

# NATIVE DIFFERENTIATION IN AN ALGEBRA

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Let  $A$  be a finite  $n$ -dimensional real algebra and let  $D \subseteq A$ , and let  $f : D \rightarrow A$ . By  $\mathbf{D}f$  we mean the total derivative of  $f$  regarded as a function  $\varphi(D) \rightarrow \mathbb{R}^n$  for some choice of basis  $\varphi : A \xrightarrow{\sim} \mathbb{R}^n$ ; we will usually use  $\phi$  implicitly.

1. Under what conditions is  $\mathbf{D}f$  represented in the algebra? I.e. when do we have

$$\mathbf{D}_{a_0}f(a) = f'(a_0)a \quad \text{or} \quad \mathbf{D}_{a_0}f(a) = af'(a_0)$$

for some function  $f' : D \rightarrow A$ ? Call such a function the (left or right) *native derivative* of  $f$ . Is this function well defined?

2. When does (1) for  $f$  imply (1) for  $f'$ ?
3. Is  $f'$  necessarily a derivation?
4. Convergence of powerseries via singularities in algebraic extension?
5. Relationship between generalized Cauchy-Riemann equations and gradients.
6. Do all gradients arise from reciprocal bases induced by a correlation?

See [1] which answers (1) and (3) mainly for unital associative algebras and touches upon (2), and see [2] which is relevant to (4).

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## Native Gradients

Our motivation is found in the *multivector derivative* from Clifford calculus, as can be found in [3]. For  $C$  a nondegenerate Clifford algebra, to any vector basis  $(e_1, \dots, e_n)$  we may associate a unique reciprocal basis  $(e^1, \dots, e^n)$  with the property that  $e^i \cdot e_j = \delta_j^i$ , where  $\delta_j^i$  is the Kronecker delta. These extend to algebra bases

$$\begin{aligned} \{e_{i_m} \wedge \dots \wedge e_{i_1} \mid 0 \leq m \leq n, 1 \leq i_1 < \dots < i_m \leq n\}, \\ \{e^{i_1} \wedge \dots \wedge e^{i_m} \mid 0 \leq m \leq n, 1 \leq i_1 < \dots < i_m \leq n\}, \end{aligned}$$

(including  $\mathbb{1}$  in the case  $m = 0$ ). Under the scalar product  $(a, b) \mapsto \langle ab \rangle$  (the scalar projection of the algebra product) these bases are reciprocal:

$$\left\langle (e^{i_1} \wedge \dots \wedge e^{i_m})(e_{j_m} \wedge \dots \wedge e_{j_1}) \right\rangle = \delta_{j_m}^{i_m} \dots \delta_{j_1}^{i_1}.$$

It is with these bases that the multivector derivative is defined as

$$\nabla f(x) = \sum_{\substack{0 \leq m \leq n \\ 1 \leq i_1 < \dots < i_m \leq n}} (e^{i_1} \wedge \dots \wedge e^{i_m}) \mathbf{D}_x f(e_{i_m} \wedge \dots \wedge e_{i_1}).$$

It can be found that this definition is independent of the initial choice of basis.

To generalize to any algebra  $A$ , let  $B = (\mathcal{B}, \mathcal{B}')$  be a pair of bases for  $A$  such that  $\mathcal{B} = (v_1, \dots, v_n)$  and  $\mathcal{B}' = (v'_1, \dots, v'_n)$ . Then define

$$\nabla^B f(x) = \sum_{i=1}^n v_i \mathbf{D}_x f(v'_i).$$

For this definition to be independent of  $B$ , consider  $C = (\mathcal{C}, \mathcal{C}')$  such that  $\mathcal{C} = (u_1, \dots, u_n)$  and  $\mathcal{C}' = (u'_1, \dots, u'_n)$  are bases. Let  $(x_1, \dots, x_n), (x'_1, \dots, x'_n), (y_1, \dots, y_n), (y'_1, \dots, y'_n)$  be the coordinates for  $\mathcal{B}, \mathcal{B}', \mathcal{C}, \mathcal{C}'$  respectively so that what we want is

$$\nabla^B f = \sum_{i=1}^n v_i \frac{\partial f}{\partial x'_i} = \sum_{j=1}^n u_j \frac{\partial f}{\partial y'_j} = \nabla^C f,$$

where we've suppressed the dependence on  $x$ . The partial derivatives between coordinates give the change-of-basis:

$$x_i = \sum_j \frac{\partial x_i}{\partial y_j} y_j, \quad v_i = \sum_j \frac{\partial y_j}{\partial x_i} u_j.$$

Together with the chain rule, this gives

$$\sum_i \sum_j u_j \frac{\partial y_j}{\partial x_i} \frac{\partial f}{\partial x'_i} = \sum_j \sum_i u_j \frac{\partial x'_i}{\partial y'_j} \frac{\partial f}{\partial x'_i} \implies \sum_i \sum_j u_j \left( \frac{\partial y_j}{\partial x_i} - \frac{\partial x'_i}{\partial y'_j} \right) \frac{\partial f}{\partial x'_i} = 0.$$

In order to be independent of  $f$ , we need  $\sum_j u_j (\partial y_j / \partial x_i - \partial x'_i / \partial y'_j) = 0$  for each  $i$ , and since  $(u_j)$  is a basis this is true iff

$$\frac{\partial y_j}{\partial x_i} - \frac{\partial x'_i}{\partial y'_j} = 0, \quad 1 \leq i, j \leq n.$$

Since the partials give the change-of-basis matrix, we may rewrite this condition as

$$(\mathcal{C} \leftarrow \mathcal{B})^T = (\mathcal{B}' \leftarrow \mathcal{C}') = (\mathcal{C}' \leftarrow \mathcal{B}')^{-1}$$

where  $(Y \leftarrow X)$  is the (coordinate) change-of-basis matrix from  $X$  to  $Y$  and  $(\_)^T$  is the transpose. It is evident that, conversely, if this condition is true then  $\nabla$  is independent of the choice of  $B$  or  $C$ . Note that when  $\mathcal{C} = \mathcal{B}'$  and  $\mathcal{C}' = \mathcal{B}$  this becomes

$$(\mathcal{B}' \leftarrow \mathcal{B})^T = (\mathcal{B}' \leftarrow \mathcal{B})$$

so that the change-of-basis matrix is symmetric.

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## Defining Native Derivatives

We find that whether or not  $A$  is unital makes a tremendous difference. First, the native derivative is not necessarily well defined: if  $f'_1, f'_2$  are both native derivatives of  $f$ , then for all  $a \in A$  we have (e.g. in the right case)

$$af'_1(a_0) = af'_2(a_0) \implies a(f'_1(a_0) - f'_2(a_0)) = 0;$$

if  $f'_1(a_0) - f'_2(a_0)$  is a nonzero *null element*, an element which multiplies with any other to 0, then  $f'_1$  and  $f'_2$  need not be equal. The case of the trivial algebra where every  $a$  is null is also possible, but seems useless in this context so we will not consider it here. Non-trivial algebras with nonzero null elements do exist:

**Example 1.** Let  $A = \mathbb{R}^2$  be an  $\mathbb{R}$ -algebra via  $(x, y)(z, w) = (xz, 0)$  for all  $x, y, z, w \in \mathbb{R}$ . It is easily verified that this defines an algebra but  $(0, y)(z, w) = (0, 0)$ , particularly when  $y \neq 0$ ; each  $(0, y)$  is a nonzero null element. Note that this is a non-unital algebra.

When  $A$  is unital there are no nonzero null elements, since if  $z$  is null then  $z = 1z = 0$ . Furthermore, we may take  $a = 1$  above to conclude that  $f'_1(a_0) = f'_2(a_0)$  so that in a unital algebra the native derivative is always well-defined when it exists. To say more about the non-unital case, we need some more facts about null elements.

**Proposition 2.** *The set of null elements  $\mathfrak{Z} \subseteq A$  forms a non-unital subalgebra; in particular, it is an algebra ideal.*

The proof of this is simple and direct so we omit it. Since  $\mathfrak{Z}$  is an ideal, we may quotient out and arrive at a definition of native derivative. In the following, let  $\pi : A \rightarrow A/\mathfrak{Z}$  be the canonical projection.

**Definition 3.** Let  $f : D \rightarrow A$  and  $a_0 \in D$ . Then  $f$  is (*left, right*) *natively differentiable at  $a_0$*  or just (*left, right*) *differentiable at  $a_0$*  if  $f$  is totally differentiable at  $a_0$  and there is  $\xi \in A$  such that for all  $a \in A$ .

$$\text{(left)} \quad \mathbf{D}_{a_0}f(a) = \xi a, \quad \text{(right)} \quad \mathbf{D}_{a_0}f(a) = a\xi$$

$f$  is (*right, left*)-differentiable if it is (*right, left*)-differentiable at each  $a_0 \in D$ . Then we define the *right derivative* of  $f$  to be the unique function  $f^R : D \rightarrow A/\mathfrak{Z}$  such that

$$\pi(\mathbf{D}_a f(b)) = \pi(b)f^R(a)$$

for all  $b \in A$ . We similarly define the *left derivative*  $f^L$ .

# Appendix

## 1 ————— Structure of Zero Divisors

Suppose  $A$  is an associative unital  $n$ -dimensional algebra. Then the (left) regular (matrix)  $\mathcal{B}$ -representation of  $A$  for  $\mathcal{B}$  a basis is

$$\mathcal{M}_{\mathcal{B}}(A) = \left\{ \left( [ae_1]_{\mathcal{B}}, [ae_2]_{\mathcal{B}}, \dots, [ae_n]_{\mathcal{B}} \right) \mid a \in A \right\}$$

for  $\mathcal{B} = (e_1, \dots, e_n)$ ; that is, the collection of all matrices with columns  $ae_1, \dots, ae_n$  represented in  $\mathcal{B}$ -coordinates. There is an isomorphism  $\Phi_{\mathcal{B}} : A \rightarrow \mathcal{M}_{\mathcal{B}}(A)$ .

For each  $a \in A$ , we may define the left and right annihilators to be

$$@_L a = \{a' \in A \mid a'a = 0\}, \quad @_R a = \{a' \in A \mid aa' = 0\}.$$

It is easy to see that these are linear subsets, and must be at most  $n-1$  dimensional since  $A$  is unital and  $\mathbb{1} \notin @_L a, @_R a$  for nonzero  $a$ . Then the set of all zero divisors may be expressed as

$$\text{zd}(A) := \{a \in A \mid \exists a' \in A \setminus \{0\}. aa' = 0 \text{ or } a'a = 0\} = \bigcup_{a \in A \setminus \{0\}} @_L a \cup @_R a.$$

This is to say that the set of zero divisors is a union of linear subspaces each with dimension  $\leq n-1$ ; however, given any basis  $\mathcal{B}$ , the zeros divisors can also be described as the set

$$\text{zd}(A) = \{a \in A \mid \det \Phi_{\mathcal{B}}(a) = 0\},$$

the roots of the polynomial  $\det \Phi_{\mathcal{B}}(a)$ . In order for  $\text{zd}(A)$  to be a union of linear subspaces, such a polynomial must have a linear factor for each such subspace. But there can only be *finitely many* such factors. It follows that  $\text{zd}(A)$  is a *finite* union of linear subspaces of  $A$  with dimension at most  $n-1$ . Note also that  $\det \Phi_{\mathcal{B}}(a)$  is a degree- $n$  polynomial and can have at most  $n$  linear factors, so  $\text{zd}(A)$  is made up of at most  $n$  such linear subspaces.

From this fact, we may construct an invertible basis for  $A$ . Let  $\text{zd}(A) = \bigsqcup_{k=1}^m Z_k$  such that  $m \leq n$  and each  $\emptyset \neq Z_k \subseteq A$  is a linear subspace. We note that  $\mathbb{1} \notin Z_k$  for any  $k$ , and so for any scalar  $c \neq 0$  and  $z \in Z_k$  it must be that  $z + c\mathbb{1} \notin Z_k$ . Then given a basis  $(\mathbb{1}, v_2, v_3, \dots, v_n)$ , we may construct an invertible basis  $(\mathbb{1}, v'_2, v'_3, \dots, v'_n)$  as follows: For any  $i$  where  $v_i$  is invertible, take  $v'_i = v_i$ . Then suppose WLOG that  $v_i \in Z_1$ . We have  $v_i + \mathbb{1} \notin Z_1$ ; if this is invertible, then we're done and take this as  $v'_i$ . Otherwise  $v_i + \mathbb{1} \in Z_2$  WLOG, and then  $v_i + \mathbb{1} + \mathbb{1} \notin Z_2$  but also still  $\notin Z_1$ ; if this is invertible, then take this as  $v'_i$ . Continue this process until  $v_i + l\mathbb{1}$  is invertible for some positive integer  $l$ ; this process terminates since there are only finitely many  $Z_k$ . Some version of this will work over any infinite field.

## 2 Reciprocal Bases

Let  $A$  be an  $n$ -dimensional  $K$ -vector space, and let  $A^*$  be its dual. Then for each basis  $\mathcal{B} = (e_1, \dots, e_n)$  of  $A$  there is a unique *dual basis*  $\mathcal{B}^* = (\varepsilon^1, \dots, \varepsilon^n)$  of  $A^*$  such that  $\varepsilon^j(e_k) = \delta_k^j$ . We use the convention that  $(\mathcal{B}^*)^* = \mathcal{B}$ . Every nondegenerate correlation  $\xi : A \rightarrow A^*$  with induced bilinear form  $\langle \_, \_ \rangle : A \times A \rightarrow K$  then gives rise to *reciprocal bases*.

**Definition 4.** The *left reciprocal basis* of  $\mathcal{B}$  is  $\mathcal{B}_L^\xi = \xi^{-1}[\mathcal{B}^*]$ , and the *right reciprocal basis* of  $\mathcal{B}$  is  $\mathcal{B}_R^\xi = \xi[\mathcal{B}]^*$ . If  $\mathcal{B}_L^\xi = (l^1, \dots, l^n)$  and  $\mathcal{B}_R^\xi = (r^1, \dots, r^n)$  then

$$\langle l^j, e_k \rangle = \langle e_k, r^j \rangle = \delta_k^j, \quad \forall j, k = 1, \dots, n.$$

**Proposition 5.**  $\mathcal{B}_L^\xi = \mathcal{B}_R^\xi$  iff  $\langle \_, \_ \rangle$  is symmetric;  $\mathcal{B}_L^\xi = -\mathcal{B}_R^\xi$  iff  $\langle \_, \_ \rangle$  is antisymmetric.

*Proof.* We give the proof for the symmetric case, and the antisymmetric case is very similar. ( $\implies$ ) Let  $e^j := l^j = r^j$ . Trivially  $\langle e^j, e_k \rangle = \langle e_k, e^j \rangle$ , from which it follows that  $\langle \_, \_ \rangle$  is symmetric. ( $\impliedby$ ) We have  $\langle e_j, l^k \rangle = \langle l^k, e_j \rangle = \delta_j^k$ , but this means that  $\mathcal{B}_L^\xi$  is the unique right reciprocal of  $\mathcal{B}$ ; thus  $\mathcal{B}_L^\xi = \mathcal{B}_R^\xi$ .  $\square$

Let  $T = (\mathcal{B}_L^\xi \leftarrow \mathcal{B})$ . Using the convention that repeated indices are summed over, we find that

$$\delta_k^j = \langle l^j, e_k \rangle = \langle l^j, l^i T_{ik} \rangle = \langle l^j, l^i \rangle T_{ik},$$

but this says that  $\langle l^j, l^i \rangle = T_{ji}^{-1}$ . Similar considerations give us the following proposition:

**Proposition 6.** Let  $T = (\mathcal{B}_L^\xi \leftarrow \mathcal{B})$  and  $U = (\mathcal{B}_R^\xi \leftarrow \mathcal{B})$ . Then

$$\begin{aligned} \langle l^j, l^k \rangle &= T_{jk}^{-1}, & \langle r^j, r^k \rangle &= U_{kj}^{-1}, \\ T_{kj} &= \langle e_j, e_k \rangle = U_{jk}, \end{aligned}$$

where the last equation is equivalent to  $T^T = U$ .

**Corollary 7.** If  $\langle \_, \_ \rangle$  is symmetric then  $(\mathcal{B}^\xi \leftarrow \mathcal{B})$  is symmetric.

## References

- [1] James S. Cook. *Introduction to  $\mathcal{A}$ -Calculus*. Aug. 4, 2017. arXiv: [1708.04135](https://arxiv.org/abs/1708.04135) [math]. URL: <http://arxiv.org/abs/1708.04135> (visited on 04/28/2021).
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- [3] Chris Doran and Anthony Lasenby. *Geometric Algebra for Physicists*. Cambridge: Cambridge University Press, 2003.