

# Discrete Elastic Rod Simulation of Shoelace Knots and Strength Progress Report

## *Mechanics of Flexible Structures and Soft Robots - MAE 263F*

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**Abstract**—In the field of flexible structure and soft robotics, analysing an elastic rod is an important fundamental subject with future potential and practical application. Recently, many innovative approaches have continued to emerge, such as [10], [6], and [1], enabling effective numerical interpretation and advanced simulation.

Among the various topics, knot tying is the most representative target for analysis, which can serve as a foundational structure with the potential to extend to various forms of soft robotics. In this project, we aim to advance our analysis beyond basic knot structures by conducting a comprehensive simulation study focusing on shoelace knots. Leveraging the Discrete Elastic Rod (D.E.R.) algorithm, we will simulate a shoelace knot structure with varying properties to compare knot strengths. We assume dynamic equilibrium under constraints for the implementation the Implicit Model Contact (IMC) model. Through this analyses and simulation, we hope to not only acquire knowledge in the simulation of flexible structures, but also a deeper understanding of the the affects of model parameters and their impacts on both simulation accuracy and speed.

### I. INTRODUCTION

Our group project goal is to simulate the tying and untying of a shoelace knot, and to study the impact of friction and elastic parameters on the strength of the knot. Our work will be based on the Discrete Elastic Rod algorithm and on the Implicit Contact Model, which are detailed below. We will implement our algorithm in C++ in order to take advantage of its computational efficiency.

The two following sections about our general simulation method and about the simulation of contact are put in this document as a reminder from our proposal report. Then we explained the assumptions we chose and our current progress. Finally, our plan for the following weeks is presented.

### II. DISCRETE ELASTIC ROD ALGORITHM

The Discrete Elastic Rod (D.E.R.) algorithm [3] is useful in 3D simulation of elastic rod. The rod discretized and made to be at dynamic equilibrium under applied constraints. My changing the constraints over time (i.e. end positions of the "shoelace"), we will be able to manipulate the rod into a tied and then subsequently untied shoelace knot. The D.E.R. algorithm is extremely useful to us because it models the rod with all elastic and external forces at equilibrium.

General implementation of the D.E.R. Algorithm can be followed in Figure 1 [8]. The first appearance of the

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#### Algorithm 1 Discrete Elastic Rods

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Require:  $\mathbf{q}(t_j), \dot{\mathbf{q}}(t_j)$  ▷ DOFs and velocities at  $t = t_j$ 
Require:  $(\mathbf{a}_1^k(t_j), \mathbf{a}_2^k(t_j), \mathbf{t}^k(t_j))$  ▷ Reference frame at  $t = t_j$ 
Require:  $\text{free\_index}$  ▷ Index of the free DOFs
Ensure:  $\mathbf{q}(t_{j+1}), \dot{\mathbf{q}}(t_{j+1})$  ▷ DOFs and velocities at  $t = t_{j+1}$ 
Ensure:  $(\mathbf{a}_1^k(t_{j+1}), \mathbf{a}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1}))$  ▷ Reference frame at  $t = t_{j+1}$ 

1: function DISCRETE_ELASTIC_RODS( $\mathbf{q}, \dot{\mathbf{q}}(t_j), (\mathbf{a}_1^k(t_j), \mathbf{a}_2^k(t_j), \mathbf{t}^k(t_j))$ )
2:   Guess:  $\mathbf{q}^{(1)}(t_{j+1}) \leftarrow \mathbf{q}(t_j)$ 
3:    $n \leftarrow 1$ 
4:   while error > tolerance do
5:     Compute reference frame  $(\mathbf{a}_1^k(t_{j+1}), \mathbf{a}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1}))^{(n)}$  using  $\mathbf{q}^{(n)}(t_{j+1})$ 
6:     Compute reference twist  $\Delta m_{k,\text{ref}}^{(n)}$  ( $k = 2, \dots, N-1$ )
7:     Compute material frame  $(\mathbf{m}_1^k(t_{j+1}), \mathbf{m}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1}))^{(n)}$ 
8:     Compute  $\mathbf{f}$  and  $\mathbb{J}$  ▷ Eqs. 7.1 and 7.2; see Algorithm 2
9:      $\mathbf{f}_{\text{free}} \leftarrow \mathbf{f}(\text{free\_index})$ 
10:     $\mathbb{J}_{\text{free}} \leftarrow \mathbb{J}(\text{free\_index}, \text{free\_index})$ 
11:     $\Delta \mathbf{q}_{\text{free}} \leftarrow \mathbb{J}_{\text{free}} \backslash \mathbf{f}_{\text{free}}$ 
12:     $\mathbf{q}^{(n+1)}(\text{free\_index}) \leftarrow \mathbf{q}^{(n)}(\text{free\_index}) - \Delta \mathbf{q}_{\text{free}}$  ▷ Update free DOFs
13:    error  $\leftarrow \text{sum}(\text{abs}(\mathbf{f}_{\text{free}}))$ 
14:     $n \leftarrow n + 1$ 
15:   end while

16:    $\mathbf{q}(t_{j+1}) \leftarrow \mathbf{q}^{(n)}(t_{j+1})$ 
17:    $\dot{\mathbf{q}}(t_{j+1}) \leftarrow \frac{\mathbf{q}(t_{j+1}) - \mathbf{q}(t_j)}{\Delta t}$ 
18:    $(\mathbf{a}_1^k(t_{j+1}), \mathbf{a}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1})) \leftarrow (\mathbf{a}_1^k(t_{j+1}), \mathbf{a}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1}))^{(n)}$ 
19:   return  $\mathbf{q}(t_{j+1}), \dot{\mathbf{q}}(t_{j+1}), (\mathbf{a}_1^k(t_{j+1}), \mathbf{a}_2^k(t_{j+1}), \mathbf{t}^k(t_{j+1}))$ 
20: end function

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Fig. 1. D.E.R. Algorithm Pseudo-code, from [8]

D.E.R. algorithm is in [2]. Since, it has been used in several research projects, especially to study knot tying [7]. The reason for that is the strength of this model on 1-D bodies in large deformation.

Kirchhoff-Love plate theory states that thin elastic plates undergo linear deformation with negligible transverse shear deformations when subjected to mechanical loads. When adapted to model rods by considering them as effectively 1-D structures, we can predict the deformation and stress distribution along the rod's length, allowing for the accurate analysis of slender structures under various loading conditions. Key elements of Kirchhoff-Love plate theory include linear elastic behavior, linear strain-displacement relations, and the assumption that shear deformations are negligible.

Since the publication of [2], which relied principally on a space-parallel transport frames to compute elastic rod twist, it was discovered that significant computational time advantages could be had if time-parallel transport was used. Hence this is the implementation that we will be using.

### III. IMPLICIT CONTACT MODEL

The D.E.R. algorithm is constructed to handle the elastic forces of the body. But in knot tying, another physical phenomenon is very important: friction. That is why a physical friction model was necessary to conduct a numerical study of knots. If a lot of numerical friction models exist, we chose the IMC model, for its physical relevancy and its computational efficiency.

The IMC model is presented in [11] and [5]. The main point of this method is to compute implicitly a forces equilibrium, including contact forces, by choosing a smoothed (and so differentiable) formulation of the contact potential. The contact potential is calculated based on an analytical segment-segment distance between the discrete edges of the rod, instead of on the distance between nodes, to obtain more realistic results. This differentiable formulation of the contact energy allows to compute its gradient and hessian, and therefore to include it in the implicit formulation. The contact potential in [11] is defined by:

$$\begin{aligned} E(\Delta, \delta) &= (2h - \Delta)^2 \quad \text{if } \Delta \in [0, 2h - \delta] \\ E(\Delta, \delta) &= \left( \frac{1}{K_1} \log(1 + e^{K_1(2h - \Delta)}) \right)^2 \quad \text{if } \Delta \in [2h - \delta, 2h + \delta] \\ E(\Delta, \delta) &= 0 \quad \text{if } \Delta \geq 2h + \delta \end{aligned}$$

In this formula,  $K_1$  is the stiffness of the energy curve,  $h$  is the radius of the rod,  $\Delta$  is the distance between the two edges in contact, and  $\delta$  is a tolerance value on the distance. Using this expression, the curves shown in figure 2 are obtained.

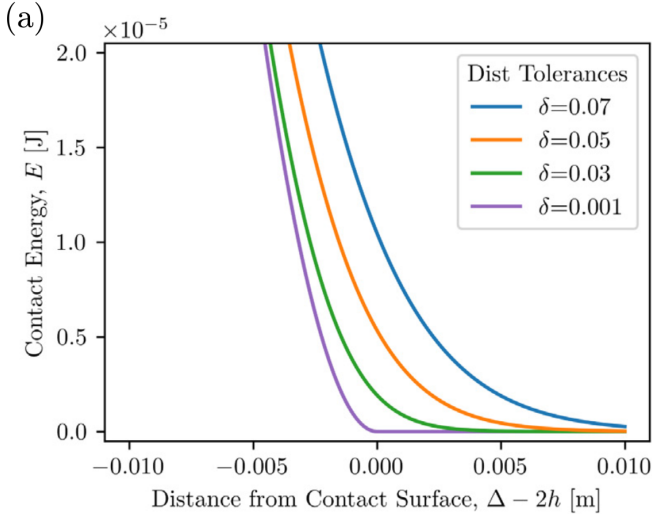


Fig. 2. Contact energy curve, this figure is extracted from [11]

The Implicit Model Contact was used in [5] to study the three-foil knot-tying phenomena. To do so, the IMC was implemented in a D.E.R. algorithm. The simulation gave very good results when compared to theoretical results or to another accurate model for contact in knot tying, the Spillmann and Teschner model (SPT). By comparison, the

Parameters	
Rod Radius	1.6[mm]
Young's Modulus, $E$	0.18 [MPa]
Poisson's Ratio	0.5
Density, $\rho$	10000 kg/m <sup>3</sup>
Time Step, $dt$	0.1 [s]
Tolerance	0.0001
Total simulation time	40 [minute]

TABLE I

PARAMETERS USED IN THE SIMULATION. NOTE LACK OF GRAVITY & FLUID FRICTION

IMC method gives smoother pull forces than SPT.

### IV. ASSUMPTIONS AND IMPLEMENTATION

In this section, we represent the specific setting of the simulation and implementation. The overall simulation flow follows the D.E.R. algorithm, the core settings are similar, and the setup values used for our simulation are shown in Table I. Here, we assume that no external force other than gravity and the force we set up is supposed to affect the system. Thus, the system modeling is assumed to be unaffected by other external forces.

We have implemented the D.E.R algorithm in C++ in order to improve speed and thus achievable quality of simulation. C++, do to its optimization advantages over MATLAB, can produce faster and more sophisticated results. Using Newton's method to simulate the system implicitly, we were able to achieve a high degree of accuracy while allowing larger time steps when compared to an explicit method.

We implemented our simulation based on DisMech [12, 4], a fully implicit penalty-based contact method for 3D elastic rod simulations. we incorporated IMC to simulate both contact and fabrication using the D.E.R. framework [3]. For external factors, we applied boundary conditions in the form of displacement instead of force. This is to ensure that the simulation can be controlled and carried out in accordance with our intentions. The approach is for the movement when a person holds onto a string and moves, where displacement restricts the structure. These conditions are sequentially applied as the structure of a shoelace changes into a certain form.

### V. CURRENT PROGRESS

Currently, we have formulated and applied the changing boundary conditions (which are a function of time) that are necessary "tie the knot". Due to the complex nature of these boundary conditions and the speed at which the simulation runs, it is a painstaking process to notice and fix small errors in boundary conditions.

Once the reef knot is successfully tied, we will release many of the fixed DOFs so that the "tying" action can be executed and simulated correctly. From here we can get knot

strength data and modify parameters so see their effects on knot strength.

A video of the current state of our system is available at the following link: <https://drive.google.com/file/d/1l6ar9xEqOrvwlKfks2ielmdALB3L6d2H/view?usp=sharing>.

## VI. FUTURE PROGRESS

After successfully getting data for the reef knot, we plan to implement the granny knot and compare the results between the two, and ultimately develop a better understanding of both knot types and which is stronger, which is more affected by parameter modification, etc. We will then be able to study the influence of the knot parameters and especially the influence of friction and elastic parameters. Each knot is indeed characterised by topological parameters, as shown in [9]. But our goal is to link this knot characterisation to a physical measurement of the knot stability and of the impact of friction and elastic parameters.

## VII. CONCLUSIONS

In this project, we will study the behavior of elastic rods using a series of simulations, with a specific focus on knot tying. Our project employs advanced simulations that closely refer to real-world scenarios, particularly with respect to the structure of shoelace knots. We leverage the Discrete Elastic Rod (D.E.R.) algorithm and the Implicit Model Contact (IMC) model to ensure a comprehensive analysis, all while considering the dynamic equilibrium under constraints. This endeavor will not only provide us with valuable practical experience in the implementation of soft robotics technology but also enable us to deepen our comprehension of the numerous factors that impact the results.

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