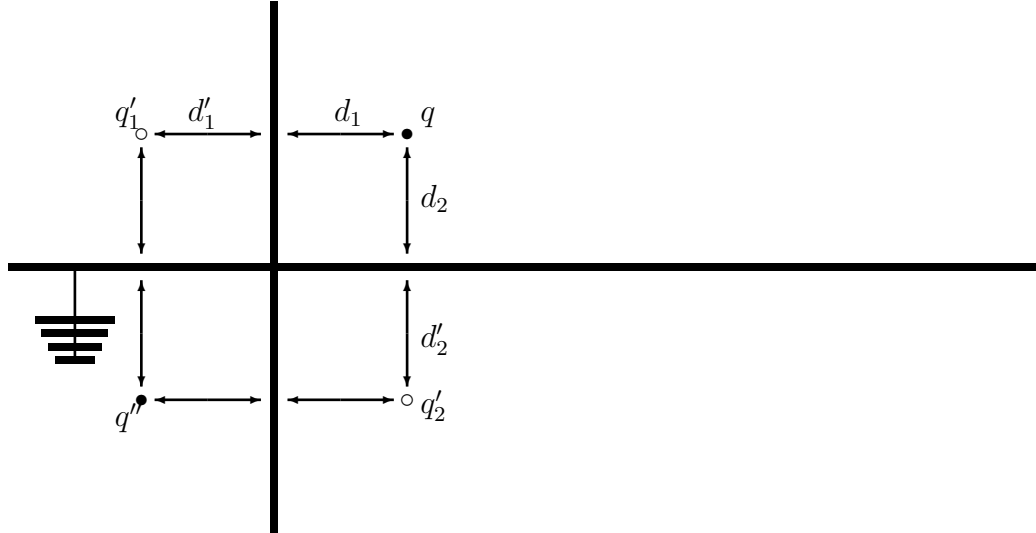


## Question 1



The infinite plates will have induced charges on them, each independently will act like an imaginary charge  $q'$  with the same charge but opposite sign as  $q$ . With both of the plates together, the induced charges on one plate will effect the charge  $q$  and also the imaginary charges  $q'_1$  and  $q'_2$ . This interaction will be like having another imaginary charge  $q''$  in the lower left corner shown in the modified diagram. This charge should be equal and opposite to the first imaginary charge, and thus  $q'' = q$ . Thus to calculate the force exerted on  $q$ , a calculation can be done as if negative charges  $q'_1$  and  $q'_2$  as well as a positive charge  $q''$  are placed in a quadrupole configuration.

The force on  $q$  is equal to the sum of the three forces:  $F_1$  from  $q'_1$ ,  $F_2$  from  $q'_2$  and  $F_3$  from  $q''$ . Using equation (1) and substituting the proper charge in for  $Q$ , the three forces can be found. Adding them together will give the total force  $\mathbf{F}_T$ .

$$\mathbf{F} = \frac{1}{4\pi\epsilon_0} \frac{qQ}{r^2} \hat{r} \quad (1)$$

$$\mathbf{F}_1 = \frac{1}{4\pi\epsilon_0} \frac{q(-q)}{(2d_1)^2} \hat{x} \quad (2)$$

$$\mathbf{F}_2 = \frac{1}{4\pi\epsilon_0} \frac{q(-q)}{(2d_2)^2} \hat{y} \quad (3)$$

$$\mathbf{F}_3 = \frac{1}{4\pi\epsilon_0} \frac{q(q)}{(2d_1)^2} \frac{\sqrt{2}}{2} \hat{x} + \frac{1}{4\pi\epsilon_0} \frac{q(q)}{(2d_2)^2} \frac{\sqrt{2}}{2} \hat{y} \quad (4)$$

$$\mathbf{F}_T = \frac{\sqrt{2} - 2}{32(d_1)^2\pi\epsilon_0} q^2 \hat{x} + \frac{\sqrt{2} - 2}{32(d_2)^2\pi\epsilon_0} q^2 \hat{y} \quad (5)$$

**Q.1 cont** I am guessing that the far field potential goes like a  $\frac{1}{r^3}$  at values of  $r \gg 2d$ , much like a quadrupole would behave.

## Question 2

(a)

Prove the Green's reciprocity theorem:

$$\int_{allspace} \rho_1 \phi_2 d^3r = \int_{allspace} \rho_2 \phi_1 d^3r$$

Evaluating the integral  $\int \mathbf{E}_1 \cdot \mathbf{E}_2 d^3r$  substituting  $\mathbf{E}_1 = -\vec{\nabla} \phi_1$  by parts:

$$\int -\vec{\nabla} \phi_1 \cdot \mathbf{E}_2 d^3r = (-\phi_1 \mathbf{E}_2)|_{allspace} - \int -\phi_1 \vec{\nabla} \cdot \mathbf{E}_2 d^3r \quad (6)$$

$$\int -\vec{\nabla} \phi_1 \cdot \mathbf{E}_2 d^3r = \frac{1}{\epsilon_0} \int \phi_1 \rho_2 d^3r \quad (7)$$

Note that:  $(-\phi_1 \mathbf{E}_2)|_{allspace} = 0$ .

Then evaluating the same integral substituting  $\mathbf{E}_2 = -\vec{\nabla} \phi_2$  again by parts:

$$\int \mathbf{E}_1 \cdot -\vec{\nabla} \phi_2 d^3r = (-\phi_2 \mathbf{E}_1)|_{allspace} - \int \phi_2 \vec{\nabla} \cdot \mathbf{E}_1 d^3r \quad (8)$$

$$\int \mathbf{E}_1 \cdot -\vec{\nabla} \phi_2 d^3r = \frac{1}{\epsilon_0} \int \phi_2 \rho_1 d^3r \quad (9)$$

Thus:

$$\int \phi_1 \rho_2 d^3r = \int \phi_2 \rho_1 d^3r \quad (10)$$

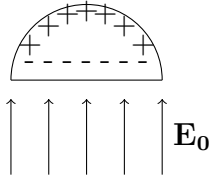
(b)

If putting a charge  $Q$  on  $A$  generates a potential  $V_{AB}$  on  $B$  and another  $V_{AA}$  on  $A$ , wouldn't generating a charge  $Q$  on  $B$  generate  $V_{BA}$  on  $A$  AND  $V_{BB}$  on  $B$ ? For the same reason? Thus, wouldn't  $V_{BB} + V_{BA} = V_{AB} + V_{AA}$ ?

## Question 3

The hemisphere is neutral to begin with, when placed in an electric field  $\mathbf{E}_0$  the hemisphere will polarize. I am not sure if the polarization will be

as displayed here, or a more difficult polarization with more negative charge towards the outer rim of the flat surface of the sphere.



The boundary conditions of the hemisphere are (i)  $V = 0$  when  $r = R$  and  $0 \leq \theta \leq \frac{\pi}{2}$ , (ii)  $V = 0$  when  $0 \leq r \leq R$  and  $\theta = \frac{\pi}{2}$  and finally (iii)  $V = -E_0 r \cos \theta + C$  when  $r \gg R$ . Unlike the sphere problem, I don't see an easy way to state that the potential is 0 for an entire plane and thus  $C$  cannot be eliminated right away.

Starting with the phi angle symmetric version of the spherical Laplace equation:

$$V(r, \theta) = \sum_{l=0}^{\infty} (A_l r^l + \frac{B_l}{r^{l+1}}) P_l(\cos \theta) \quad (11)$$

Boundary conditions are used to solve for  $A_l$  and  $B_l$ .  $P_l(\cos \theta)$  is the Legendre polynomial for  $\cos \theta$ .

From boundary condition (i):

$$0 = A_l R^l + \frac{B_l}{R^{l+1}} \text{ for } 0 \leq \theta \leq \frac{\pi}{2} \quad (12)$$

So in the northern hemisphere:

$$B_l = -A_l R^{2l+1} \quad (13)$$

$$V(r, \theta) = \sum_{l=0}^{\infty} A_l (r^l - \frac{R^{2l+1}}{r^{l+1}}) P_l(\cos \theta) \quad (14)$$

From boundary condition (iii)  $r \gg R$  the second term is very small and thus:

$$\sum_{l=0}^{\infty} A_l r^l P_l(\cos \theta) = -E_0 r \cos \theta + C \quad (15)$$

Since the result has  $\cos \theta$  and  $C$ , only  $l = 0$ ,  $l = 1$  matter.  $A_0 = C$  and  $A_1 = -E_0$

$$V(r, \theta) = C \frac{R}{r} - E_0 (r - \frac{R^3}{r^2}) \cos \theta \quad (16)$$

To eliminate C, I will simply drop it and test the boundary conditions. Does (17) satisfy my boundary conditions?

$$V(r, \theta) = -E_0(r - \frac{R^3}{r^2})\cos \theta \quad (17)$$

$$(i) \ r = R: \quad 0 \stackrel{?}{=} -E_0(R - \frac{R^3}{R^2})\cos \theta \quad (18)$$

$$0 \stackrel{\checkmark}{=} R - R \quad (19)$$

$$(ii) \ \theta = \frac{\pi}{2}: \quad 0 \stackrel{?}{=} -E_0(r - \frac{R^3}{r^2})\cos \frac{\pi}{2} \quad (20)$$

$$0 \stackrel{\checkmark}{=} \cos \frac{\pi}{2} \quad (21)$$

$$(iii) \ r \gg R: \quad V(r, \theta) \stackrel{\checkmark}{\rightarrow} -E_0 r \cos \theta : C = 0 \quad (22)$$

The conductor will look like a dipole far away.