

# 2D $S_N$ Radiation Transport with Diffusion Acceleration

MATH 676 – Final Presentation

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# Introduction

## Background

- My Ph.D. work involves various acceleration techniques for method of characteristics (MOC) radiation transport
- Some acceleration methods for discrete ordinates ( $S_N$ ) transport are similar to those in MOC transport
- In specific, the focus is diffusion synthetic acceleration (DSA), which attenuates the errors most poorly attenuated by source iteration

## Goals

- Develop a one-group, DGFEM,  $S_N$  radiation transport code
- Develop a one-group, DGFEM, diffusion radiation transport code
- Accelerate the  $S_N$  source iterations with DSA using the diffusion code
- More... if time permits (which it did!)

# One-group Linear Boltzmann Equation

Start with the one-group  $S_N$  transport equation for a single direction  $d$  (neglecting boundary conditions for simplicity), as

$$\mathbf{\Omega}_d \cdot \nabla \psi_d(\mathbf{x}) + (\sigma_a(\mathbf{x}) + \sigma_s(\mathbf{x})) \psi_d(\mathbf{x}) - \frac{\sigma_s(\mathbf{x})}{2\pi} \sum_{d=1}^{N_\Omega} \omega_d \psi_d(\mathbf{x}) = q(\mathbf{x}), \quad (1)$$

Let  $\mathbb{T}_h$  be the set of all cells of the triangulation in a discontinuous approximation space. The DG weak form with test function  $v_d$  is

$$\sum_{K \in \mathbb{T}_h} \left[ (-\mathbf{\Omega}_d \cdot \nabla v_d, \psi_d)_K + (\psi_d^+ \mathbf{\Omega}_d \cdot \mathbf{n}, v_d)_{\delta K} + (\sigma_t \psi_d, v_d)_K - (\sigma_s \phi, v_d)_K = (q, v_d)_K \right], \quad (2)$$

where  $\phi$  is the *scalar flux*,  $\phi = \sum_{d=1}^{N_\Omega} \omega_d \psi_d$ , and  $\psi_d^+$  is the upwind value of  $\psi_d$  (the value from the side of the face in which  $\mathbf{\Omega} \cdot \mathbf{n} \geq 0$ ).

# Source Iteration

To solve, cast Eq. (1) with iterative index  $\ell$  as

$$\mathbf{\Omega}_d \cdot \nabla \psi_d^{(\ell+1)} + \sigma_t \psi_d^{(\ell+1)} = \sigma_s \phi^{(\ell)} + q, \quad (3)$$

where  $\ell$  is the iterative index,  $\psi_d^{(0)} = \phi^{(0)} = \vec{0}$ . After solving each direction,  $d$ , for an iteration  $\ell$  in Eq. (3), update the scalar flux with

$$\phi^{(\ell+1)} = \frac{1}{2\pi} \sum_{d=1}^{N_\Omega} w_d \psi_d^{(\ell+1)}.$$

As  $\sigma_s/\sigma_t \rightarrow 1$ , particles scatter more before they are absorbed  $\rightarrow$  **the number of source iterations becomes significant!**

# Diffusion Acceleration

Simple algebraic manipulations can show that the error in  $\psi^{\ell+1}$  satisfies the transport equation with a source equal to:

$$R^{\ell+1} = \frac{\sigma_s}{2\pi}(\phi^{\ell+1} - \phi^\ell).$$

Fourier analysis shows that the angular flux error has a linearly anisotropic angular dependence. The diffusion approximation is exact for such a dependence  $\rightarrow$  cast the diffusion problem to form an error equation with the diffusion approximation that will attenuate the errors most poorly attenuated by the transport solve.

Said approximation is cast as:

$$-\nabla \cdot D \nabla \delta e^{\ell+1} + \sigma_a \delta e^{\ell+1} = \sigma_s (\phi^{\ell+1} - \phi^\ell), \quad (4)$$

where  $e^{\ell+1}$  is the approximated error in  $\phi^{\ell+1}$  and  $D = 1/3\sigma_t$ .

## Diffusion Acceleration (cont.)

Casting Equation (4) in the same DG space with interior edges  $\mathcal{E}_h^i$  and boundary edges  $\mathcal{E}_h^b$  using a modified interior penalty method for the face terms we obtain

$$\begin{aligned} \int_{\mathbb{T}_h} (D \nabla e \cdot \nabla v + \sigma_a e v) + \int_{\mathcal{E}_h^i} (\{ \{ D \delta_n e \} \} [v] + \{ \{ D \delta_n v \} \} [e] + \kappa [e] [v]) \\ + \int_{\mathcal{E}_h^b} (\kappa e v - D v \delta_n e - D e \delta_n v) = \int_{\mathbb{T}_h} (\phi^{\ell+1} - \phi^\ell) v, \quad (5) \end{aligned}$$

where

$$\{ \{ u \} \} \equiv \frac{u^+ + u^-}{2} \quad \text{and} \quad [u] \equiv u^+ - u^-,$$

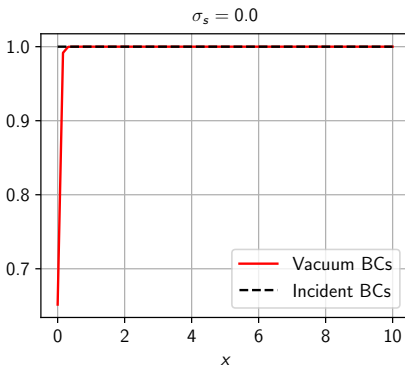
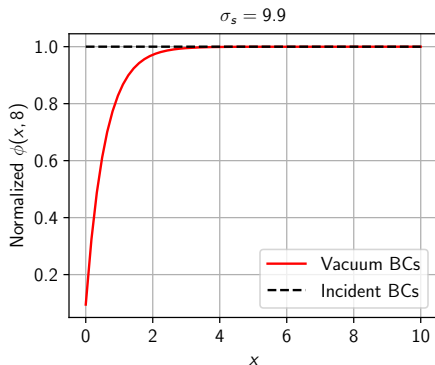
in which the penalty coefficient is

$$\kappa = \begin{cases} 2 \left( \frac{D^+}{h_\perp^+} + \frac{D^-}{h_\perp^-} \right) & \text{for interior edges,} \\ 8 \frac{D^-}{h_\perp^-} & \text{for boundary edges,} \end{cases}$$

and  $h_\perp^\pm$  is a characteristic length of the cell in the direction orthogonal to the edge.

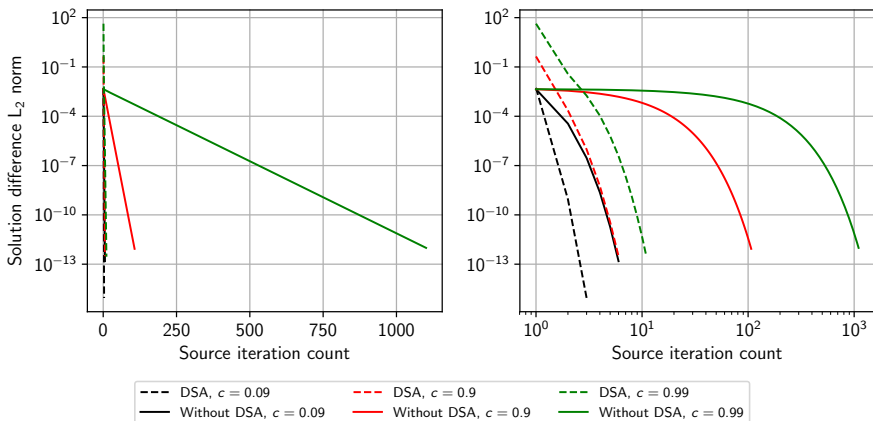
# Constant Solution Verification

Consider  $\mathcal{D} = [0, 10]^2$ ,  $N_\Omega = 4$ ,  $q = 1$ ,  $\sigma_t = 10$ , and  $64^2$  elements. Top and right boundary conditions are reflective. Bottom and left boundary conditions are either vacuum or incident isotropic flux of  $q/\sigma_a$ .



# Diffusion Acceleration Results

Consider  $\mathcal{D} = [0, 10]^2$ ,  $N_\Omega = 20$ ,  $q = 1$ ,  $\sigma_a + \sigma_s = \sigma_t = 100$ , and  $16^2$  elements. Increase the scattering ratio,  $c = \sigma_s/\sigma_t$ , with and without diffusion acceleration and observe the results:





# Additional Goals Completed

The primary implementation (transport with diffusion acceleration) was completed earlier than expected. Therefore, additional goals were added (and completed!) to round out the project:

## Parallel support

- Supports parallel solves using MPI and Trillinos wrappers, completed primarily by following step-40
- Transport is solved with GMRES and the AMG preconditioner
- Diffusion is solved with CG and the AMG preconditioner

## Reflecting boundary conditions

- Reflecting boundaries require storing the outgoing flux on the boundaries and then reflecting on the incoming boundaries (bit of a pain, but it works)
- Also supported in the diffusion acceleration scheme through adding an additional source term for boundary flux error

# Whoop

