EFT Wilson Coefficent Conversions

Lukas Lechner, Robert Schöfbeck

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1 Conversions in dim6top

1.1 C_{tB} to C_{tZ} conversion

Basis transformation of the Wilson coefficient C_{tB} to C_{tZ} using the Weinberg angle θ_W and $\sin \theta_W \approx 0.471$

$$C_{tZ} = \operatorname{Re}\left(-\sin\theta_W C_{tB} + \cos\theta_W C_{tW}\right),\tag{1}$$

where only one Wilson coefficient at a time is used, thus for the conversion of C_{tB} to C_{tZ} , the coefficient C_{tW} is set to 0.

$$C_{\rm tZ} = \operatorname{Re}\left(-\sin\theta_W C_{\rm tB}\right) \approx -0.471 C_{\rm tB} \tag{2}$$

C_{tB} limits translate to C_{tZ} limits with a factor of \approx -0.471

1.2 $C_{\phi Q}$ conversion

In the FeynRules model dim6top, a linear combination of Wilson coefficients is used to parametrize the Z boson coupling. The coefficient cpQM is given as

$$C_{\phi Q}^{-} = C_{\phi q}^{1(33)} - C_{\phi q}^{3(33)}, \tag{3}$$

where additionally cpQ3 is defined as $C_{\phi q}^{3(33)}$. If cpQM is the only Wilson coefficient set to non-zero, its definition reduces to

$$C_{\phi Q}^{-} = C_{\phi q}^{1(33)} \tag{4}$$

which removes effects of cpQM on the tWb vertex, but adds effects on the bbZ coupling. This can be seen in the operator definitions of the processes

$$O_{\phi a}^{1(ij)} = (\phi^{\dagger} i \overleftrightarrow{D}_{\mu} \phi)(\bar{q}_i \gamma^{\mu} q_j) \tag{5}$$

$$O_{\phi q}^{3(ij)} = (\phi^{\dagger} i \overleftrightarrow{D_{\mu}}^{I} \phi) (\bar{q}_{i} \gamma^{\mu} \tau^{I} q_{j})$$

$$(6)$$

where τ^I are the Pauli matrices and thus, $O_{\phi q}^{3(ij)}$ introduces anomalous couplings to the W boson.

2 Conversion of ATLAS Wilson coefficients

ATLAS presented ttZ limits on EFT in the talk of Ref. [2]. However different definitions of the operators \mathcal{O} are used in ATLAS [3] compared to CMS [4], leading to a scale factor for the Wilson coefficients.

2.1 C_{tB} conversion: ATLAS vs. CMS

To convert the definitions, the Lagrangians have to be equal.

$$C_{\rm tB}^{\rm ATLAS} \mathcal{O}_{\rm tB}^{\rm ATLAS} \stackrel{!}{=} C_{\rm tB}^{\rm CMS} \mathcal{O}_{\rm tB}^{\rm CMS}$$
 (7)

$$\mathcal{O}_{\rm tB}^{\rm CMS} = (\bar{Q}\sigma^{\mu\nu}t)\,\tilde{\phi}B_{\mu\nu} \tag{8}$$

$$\mathcal{O}_{\rm tB}^{\rm ATLAS} = y_t \, g_Y \, \left(\bar{Q} \sigma^{\mu\nu} t \right) \, \tilde{\phi} B_{\mu\nu} = y_t \, g_Y \, \mathcal{O}_{\rm tB}^{\rm CMS}. \tag{9}$$

Using Eq. 7 and Eq. 9 leads to

$$C_{\rm tB}^{\rm CMS} = y_t g_Y C_{\rm tB}^{\rm ATLAS} \tag{10}$$

where $y_t = \frac{\sqrt{2}m_t}{v}$. Using the definitions from Ref. [3]

$$m_{t} = 173.3 \text{ GeV}$$

$$\alpha_{\text{EW}} = \frac{1}{127.9}$$

$$v = 246 \text{ GeV}$$

$$\tan \theta_{W} = \frac{g_{1}}{g_{2}} = 0.535$$

$$\alpha = \frac{1}{4\pi} \frac{(g_{1}g_{2})^{2}}{g_{1}^{2} + g_{2}^{2}} = \frac{g_{1}^{2}}{4\pi(1 + \tan^{2}\theta_{W})} = \frac{g_{1}^{2}\cos^{2}\theta_{W}}{4\pi} = \frac{1}{127.9}$$

$$g' = g_{1} = g_{Y} = \sqrt{4\pi\alpha(1 + \tan^{2}\theta_{W})} = \frac{\sqrt{4\pi\alpha}}{\cos\theta_{W}} \approx 0.34$$

$$g = g_{2} = g_{w} = \frac{g_{1}}{\tan\theta_{W}} \approx 0.636$$

$$y_{t} = \frac{\sqrt{2}m_{t}}{v} \approx 0.996$$

leads to a conversion factor for C_{tB} from ATLAS to CMS of

$$C_{\rm tB}^{\rm CMS} = y_t \, g_Y C_{\rm tB}^{\rm ATLAS} \approx 0.34 \, C_{\rm tB}^{\rm ATLAS} \tag{11}$$

The ATLAS C_{tB} limits translates to limits used in CMS with a factor of ≈ 0.34

2.1.1 Total conversion from C_{tB}^{ATLAS} to C_{tZ}^{CMS} limits

To convert from C_{tB}^{ATLAS} to C_{tZ}^{CMS} , an additional conversion factor for C_{tB} to C_{tZ} from Sec. 1.1 has to be applied.

$$C_{\rm tZ}^{\rm CMS} = -\sin\theta_W \, y_t \, g_Y C_{\rm tB}^{\rm ATLAS} \approx -0.16 \, C_{\rm tB}^{\rm ATLAS} \tag{12}$$

The ATLAS C_{tB} limits translates to CMS C_{tZ} limits with a factor of \approx -0.16

2.1.2 Conversion of ATLAS ttZ C_{tB}^{ATLAS} limits

The presented ttZ ATLAS C_{tB}^{ATLAS} limits of Ref. [2] convert to CMS C_{tZ}^{CMS} limits as

$$C_{\rm tZ}^{\rm ATLAS \to CMS}(36.1 \text{ fb}^{-1}, 95\%, \text{x-sec only}) : [-2.40, +2.40]$$
 (13)

2.2 C_{tW} conversion: ATLAS vs. CMS

To convert the definitions, the Lagrangians have to be equal.

$$C_{\text{tW}}^{\text{ATLAS}} \mathcal{O}_{\text{tW}}^{\text{ATLAS}} \stackrel{!}{=} C_{\text{tW}}^{\text{CMS}} \mathcal{O}_{\text{tW}}^{\text{CMS}}$$
(14)

$$\mathcal{O}_{tW}^{CMS} = (\bar{q}_i \sigma^{\mu\nu} \tau^I u_i) \tilde{\phi} W_{\mu\nu}^I \tag{15}$$

$$\mathcal{O}_{\text{tW}}^{\text{ATLAS}} = y_t \, g_w \, (\bar{q}_i \sigma^{\mu\nu} \tau^I u_j) \tilde{\phi} W_{\mu\nu}^I = y_t \, g_w \, \mathcal{O}_{\text{tW}}^{\text{CMS}}$$
(16)

Using Eq. 14 and Eq. 16 leads to

$$C_{\rm tW}^{\rm CMS} = y_t g_w C_{\rm tW}^{\rm ATLAS} \approx 0.633 C_{\rm tW}^{\rm ATLAS}$$
(17)

where $y_t = \frac{\sqrt{2}m_t}{v} \approx 0.996$ and $g_w \approx 0.636$ (see Sec. 2.1).

The ATLAS C_{tW} limits translates to limits used in CMS with a factor of ≈ 0.633

2.3 $C_{\phi t}$ and $C_{\phi Q}$ conversion: ATLAS vs. CMS

To convert the definitions, the Lagrangians have to be equal.

$$C_{\phi t}^{\text{ATLAS}} \mathcal{O}_{\phi t}^{\text{ATLAS}} \stackrel{!}{=} C_{\phi t}^{\text{CMS}} \mathcal{O}_{\phi t}^{\text{CMS}}$$

$$\tag{18}$$

$$\mathcal{O}_{\phi t}^{\text{CMS}} = (\phi^{\dagger} i \overrightarrow{D}_{\mu} \phi)(\bar{t} \gamma^{\mu} t) \tag{19}$$

$$\mathcal{O}_{\phi t}^{\text{ATLAS}} = \frac{1}{2} y_t^2 \left(\phi^{\dagger} i \overleftrightarrow{D_{\mu}} \phi \right) (\bar{t} \gamma^{\mu} t) = \frac{1}{2} y_t^2 \, \mathcal{O}_{\phi t}^{\text{CMS}}$$
 (20)

However, the factor $\frac{1}{2}$ cancels, as the definition of ATLAS [3] includes the hermitian conjugate state in the Lagrangian, where the CMS definitions [4] uses the fact that these operators are hermitian. Thus, using Eq. 18 and Eq. 20 leads to

$$C_{\phi t}^{\rm CMS} = y_t^2 C_{\phi t}^{\rm ATLAS} \approx 0.992 C_{\phi t}^{\rm ATLAS}$$
(21)

where $y_t = \frac{\sqrt{2}m_t}{v} \approx 0.996$ (see Sec. 2.1). The same pre-factor is applied to $C_{\phi Q}^{(1)}$ and $C_{\phi Q}^{(3)}$, thus $C_{\phi Q}$.

The ATLAS $C_{\phi t}$, $C_{\phi Q}^{(1)}$, $C_{\phi Q}^{(3)}$ and $C_{\phi Q}$ limits translate to limits used in CMS with a factor of 0.992.

3 Conversion of HEL UFO Wilson coefficients

Comparison of the Wilson coefficients using couplings of Lagrangians from the couplings.py files in the FeynRules directories for the dim6top- and HEL-UFO models. Here, only an example of one coupling for one Lagrangian is used to show the conversion factor, where the Lagrangians are labeled similar to the definitions in the couplings.py file (GC #).

3.1 C_{tB} to C_{tZ} conversion

$$g(\mathcal{L}_{\text{dim6}}, \text{GC 664}) = \frac{C_{\text{tZ}}}{\Lambda^2 \sqrt{2}}$$
(22)

$$g(\mathcal{L}_{\text{HEL}}, \text{GC 2191}) = \frac{\sqrt{4\pi\alpha_{\text{EW}}} y_t}{M_W^2 \sqrt{2}} \left(C_{\text{uB}} \frac{\sin \theta_W}{\cos \theta_W} - C_{\text{uW}} \frac{\cos \theta_W}{\sin \theta_W} \right)$$
(23)

$$g(\mathcal{L}_{\text{dim6}}) \stackrel{!}{=} g(\mathcal{L}_{\text{HEL}})$$
 (24)

but only one Wilson coefficient at a time is used, thus $C_{uW} = 0$, leading to a direct comparison of C_{tZ} and C_{uB} (see Sec. 1.1):

$$\frac{C_{\rm tZ}}{\Lambda^2} = -C_{\rm uB} \frac{\sqrt{4\pi\alpha_{\rm EW}} \, y_t}{M_W^2} \frac{\sin\theta_W}{\cos\theta_W} \tag{25}$$

A second example is the comparison of the Lagrangians $g(\mathcal{L}_{\text{dim6}}, \text{GC 948})$ and $g(\mathcal{L}_{\text{HEL}}, \text{GC 2187})$, leading to the same results. Even though it may look reasonable to compare the Lagrangian in switched order (DIM6: GC 664 \leftrightarrow HEL: GC 2187, DIM6: GC 948 \leftrightarrow HEL: GC 2187), the results are then not consistant.

3.1.1 Comparison of Limits in Top 17-005

In Ref. [1], a factor k is applied to the presented limits, where k is taken from Ref. [5].

$$k = \frac{\sqrt{4\pi\alpha_{\rm EW}}\,y_t}{2\,M_W^2\,\cos\theta_W}\tag{26}$$

$$\bar{C}_{\text{uB}} = C_{\text{uB}} k = C_{\text{uB}} \frac{\sqrt{4\pi\alpha_{\text{EW}}} y_t}{2M_W^2 \cos\theta_W}$$
(27)

To compare the limits to current efforts using the dim6top model, the conversion factor is thus

$$\rightarrow \frac{C_{\rm tZ}}{\Lambda^2} = -2\sin\theta_W \,\bar{C}_{\rm uB} = -0.96 \,\bar{C}_{\rm uB} \tag{28}$$

The HEL UFO C_{uB} limits in Ref. [1] translate to limits of C_{tZ} with a factor of -0.96

3.1.2 Conversion of Top 17-005 \bar{C}_{uB} limits

Thus, the presented ttZ C_{uB} limits of Ref. [1] convert to dim6top limits as

$$C_{tZ}^{Top\ 17-005\to\ dim6top}(36.1\ fb^{-1}, 95\%, x\text{-sec only}): [-2.02, +2.02]$$
 (29)

3.2 C_{Hu} to $C_{\phi t}$ conversion

$$g(\mathcal{L}_{\text{dim6}}, \text{GC 1112}) = -i \frac{C_{\phi t} \cos \theta_W \sqrt{4 \alpha_{\text{EW}} \pi} v}{\Lambda^2 \sin \theta_W} - i \frac{C_{\phi t} \sqrt{4 \alpha_{\text{EW}} \pi} \sin \theta_W v}{\cos \theta_W \Lambda^2}$$

$$g(\mathcal{L}_{\text{HEL}}, \text{GC 1136}) = -i \frac{C_{\text{Hu}} \cos \theta_W \sqrt{4 \alpha_{\text{EW}} \pi}}{\sin \theta_W v} - i \frac{C_{\text{Hu}} \sqrt{4 \alpha_{\text{EW}} \pi} \sin \theta_W}{\cos \theta_W v}$$
(30)

$$g(\mathcal{L}_{\text{HEL}}, \text{GC 1136}) = -i \frac{C_{\text{Hu}} \cos \theta_W \sqrt{4 \alpha_{\text{EW}} \pi}}{\sin \theta_W v} - i \frac{C_{\text{Hu}} \sqrt{4 \alpha_{\text{EW}} \pi} \sin \theta_W}{\cos \theta_W v}$$
(31)

$$g(\mathcal{L}_{\text{dim6}}) \stackrel{!}{=} g(\mathcal{L}_{\text{HEL}})$$
 (32)

3.2.1Comparison of Limits in Top 17-005

In Ref. [1], a factor k is applied to the presented limits, where k is taken from Ref. [5].

$$k = \frac{1}{2v^2} \tag{34}$$

$$\bar{C}_{\rm Hu} = C_{\rm Hu} \, k = \frac{C_{\rm Hu}}{2 \, v^2}$$
 (35)

To compare the limits to current efforts using the dim6top model, the conversion factor is thus

$$\rightarrow \frac{C_{\phi t}}{\Lambda^2} = 2\,\bar{C}_{Hu} \tag{36}$$

The HEL UFO \bar{C}_{Hu} limits in Ref. [1] translate to limits of $C_{\phi t}$ with a factor of 2.0

Conversion of Top 17-005 \bar{C}_{Hu} limits

Thus, the presented ttZ \bar{C}_{Hu} limits of Ref. [1] convert to dim6top limits as

$$C_{\phi t}^{\text{Top 17-005} \to \text{dim6top}}(36.1 \text{ fb}^{-1}, 95\%, \text{x-sec only}) : [-20.2, +4.0]$$
 (37)

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