

Basic Tensor

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不变

From vector to tensor:

Vector and contravariant:

$$\vec{V} = \sum_i \vec{e}_i V^i$$

Following Einstein summation rule, we have

$$\vec{V} = \vec{e}_i V^i$$

Written in explicit:

$$\vec{V} = \left[\vec{e}_1, \vec{e}_2, \dots, \vec{e}_n \right] \begin{bmatrix} V^1 \\ V^2 \\ \dots \\ V^n \end{bmatrix}$$

To obtain the following form,

$$\begin{bmatrix} V^1 \\ V^1 \\ \dots \\ V^n \end{bmatrix} = B \begin{bmatrix} V^1 \\ V^1 \\ \dots \\ V^n \end{bmatrix} \quad \text{B for backward}$$

Written in explicit,

$$V^i = \underline{B_j^i} V^j$$

Note the order here

We define the transformation rule of basis

$$\begin{bmatrix} \tilde{e}_1, \tilde{e}_2, \dots, \tilde{e}_n \end{bmatrix} = \begin{bmatrix} \vec{e}_1, \vec{e}_2, \dots, \vec{e}_n \end{bmatrix} F \quad \text{F for forward}$$

Written in explicit,

$$\tilde{e}_i = \underline{\vec{e}_j} F_j^i$$

Note the order here

Since F and B are contra to each other,
this is the reason of contravariant

The invariant of vector gives

$$\vec{V} = \begin{bmatrix} \vec{e}_1, \vec{e}_2, \dots, \vec{e}_n \end{bmatrix} \begin{bmatrix} V^1 \\ V^1 \\ \dots \\ V^n \end{bmatrix} = \begin{bmatrix} \vec{e}_1, \vec{e}_2, \dots, \vec{e}_n \end{bmatrix} FB \begin{bmatrix} V^1 \\ V^1 \\ \dots \\ V^n \end{bmatrix} = \begin{bmatrix} \tilde{e}_1, \tilde{e}_2, \dots, \tilde{e}_n \end{bmatrix} \begin{bmatrix} V^1 \\ V^1 \\ \dots \\ V^n \end{bmatrix}, FB = 1$$

Covector

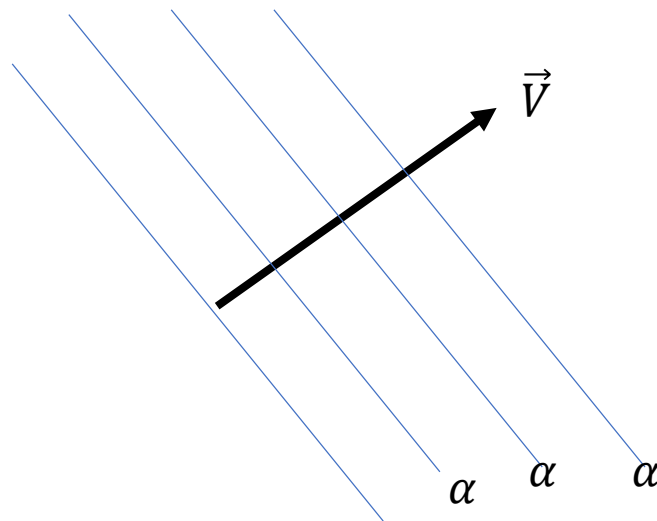
Map a vector to a scalar:

$$\alpha : V \rightarrow R$$

Geometric illustration:

Consider the fact that the inner product of two vectors is a scalar,

If we regard vector as arrow, then covector is stack, which is living in the mirror world (dual vector space)



This is because:

$$\alpha \vec{V} = aV_1 + bV_2$$

is a set of parallel lines (stack)

Define dual vector basis

$$\vec{\mathcal{E}}^1, \vec{\mathcal{E}}^2, \dots, \vec{\mathcal{E}}^n$$

Which satisfies:

$$\vec{\mathcal{E}}^i \cdot \vec{e}_j = \delta_j^i$$

So we can express any covector as:

$$\vec{W} = \vec{\mathcal{E}}^i W_i$$

It's easy to show

$$\vec{\mathcal{E}}^i = \sum_j \vec{\mathcal{E}}^j B_i^j$$

$$W_i = \sum_j W_j F_i^j$$

Summary here:

$$\vec{\tilde{e}}_i = \sum_j \vec{e}_j F_i^j$$

covariant

把指标放在下标表示做transformation的时候使用F，即表示协变covariant

$$\vec{\mathcal{E}}^i = \sum_j B_j^i \vec{\mathcal{E}}^j$$

contravariant

把指标放在上标表示做transformation的时候使用B，即表示逆变contravariant

$$V^i = \sum_j B_j^i V^j$$

contravariant

把指标放在上标表示做transformation的时候使用B，即表示逆变contravariant

$$W_i = \sum_j W_j F_i^j$$

covariant

把指标放在下标表示做transformation的时候使用F，即表示协变covariant

Matric tensor

Now we define length.

First, we define metric tensor, so we can separate the dependence of basis apart:

$$g_{ij} = \vec{e}_i \vec{e}_j$$

Written in explicit:

$$g = \begin{bmatrix} \vec{e}_1 \vec{e}_1 & \vec{e}_1 \vec{e}_2 & \dots & \vec{e}_1 \vec{e}_n \\ \vec{e}_2 \vec{e}_1 & \vec{e}_2 \vec{e}_2 & \dots & \vec{e}_2 \vec{e}_n \\ \dots & \dots & \dots & \dots \\ \vec{e}_m \vec{e}_1 & \vec{e}_m \vec{e}_2 & \dots & \vec{e}_m \vec{e}_n \end{bmatrix}$$

And the length can be written as:

$$\|\vec{V}\| = \sum_{ij} V^i V^j \vec{e}_i \vec{e}_j = \sum_{ij} V^i V^j g_{ij}$$

And the metric tensor follows the transformation rule:

$$g_{ij} = \vec{e}_i \vec{e}_j = \vec{e}_k F_i^k \vec{e}_m F_j^m = \vec{e}_k \vec{e}_m F_i^k F_j^m$$

Metric tensor as a map:

$$g : V \times V \rightarrow R$$

Metric tensor maps two vectors into a scalar:

$$g(\vec{v}, \vec{w}) = v^i w^j g_{ij}$$

Multilinear maps

Let V_1, V_2, \dots, V_N be vector spaces over the field R , a map

$$T : V_1 \times V_2 \times \dots \times V_N \rightarrow R$$

$$T(v_1, v_2, \dots, v_n) = a$$

is said to be multilinear if it is linear in each argument separately:

$$\begin{aligned} & T(v_1, v_2, \dots, av_i + bv'_i, \dots, v_n) \\ &= aT(v_1, v_2, \dots, v_i, \dots, v_n) + bT(v_1, v_2, \dots, v'_i, \dots, v_n) \end{aligned}$$

Multilinear maps can be added and multiplied by scalars in the usual fashion

$$\begin{aligned} & (aT + bS)(v_1, v_2, \dots, v_n) \\ &= aT(v_1, v_2, \dots, v_n) + bS(v_1, v_2, \dots, v_n) \end{aligned}$$

Therefore, multilinear map forms a vector space, denoted by

$$V_1^* \times V_2^* \times \dots \times V_N^* \quad \text{我们总是需要covector去吃掉vector}$$

called the tensor product of the dual spaces $V^*_1, V^*_2, \dots, V^*_N$

Tensor spaces of type (r, s)

Let V be a vector space of dimension n over the field R . Setting

$$V_1 = V_2 = \dots = V_r = V^*, V_{r+1} = V_{r+2} = \dots = V_{r+s} = V$$

We refer to any multilinear map

$$T : \underbrace{V^* \times V^* \times \dots \times V^*}_r \times \underbrace{V \times \dots \times V}_s \rightarrow R$$

contravariant degree covariant degree

as a tensor of type (r, s) on V .

因为 T 可以吃掉 r 个 covector 以及 s 个 vector, 所以 T 的基可以写成

$$V^{(r,s)} = \underbrace{V \otimes V \otimes \dots \otimes V}_r \otimes \underbrace{V^* \otimes \dots \otimes V^*}_s$$

Basis representation

$$V = \{\vec{e}_1, \dots, \vec{e}_n\}$$

$$V^* = \{\vec{\varepsilon}^1, \dots, \vec{\varepsilon}^n\}$$

where the dual basis is defined by

$$\vec{\varepsilon}^i(\vec{e}_j) = \delta_j^i$$

And the basis for (r, s)-tensor is

$$V^{(r,s)} = \{\vec{e}_{i_1} \otimes \vec{e}_{i_2} \otimes \dots \otimes \vec{e}_{i_r} \otimes \vec{\varepsilon}^{j_1} \otimes \dots \otimes \vec{\varepsilon}^{j_s}\}$$

And any tensor can be expanded as

$$T^{(r,s)} = T_{j_1 j_2 \dots j_s}^{i_1 i_2 \dots i_r} \vec{e}_{i_1} \otimes \vec{e}_{i_2} \otimes \dots \otimes \vec{e}_{i_r} \otimes \vec{\varepsilon}^{j_1} \otimes \dots \otimes \vec{\varepsilon}^{j_s}$$

And the coefficient is:

$$T_{j_1 j_2 \dots j_s}^{i_1 i_2 \dots i_r} = T^{(r,s)}(\vec{\varepsilon}^{j_1}, \dots, \vec{\varepsilon}^{j_s}, \vec{e}_{i_1}, \dots, \vec{e}_{i_r})$$

Example: (0, 2)-tensor

$$V^{(0,2)} = \mathbf{V}^* \otimes \mathbf{V}^* \quad \omega \otimes \rho$$

And u, v are two vector, so we have

$$\omega \otimes \rho(u, v) = \omega(u)\rho(v)$$

Into basis

$$\mathbf{e}^i \otimes \mathbf{e}^j$$

And we can represent a (0, 2)-tensor $\omega \otimes \rho$ as

$$\omega \otimes \rho = \omega_i \mathbf{e}^i \otimes \rho_j \mathbf{e}^j = \omega_i \rho_j \mathbf{e}^i \otimes \mathbf{e}^j$$

$$T = T_{ij} \mathbf{e}^i \otimes \mathbf{e}^j$$

So we have

$$T(u, v) = T_{ij} u^i v^j$$

We can express T as a matrix

$$T = \begin{bmatrix} T_{11} & T_{12} & \dots & T_{1n} \\ T_{21} & T_{22} & \dots & T_{2n} \\ \dots & \dots & \dots & \dots \\ T_{n1} & T_{n2} & \dots & T_{nn} \end{bmatrix}$$

For example, the metric tensor is a (2,0)-type tensor:

$$g = \{ g | g_{ij} = \vec{e}_i \vec{e}_j \}$$

So it gives the expected length:

$$g(\vec{v}, \vec{w}) = g_{ij} v^i w^j$$

For example, the metric tensor is a (0,2)-type tensor:

$$G = \{ G | G^{ij} = \vec{\varepsilon}^i \vec{\varepsilon}^j \}$$

For example, the linear map is a (1,1)-type tensor:

$$L = L_j^i \vec{e}_i \vec{\varepsilon}^j$$

So the linear map gives the expected transformation rule:

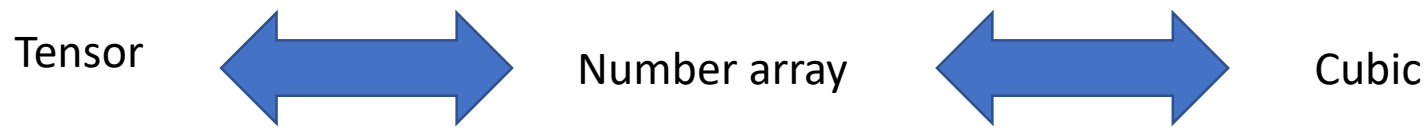
$$L(\vec{v}) = L_j^i \vec{e}_i \vec{\varepsilon}^j v = L_j^i \vec{e}_i v^j = w^i \vec{e}_i = \vec{w}$$

它使用一个covector收缩掉一个vector，剩下一个vector

Written in matrix form, we have

$$L = \begin{bmatrix} L_1^1 & L_2^1 & \dots & L_m^1 \\ L_1^2 & L_2^2 & \dots & L_m^2 \\ \dots & \dots & \dots & \dots \\ L_1^n & L_2^n & \dots & L_m^n \end{bmatrix}$$

How to picture a tensor with cubic?



(1, 0)

$$\begin{bmatrix} V^1 \\ V^2 \\ \dots \\ V^n \end{bmatrix}$$



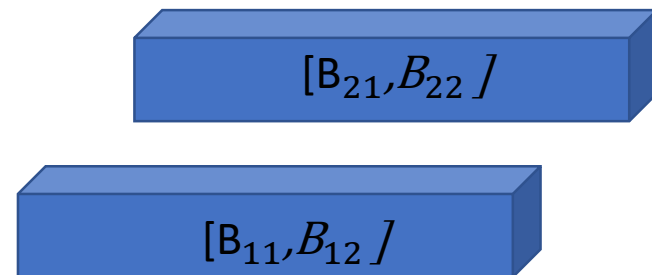
(0, 1)

$$[W_1 \quad W_2 \quad \dots \quad W_n]$$



(0, 2)

$$B = [[B_{11}, B_{12}], [B_{21}, B_{22}]]$$



(1, 1)

$$\begin{bmatrix} L_1^1 & L_2^1 & \dots & L_m^1 \\ L_1^2 & L_2^2 & \dots & L_m^2 \\ \dots & \dots & \dots & \dots \\ L_1^n & L_2^n & \dots & L_m^n \end{bmatrix}$$



一个立方体最多可可视化三阶张量，更高阶张量无法可视化。

Tensor equation

如果我们把所有的方程都写成张量方程的形式，我们知道不管在什么样的坐标系下，方程的形式都是一样的。这个就是张量方程的好处。

Tensor and quantum mechanics

量子力学里面的算符都是线性映射，都可以用张量进行表示。其中Dirac标记里面的bra和ket分别对应偶空间和空间，所以量子力学是具有张量结构的。

因为单体问题是（0+1）维的问题，所以我们只需要矩阵就可以描述问题；

但是多体问题是（3+1）维的问题，所以矩阵已经不足够描述问题了，必须启用张量。

