

介绍



教程简介:

• 面向对象:量子计算初学者

• 依赖课程:线性代数,量子力学(非必需)

知乎专栏:

https://www.zhihu.com/column/c_1501138176371011584

Github & Gitee 地址:

https://github.com/mymagicpower/qubits https://gitee.com/mymagicpower/qubits

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初态制备



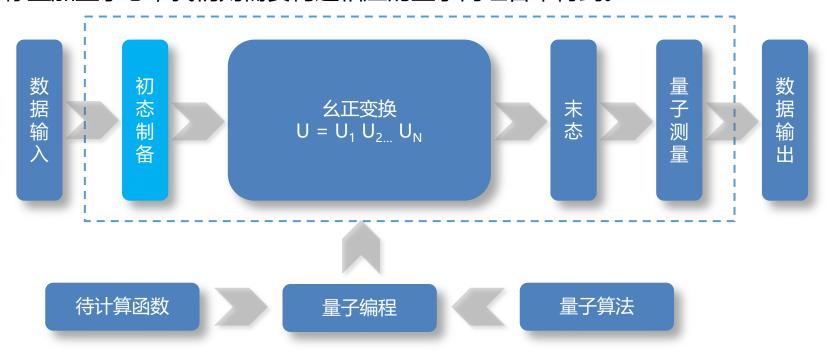
初态制备

指的是量子计算中初始量子态的构造,是量子计算的初始步骤。

以单比特为例:

从基态 |0> 出发制备任给目标叠加态的过程称为初态制备。

在实际量子运算中,我们可以直接得到的默认量子态是基态 $|0\rangle$,通过 X 门可以得到基态 $|1\rangle$ 。 对于任给的目标叠加量子态,我们则需要构造相应的量子门组合来得到。



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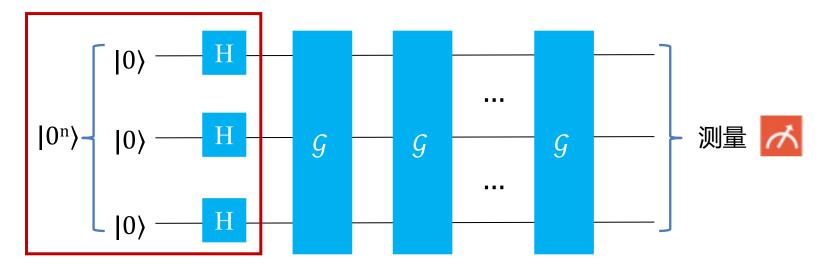


以两比特态空间为例,从 |0⟩^{⊗2} (|0⟩⊗|0⟩) 出发,对每个量子比特进行 H 门操作可以得到两比特空间中所有基态的均匀叠加。

类似地,在任意维态空间中,均可以借助 H 门从多维的 |0> 基态出发,得到所有基态均匀线性组合的量子态。这种量子态称为最大叠加态,很多初始状态要求为最大叠加态,量子计算的并行性也有赖于此。

$$(H|0\rangle)^{\otimes n} = \frac{1}{\sqrt{2^n}} \sum_{x=0}^{2^{n-1}} |x\rangle$$

量子态准备 - 最大叠加态



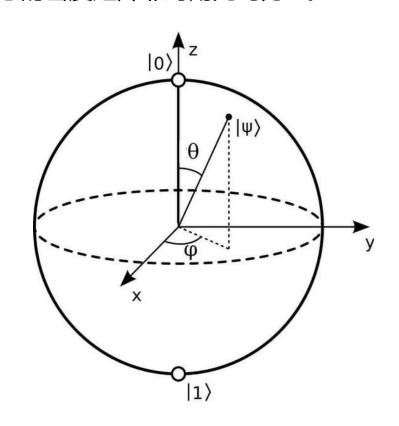
纯态和混态



非基态的量子态都为叠加态。

叠加态又可以分为相干叠加和非相干叠加,分别称为纯态和混态。

如将态空间与Bloch球关联,球面上量子态为纯态,球体内的量子态为混态。另一种重要的区分方式为密度矩阵,混态的密度矩阵非对角元均为 0。



密度矩阵:

$$\rho = |\psi\rangle\langle\psi| = \begin{bmatrix} \alpha \\ \beta \end{bmatrix} \begin{bmatrix} \overline{\alpha} & \overline{\beta} \end{bmatrix} = \begin{bmatrix} \alpha\overline{\alpha} & \alpha\overline{\beta} \\ \beta\overline{\alpha} & \beta\overline{\beta} \end{bmatrix} = \begin{bmatrix} |\alpha|^2 & \alpha\overline{\beta} \\ \beta\overline{\alpha} & |\beta|^2 \end{bmatrix}$$

其中
$$\langle \psi | = |\psi \rangle^{\dagger}$$
 ,且矩阵迹为: $\mathrm{tr}(\rho) = |\alpha|^2 + |\beta|^2 = 1$

量子纠缠



对于一个量子系统的量子态 |ψ⟩:

直积态:可以表示成形如 $|\Psi\rangle = |\Psi_0\rangle \otimes |\Psi_1\rangle$ 的两个量子系统的张量积形式

纠缠态:不能进行这种直积分解的量子态

例如,对两比特的 Bell 态 $\frac{1}{\sqrt{2}} \mid 00 \rangle + \frac{1}{\sqrt{2}} \mid 11 \rangle$:它不能写成两个单比特量子态的直积(张量积)形式。

量子纠缠态有超越经典关联的量子关联。 为了发挥量子计算的并行性和高效性,量子计算的量子比特之间应当有着纠缠关联。



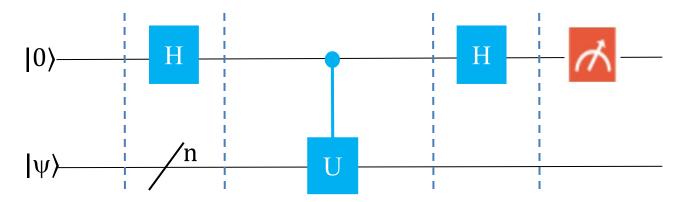
Hadamard Test

Hadamard Test量子线路的主要作用:

对任给的幺正变换 U 和量子态 $|\psi\rangle$,可以给出该**幺正算符 U 在量子态上的投影期望** $\langle\psi|U|\psi\rangle$ 。即可以通过测量一个辅助比特(ancilla qubit)来方便地得到一个**幺正算符 U 对于一个量子态** $|\psi\rangle$ **的平均值**。

$$\langle \psi | U | \psi \rangle = \text{Re } \langle \psi | U | \psi \rangle + \text{Im } \langle \psi | U | \psi \rangle \text{ i}$$
 实部 虚部

其中的实部对应的量子线路图为:



整个量子线路可以视为,对两个寄存器中量子比特组成的一个n+1维量子态 $|0\rangle|\psi\rangle$,进行量子门操作组合: $Q=(H\otimes I^{\otimes n})(\operatorname{Ctrl} - U)(H\otimes I^{\otimes n})$ 其中 $\operatorname{Ctrl} - U$ 表示基于幺正算符 U 的受控门





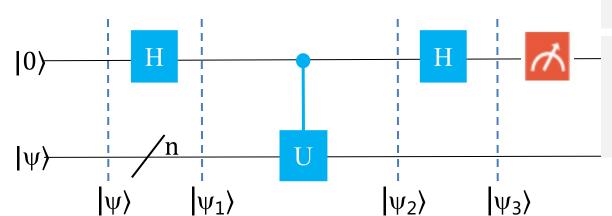
$$|\psi_1\rangle = H|0\rangle \otimes |\psi\rangle$$

$$= \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \otimes |\psi\rangle$$

$$= \frac{1}{\sqrt{2}}(|0\rangle|\psi\rangle + |1\rangle|\psi\rangle)$$

$$|\psi_2\rangle = (\operatorname{Ctrl} - U) |\psi_1\rangle$$

= $\frac{1}{\sqrt{2}} (\operatorname{Ctrl} - U) (|0\rangle |\psi\rangle + |\mathbf{1}\rangle |\psi\rangle)$



$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

H 门作用在基态:

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

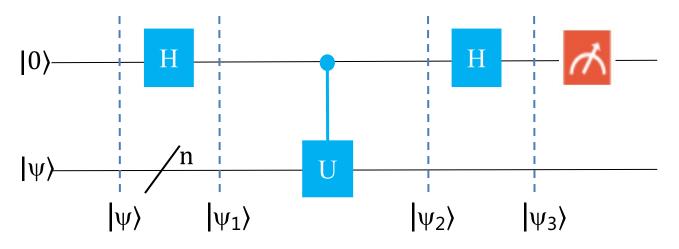
$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

因为 Ctrl -U 表示基于幺正变换 U 的受控门,只有控制位为 $|1\rangle$ 时,才会作用于 $|\psi\rangle$,所以有:

$$|\psi_2\rangle = \frac{1}{\sqrt{2}}(|0\rangle|\psi\rangle + |\mathbf{1}\rangle U|\psi\rangle)$$

Hadamard Test - 测量前的状态(实部)





$$|\psi_{3}\rangle = H|\psi_{2}\rangle$$

$$= \frac{1}{\sqrt{2}}(H|0\rangle|\psi\rangle + H|1\rangle U|\psi\rangle)$$

$$= \frac{1}{\sqrt{2}}(\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)|\psi\rangle + \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)U|\psi\rangle)$$

$$= |0\rangle \frac{|\psi\rangle + U|\psi\rangle}{2} + |1\rangle \frac{|\psi\rangle - U|\psi\rangle}{2}$$

$$= |0\rangle \frac{I+U}{2}|\psi\rangle + |1\rangle \frac{I-U}{2}|\psi\rangle$$

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

H 门作用在基态:

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$



Hadamard Test - 测量计算(实部)

$$|\psi_3\rangle = |0\rangle \frac{I+U}{2} |\psi\rangle + |1\rangle \frac{I-U}{2} |\psi\rangle$$

测量 |0>:

$$(|0\rangle\langle 0| \otimes I) |\psi_{3}\rangle = (|0\rangle\langle 0| \otimes I)(|0\rangle \frac{I+U}{2} |\psi\rangle + |1\rangle \frac{I-U}{2} |\psi\rangle)$$

$$= |0\rangle\langle 0|0\rangle \frac{I+U}{2} |\psi\rangle + |0\rangle\langle 0|1\rangle \frac{I-U}{2} |\psi\rangle$$

$$= \frac{1}{2} |0\rangle \otimes (I+U) |\psi\rangle$$

测量的概率:

$$\begin{aligned} \operatorname{Prob}(0) &= ||\frac{1}{2}|0\rangle \otimes (I+U)|\psi\rangle ||^2 = \frac{1}{4}||0\rangle||^2 ||(I+U)|\psi\rangle ||^2 \\ &= \frac{1}{4}\langle \psi|(I+U^{\dagger})(I+U)|\psi\rangle = \frac{1}{4}(\langle \psi| + \langle \psi|U^{\dagger})(|\psi\rangle + U|\psi\rangle) \\ &= \frac{1}{4}(\langle \psi|\psi\rangle + \langle \psi|U|\psi\rangle + \langle \psi|U^{\dagger}|\psi\rangle + \langle \psi|U^{\dagger}U|\psi\rangle) \\ &= \frac{1}{4}(2 + \langle \psi|U|\psi\rangle + \langle \psi|U^{\dagger}|\psi\rangle) = \frac{1}{4}(2 + \langle \psi|U|\psi\rangle + \langle \psi|U|\psi\rangle^*) \\ &= \frac{1+Re(\langle \psi|U|\psi\rangle)}{2} \end{aligned}$$

公式:
$$\langle 0|0\rangle = 1 \quad \langle 0|1\rangle = 0$$

$$(A \otimes B) \quad (C \otimes D) = (AC) \otimes (BD)$$

$$||\psi\rangle|^2 = |\psi\rangle^{\dagger}|\psi\rangle = \langle \psi|\psi\rangle$$

$$(A \otimes B)^{\dagger} = A^{\dagger} \otimes B^{\dagger}$$

$$\langle e_j|A|e_k\rangle = \langle e_k|A^{\dagger}|e_j\rangle^*$$

$$\langle u|A|v\rangle = \langle A^{\dagger}u|v\rangle = \langle v|A^{\dagger}|u\rangle^*$$

$$\langle \psi | U | \psi \rangle = \text{Re } \langle \psi | U | \psi \rangle + \text{Im } \langle \psi | U | \psi \rangle \text{ i}$$
 实部 虚部

 $Re(\langle \psi | U | \psi \rangle)$ 幺正算符 U 在量子态 ψ 上投影期望的实部



Hadamard Test – 测量计算(实部)

那么:

$$Prob(1) = 1 - Prob(0) = \frac{1 - Re(\langle \psi | U | \psi \rangle)}{2}$$

经量子线路,如果测量结果为 |0>,则让输出为1,如果测量结果为 |1>,则让输出为-1,那么期望值为:

$$E(M) = \sum_{m} mP(m)$$

$$= 1 * Prob(0) + (-1) * Prob(1)$$

$$= \frac{1 + Re(\langle \psi | U | \psi \rangle)}{2} - \frac{1 - Re(\langle \psi | U | \psi \rangle)}{2}$$

$$= Re(\langle \psi | U | \psi \rangle)$$

其为幺正算符 U 在量子态 ψ 上投影期望的实部。





$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \qquad S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

$$S = \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix}$$

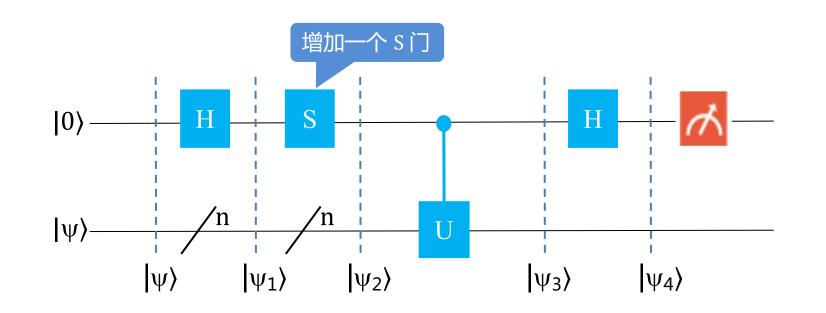
$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

所以有:

$$|\psi_1\rangle = H|0\rangle \otimes |\psi\rangle$$

$$\begin{aligned} |\psi_{2}\rangle &= \mathrm{SH}|0\rangle \otimes |\psi\rangle \\ &= \begin{bmatrix} 1 & 0 \\ 0 & i \end{bmatrix} \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 0 \end{bmatrix} \otimes |\psi\rangle \\ &= \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix} \otimes |\psi\rangle \\ &= \frac{1}{\sqrt{2}} (|0\rangle \otimes |\psi\rangle + i |1\rangle \otimes |\psi\rangle) \end{aligned}$$





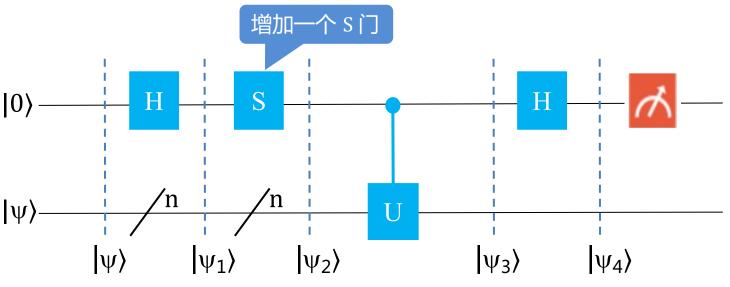
Hadamard Test - 测量前的状态(虚部)

因为:

$$|\psi_2\rangle = \frac{1}{\sqrt{2}} (|0\rangle \otimes |\psi\rangle + i |1\rangle \otimes |\psi\rangle)$$

所以有:

$$\begin{aligned} |\psi_3\rangle &= (\mathsf{Ctrl} - U) \ |\psi_2\rangle \\ &= \frac{1}{\sqrt{2}} (\mathsf{Ctrl} - U) (|0\rangle \otimes |\psi\rangle + i \ |\mathbf{1}\rangle \otimes |\psi\rangle \) \end{aligned}$$



因为 Ctrl -U 表示基于幺正变换 U 的受控门,只有控制位为 $|1\rangle$ 时,才会作用于 $|\psi\rangle$,所以有:

$$|\psi_3\rangle = \frac{1}{\sqrt{2}}(|0\rangle|\psi\rangle + i|\mathbf{1}\rangle \otimes \mathbf{U}|\psi\rangle)$$



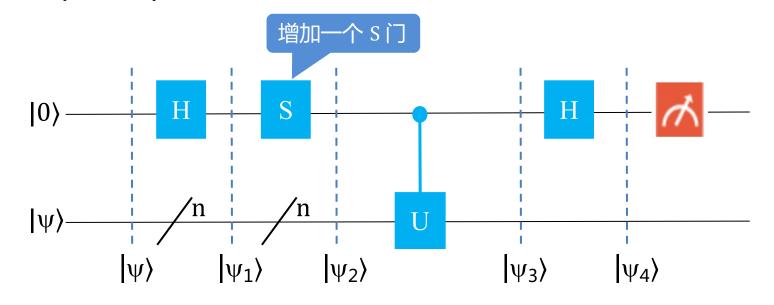
Hadamard Test - 测量前的状态(虚部)

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

H 门作用在基态:

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$



$$\begin{aligned} |\psi_{4}\rangle &= H|\psi_{3}\rangle \\ &= \frac{1}{\sqrt{2}} \left(H|0\rangle \otimes |\psi\rangle + iH|1\rangle \otimes U|\psi\rangle \right) \\ &= \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \otimes |\psi\rangle + i \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \otimes U|\psi\rangle \right) \\ &= \frac{1}{2} (|0\rangle + |1\rangle) \otimes |\psi\rangle + i (|0\rangle - |1\rangle) \otimes U|\psi\rangle \right) \end{aligned}$$



Hadamard Test - 测量计算(虚部)

$$|\psi_4\rangle = \frac{1}{2}(|0\rangle + |1\rangle) \otimes |\psi\rangle + i(|0\rangle - |1\rangle) \otimes U|\psi\rangle)$$

测量 |0>:

$$(|0\rangle\langle 0| \otimes I) |\psi_{4}\rangle = (|0\rangle\langle 0| \otimes I) (\frac{1}{2}(|0\rangle + |1\rangle) \otimes |\psi\rangle + i (|0\rangle - |1\rangle) \otimes U|\psi\rangle))$$

$$= \frac{1}{2}(|0\rangle\langle 0|0\rangle + |0\rangle\langle 0|1\rangle) \otimes |\psi\rangle + i (|0\rangle\langle 0|0\rangle - |0\rangle\langle 0|1\rangle) \otimes U|\psi\rangle)$$

$$= \frac{1}{2}(|0\rangle \otimes |\psi\rangle + i|0\rangle \otimes U|\psi\rangle)$$

测量的概率:

$$\begin{aligned} & \operatorname{Prob}(0) = ||\frac{1}{2}(|0\rangle \otimes |\psi\rangle + i|0\rangle \otimes U|\psi\rangle) \,||^2 \\ & = \frac{1}{4}(|0\rangle \otimes |\psi\rangle + i|0\rangle \otimes U|\psi\rangle)^{\dagger}(|0\rangle \otimes |\psi\rangle + i|0\rangle \otimes U|\psi\rangle) \\ & = \frac{1}{4}(\langle 0| \otimes \langle \psi| - i \langle 0| \otimes \langle \psi|U^{\dagger}) (|0\rangle \otimes |\psi\rangle + i|0\rangle \otimes U|\psi\rangle) \\ & = \frac{1}{4}(\langle 0| \otimes \langle \psi|)(|0\rangle \otimes |\psi\rangle) + (\langle 0| \otimes \langle \psi|)(i|0\rangle \otimes U|\psi\rangle) + (-i \langle 0| \otimes \langle \psi|U^{\dagger})(|0\rangle \otimes |\psi\rangle) + (-i \langle 0| \otimes \langle \psi|U^{\dagger})(i|0\rangle \otimes U|\psi\rangle)) \\ & = \frac{1}{4}(\langle 0|0\rangle \otimes \langle \psi|\psi\rangle + i \langle 0|0\rangle \otimes \langle \psi|U|\psi\rangle - i \langle 0|0\rangle \otimes \langle \psi|U^{\dagger}|\psi\rangle + \langle 0|0\rangle \otimes \langle \psi|U^{\dagger}U|\psi\rangle) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U^{\dagger}|\psi\rangle + 1) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U^{\dagger}|\psi\rangle + 1) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U^{\dagger}|\psi\rangle + 1) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U^{\dagger}|\psi\rangle + 1) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U^{\dagger}|\psi\rangle + 1) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle - i \langle \psi|U|\psi\rangle^*) \\ & = \frac{1}{4}(\langle 0+i \langle \psi|U|\psi\rangle) \\ & = \frac{$$

公式:

- $\langle 0|0\rangle = 1 \ \langle 0|1\rangle = 0$
- $(A \otimes B) (C \otimes D) = (AC) \otimes (BD)$
- $| |\psi\rangle |^2 = |\psi\rangle^{\dagger} |\psi\rangle = \langle \psi |\psi\rangle$
- $(A \otimes B)^{\dagger} = A^{\dagger} \otimes B^{\dagger}$
- $(cA)^{\dagger} = c^*A^{\dagger}$
- $(A + B)^{\dagger} = A^{\dagger} + B^{\dagger}$
- $(AB)^{\dagger} = B^{\dagger}A^{\dagger}$
- $\langle e_i | A | e_k \rangle = \langle e_k | A^{\dagger} | e_i \rangle^*$
- $\langle u|A|v\rangle = \langle A^{\dagger}u|v\rangle = \langle v|A^{\dagger}|u\rangle^*$

 $Im(\langle \psi | U | \psi \rangle)$ 幺正算符 U 在量子态 ψ 上投影期望的虚部

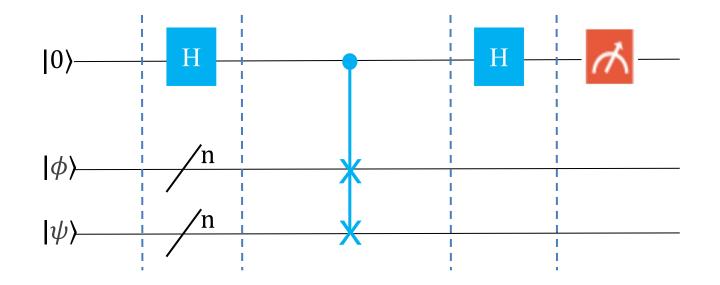
 $\langle \psi | U | \psi \rangle = a + b i$ 则: $b = Im(\langle \psi | U | \psi \rangle)$ $\langle \psi | U | \psi \rangle^* = a - b i$ $(\langle \psi | U | \psi \rangle - \langle \psi | U | \psi \rangle^*) i = 2b i * i = -2b$



SWAP Test

Hadamard Test有着多种形式和广泛用途,其中一种特殊形式是基本量子线路SWAP Test。任给两个维数相同的量子态,通过SWAP Test线路,可以得到两个量子态的保真度,反应了它们的重叠情况。两个量子态 $|\phi\rangle$, $|\psi\rangle$ 的保真度是指量子态内积范数的平方 $|\langle\phi|\psi\rangle|^2$

SWAP Test的量子线路图结构:



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SWAP Test - 测量前的状态



$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

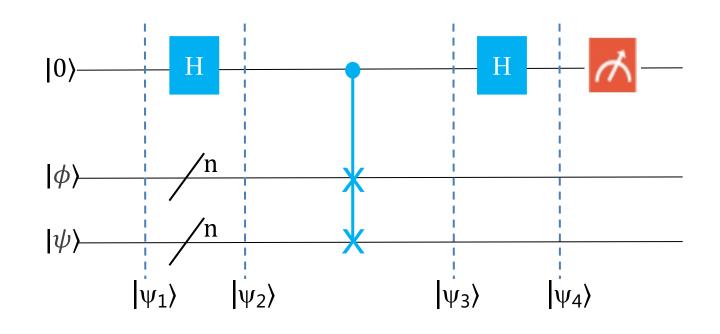
$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

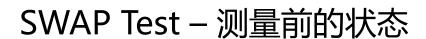
$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

1. 输入态 $|\psi_1\rangle = |0\rangle \otimes |\phi\rangle \otimes |\psi\rangle$

2. 输入态经过第一个 H 门
$$|\psi_{2}\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) |\phi\rangle |\psi\rangle = \frac{1}{\sqrt{2}}(|0\rangle |\phi\rangle |\psi\rangle + |1\rangle |\phi\rangle |\psi\rangle)$$

3. 再经过一个 swap 门 ($|1\rangle$ 时交换) $|\psi_3\rangle = \frac{1}{\sqrt{2}}(|0\rangle|\phi\rangle|\psi\rangle + |1\rangle|\psi\rangle|\phi\rangle)$







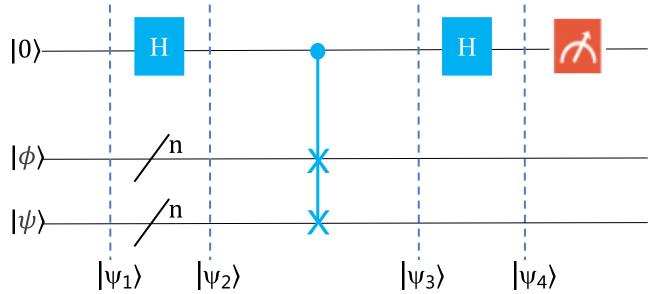
$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

$$H|0\rangle = \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle)$$

$$H|1\rangle = \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle)$$

4. 再经过最后一个 H 门

$$\begin{split} |\psi_{4}\rangle &= H|\psi_{3}\rangle \\ &= \frac{1}{\sqrt{2}}(H|0\rangle|\phi\rangle|\psi\rangle + H|1\rangle|\psi\rangle|\phi\rangle) \\ &= \frac{1}{\sqrt{2}}(\frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)|\phi\rangle|\psi\rangle + \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)|\psi\rangle|\phi\rangle) \\ &= \frac{1}{2}[|0\rangle(|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle) + |1\rangle(|\phi\rangle|\psi\rangle - |\psi\rangle|\phi\rangle)] \end{split}$$





SWAP Test - 测量

$$|\psi_{4}\rangle = \frac{1}{\sqrt{2}} \left(\frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) |\phi\rangle |\psi\rangle + \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) |\psi\rangle |\phi\rangle \right)$$
$$= \frac{1}{2} [|0\rangle (|\phi\rangle |\psi\rangle + |\psi\rangle |\phi\rangle) + |1\rangle (|\phi\rangle |\psi\rangle - |\psi\rangle |\phi\rangle)]$$

测量 |0>:

$$(|0\rangle\langle 0| \otimes I) |\psi_{4}\rangle = (|0\rangle\langle 0| \otimes I) \frac{1}{2}[|0\rangle\langle (|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle) + |1\rangle\langle (|\phi\rangle|\psi\rangle - |\psi\rangle|\phi\rangle)]$$

$$= \frac{1}{2}[|0\rangle\langle 0|0\rangle\langle (|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle) + |0\rangle\langle 0|1\rangle\langle (|\phi\rangle|\psi\rangle - |\psi\rangle|\phi\rangle)]$$

$$= \frac{1}{2}[|0\rangle\langle 0|\psi\rangle + |\psi\rangle|\phi\rangle)]$$

那么测量的概率:

$$\begin{aligned} \operatorname{Prob}(0) &= ||\frac{1}{2}[|0\rangle \otimes (|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle)] ||^2 = \frac{1}{4}||0\rangle||^2 ||(|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle) ||^2 \\ &= \frac{1}{4}(|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle)^{\dagger} (|\phi\rangle|\psi\rangle + |\psi\rangle|\phi\rangle) \\ &= \frac{1}{4}(\langle\psi|\langle\phi| + \langle\phi|\langle\psi| \rangle|\psi\rangle + |\psi\rangle|\phi\rangle) \\ &= \frac{1}{4}(\langle\psi|\langle\phi|\phi\rangle|\psi\rangle + \langle\phi|\langle\psi|\phi\rangle|\psi\rangle + \langle\phi|\langle\psi|\phi\rangle|\psi\rangle + \langle\phi|\langle\psi|\psi\rangle|\phi\rangle) \\ &= \frac{1}{4}(1 + \langle\phi|\langle\psi|\phi\rangle|\psi\rangle + \langle\phi|\langle\psi|\phi\rangle|\psi\rangle + 1) \\ &= \frac{1}{2}(1 + ||\langle\psi|\phi\rangle||^2) \end{aligned}$$

公式:

- $\langle 0|1\rangle = 1$ $\langle 0|1\rangle = 0$
- $(A \otimes B) (C \otimes D) = (AC) \otimes (BD)$
- $| | \psi \rangle |^2 = | \psi \rangle^{\dagger} | \psi \rangle = \langle \psi | \psi \rangle$
- $(A \otimes B)^{\dagger} = A^{\dagger} \otimes B^{\dagger}$
- $(cA)^{\dagger} = c^*A^{\dagger}$
- $(A + B)^{\dagger} = A^{\dagger} + B^{\dagger}$
- $(AB)^{\dagger} = B^{\dagger}A^{\dagger}$
- $\langle e_i | A | e_k \rangle = \langle e_k | A^{\dagger} | e_i \rangle^*$
- $\langle u|A|v\rangle = \langle A^{\dagger}u|v\rangle = \langle v|A^{\dagger}|u\rangle^*$





对SWAP Test的公式推导验证过程完全类似于Hadamard Test,结果量子态的第一个寄存器测量得到 $|0\rangle$, $|1\rangle$ 的概率均与给定的两个量子态的保真度相关。也就是可以多次测量,判断两个量子态 $|\phi\rangle$, $|\psi\rangle$ 具体区别有多大。

$$P_0 = \frac{1 + |\langle \psi | \phi \rangle|^2}{2}, P_1 = 1 - P_0$$

SWAP Test作为Hadamard的一种特殊形式,它对两个给定量子态给出了其保真度相关的测量结果,具有重要应用意义。在量子态的内积相关研究中有着重要作用。

