

Physics 195 Problem Set 4

Problem 2 (Problem 2.5 of Ryden.)

Consider blackbody radiation at a temperature T . Show that for an energy threshold $E_0 \gg kT$, the fraction of the blackbody photons that have energy $hf > E_0$ is

$$\frac{n(hf > E_0)}{n_\gamma} \approx 0.42 \left(\frac{E_0}{kT} \right)^2 \exp \left(-\frac{E_0}{kT} \right)$$

The cosmic background radiation is currently called the “cosmic *microwave* background.” However, photons with $\lambda < 1$ mm actually lie in the *far infrared* range of the electromagnetic spectrum. It’s time for truth in advertising: what fraction of the photons in today’s “cosmic microwave background” are actually far infrared photons?

Solution:

The number density of photons is given by

$$n(f)df = \frac{\varepsilon(f)df}{hf} = \frac{8\pi}{c^3} \frac{f^2 df}{e^{hf/kT} - 1} \quad (1)$$

for a given frequency range f and $f + df$.

To get the number of the blackbody photons that have energy $hf > E_0$, we integrate this from E_0/h to ∞ :

$$n(hf > E_0) = \frac{8\pi}{c^3} \int_{E_0/h}^{\infty} \frac{f^2 df}{e^{hf/kT} - 1} \quad (2)$$

Using the substitution

$$u = f - \frac{E_0}{h} \quad (3)$$

we can transform the integral into

$$n(hf > E_0) = \frac{8\pi}{c^3} \int_0^{\infty} \frac{(u + E_0/h)^2 du}{\exp \left(\frac{hu}{kT} + \frac{E_0}{kT} \right) - 1} \quad (4)$$

Multiplying the numerator and denominator by $(h/kT)^2$,

$$n(hf > E_0) = \frac{8\pi}{c^3} \left(\frac{kT}{h} \right)^3 \int_0^{\infty} \left(\frac{hu}{kT} + \frac{E_0}{kT} \right)^2 \left[\exp \left(\frac{hu}{kT} + \frac{E_0}{kT} \right) - 1 \right]^{-1} du \left(\frac{h}{kT} \right) \quad (5)$$

Let the whole integral be I such that

$$n(hf > E_0) = \frac{8\pi}{c^3} \left(\frac{kT}{h} \right)^3 I \quad (6)$$

The number density of photons in blackbody radiation is

$$n_\gamma = \beta T^3 \quad (7)$$

where

$$\beta = \frac{2.4041}{\pi^2} \frac{k^3}{h^3 c^3} \quad (8)$$

Thus, we can expand n_γ into

$$n_\gamma = \frac{8\pi k^3}{h^3 c^3} (2.4041) T^3 \quad (9)$$

The fraction of the blackbody photons that have energy $hf > E_0$ will be

$$\frac{n(hf > E_0)}{n_\gamma} = \frac{\frac{8\pi}{c^3} \left(\frac{kT}{h}\right)^3}{\frac{8\pi k^3 T^3}{h^3 c^3}} \cdot \frac{1}{2.4041} I = 0.42I \quad (10)$$

Going back to the integral, we have defined I to be

$$I = \int_0^\infty \left(\frac{hu}{kT} + \frac{E_o}{kT} \right)^2 \left[\exp \left(\frac{hu}{kT} + \frac{E_o}{kT} \right) - 1 \right]^{-1} du \left(\frac{h}{kT} \right) \quad (11)$$

Using another change of variables

$$x = \frac{hu}{kT}, \quad (12)$$

we transform the integral into

$$I = e^{-E_o/kT} \int_0^\infty \left(x + \frac{E_o}{kT} \right)^2 e^{-x} dx \quad (13)$$

This expands to

$$I = e^{-E_o/kT} \left\{ \int_0^\infty x^2 e^{-x} dx + \frac{2E_o}{kT} \int_0^\infty x e^{-x} dx + \left(\frac{E_o}{kT} \right)^2 \int_0^\infty e^{-x} dx \right\} \quad (14)$$

We can evaluate the integrals using the gamma function and get

$$I = e^{-E_o/kT} \left[2 + \frac{2E_o}{kT} + \left(\frac{E_o}{kT} \right)^2 \right] \quad (15)$$

Note that for $E_0 \gg kT$, we can ignore the first two terms:

$$I \approx \left(\frac{E_0}{kT} \right)^2 \exp \left(-\frac{E_0}{kT} \right) \quad (16)$$

Therefore, the fraction of the blackbody photons that have energy $hf > E_0$ is

$$\frac{n(hf > E_0)}{n_\gamma} \approx 0.42 \left(\frac{E_0}{kT} \right)^2 \exp \left(-\frac{E_0}{kT} \right) \quad (17)$$