

Lecture 9: Nonlinear Partial Differential Equations (PDEs)

COMP5930M Scientific Computation

Today

Example PDEs

Numerical modelling of PDEs

- Approximation in space

- Approximation in time

- Nonlinear system

- Solution algorithm

Summary

Differential equations

A **differential equation** is a type of equation where the unknown being sought is a function $u(t)$ of a free variable t , and the equation involves one of the derivatives of u with respect to t .

Example 1: The differential equation

$$\frac{du}{dt} = \cos(t)$$

has general solution $u(t) = \sin(t) + C$, where C is any constant.

To determine C , we need to fix the value of $u(t)$ in at least one point, say $u(0) = u_0$. If we fix the value of $u(t)$ at time $t = 0$ and look for a solution for $t > 0$, we say it is an **initial-value problem**.

Partial differential equations

A **partial differential equation** is a type of equation where the unknown being sought is a function $u(x, t)$ of two (or more) free variables x and t , and the equation involves the **partial derivatives** of u with respect to both x and t .

Example 1: The partial differential equation

$$\frac{\partial u}{\partial t} = \cos(t) \frac{\partial u}{\partial x}$$

includes partial derivatives w.r.t to both x and t .

Typically, we will consider the variable t as **time** and treat the problem as in initial value problem in time i.e. the value of the function at time t_0 , $u(x, t_0) = u_0(x)$, is a known function.

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We will typically treat x as a variable representing a **spatial coordinate**. For such variables, a natural condition is a **boundary condition**. If $x \in [a, b]$ then we fix the values of u at the interval end-points: $u(a, t) = \alpha$, $u(b, t) = \beta$, for all $t > 0$.

Example: 1d viscous Burger's Equation

Find $u(x, t)$ satisfying

$$\frac{\partial u}{\partial t} + R u \frac{\partial u}{\partial x} = \frac{\partial^2 u}{\partial x^2}$$

on $x \in [X_1, X_2]$, and $t > 0$,

with boundary conditions $u(X_1, t) = u_1(t)$, $u(X_2, t) = u_2(t)$

and initial conditions $u(x, 0) = u_0(x)$.

$R > 0$ is a known constant

- ▶ Prototype model for fluid dynamics
- ▶ Time dependent
- ▶ Nonlinear convection

Behaviour of solutions to Burger's equation

Presence of shock fronts in 1-D (for $R \gg 1$):

https://www.youtube.com/watch?v=FAY_N1a-LYQ

Example: Cellular action potential models

Find $u(x, y, t)$ and $v(x, y, t)$ satisfying

$$\begin{aligned}\frac{\partial u}{\partial t} + g(u, v) &= \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + f(t) \\ \frac{\partial v}{\partial t} &= h(u, v)\end{aligned}$$

on $(x, y) \in \Omega$, with $u(t) = u_b(x, y, t)$ on the boundary $\partial\Omega$.

Initial conditions $u(x, y, 0)$ and $v(x, y, 0)$ also need to be specified.

- ▶ The function $u(x, y, t)$ describes the transmembrane potential across the cell membranes (e.g. nervous system or the heart)
- ▶ The function $v(x, y, t)$ describes the opening/closing of ionic channels that regulate the cell membranes
- ▶ The functions g, h are nonlinear in u, v .

Behaviour of solutions to action potential models

Nonlinear dynamics exhibiting chaotic behaviour:

<https://www.youtube.com/watch?v=S9lHkUCImcs>

Example: Navier-Stokes Equations

Find $\mathbf{u}(\mathbf{x}, t)$ and $\rho(\mathbf{x}, t)$ satisfying

$$\begin{aligned}\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p &= \frac{1}{Re} \nabla \cdot \nabla \mathbf{u} \\ \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) &= 0\end{aligned}$$

on $\mathbf{x} \in \Omega$ with *appropriate* boundary and initial conditions.

- ▶ Fundamental model for fluid dynamics at the continuum level
- ▶ Time dependent
- ▶ Diffusive terms
- ▶ Nonlinear convection

Behaviour of solutions to Navier-Stokes equations

Complex instabilities can arise even in simple cases:

<https://www.youtube.com/watch?v=Nh1dX7MrukA>

Numerical modelling approach

The method of lines:

1. Approximation in space only
 - ▶ Semi-discrete system of ordinary differential equations (ODEs)

$$\mathbf{U}'(t) = \mathbf{F}(\mathbf{U}(t))$$

2. Approximation in time
 - ▶ Fully discrete (nonlinear) system, e.g.

$$\mathbf{U}_{k+1} = \mathbf{U}_k + \Delta t \mathbf{F}(\mathbf{U}_{k+1})$$

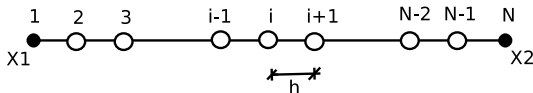
3. Discrete solution algorithm for \mathbf{U}_{k+1} at each step
 - ▶ Numerical solution

Approximations in 1d: Finite Difference Methods

Define a grid (mesh) for the spatial domain x

For FDM approximations define a uniform spacing h

$$x_i = X_1 + (i - 1)h, \quad i = 1, 2, \dots, n$$



At point (node) i approximate all spatial terms of the PDE

Notation: $u(x_i, t) \equiv u_i(t)$, $u(x_i, t^n) \equiv u_i^n$, $\frac{\partial u}{\partial t} \equiv \dot{u}$

Spatial approximation: FDM

Some standard difference approximations

$$\frac{\partial u}{\partial x}(x_i) \approx \frac{u_i - u_{i-1}}{h} \quad (\text{upwind difference, 1st order})$$

$$\frac{\partial u}{\partial x}(x_i) \approx \frac{u_{i+1} - u_{i-1}}{2h} \quad (\text{central difference, 1st order})$$

$$\frac{\partial^2 u}{\partial x^2}(x_i) \approx \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} \quad (\text{central difference, 2nd order})$$

We say a FDM is **accurate of order p** if the error of the approximation is $o(h^p)$. Thus the higher the order, the better.

Spatial approximation: FDM

Some standard difference approximations

$$\frac{\partial u}{\partial x}(x_i) \approx \frac{u_{i+1} - u_i}{h} \quad (\text{downwind difference, 1st order})$$

$$\frac{\partial u}{\partial x}(x_i) \approx \frac{u_{i+1} - u_{i-1}}{2h} \quad (\text{central difference, 1st order})$$

$$\frac{\partial^2 u}{\partial x^2}(x_i) \approx \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2} \quad (\text{central difference, 2nd order})$$

Many other variants are possible, also of higher orders.

FDM for 1d Burger's Equation

At node i of the grid, the equation is:

$$\frac{\partial u}{\partial t}(x_i) + R u(x_i) \frac{\partial u}{\partial x}(x_i) = \frac{\partial^2 u}{\partial x^2}(x_i)$$

Semi-discrete form:

$$\dot{u}_i + R u_i \left(\frac{u_i - u_{i-1}}{h} \right) = \frac{u_{i+1} - 2u_i + u_{i-1}}{h^2}$$

- ▶ Semi-discrete ODE system defined for $i = 2, 3, \dots, n - 1$
- ▶ Upwind/downwind approximation for $\frac{\partial u}{\partial x}$ chosen depending on the direction in which the shock wave travels

Options for nonlinear terms

There are usually alternatives, eg. $u \frac{\partial u}{\partial x} \equiv \frac{1}{2} \frac{\partial(u^2)}{\partial x}$

$$\begin{aligned} u \frac{\partial u}{\partial x}(x_i) &\approx u_i \left(\frac{u_i - u_{i-1}}{h} \right) \\ \frac{1}{2} \frac{\partial u^2}{\partial x}(x_i) &\approx \frac{1}{2} \left(\frac{u_i^2 - u_{i-1}^2}{h} \right) \\ &= \frac{1}{2} (u_i + u_{i-1}) \left(\frac{u_i - u_{i-1}}{h} \right) \end{aligned}$$

- ▶ Will lead to differences in Jacobian terms and numerics
- ▶ Both are consistent and formally of equal accuracy

Approximating in time

The continuous problem is reduced to a semi-discrete system of ODEs for $\mathbf{u}(t)$: $\dot{\mathbf{u}} = f(\mathbf{u})$

Initial-value problem: We start from a given \mathbf{u}^0 and find a formula that computes \mathbf{u}^{k+1} given \mathbf{u}^k , for $k = 1, 2, 3, \dots$. Here k is the time index, i.e. the solution at time t_k is \mathbf{u}^k .

This procedure is called **time-stepping** or **time-integration**. Unlike the spatial variable, we do not need to solve all the unknowns at once but proceed one time-step at a time.

Approximating in time

The continuous problem is reduced to a semi-discrete system of ODEs for $\mathbf{u}(t)$: $\dot{\mathbf{u}} = \mathbf{f}(\mathbf{u})$

- Explicit Euler

$$\frac{\mathbf{u}^{k+1} - \mathbf{u}^k}{\Delta t} = \mathbf{f}(\mathbf{u}^k)$$

No actual equation to solve (very efficient), but typically result is unstable unless $\Delta t \ll 1$...

- Implicit Euler

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Fully nonlinear equation for \mathbf{u}^{k+1} to solve, but usually can take larger time steps Δt .

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1d Burger's Equation

Use Implicit Euler. At node i of the grid:

$$\frac{u_i^{k+1} - u_i^k}{\Delta t} + R u_i^{k+1} \frac{u_i^{k+1} - u_{i-1}^{k+1}}{h} = \frac{u_{i+1}^{k+1} - 2u_i^{k+1} + u_{i-1}^{k+1}}{h^2}$$

- ▶ Fully discrete nonlinear system for the unknown u_i^{k+1} , $i = 2, \dots, n-1$
- ▶ The equations are coupled to the neighbouring nodes through the FDM approximation

Nonlinear system structure

Equations are formed at each internal node $i = 2, \dots, n - 1$

- ▶ At node $i - 1$

$$F_{i-1}(u_{i-2}^{k+1}, u_{i-1}^{k+1}, u_i^{k+1}) = 0$$

- ▶ At node i

$$F_i(u_{i-1}^{k+1}, u_i^{k+1}, u_{i+1}^{k+1}) = 0$$

- ▶ At node $i + 1$

$$F_{i+1}(u_i^{k+1}, u_{i+1}^{k+1}, u_{i+2}^{k+1}) = 0$$

Each equation depends only on local, neighbouring information

Time-stepping approach

- ▶ Initial conditions are specified at $t = t_0$ as $u = U(x, t_0)$
 - ▶ Sets the discrete solution $u_i^0 = U(x_i, t_0)$
- ▶ For each time step $k = 0, 1, 2, \dots$
we solve a nonlinear system to find \mathbf{u}_{k+1} :
 - ▶ $\mathbf{F}(\mathbf{u}^{k+1}, \mathbf{u}^k) = \mathbf{0}$
 - ▶ Advances u_i^k to u_i^{k+1} , $i = 2, 3, \dots, n-1$

Summary

- ▶ Nonlinear PDEs can exhibit complex spatiotemporal dynamics
- ▶ The method of lines for numerically approximating PDEs
 - ▶ Separate approximations in space and time
- ▶ Approximation in space: Finite difference method
- ▶ Approximation in time: Implicit Euler's method
- ▶ The final discrete system can be solved with Newton's method at each timestep