



# Studies of covariance and nonlocality in Bohmian mechanics

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Group: Foundations of Quantum Mechanics

- The principle of general covariance: *coordinates are mere event markers upon which the ultimate laws of physics should not rely.*
- Entanglement is *not invariant* upon changes of quantum reference frames<sup>1</sup>.
  - For pure states, nonlocality (NL) is a frame-dependent quantity.
- Our proposal:
  1. Explore a broader notion of NL within Bohmian mechanics
  2. Quantify NL by analysing the particle's motion.

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1. R. M. Angelo, N. Brunner, S. Popescu, A. Short, and P. Skrzypczyk, J. Phys. A: Math. Theor. 44, 145304 (2011),
2. R. M. Angelo and A. D. Ribeiro, J. Phys. A: Math. Theor. 45, 465306 (2012),
3. F. Giacomini, E. Castro-Ruiz, and C. . Brukner, Nat. Commun.10, 494 (2019).

Bohmian mechanics (BM):

- The universe consists of **particles** with well-defined **trajectories**; Generalized state  $(\mathbf{x}, \psi)$ .
- Two-particle system ( $a$  and  $b$ ): The motion of  $b$  is given by

$$m_b \ddot{x}_b = F_b^{[\psi]}(x_a, x_b, t) = -\partial_{x_b} \left( \mathcal{V} + Q^{[\psi]} \right),$$

$$Q^{[\psi]} = \underbrace{-\frac{\hbar^2}{2|\psi|} \left( \frac{\partial_{x_a}^2}{m_a} + \frac{\partial_{x_b}^2}{m_b} \right) |\psi|}_{\text{Quantum potential}}$$

- Physical state at  $t$ :  $\left( \underbrace{x_a, x_b}_{\text{System configuration}}, \underbrace{\psi}_{\text{Wave function}} \right)$
- $Q^{[\psi]}$  is one of the ingredients needed to quantify NL.

In BM one needs to consider two types of NL.

- Einstein's principle of local action (**LA**): *externally influencing A has no immediate influence on B*
- In Newton's mechanics **x**: *disturbances in A's position cannot instantaneously alter B's acceleration*
- In quantum mechanics  **$\psi$** : *a measurement result on B is unaffected by operations on A*
- The **LA** is formulated according to a given definition of state
- In BM  **$(\mathbf{x}, \psi)$** : the system's motion can violate **LA** for reasons concerning its particle **and** wave aspects.

Next: How to assess this violation?

- The conditional average of  $F_b^{[\psi]}$  includes the systems preparation.

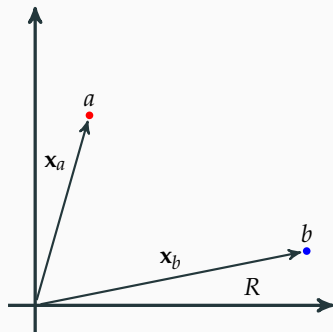
- For a system with two particles, we proposed

$$\begin{aligned}\langle F_b^{[\psi]} \rangle_{x_b} &:= \int F_b^{[\psi]} \rho_{x_a|x_b} dx_a \\ &= \int F_b^{[\psi]}(x_a, x_b) \rho_{x_a} dx_a + \underbrace{\int F_b^{[\psi]}(x_a, x_b) (\rho_{x_a|x_b} - \rho_{x_a}) dx_a}_{\eta},\end{aligned}$$

with

- $\rho_{x_a, x_b} = |\psi|$  is the joint probability;
  - $\rho_{x_{a(b)}} = \int dx_{b(a)} \rho_{x_a, x_b}$  are the marginals probabilities;
  - $\rho_{x_a|x_b} := \rho_{x_a, x_b} / \rho_{x_b}$  the conditional probabilities.
- $\eta$  truly includes the system's preparation into the dynamics.
  - it explicitly depends on the correlation between the variables: if  $x_a$  is independent, then  $\rho_{x_a|x_b} = \rho_{x_a}$  and  $\eta = 0$ .

## Two free particles: analysis from a classical inertial reference frame $R$



Assume:

- a separable gaussian wave function  
 $\psi = \psi_a \psi_b$ ;  
 $\psi_k = \psi_k(x_k, t) = \sqrt{G_{\delta_k}}(x_k - \bar{x}_k)$ ,  
with mean value  $\bar{x}_k$  and deviation  $\delta_k$  and  $k = a, b$ .

Then:

- The acceleration of  $b$  is

$$\ddot{x}_b = \frac{x_b - \bar{x}_b}{t_b^2},$$

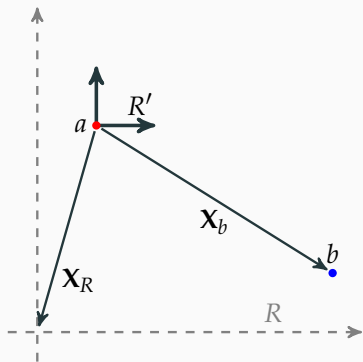
with  $t_b := t_{b,0} \sqrt{1 + (t/t_{b,0})^2}$  and  $t_{b,0} = 2m_b \delta_{b,0}^2 / \hbar$  the **Ehrenfest's scale** of time for  $b$ .

- The motion is entirely local: no entanglement + inertial  $R$ . Thus, our formalism is innocuous,

$$\langle \ddot{x}_b \rangle_{x_b} = \ddot{x}_b;$$

**there is no NL to be detected.**

## Two free particles: analysis from a quantum reference frame $R'$



Consider:

- Particle  $b$  is entangled with  $R$ :

$$\begin{aligned}\Psi(X_R, X_b) &= \psi_a(-X_R)\psi_b(X_b - X_R) \\ &\neq \Psi_R(X_R)\Psi_b(X_b)\end{aligned}$$

- Acceleration of  $b$

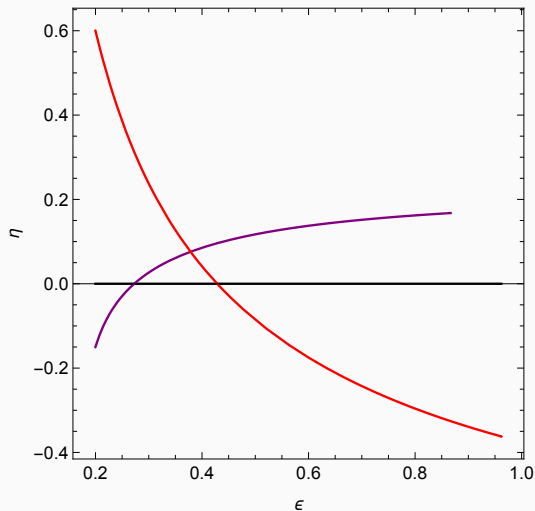
$$\ddot{X}_b = \frac{X_b - \bar{X}_b}{t_b^2} - \frac{X_R - \bar{X}_R}{t_b^2} + \frac{X_R - \bar{X}_R}{t_a^2},$$

is in general non-local (except for the special case  $t_a = t_b$ ).

- Therefore, we must identify a non-local influence caused by entanglement in the averaged motion.
- Local average of  $b$

$$\langle \ddot{X}_b \rangle_{X_b} = \underbrace{\frac{X_b - \bar{X}_b}{t_b^2}}_{\text{Local acceleration}} \left\{ 1 + \underbrace{\left[ \left( \frac{t_b}{t_a} \right)^2 - 1 \right] \lambda^2}_{\eta} \right\},$$

which  $\lambda$  is the statistical correlation between  $X_b$  and  $X_R$ .



- It indicates a **monotonic** and a **counter monotonic** relation between  $\eta$  and entanglement ( $\epsilon$ ).
- It is **monotonic** whenever Ehrenfest's scale associated with the **quantum reference frame** is bigger than the scale associated with the **system**:  $t_{a,0} > t_{b,0}$ . And **counter monotonic**, otherwise.
- Conclusion
  1. Broader notion of nonlocality (NL) within Bohmian mechanics;
  2. NL is intimately related to the physical state definition;
  3. Our proposal: NL quantifier based on the averaged motion;
  4. Case study: 2 free particles.  $\eta$  is a promising NL quantifier.