i2 = (y[0] + 4*y[1] + 2*y[2] + 4*y[3] + y[4])/3*dxmyerr=np.abs(i1-i2) if myerr<tol:</pre> return i2 else: mid=(a+b)/2int1=Ez(fun, z, a, mid, to1/2)int2=Ez(fun,z,mid,b,tol/2) return int1+int2 In [84]: def compute field(zs=np.linspace(0,10,101)): #computes the electruc field for an array of z values R = E(1, 1, return r=True)E field = np.zeros(len(zs)) for i in tqdm(range(len(zs))): z=zs[i]if z == R: #if we encounter the value of R in the z array, we have a singularity, and we can't solve the integral print('Skipped evaluation at \$z=R\$') continue #skip this value if this happens E field[i] = Ez(E,z,0,np.pi,1)E field/=np.max(E field) fig, ax = plt.subplots() ax.plot(zs,E field) ax.set xlabel('\$Z\$') ax.set_ylabel('\$E/E_{max}\$') return E field E field = compute field() | 20/101 [00:00<00:01, 45.50it/s] Skipped evaluation at \$z=R\$ 100%| 101/101 [00:00<00:00, 110.22it/s] 1.0 8.0 0.6 0.4 0.2 0.0 10 6 Ζ In [85]: #now compare with the scipy quad function zs = np.linspace(0, 10, 101)E field = np.zeros(len(zs)) E_field_err = np.zeros(len(zs)) for i in tqdm(range(len(zs))): z=zs[i]E_field[i], E_field_err[i] = quad(E,0,np.pi, args = z) E field err/=np.max(E field) E field/=np.max(E field) 0%| | 0/101 [00:00<?, ?it/s]<ipython-input-85-6baacf238d6e>:10: IntegrationWarn ing: The occurrence of roundoff error is detected, which prevents the requested tolerance from being achieved. The error may be underestimated. E_field[i], E_field_err[i] = quad(E,0,np.pi, args = z) 101/101 [00:00<00:00, 1722.54it/s] In [89]: fig, ax = plt.subplots(2,1, gridspec_kw={'height_ratios': [3, 1]}) ax[0].plot(zs,E field) ax[1].set xlabel('\$z\$')ax[0].set_ylabel('\$E/E_{max}\$') ax[1].plot(zs, E field err) ax[1].set ylabel('Error\$/E {max}\$') 1.0 0.8 0.6 0.4 0.2 0.0 2 $1e^{0}_{-9}$ 6 10 Error/E_{max} 10 Out[89]: Text(0, 0.5, 'Error\$/E_{max}\$') We can see that both the integrator from the lectures and the scipy.integrate.quad integrator yield the expected results. I.e., the field is 0 inside the shell and goes down as $1/z^2$ when we hit z=R. However, there is a singularity at z=R where the electric field is undefined. Quad does not care whether or not there is a z value that is equal to R, but our integrator has trouble dealing with that value because it can never reach the desired tolerance. This is why we have to skip that value for our integrator. **Problem 2** The integrator made in class was making way too many function evaluations. When dividing the interval in 2 parts, we were recalculating all 5 points in the new interval. However, the 5 points form the previous interval are all re-used in the two new smaller intervals (i.e., the first and third point of a large interval bounds the first smaller interval, with the 2nd point of the original interval being the midpoint of the first smaller interval, and similarly for the second small interval). Hence, each time we divide our interval in 2, we only need to do 4 new function evaluations instead of 10. When we want this to happen, we need to have an argument in the function where if we feed it the points from the previous iteration, it uses and only computes the remaining points. If not, we proceed as usual. Each time the extra information is used, we save 3 function calls, hence the counter goes up by three in each block where extra exist. This is what is shown below. In [91]: def integrate adaptive(fun,a,b,tol, extra = None): x=np.linspace(a,b,5)dx = x[1] - x[0]if isinstance (extra, tuple): #if extra actually contains information ys = extra[0] # the old reusable y values counts = extra[1]+3 #counting the saved calls y new = fun(x[1::2]) #evaluating the function only where we dont know it already $y = np.array([ys[0], y_new[0], ys[1], y_new[1], ys[2]])$ #combine the new and old yval ues else: #if it is the first function call, proceed as before y=fun(x)counts = 0#do the 3-point integral i1 = (y[0] + 4*y[2] + y[4])/3*(2*dx)i2 = (y[0] + 4*y[1] + 2*y[2] + 4*y[3] + y[4]) / 3*dxmyerr=np.abs(i1-i2) if myerr<tol:</pre> return i2, counts else: mid=(a+b)/2int1, counts1=integrate adaptive(fun, a, mid, to1/2, (y[:3], counts)) #pass the first 3 points to the first smaller interval int2, counts2=integrate_adaptive(fun, mid, b, to1/2, (y[2:], counts)) #pass the last 3 points to the second smaller inteval return int1+int2, counts1+counts2 We can test the improved integrator below for a couple of functions In [70]: start = time.time() ans, saved_calls=integrate_adaptive(np.exp,0,10,1e-8) stop = time.time() print('Answer was ',ans, 'with error',ans-(np.exp(10)-np.exp(0)), 'saving', saved_calls, 'calls. Ran in ',stop-start,'seconds') Answer was 22025.465794806933 with error 2.1464074961841106e-10 saving 52002 calls. Ran in 0.13402295112609863 seconds In [71]: start = time.time() ans, saved calls=integrate adaptive (np.sin, 0, 15, 1e-8) stop = time.time() print('Answer was ',ans, 'with error',ans-(-np.cos(15)+np.cos(0)), 'saving', saved_calls, 'calls. Ran in ',stop-start,'seconds') Answer was 1.7596879128806284 with error 2.1807000649687325e-11 saving 12987 calls. Ran in 0.04883980751037598 seconds **Problem 3** To model our function with a Chebyshev polynomial, we want $y_i = \sum_j c_j T_j(x_i)$, which can be expressed in matrix form as y=Ax=Ac, where c is the vector of coefficients, and A is an m imes n matrix where m is the number of points we are fitting to, and n is the order of the polynomial we want to fit. Each column is the j^{th} Chebyshev polynomial evaluated at x_i . i.e, the 3rd column of the matrix has entries corresponding to $T_3(x_i)=4x_i^3-3x_i$. Doing a least squares fit, we can find the coefficients by the following operation $c=(A^TA)^{-1}A^Ty$. We can then evaluate this polynomial for $y=log_2(x)$ between 0.5 and 1. The following function uses np.polynomial.chebyshev.Chebyshev.fit, which returns a Chebyshev object, which is itself a function, so we can evaluate the polynomial directly. It performs the least-squares fit. It also has a trim method that trims the polynomial to the desired accuracy by removing all of the terms with a smaller order of magnitude than the tolerance. We can do this because the error is only as big as the sum of the coefficients. The image of the chebyshev polynomials is bounded between -1 and 1, so the "worst possible error" is if all points happen to evaluate at ± 1 in the polynomials, so the total error is the sum of the coefficients we drop. In [72]: **def** model_log2_cheb(a=0.5,b=1,ord = 150,tol=1e-6): x = np.linspace(a,b,1001) #create sufficiently large x array y = np.log2(x)cheb = np.polynomial.chebyshev.Chebyshev.fit(x,y,deg=ord, domain = (a,b)) #perform le ast squares fit cheb = cheb.trim(tol) #remove coeffecients below tolerance ncoeffs = len(cheb) cheb eval = cheb(x) #evaluate the polynomial for all x xy = np.vstack((x,cheb_eval)).T #the full curve return xy, ncoeffs, cheb #return the full x-y points, the number of coefficients and t he chebyshev object data,ncoeffs,cheb = model_log2_cheb() err = np.log2(data[:,0])-data[:,1] fig,ax = plt.subplots(2,1,gridspec kw={'height ratios': [3,1]}) ax[0].plot(data[:,0],data[:,1], linewidth = 2, label = 'Chebyshev series of \$log 2(x)\$ wi th {} coefficients'.format(ncoeffs)) $ax[0].plot(data[:,0],np.log2(data[:,0]),'--', linewidth = 1, color = 'r', label = '<math>log_2$ ' (x)\$') ax[1].plot(data[:,0],err) ax[1].set xlabel('\$x\$') ax[1].set ylabel('Error') $ax[0].set_ylabel('$f(x)$')$ ax[0].legend(loc = 'lower right') fig.tight layout() 0.0 -0.2-0.4-0.6-0.8Chebyshev series of $log_2(x)$ with 8 coefficients $log_2(x)$ -1.00.7 0.5 0.6 0.8 0.9 1.0 2.5 0.0 -2.50.5 0.6 0.7 8.0 0.9 1.0 Х To compute the natural log of any number, we can use our previous function, using the log_2 to get to the ln. When splitting a number into its mantissa and exponent using np.frexp, it expresses it as $a imes 2^b$, where a is between 0.5 and 1. We know that $log_2(a imes 2^b)=log_2(a)+b$, so we can use our previous function to compute $log_2(a)$. We can then change our basis to find the ln: $ln(x)=rac{log_2(x)}{log_2(e)}$. We need to express ein the same way explained above before taking its log_2 , so that our function can handle it. In [78]: def mylog2(num, tol=1e-6, return err = False): man num, exp num = np.frexp(num) #split the number and e in exp-mantissa man e,exp e = np.frexp(np.e) cheb = model log2 cheb(tol = tol)[2] #get the chebyshev polynomial for the desired to

lerance

log2 e = cheb(man e) + exp e

err = np.log(num) - ln

return ln, err

We can test this with an array of values:

return ln

x = np.linspace(1, 101, 1001)

ax[1].plot(x, lns[1])

4

3

1

0

Out[92]: Text(0, 0.5, '\$f(x)\$')

1e-7

20

20

ax[1].set_xlabel('\$x\$')
ax[1].set_ylabel('Error')
ax[0].set_ylabel('\$f(x)\$')

lns = mylog2(x,return err=True)

ax[0].legend(loc = 'lower right')

if return err:

else:

In [92]:

ln=log2 num/log2 e #change of basis

log2 num = cheb(man num) + exp num # compute both log 2 as explained above

ax[0].plot(x,lns[0],linewidth = 2, label = 'Natural log using Chebyshev evaluation')

Natural log using Chebyshev evaluation

80

80

100

100

--- True Natural log

60

60

40

40

ax[0].plot(x,np.log(x), linewidth = 1,linestyle='--',label='True Natural log')

fig, ax = plt.subplots(2,1,gridspec_kw={'height_ratios': [3,1]})

In [1]: import numpy as np

import time

Problem 1

we can write

import matplotlib.pyplot as plt
from scipy.integrate import quad
import astropy.constants as c

If we split the charged sphere in little charge elements dq, each dq contributes by $dE=rac{dq}{4\pi\epsilon_0 r^2}$, where r is

the distance between the dq element and the point where we want to compute the field. We can rewrite r by using the law of cosines $r^2=z^2+R^2-2zRcos(\theta)$. Where z is the distance from the center of the sphere

However, only the electric field in the z direction will survive, the others get cancelled out by symmetry, hence

 $dE_z=rac{dq}{4\pi\epsilon_0(z^2+R^2-2zRcos(heta))}cos(eta)$, where eta is the angle between the z axis and the segment from dq to

def E(theta,z,R = 1,sigma = 0.01,return r=False):#defining the function to integrate

In [83]: def Ez(fun,z,a,b,tol): #integrator used in the lectures modified such that it evaluates at

return R**2*sigma/(2*c.eps0.value)*((z-R*np.cos(theta))*np.sin(theta))/(R**2+z**2-2*R

and R the radius of the sphere. We can plug this r^2 in the expression for dE.

from tqdm import tqdm

%matplotlib notebook

 $dE = rac{dq}{4\pi\epsilon_0(z^2 + R^2 - 2zRcos(heta))}$

the point we evaluate the field at.

Hence $dE_z=rac{(z-Rcos(heta))~dq}{4\pi\epsilon_0(z^2+R^2-2zRcos(heta))^{3/2}}$

and $dq=\sigma dA=\sigma R^2 sin(heta)d heta d\phi$

 $E(z)=rac{R^2\sigma}{2\epsilon_0}\int_0^\pirac{(z-Rcos(heta))sin(heta)d heta}{(z^2+R^2-2zRcos(heta))^{3/2}}$,

Integrating both sides, we get

if return r:

dx=x[1]-x[0]y=fun(x,z)

a z value

return R

*z*np.cos(theta))**(3/2)

x=np.linspace(a,b,5)

#do the 3-point integral

i1 = (y[0] + 4*y[2] + y[4])/3*(2*dx)

In [82]:

and $cos(eta) = rac{z - Rcos(heta)}{r} = rac{z - Rcos(heta)}{(z^2 + R^2 - 2zRcos(heta))^{1/2}}$

Combining all of this, $dE_z=rac{(z-Rcos(heta))\ \sigma R^2 sin(heta)d heta d\phi}{4\pi\epsilon_0(z^2+R^2-2zRcos(heta))^{3/2}}$

We can solve this integral using numerical methods as below