-of-fit between theory and numerical results. It is a non-parametric test that measures how far apart two distributions are, and it returns a KS-statistic that can be converted into a p-value. A small p-value implies that the model is unlikely to have been the generating function for the data and should be rejected, while a large p-value implies that the differences between data and model may be attributed to statistical differences.

The table below shows the results of the KS test on m = 1 to m = 32:

m	p-value
1	1.000
2	0.999
4	0.992
8	0.917
16	0.514
32	0.216

However, these values are meaningless unless we know what kind of deviation from theoretical is considered acceptable, and at what value we should reject the null hypothesis. To answer that, we generate synthetic datasets governed by the theoretical distribution in Equation 17 to measure how far they fluctuate from the reference theoretical distribution in a similar way as described by Clauset et al. (2009), and compare the results with the simulated data. The synthetic datasets were generated with a lower bound of m and an upper bound given by the maximum degree in the simulated dataset. We can then measure how far the synthetic datasets deviate from the theoretical distribution. If the simulated data is much further from the theoretical distribution than the typical synthetic data, then we would have grounds to reject the null hypothesis.

For each of the synthetic datasets, we compare it to the theoretical distribution and calculate its KS-statistic. Then we define our p-value as the fraction of the time the resulting statistic is larger than the value for the numerical data. We can then reject the null hypothesis if $p \leq 0.1$ (Clauset et al., 2009), that is, if there is a 1 in 10 probability

or less that we would, by chance, get data that agree as poorly with the model as the current data.

Another issue to consider is the number of synthetic datasets to generate. Again, Clauset et al. (2009) suggests a useful rule of thumb: to have p-values accurate to within about ϵ , we need at least $\frac{1}{4}\epsilon^{-2}$ datasets. In this project,

1.2 Largest expected degree

1.2.1 Theory

The finite size of the system imposes a structural cutoff on the largest expected degree. For scale free networks, Aiello et al. (2001) defined the maximum degree to be approximately the value above which there is less than one vertex of that degree in the graph on average, that is, $N \sum_{k=k_1}^{\infty} p_{\infty}(k) = 1$.

Generally, it is shown (Boguñá et al., 2004) that for a scale free network with $p_{\infty}(k) \propto k^{\gamma}$, the largest expected degree will be

$$k_1(N) \sim N^{1/(\gamma - 1)}$$
. (18)

In our case, this can be easily verified. Starting with the equation

$$N\sum_{k=k_1}^{\infty} p_{\infty}(k) = 1, \tag{19}$$

we can see that this is almost identical to Equation 13, just with a different factor and lower limit. Hence we have

$$2m(m+1)\frac{1}{2k_1(k_1+1)} = \frac{1}{N}. (20)$$

We can then rearrange this to give us an expression for k_1 :

$$k_1 = \frac{-1 + \sqrt{1 + 4Nm(m+1)}}{2} \tag{21}$$

where the other negative solution is rejected as it is unphysical, and confirming that verifying that $k \propto N^{0.5}$.

1.2.2 Numerical analysis

As can be seen from the ??, there seems to be a constant offset between the numerical and theoretical k_1 values. By looking at the ratios of $k_1(\text{numerical})/k_1(\text{theoretical})$ as shown in ??, we can see that deviations are generally constant.

To produce a data collapse, we need to find the function f such that

$$p_N(k) = f(k)\mathcal{G}(k/k_1) \tag{22}$$

To find f(k), we know that in the limit of $N \to \infty$, p must have no dependence on L, and in the large N limit, $p_N(k) = f(k)$ which is just $p_\infty(k)$. This implies

$$\frac{p_N(k)}{p_\infty(k)} = G(k/k_1) \tag{23}$$

which means that plotting $p_N(k)/p_\infty(k)$ against k/k_1 will produce a data collapse, as shown in ??.

2 Pure random attachment

2.1 Degree distribution

2.1.1 Theory

The pure random attachment model can be seen as a limiting case of the BA model. In this model, all existing vertices are chosen with equal probability, i.e. $\Pi = \Pi_{rnd} \propto 1$. This preserves growth but removes preferential attachment.

Similar to the previous section, we start from the master equation in Equation 1, and instead of $\Pi = k/2E(t)$ as in preferential attachment, we use $\Pi_{rnd} = 1/N(t)$. Again, we consider the long-time ansatz $n(k,t) \to N(t)p_{\infty}(k)$. Substituting these terms into Equation 1, we have

$$p_{\infty}(k) = mp_{\infty}(k-1) - mp_{\infty}(k) + \delta_{k,m}. \tag{24}$$

Considering the case of k > m, we obtain the recurrence relation

$$p_{\infty}(k) = \left(\frac{m}{m+1}\right) p_{\infty}(k-1) = \dots = \left(\frac{m}{m+1}\right)^{k-m} p_{\infty}(m)$$
 (25)

Now we consider k = m. Substituting k = m into Equation 24 and remembering that $p_{\infty}(k < m) = 0$, we get

$$p_{\infty}(m) = -mp_{\infty} + 1, \tag{26}$$

giving us

$$p_{\infty}(m) = \frac{1}{m+1}.\tag{27}$$

Combining this result with Equation 25, we get the following formula for $p_{\infty}(k)$:

$$p_{\infty}(k) = \frac{1}{m+1} \left(\frac{m}{m+1}\right)^{k-m}.$$
 (28)

For normalization, we need to check that

$$\sum_{k=m}^{\infty} p_{\infty}(k) = \frac{1}{m+1} \sum_{k=m}^{\infty} \left(\frac{m}{m+1}\right)^{k-m} = 1.$$
 (29)

The terms in the summation form a converging geometric series, with the starting term being zero and common ratio being m/(m+1). Hence we have

$$\sum_{k=m}^{\infty} \left(\frac{m}{m+1} \right)^{k-m} = \frac{1}{1 - [m/(m+1)]}.$$
 (30)

By substituting this back into the Equation 29, we can see that normalization is satisfied.

As we can see from Equation 28, the resulting degree distribution in this limit is geometric (Peköz et al., 2013), indicating that growth alone is not sufficient to produce a scale free structure.

2.1.2 Numerical analysis

Numerical simulations confirmed that growth alone is not sufficient to produce a scale free structure.

2.2 Largest expected degree

2.2.1 Theory

Using the same definition from the previous section for largest expected degree, we need

$$N\sum_{k=k_1}^{\infty} p_{\infty}(k) = N\sum_{k=k_1}^{\infty} \frac{1}{m+1} \left(\frac{m}{m+1}\right)^{k-m} = 1.$$
 (31)

Again, the summation is similar to the previous Equation 30 just with a different lower limit.

3 Random walks and preferential attachment

3.1 Theoretical degree distribution

One of the weaknesses of the BA model and its generalizations is that this implicitly requires a knowledge of the total degree and a calculation across existing vertices on the graph. This requirements then destroys the potential for this model to exhibit emergent properties based on local behaviour. The attachment by performing a random walk is a solution proposed by Saramäki and Kaski (2004). In this model, a vertex is chosen at random from existing vertices and then executes a random walk of length L from that vertex. The new vertex then attaches to the destination vertex.

This model was thought to be able to reproduce the BA degree distribution even for L=1 (Saramäki and Kaski, 2004; J.P.Saramaki and T.S.Evans, 2004). While this is the case for large L, Cannings and Jordan (2013) later showed that the L=1 degree sequence converges to a degenerate limiting solution in which almost every vertex has degree 1, instead of a power law distribution, and demonstrated that this model is fundamentally

different from the BA model, unless we allow an indefinite length for the random walk. For L=0, this reduces to the random attachment model.

For L=1: likelihood ratio tests, clauset

References

- Aiello, William et al. (2001). "A Random Graph Model for Power Law Graphs". In: Experimental Mathematics 10.1, pp. 53-66. ISSN: 1058-6458. DOI: 10.1080/10586458. 2001.10504428. URL: http://www.tandfonline.com/doi/abs/10.1080/10586458. 2001.10504428.
- Boguñá, M., R. Pastor-Satorras, and A. Vespignani (2004). "Cut-offs and finite size effects in scale-free networks". In: *European Physical Journal B* 38.2, pp. 205–209. ISSN: 14346028. DOI: 10.1140/epjb/e2004-00038-8. arXiv: 0311650 [cond-mat].
- Cannings, Chris and Jonathan Jordan (2013). "Random walk attachment graphs". In: Electronic Communications in Probability 18, pp. 1–8. ISSN: 1083589X. DOI: 10.1214/ECP.v18-2518. arXiv: 1303.1052.
- Clauset, Aaron, Cosma Rohilla Shalizi, and M E J Newman (2009). "Power-Law Distributions in Empirical Data". In: *SIAM Review* 51.4, pp. 661–703. ISSN: 19417330. DOI: 10.1214/13-AOAS710. arXiv: arXiv:0706.1062v2.
- J.P.Saramaki and T.S.Evans (2004). "Scale Free Networks from Self-Organisation". In: *Physical Review E* 72.2, pp. 1–33. ISSN: 1539-3755. DOI: 10.1103/PhysRevE.72. 026138. arXiv: 0411390v2 [cond-mat]. URL: http://link.aps.org/doi/10.1103/PhysRevE.72.026138.
- Peköz, Erol A., Adrian Röllin, and Nathan Ross (2013). "Total variation error bounds for geometric approximation". In: *Bernoulli* 19.2, pp. 610–632. ISSN: 1350-7265. DOI: 10.3150/11-BEJ406. URL: http://projecteuclid.org/euclid.bj/1363192040.
- Saramäki, Jari and Kimmo Kaski (2004). "Scale-free networks generated by random walkers". In: *Physica A: Statistical Mechanics and its Applications* 341.1-4, pp. 80–86. ISSN: 03784371. DOI: 10.1016/j.physa.2004.04.110. arXiv: 0404088 [cond-mat].