



# Outline

- 1 Motivation
- 2 ARK Methods
- 3 Example Results
- 4 Conclusions

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# Multiphysics Problems

“Multiphysics problems” typically involve a variety of interacting processes:

- System of multiple components coupled in the bulk:
  - Cosmology: radiation + (magneto)hydrodynamics + chemistry + gravity
  - Combustion/subsurface flow: reaction + transport
- System of multiple components coupled across interfaces:
  - Climate: ocean + atmosphere + sea ice
  - Tokamak: fluid core + kinetic edge

In this talk, we'll consider the prototypical multiphysics form,

$$\partial_t u = f(t, u) = f_1(t, u) + f_2(t, u),$$

where  $u$  comprises all of our unknowns, and  $f_i(t, u)$  are different physical processes that each act on all or part of  $u$ .

A primary difficulty with multiphysics problems is that each process may evolve at a different speed, e.g.  $f_1(t, u)$  is “slow” while  $f_2(t, u)$  is “fast”.



# Multiphysics Often Means “Multirate”

A single time scale is ideal for explicit-time methods, allowing for simpler algorithms, high-order accuracy, and predictable stability.

Wide temporal disparities can be analytically reformulated, but only for the scale of interest.

True multi-rate problems, however, require something more:

- Fully implicit methods are valid for stiff problems, but may require adaptation of solvers for all physical components.
- Operator-split methods are often chosen to match methods with physics.

Unfortunately, “standard” splitting approaches suffer from:

- *Low Accuracy* – even fractional-step methods may not achieve asymptotic  $\mathcal{O}(h^2)$  accuracy until  $h$  is very small, since error terms are dominated by inter-process interactions [Ropp, Shadid,& Ober 2005].
- *Low Stability* – numerical stability isn't guaranteed *even if  $h$  is stable for each component* [Estep et al., 2007].



# Fixing Operator Splitting

Although potentially dangerous, splittings are pervasive in scientific computing:

- Reuse existing/legacy software,
- Allow incorporation of domain-specific knowledge,
- No monolithic solvers for complex (and often non-differentiable) physics,
- Results “look reasonable,” especially once the time stepping parameters have been tweaked.

Can we enhance splitting’s stability & accuracy while retaining these benefits?

- The primary problem with basic splittings is that the component solvers are derived in isolation, with no concern for the coupling error.
- What if we instead derived new splitting approaches that explicitly account for inter-component coupling?

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# Additive Runge-Kutta Methods [Ascher et al. 1997; Araújo et al. 1997]

Although ARK methods may be derived for arbitrary splittings, here we consider splittings into two components: *explicit* and *implicit*,

$$\partial_t u = f_E(t, u) + f_I(t, u), \quad t \in [0, T], \quad u(0) = u_0,$$

We combine two  $s$ -stage methods (ERK + DIRK). Denoting  $t_{n,j} = t_n + c_j h$ ,

$$z_i = u_n + h \sum_{j=1}^{i-1} a_{i,j}^E f_E(t_{n,j}, z_j) + h \sum_{j=1}^i a_{i,j}^I f_I(t_{n,j}, z_j), \quad i = 1, \dots, s,$$

$$u_{n+1} = u_n + h \sum_{j=1}^s b_j (f_E(t_{n,j}, z_j) + f_I(t_{n,j}, z_j)) \quad [\text{solution}]$$

$$\tilde{u}_{n+1} = u_n + h \sum_{j=1}^s \tilde{b}_j^E (f_E(t_{n,j}, z_j) + f_I(t_{n,j}, z_j)) \quad [\text{embedding}]$$

We therefore have two Butcher tables to work with:

$c_1$	0	0	$\cdots$	0
$c_2$	$a_{2,1}^E$	0	$\cdots$	0
$\vdots$	$\vdots$	$\ddots$	$\ddots$	$\vdots$
$c_s$	$a_{s,1}^E$	$\cdots$	$a_{s,s-1}^E$	0
	$b_1$	$\cdots$	$b_{s-1}$	$b_s$
	$\tilde{b}_1$	$\cdots$	$\tilde{b}_{s-1}$	$\tilde{b}_s$

$c_1$	$a_{1,1}^I$	0	$\cdots$	0
$c_2$	$a_{2,1}^I$	$a_{2,2}^I$	$\cdots$	0
$\vdots$	$\vdots$	$\ddots$	$\ddots$	$\vdots$
$c_s$	$a_{s,1}^I$	$\cdots$	$a_{s,s-1}^I$	$a_{s,s}^I$
	$b_1$	$\cdots$	$b_{s-1}$	$b_s$
	$\tilde{b}_1$	$\cdots$	$\tilde{b}_{s-1}$	$\tilde{b}_s$





# ARK Coefficients

We have  $s^2 + 3s$  total free parameters  $(c_i, b_j, \tilde{b}_j, a_{i,j}^E, a_{i,j}^I)$ . As with traditional RK methods, these are chosen to satisfy desired constraints:

- Maximize the order of accuracy for each elemental method,
- Maximize the stability of each elemental method,
- Simplify repeated implicit solves (e.g. SDIRK or ESDIRK),
- Enable accurate/stable embeddings for temporal error estimation,
- Conservation of certain integrals (linear & quadratic first integrals).

Additionally, ARK methods must also satisfy coupling conditions *between* the methods, to the same accuracy as each elemental method.

# ARK Solution Algorithm – Tables

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0
0.87	0.87	0	0	0
0.6	0.53	0.07	0	0
1.0	0.40	-0.44	1.04	0
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

0	0	0	0	0
0.87	0.44	0.44	0	0
0.6	0.26	-0.09	0.44	0
1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

# ARK Solution Algorithm – Stage 1

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0
0.87	0.87	0	0	0
0.6	0.53	0.07	0	0
1.0	0.40	-0.44	1.04	0
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

0	0	0	0	0
0.87	0.44	0.44	0	0
0.6	0.26	-0.09	0.44	0
1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

Stage 1:  $z_1 = u_n$ ,

# ARK Solution Algorithm – Stage 2

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0	0	0	0	0	0
0.87	0.87	0	0	0	0.87	0.44	0.44	0	0
0.6	0.53	0.07	0	0	0.6	0.26	-0.09	0.44	0
1.0	0.40	-0.44	1.04	0	1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44		0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40		0.21	-0.49	0.87	0.40

Stage 1:  $z_1 = u_n$ ,

Stage 2:  $z_2 - 0.44hf_I(t_{n,2}, z_2) = u_n + h \left( a_{2,1}^E f_E(t_{n,1}, z_1) + a_{2,1}^I f_I(t_{n,1}, z_1) \right),$

# ARK Solution Algorithm – Stage 3

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0
0.87	0.87	0	0	0
0.6	0.53	0.07	0	0
1.0	0.40	-0.44	1.04	0
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

0	0	0	0	0
0.87	0.44	0.44	0	0
0.6	0.26	-0.09	0.44	0
1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40

Stage 1:  $z_1 = u_n$ ,

Stage 2:  $z_2 - 0.44hf_I(t_{n,2}, z_2) = u_n + h \left( a_{2,1}^E f_E(t_{n,1}, z_1) + a_{2,1}^I f_I(t_{n,1}, z_1) \right),$

Stage 3:  $z_3 - 0.44hf_I(t_{n,3}, z_3) = u_n + h \sum_{j=1}^2 \left( a_{3,j}^E f_E(t_{n,j}, z_j) + a_{3,j}^I f_I(t_{n,j}, z_j) \right),$

# ARK Solution Algorithm – Stage 4

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0	0	0	0	0	0
0.87	0.87	0	0	0	0.87	0.44	0.44	0	0
0.6	0.53	0.07	0	0	0.6	0.26	-0.09	0.44	0
1.0	0.40	-0.44	1.04	0	1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44		0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40		0.21	-0.49	0.87	0.40

Stage 1:  $z_1 = u_n$ ,

Stage 2:  $z_2 - 0.44hf_I(t_{n,2}, z_2) = u_n + h \left( a_{2,1}^E f_E(t_{n,1}, z_1) + a_{2,1}^I f_I(t_{n,1}, z_1) \right)$ ,

Stage 3:  $z_3 - 0.44hf_I(t_{n,3}, z_3) = u_n + h \sum_{j=1}^2 \left( a_{3,j}^E f_E(t_{n,j}, z_j) + a_{3,j}^I f_I(t_{n,j}, z_j) \right)$ ,

Stage 4:  $z_4 - 0.44hf_I(t_{n,4}, z_4) = u_n + h \sum_{j=1}^3 \left( a_{4,j}^E f_E(t_{n,j}, z_j) + a_{4,j}^I f_I(t_{n,j}, z_j) \right)$ ,

# ARK Solution Algorithm – Finish

To better understand the workings of an ARK time step, let's consider the ERK/ESDIRK pair, ARK3(2)4L[2]SA [Kennedy & Carpenter, 2001],

0	0	0	0	0	0	0	0	0	0
0.87	0.87	0	0	0	0.87	0.44	0.44	0	0
0.6	0.53	0.07	0	0	0.6	0.26	-0.09	0.44	0
1.0	0.40	-0.44	1.04	0	1.0	0.19	-0.60	0.97	0.44
	0.19	-0.60	0.97	0.44		0.19	-0.60	0.97	0.44
	0.21	-0.49	0.87	0.40		0.21	-0.49	0.87	0.40

Stage 1:  $z_1 = u_n$ ,

Stage 2:  $z_2 - 0.44hf_I(t_{n,2}, z_2) = u_n + h \left( a_{2,1}^E f_E(t_{n,1}, z_1) + a_{2,1}^I f_I(t_{n,1}, z_1) \right)$ ,

Stage 3:  $z_3 - 0.44hf_I(t_{n,3}, z_3) = u_n + h \sum_{j=1}^2 \left( a_{3,j}^E f_E(t_{n,j}, z_j) + a_{3,j}^I f_I(t_{n,j}, z_j) \right)$ ,

Stage 4:  $z_4 - 0.44hf_I(t_{n,4}, z_4) = u_n + h \sum_{j=1}^3 \left( a_{4,j}^E f_E(t_{n,j}, z_j) + a_{4,j}^I f_I(t_{n,j}, z_j) \right)$ ,

Finish:  $u_{n+1} = u_n + h \sum_{j=1}^4 b_j (f_{E,j} + f_{I,j})$  and  $\tilde{u} = u_n + h \sum_{j=1}^4 \tilde{b}_j (f_{E,j} + f_{I,j})$ .





# The ARKode Library

As a part of the FASTMath SciDAC Institute, we are constructing a library comprising these solvers named *ARKode*, that will be released as a new component solver within SUNDIALS.

- Nearly identical user interface as CVODE, albeit with separate user-specified  $f_E(t, y)$  and  $f_I(t, y)$  routines.
- Data structure agnostic – as long as the basic vector kernels are supplied this works with your native data structures.
- High-order accurate dense output, allowing efficient interpolation of results between integration steps, and reliable implicit predictors.
- Parameters optimized for iterative solvers and large-scale parallelism.
- Exhaustive suite of example and regression test problems.

# The ARKode Library

Key differences between ARKode and CVODE include:

- Support for disabling either  $f_E$  or  $f_I$ , allowing adaptive DIRK or ERK.
- Optional accelerated fixed-point nonlinear solver,
- Optional PCG, FGMRES Krylov solvers,
- Support for non-identity mass matrix,  $Mu' = f_E(t, u) + f_I(t, u)$ ,
- “Hot restart” support for problems with spatial adaptivity.
- “Set routines” allowing complete control over: Butcher table coefficients, time step adaptivity algorithm, temporal error estimation algorithm, implicit predictor algorithm, all internal solver parameters.

Plans:

- Fall 2013 public release; “friendly-user” release available now.
- Future support for partitioned symplectic methods (Hamiltonian systems).

# Transitioning Between CVODE and ARKode

Consistent API simplifies experimentation:

```
#include <cvode/cvode.h>
#include <cvode/cvode_dense.h>
#include <sundials/sundials_types.h>

/* User-supplied Functions */
static int f(realtype t, N_Vector y,
             N_Vector ydot, void *udata);
static int Jac(long int N, realtype t, N_Vector y,
               N_Vector fy, DlsMat J, void *udata,
               N_Vector t1, N_Vector t2, N_Vector t3);

int main() {

    ...

    void *cvode_mem = CVodeCreate(CV_BDF, CV_NEWTON);
    CVodeInit(cvode_mem, f, T0, y);
    CVodeSStolerances(cvode_mem, reltol, abstol);
    CVDense(cvode_mem, NEQ);
    CVDlsSetDenseJacFn(cvode_mem, Jac);

    for (iout=0; iout<Nt; iout++)
        CVue(cvode_mem, tout, y, &t, CV_NORMAL);
}
```

```
#include <arkode/arkode.h>
#include <arkode/arkode_dense.h>
#include <sundials/sundials_types.h>

/* User-supplied Functions */
static int fe(realtype t, N_Vector y,
              N_Vector ydot, void *udata);
static int fi(realtype t, N_Vector y,
              N_Vector ydot, void *udata);
static int Jac(long int N, realtype t, N_Vector y,
               N_Vector fy, DlsMat J, void *udata,
               N_Vector t1, N_Vector t2, N_Vector t3);

int main() {

    ...

    void *arkode_mem = ARKodeCreate();
    ARKodeInit(arkode_mem, fe, fi, T0, y);
    ARKodeSStolerances(arkode_mem, reltol, abstol);
    ARKDense(arkode_mem, NEQ);
    ARKDlsSetDenseJacFn(arkode_mem, Jac);

    for (iout=0; iout<Nt; iout++)
        ARKode(arkode_mem, tout, y, &t, ARK_NORMAL);
}
```

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# PDE Cosmology Model Problem – Radiating Ionization Front

Find  $u(\mathbf{x}, t)$ ,  $v(\mathbf{x}, t)$  s.t. for  $(\mathbf{x}, t) \in [-1, 1]^3 \times [0, 1]$ ,

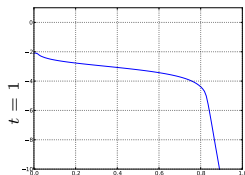
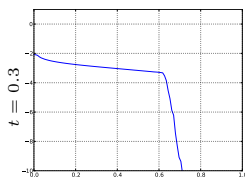
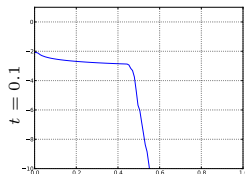
$$\partial_t u = \nabla \cdot (\beta u) + \nabla \cdot (\mu \nabla u) + f_u(u, v),$$

$$\partial_t v = \nabla \cdot (\beta v) + f_v(u, v),$$

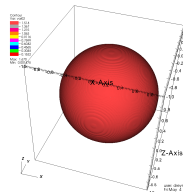
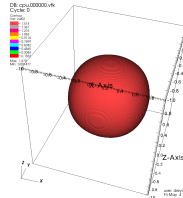
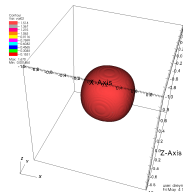
Where

- $u(\mathbf{x}, 0) = 10^{-8}$ ,  $v(\mathbf{x}, 0) = 1.67$ .
- $\nabla u \cdot \mathbf{n}|_{\partial\Omega} = 0$ ,  $\nabla v \cdot \mathbf{n}|_{\partial\Omega} = 0$ .
- $f_u(u, v) = 25000 \delta_0(\mathbf{x}) - 1800000 uv$ .
- $f_v(u, v) = a(c - v)^2 - bv(c - v) + 567000 uv$ .
- $\beta(u, v, \mathbf{x}) = -\frac{\alpha u(c - v)}{\|\mathbf{x}\|} \mathbf{x}$ ,  $\alpha = \{10, 100, 250\}$ ,
- $\mu(v) = \frac{100}{v}$ ,
- $a = 2.445$ ,  $b = 0.01118$ ,  $c = 1.673$ .

$\log_{10}(u)$  Profiles



Ionized region



# Radiating Ionization Test Solvers

## Discretizations:

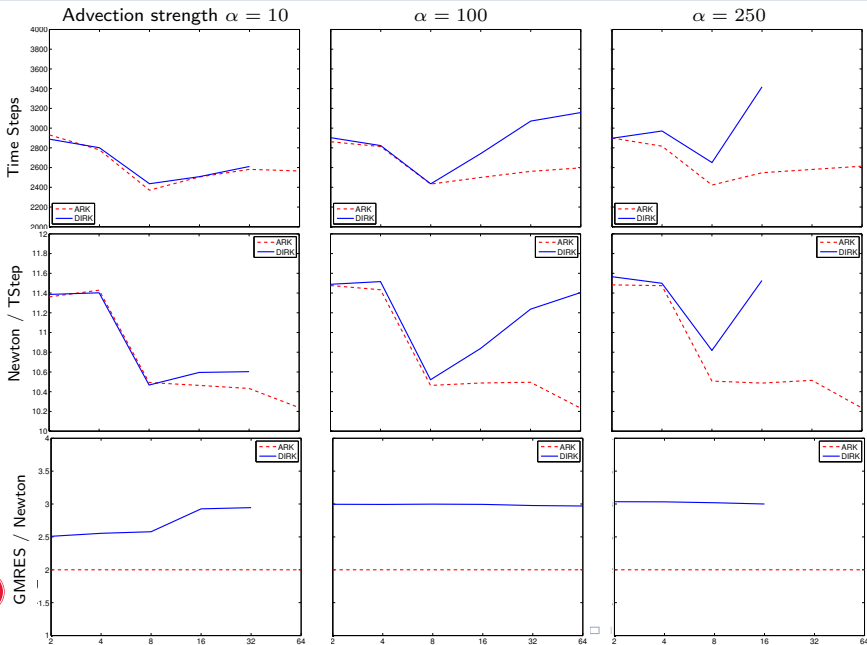
- FV grid of size  $N_x \times N_y \times N_z$ , with a fixed  $32 \times 32 \times 32$  grid per MPI task.
- Process sizes:  $2(2 \times 1 \times 1)$ ,  $4(2 \times 2 \times 1)$ ,  $8(2 \times 2 \times 2)$ ,  $16(4 \times 2 \times 2)$ ,  $32(4 \times 4 \times 2)$ ,  $64(4 \times 4 \times 4)$ .
- $\nabla \cdot (\mu \nabla u)$  discretized using 2nd-order centered differences.
- $\nabla \cdot (\beta u)$  discretized using CD or  $\mathcal{O}(\Delta x)$  upwind, depending on  $\|\beta\| > 10^{-3}$ .
- Time:  $\mathcal{O}(\Delta t^4)$  ARK and DIRK methods.

## Solvers:

- Nonlinear: inexact Newton, `maxit` = 5.
- Linear: matrix-free GMRES, `maxit` = 50.
- Precond: AMG [HYPRE], Rx-Diff only (no advection).
- PID controller [Kennedy & Carpenter, 2001], with error estimate

$$e = \left( \frac{1}{N} \sum_i \left( \frac{u_i - \tilde{u}_i}{r_{tol} u_i + a_{tol}} \right)^2 \right)^{1/2}, \quad r_{tol} = 10^{-3}, \quad a_{tol} = 10^{-9}.$$

# Radiating Ionization Results – Iteration Weak Scaling



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# Conclusions

ARK methods allow accurate/stable splitting of multi-rate problems:

- IMEX couplings are properly handled,
- No Dahlquist barrier – high accuracy and stability simultaneously possible.

RK under-pinnings allow robust and theoretically rigorous methods for error estimation and time adaptivity.

Works well with spatial adaptivity (due to one-step approach), with tunings for vector resizing and finite element discretizations.

Flexible infrastructure:

- Allows adaptive ERK or IRK methods alone,
- Enables domain-specific knowledge into the formulation.

“Convenient” preconditioners need not be chastised, since many solver difficulties (non-differentiability, nonlinearity, complex discretizations) often lie in slow components that can be treated explicitly.



# Thanks & Acknowledgements

## Collaborators/Students:

- Carol S. Woodward [LLNL]
- Alan C. Hindmarsh [LLNL]
- David J. Gardner [SMU, PhD]
- Jean M. Sexton [SMU, MS]



## Current Grant/Computing Support:

- DOE SciDAC & INCITE Programs
- NSF AST & XSEDE Programs
- DOD DURIP Program
- SMU Center for Scientific Computation



## Software:

- ARKode – <http://faculty.smu.edu/reynolds/arkode>
- SUNDIALS – <https://computation.llnl.gov/casc/sundials>
- Enzo (cosmology) – <http://enzo-project.org>



SMU



## Outline

## 5 Extra Slides



# First-Order Splittings

Denote  $S_i(h, u(t_n))$  as a solver for the component  $\partial_t u = f_i(t, u)$  over a time step  $t_n \rightarrow t_n + h \equiv t_{n+1}$ , with initial condition  $u(t_n)$ .

To evolve  $u(t_n) \rightarrow u(t_{n+1})$ , we can use different solvers at the same  $h$ ,

$$\begin{aligned}\hat{u} &= S_1(h, u(t_n)), \\ u(t_{n+1}) &= S_2(h, \hat{u}),\end{aligned}$$

or we may subcycle time steps for individual components,

$$\begin{aligned}\hat{u}_{j+1} &= S_1\left(\frac{h}{m}, \hat{u}_j\right), \quad j = 0, \dots, m, \quad \hat{u}_0 = u(t_n), \\ u(t_{n+1}) &= S_2(h, \hat{u}_m),\end{aligned}$$

Unless the  $S_i$  commute [i.e.  $S_1(h, S_2(h, u)) = S_2(h, S_1(h, u))$ ] or the splitting is symmetric, these methods are at best  $O(h)$  accurate (*no matter the accuracy of the individual solvers*).

# Fractional Step (Strang) Splitting [Strang 1968]

“Strang splitting” attempts to achieve a higher-order method using these separate component solvers, through manually symmetrizing the operator:

$$\hat{u}_1 = S_1 \left( \frac{h}{2}, u(t_n) \right),$$

$$\hat{u}_2 = S_2 (h, \hat{u}_1),$$

$$u(t_{n+1}) = S_1 \left( \frac{h}{2}, \hat{u}_2 \right).$$

This approach is  $O(h^2)$  as long as each  $S_i$  is  $O(h^2)$ .

*However:*

- This asymptotic accuracy may not be achieved until  $h$  is very small, since error terms are dominated by inter-process interactions [Ropp, Shadid,& Ober 2005].
- Numerical stability isn't guaranteed *even if  $h$  is stable for each component* [Estep et al., 2007].

## Operator-Splitting Issues – Accuracy [Ropp, Shadid, &amp; Ober 2005]

Coupled systems can admit destabilizing modes not present in either component, due to *numerical resonance instabilities* [Grubmüller 1991].

Brusselator Example (Reaction-Diffusion):

$$\partial_t T = \frac{1}{40} \nabla^2 T + 0.6 - 3T + T^2 C,$$

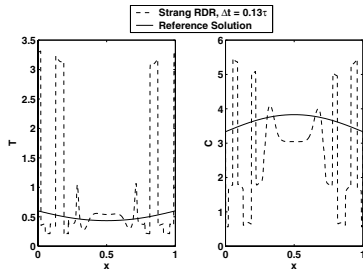
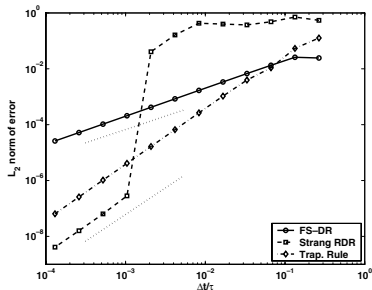
$$\partial_t C = \frac{1}{40} \nabla^2 C + 2T - T^2 C,$$

Three solvers:

- Basic split: **D** (trap.) then **R** (subcycled BDF).
- Strang:  $\frac{h}{2}\mathbf{R}$ ,  $h\mathbf{D}$ ,  $\frac{h}{2}\mathbf{R}$ ,
- Fully implicit trapezoidal rule,

Results:

- is stable but inaccurate for all tests;
- unusable until  $h$  is “small enough”.



# Operator Splitting Issues – Accuracy [Estep 2007]

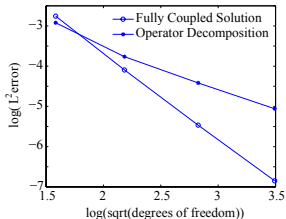
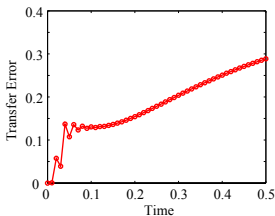
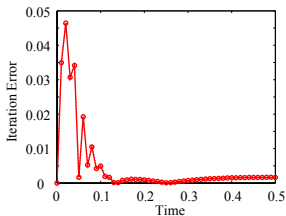
Consider  $\Omega = \Omega_1 \cup \Omega_2$  where the subdomains share a boundary  $\Gamma = \partial\Omega_1 \cap \partial\Omega_2$ :

$$\partial_t u_1 = \nabla^2 u_1, \quad x \in \Omega_1, \quad \partial_t u_2 = \frac{1}{2} \nabla^2 u_2, \quad x \in \Omega_2,$$

$$u_1 = u_2, \quad \nabla u_1 \cdot n = \nabla u_2 \cdot n, \quad \text{for } x \in \Gamma.$$

Solved using one Gauss-Seidel iteration:  $S_1$  on  $\Omega_1$ , then  $S_2$  on  $\Omega_2$  (both trapezoidal). Errors from not iterating to convergence, and from error transfer between subdomains.

Using adjoints, they measured these errors separately:



- Error from incomplete iteration decreased with time.
- Transfer error accumulated and became dominant with time.
- While each  $S_i$  was  $O(h^2)$ , the coupled method was only  $O(h)$ .

# Operator-Splitting Issues – Stability [Estep et al., 2007]

Second Reaction-Diffusion Example (split subcycling; exact solvers):

$$\partial_t u = -\lambda u + u^2, \quad u(0) = u_0, \quad t > 0.$$

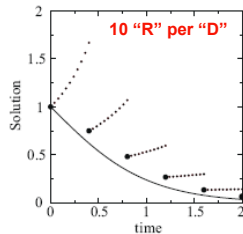
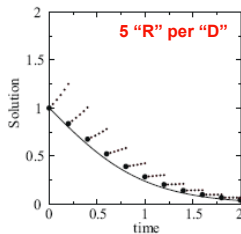
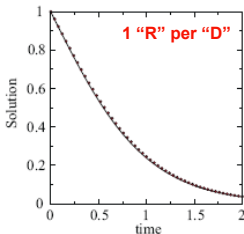
Phase 1 (R):  $\partial_t u_r = u_r^2, \quad u_r(t_n) = u_n, \quad t \in [t_n, t_{n+1}],$

Phase 2 (D):  $\partial_t u_d = -\lambda u_d, \quad u_d(t_n) = u_r(t_{n+1}), \quad t \in [t_n, t_{n+1}].$

True solution,  $u(t) = \frac{u_0 e^{-\lambda t}}{1 + \frac{u_0}{\lambda} (e^{-\lambda t} - 1)},$  is well-defined  $\forall t$  if  $\lambda > u_0$ .

Split solution,  $u(t_{n+1}) = \frac{u(t_n) e^{-\lambda h}}{1 - u(t_n) h},$  can blow up in finite time.

Results using 50  
time steps, with  
varying amounts  
of subcycling.





# Time Step Selection

## Stability:

If IRK portion is  $A$ -stable, linear instability can only arise due to the ERK, so time steps must satisfy the explicit stability restriction,

$$h \leq h_{exp}.$$

## Accuracy:

For an order  $q$  method with order  $p$  embedding (typically  $q = p + 1$ ),

$$\begin{aligned} \|u(t_{n+1}) - u_{n+1}\| &\leq \|\tilde{u} - u_{n+1}\| + \|u(t_{n+1}) - \tilde{u}\| \\ &\leq \|\tilde{u} - u_{n+1}\| + \mathcal{O}(h^{p+1}). \end{aligned}$$

Hence  $h$  adaptivity may utilize  $\|\tilde{u} - u_{n+1}\|$  for error estimation, e.g.

$$h_{n+1} = c h_n \left( \frac{r_{tol}}{\|\tilde{u} - u_{n+1}\|} \right)^{1/(p+1)}$$

In practice, a variety of adaptivity methods may be used, and are often based on control-theoretic techniques (e.g. I, PI, PID) [Gustafsson 1991 & 1994; Söderlind 2003].

# Implicit Predictors and Dense Output

Accurate initial guesses are critical for nonlinear solver efficiency & robustness.

Each stage approximates the solution at a different time, i.e.  $z_i \approx y(t_n + c_i h)$ , giving rise to an interpolating polynomial (a.k.a. *dense output*),

$$u_d(t) = u_n + h \sum_{i=1}^s \left( \left( \sum_{j=1}^{p^*} b_{i,j}^* \left( \frac{t-t_n}{h} \right)^j \right) f(t_{n,i}, z_i) \right).$$

- Coefficients  $b_{i,j}^*$  may be determined *a priori* for each method.
- $u_d(t)$  accuracy is limited by order of each  $z_i$  (typically lower than  $q$ ).
- Extrapolation accuracy rapidly diminishes for  $t > t_{n+1}$ . Hence predictors are most helpful for early stages of next step.
- Later stages  $z_i$  may instead be predicted by previous stages,  $z_1, \dots, z_{i-1}$ .