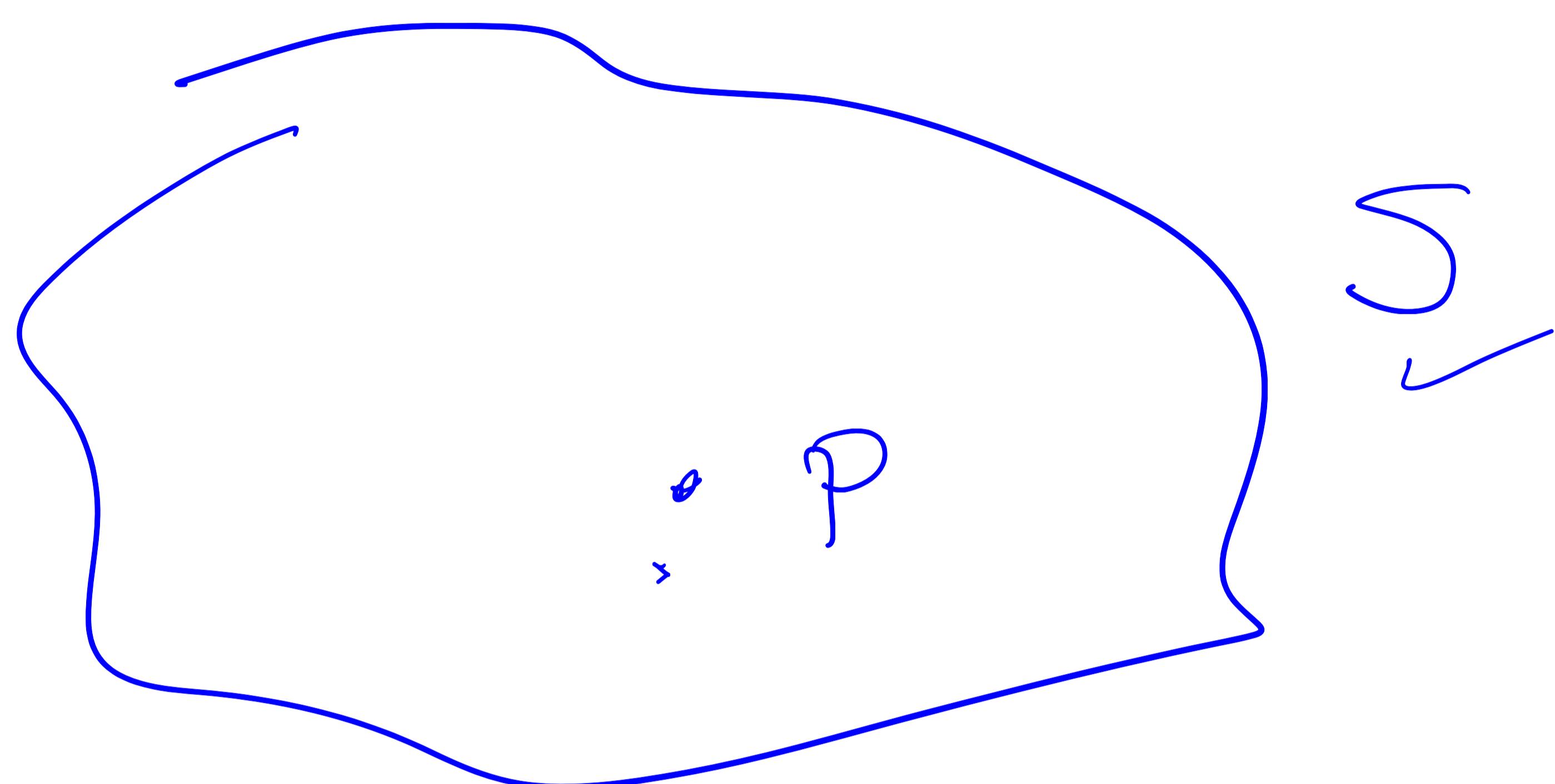


Kirchoff's theory of diffraction



$$(1) \quad V(x, y, z, t) = \underbrace{U(x, y, z)}_{(2)} e^{-i\omega t} \quad \nabla^2 + k^2 U = 0 \quad \checkmark$$

$$(2) \quad \int_V \bar{\nabla} \cdot \bar{A} d^3 \bar{x} = \oint_S \bar{A} \cdot \hat{n} dS$$

$$\bar{A} = U \nabla U'$$

$$\bar{\nabla} \cdot \bar{A} = \bar{\nabla} \cdot (U \nabla U') = U \nabla^2 U' + \bar{\nabla} U \cdot \nabla U'$$

$$\int_V U \nabla^2 U' + \underline{\nabla U \cdot \nabla U'} = \oint_S U \frac{\partial U'}{\partial n} dS \quad (1)$$

$$U \rightarrow U'$$

$$\int_V U' \nabla^2 U + \nabla U \cdot \nabla U' = \oint_S U' \frac{\partial U}{\partial n} dS \quad (2)$$

$$\int_V \left(U \nabla^2 U' - U' \nabla^2 U \right) d^3\bar{x} = \int_S \left(U \frac{\partial U'}{\partial n} - U' \frac{\partial U}{\partial n} \right) ds$$

Extra condition

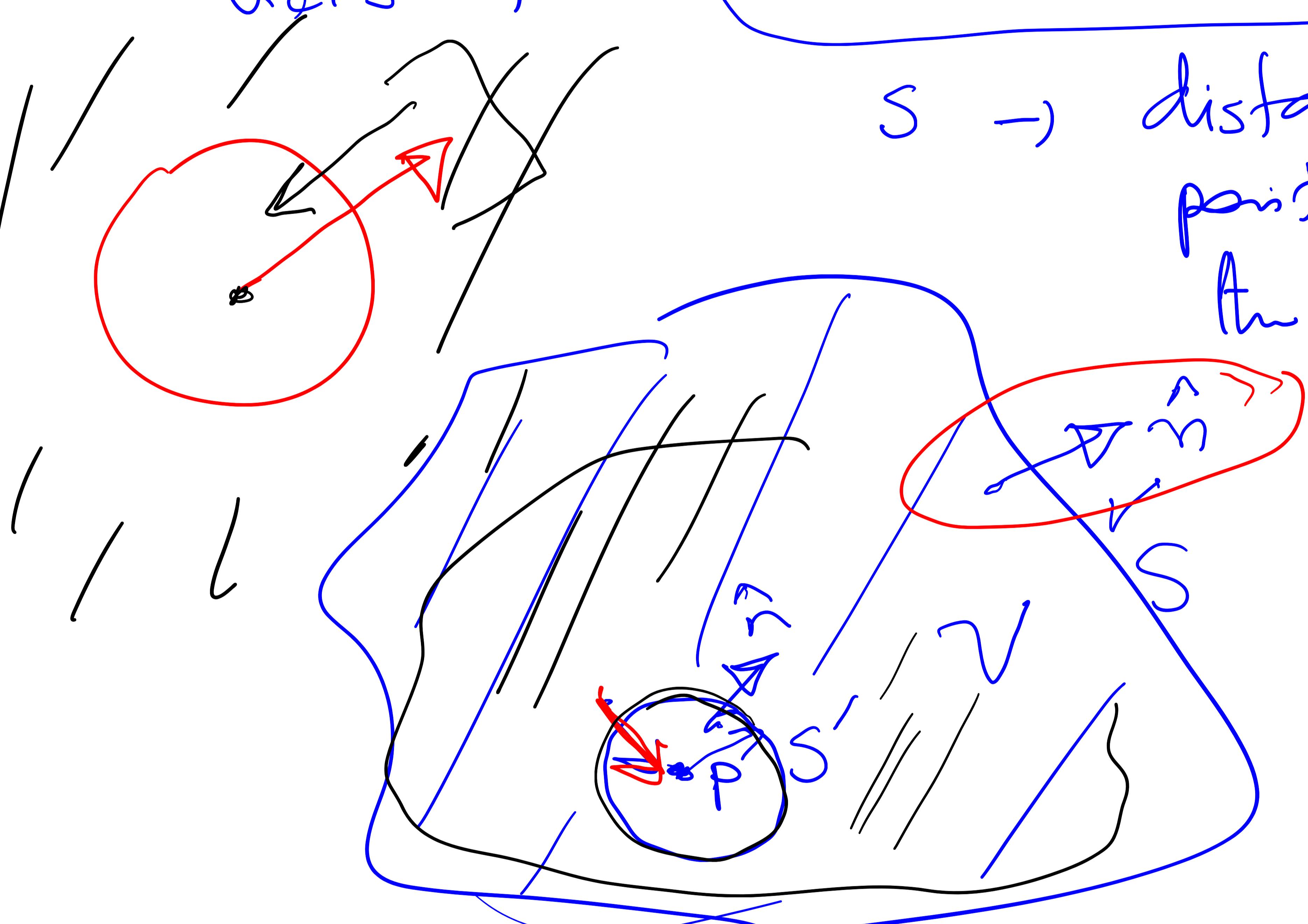
$$(\nabla^2 + k^2) U = 0, \quad (\nabla^2 + k^2) U' = 0$$

LHS = 0

$$\int_S \left(U \frac{\partial U'}{\partial n} - U' \frac{\partial U}{\partial n} \right) ds = 0 \quad \text{--- (3)}$$

Let's take $U' = e^{iks}/s$

$s \rightarrow$ distance between observation point and any point of the surface S



Excl. $s = 0$

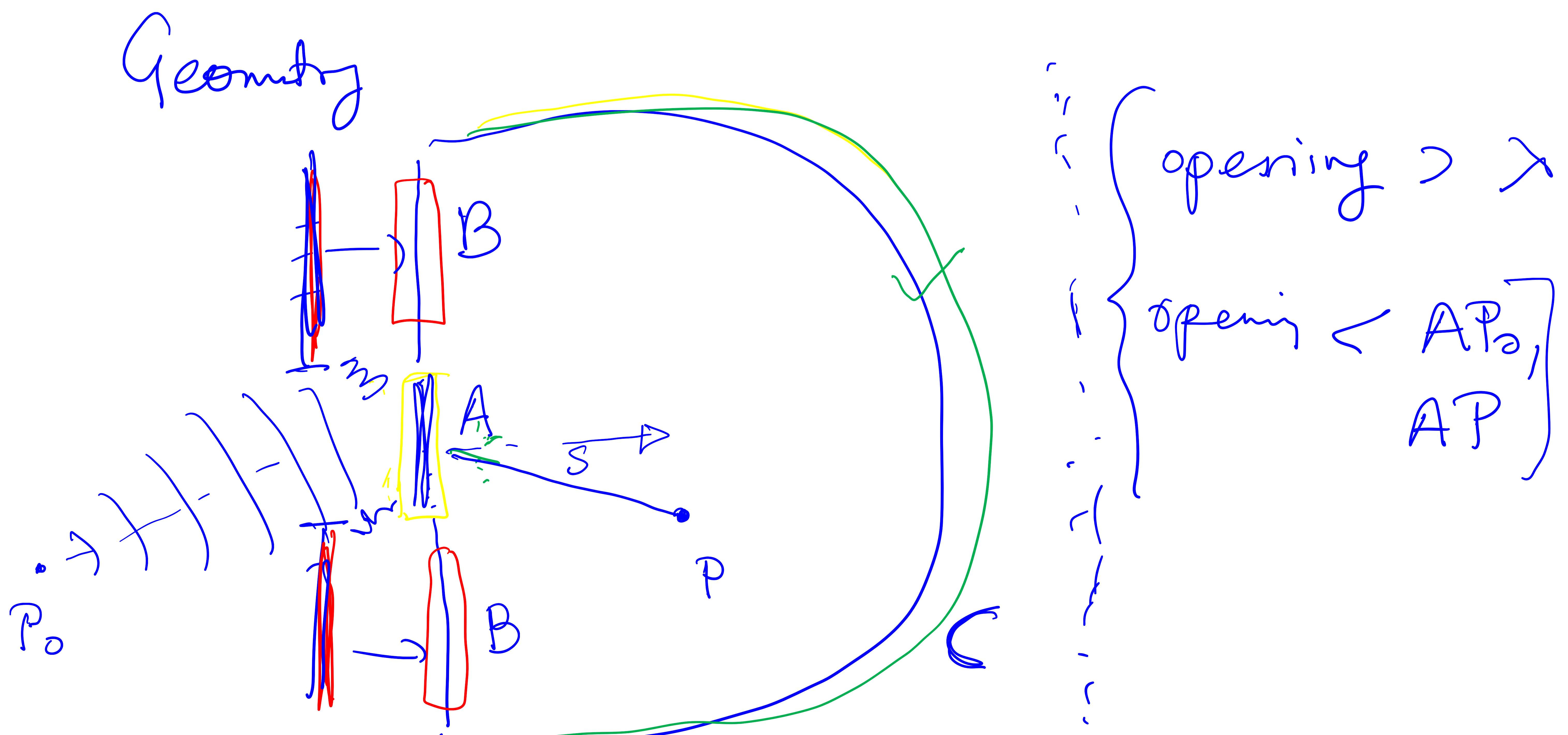
singularity

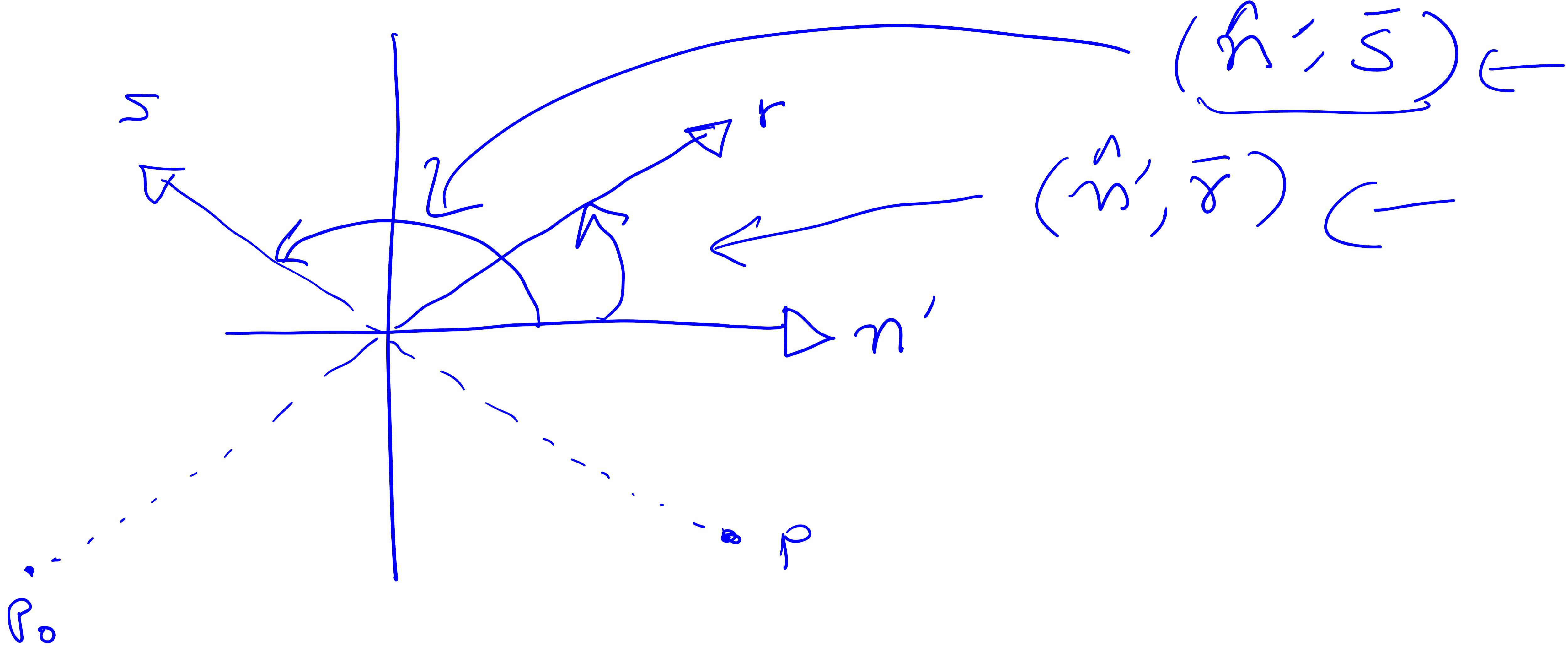
$$e^{ikr}$$

$$\left(\int_S + \int_{S'} \right) \underbrace{\left(U \frac{\partial}{\partial n} \left(\frac{e^{iks}}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n} \right)}_{ds} = 0$$

$$\begin{aligned}
 & \int_S \left(U \frac{\partial}{\partial n} \left(\frac{e^{iks}}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n} \right) dS \\
 &= \int_{S'} \left(U \frac{e^{iks}}{s} \left(ik - \frac{1}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n} \right) dS' \\
 &= \int_{S'} \left(U \frac{e^{ik\varepsilon}}{\varepsilon} \left(ik - \frac{1}{\varepsilon} \right) - \frac{e^{ik\varepsilon}}{\varepsilon} \frac{\partial U}{\partial n} \right) dS' \\
 &\stackrel{\varepsilon \rightarrow 0}{=} -4\pi U(P)
 \end{aligned}$$

$$\boxed{U(P) = -\frac{1}{4\pi} \iint_S \left(U \frac{\partial}{\partial n} \left(\frac{e^{iks}}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n} \right) dS}$$





Apply Kirchoff

$$U(P) = -\frac{1}{4\pi} \left(S_A + \cancel{S_B} + \cancel{S_C} \right) \left[U \frac{\partial}{\partial n} \left(\frac{e^{iks}}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n} \right] dS$$

"Kirchoff's assumptions"

- On A : $U = U^{(i)}$, $\frac{\partial U}{\partial n} = \frac{\partial U^{(i)}}{\partial n}$

$$\frac{A e^{ikr}}{\gamma}$$

$$- \frac{A e^{ikr}}{\gamma} (ik - \frac{1}{2})$$

- On B : $U = 0$, $\frac{\partial U}{\partial n} = 0$ (converges)

- On C : Large R, $U, \frac{\partial U}{\partial n} \rightarrow 0$

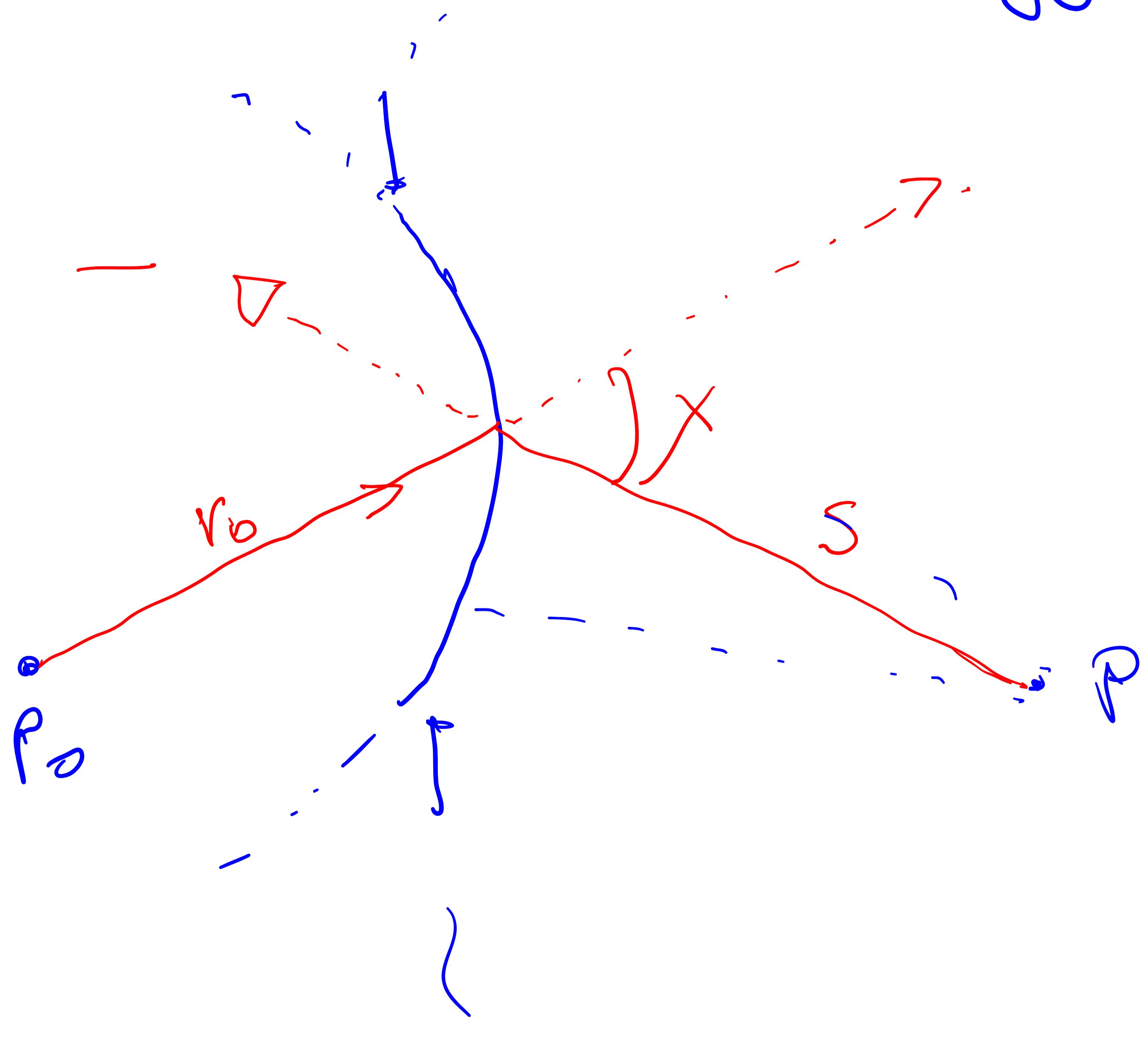
$$\Rightarrow U(P) = \frac{1}{4\pi} \iint_A \left(U \frac{\partial}{\partial n'} \left(\frac{e^{iks}}{s} \right) - \frac{e^{iks}}{s} \frac{\partial U}{\partial n'} \right) dS$$

$$= \frac{1}{4\pi} A \iint_A \frac{e^{ik(\gamma+s)}}{\gamma s} \left[\left(ik - \frac{\gamma}{s} \right) \cos(n', s) - \left(ik - \frac{\gamma}{s} \right) \cos(n', \gamma) \right]$$

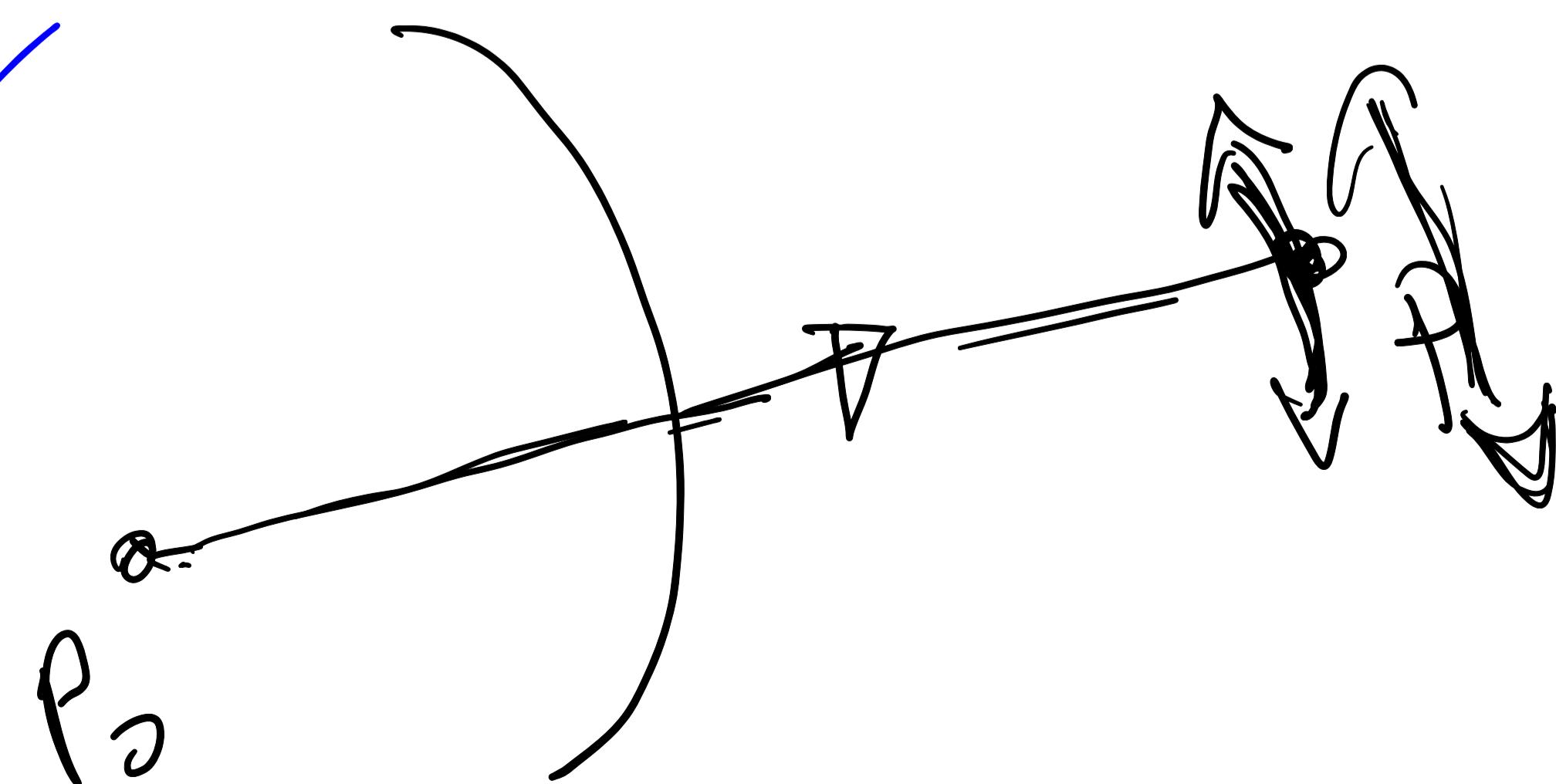
$r, s \rightarrow$ large compared to λ

$$U(P) = \frac{iA}{2\lambda} \iint_A \frac{e^{ik(\gamma+s)}}{\gamma s} \left[\cos(n', s) - \cos(n', \gamma) \right] dS$$

Connections to Huygen Fresnel theory



$$x = \pi - (\gamma_0, s)$$



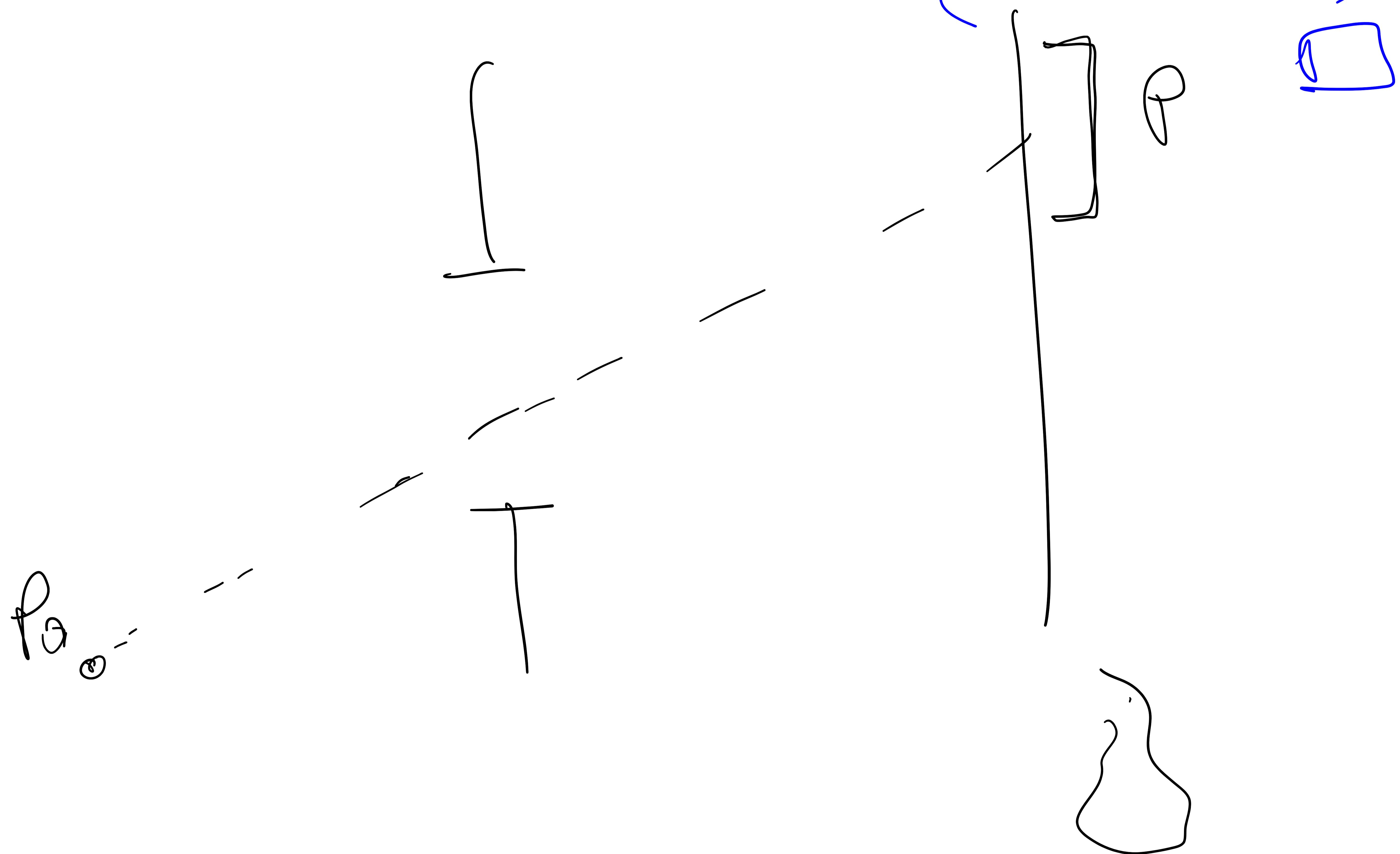
$$U(P) = \frac{-iA}{2\lambda} \frac{e^{ikr_0}}{r_0} \iint_A \frac{e^{iks}}{s} (1 + \cos x)$$

Huygen - Fresnel

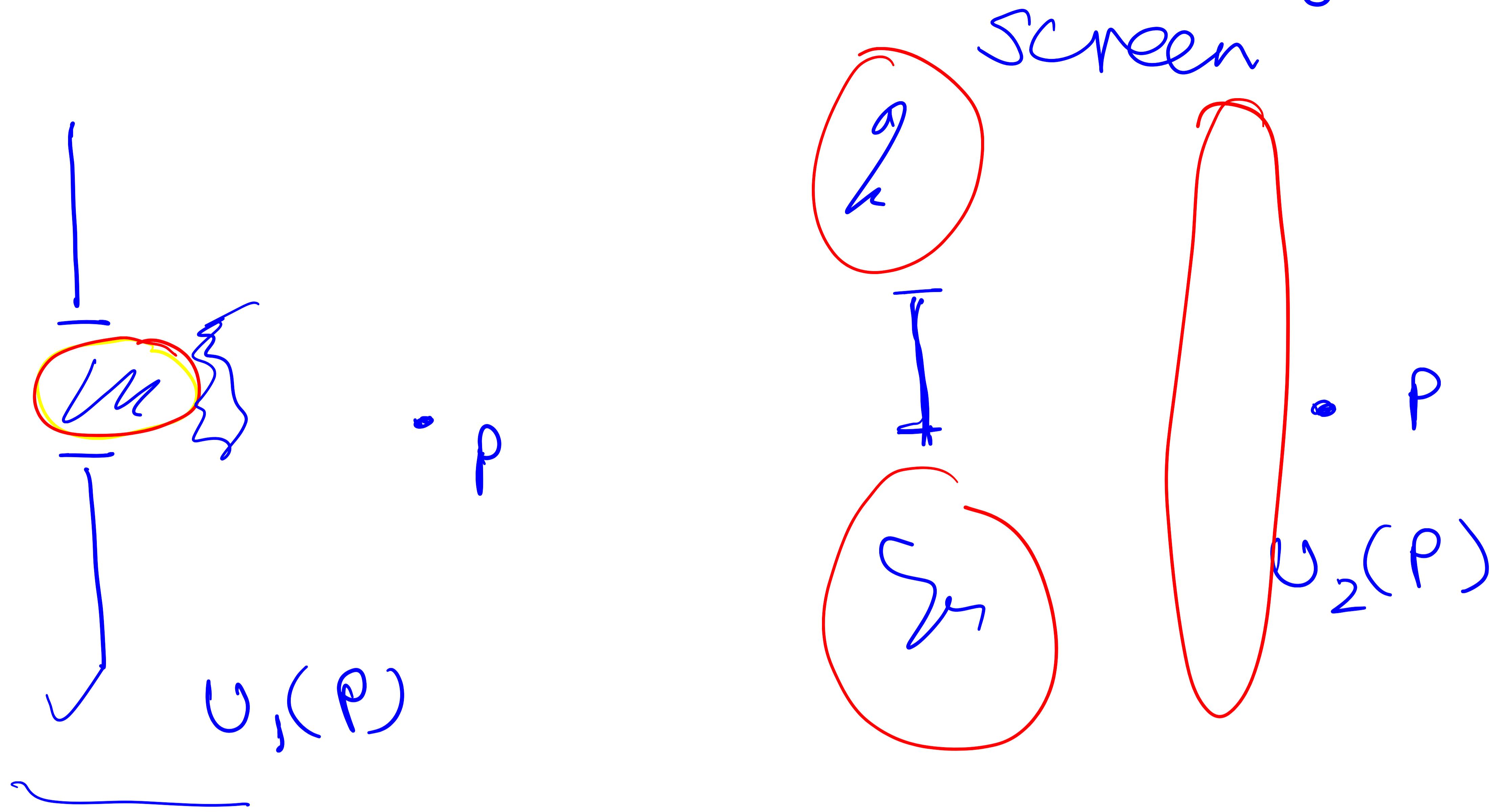
$$U(P) = \frac{A e^{ikr_0}}{r_0} \iint_A \frac{e^{iks}}{s} K(x) dS$$

$$K(x) = \frac{-i}{2x} (1 + \cos x) \quad \checkmark$$

(induction)



Babinet principle : "Complementary"



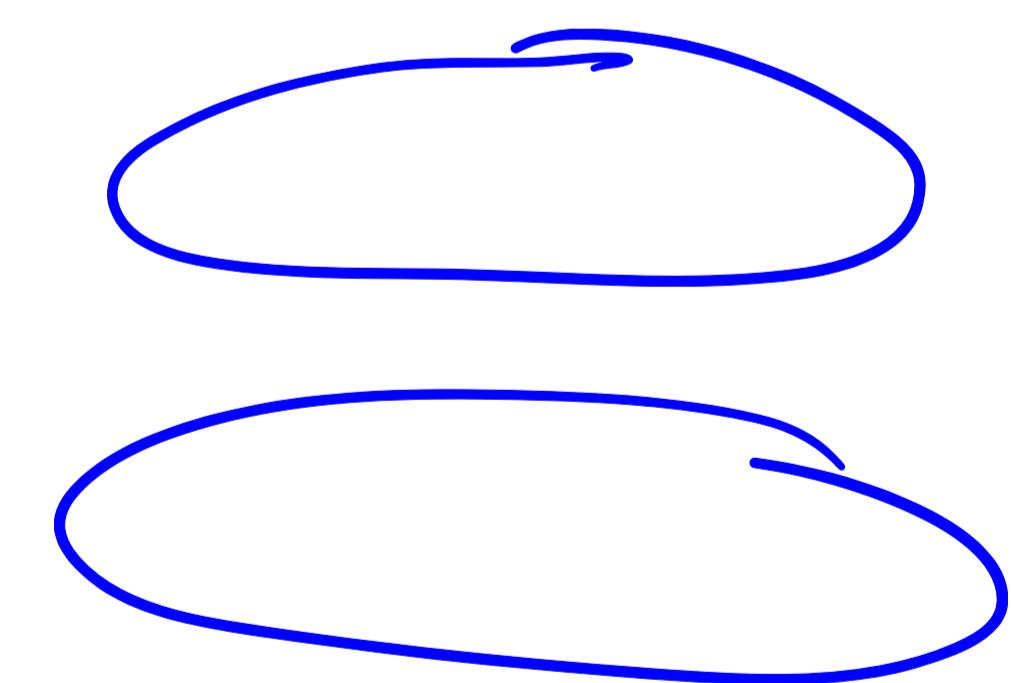
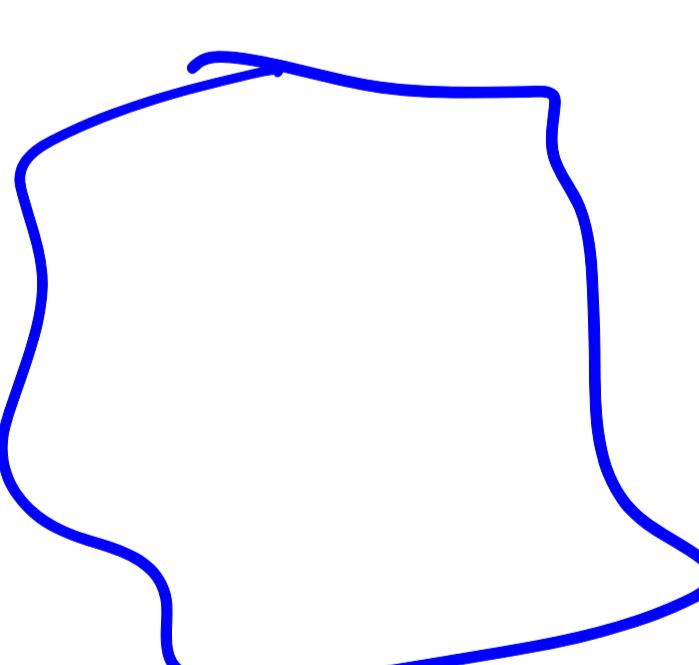
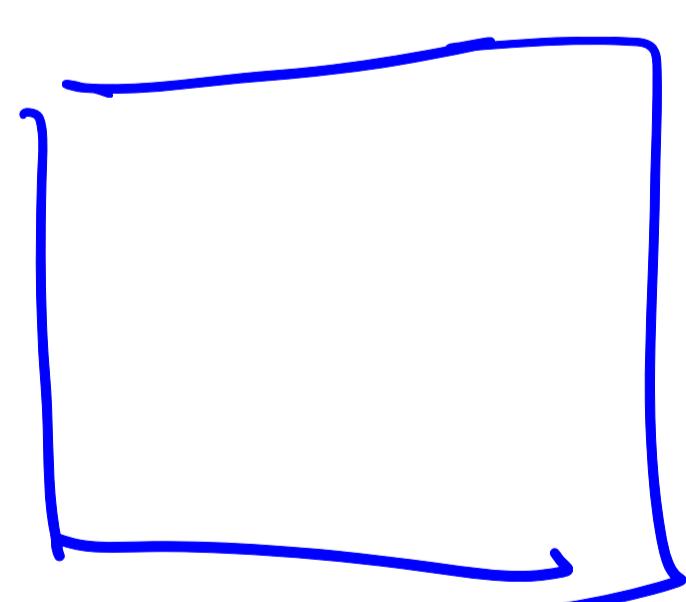
If $U(P)$ is the field at P when
no screen is present,

$$U(P) = \boxed{U_1(P)} + \boxed{U_2(P)}$$

Babinet

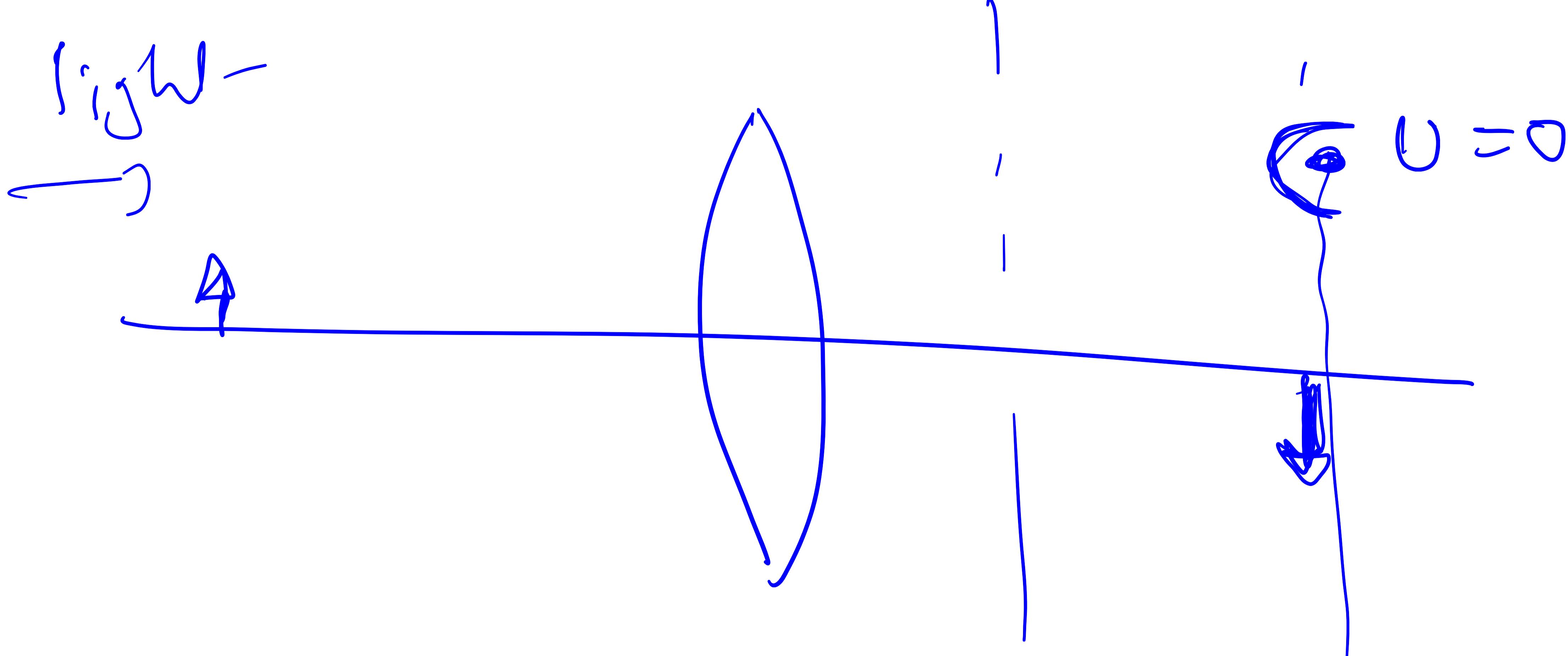
Implications

(1) If $U_1 = 0$; $U_2 = U$



(2) If $U = 0 \Rightarrow U_1 = -U_2$

Ex 1



$$U_1 = -U_2$$

Ex 2

