

# Signal To Noise

Often abbreviated as S/N or SN

Assume:

- photon counting detector
- not the same as integrating energy of arriving photons!
- sources of noise are independent, so they add in quadrature:

$$N_{total} = \sqrt{N_1^2 + N_2^2 + N_3^2 + \dots}$$

In the optical and shorter wavelengths, photon arrival obeys Poisson statistics

Probability of observing  $n$  photons in time  $t$ :

$$\text{Prob}(n) = \frac{(\lambda t)^n}{n!} e^{-\lambda t}$$

For photon arrival rate  $\lambda$ , which is related to the flux from the object and the sensitivity of the system you are observing with.

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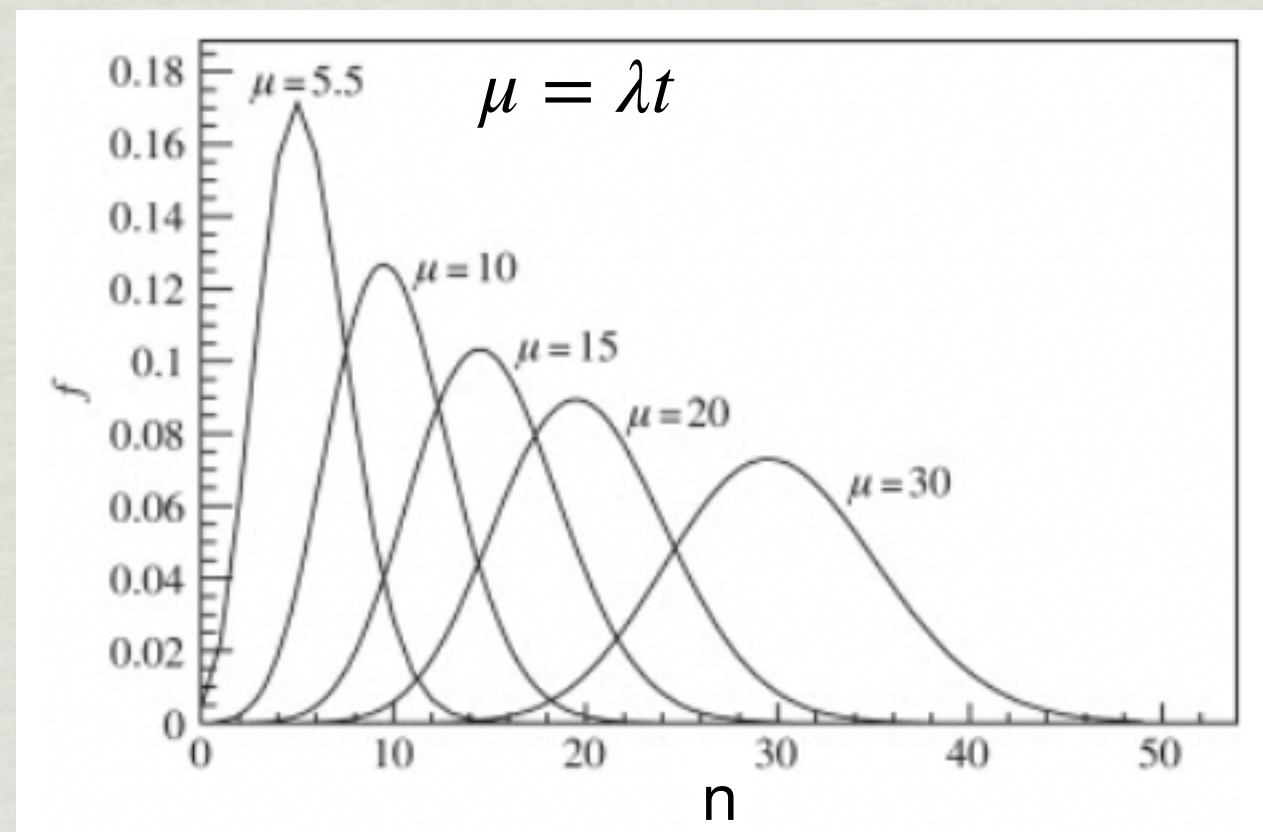
For photon arrival rate  $\lambda$ , which is related to the luminosity of the object and the sensitivity of the system you are observing with.

Even if you know  $\lambda$  perfectly, the distribution of the number of photons that arrive in time  $t$  will follow the Poisson distribution

How to reduce this uncertainty?

Measure  $N$  photons collected in many pixels to measure a uniform background, like the sky

Measure  $N$  photons in a pixel in many exposures of your source, or use a longer exposure time so  $N$  is larger

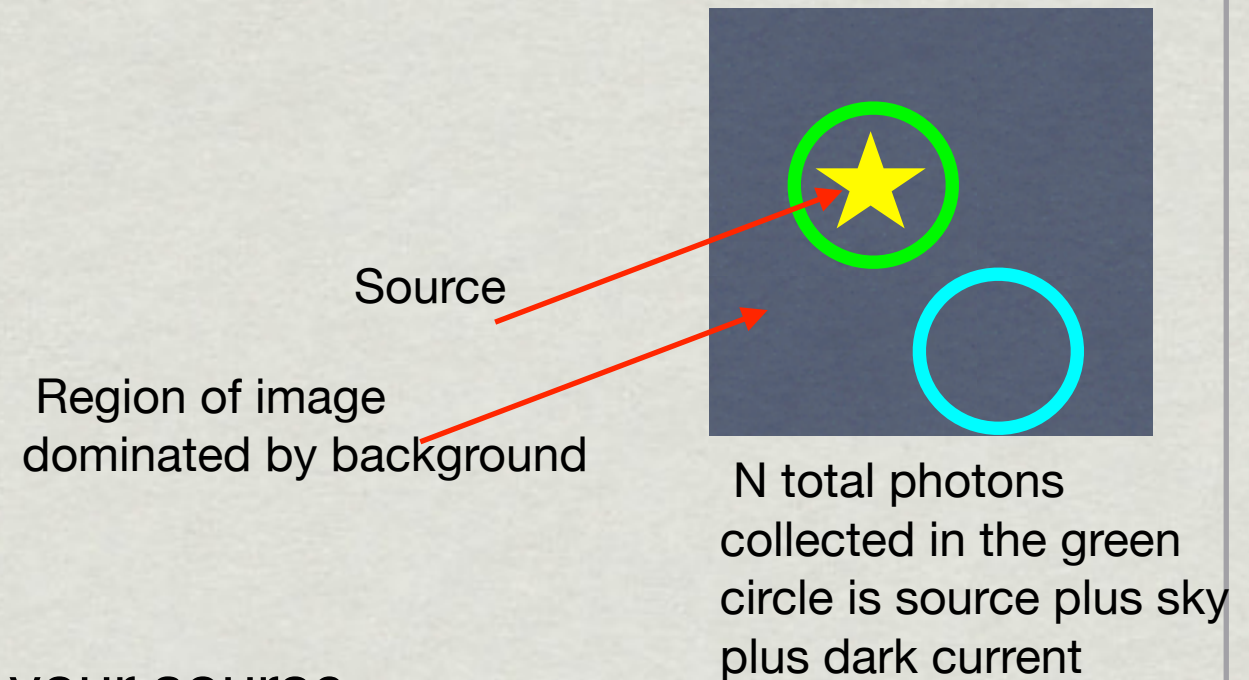




# Signal To Noise

Signal to Noise ratio for an unresolved source:

$$\frac{S_{\gamma}t}{\sqrt{S_{\gamma}t + B_{\gamma}tA_{\Omega} + RN_e^2N_{pix}}}$$



$S_{\gamma}t$  = your signal, the number of photons from your source

Written as the product of the rate of photon arrival from the source  $S_{\gamma}$  in counts/second and the exposure time  $t$  you spent collecting the photons

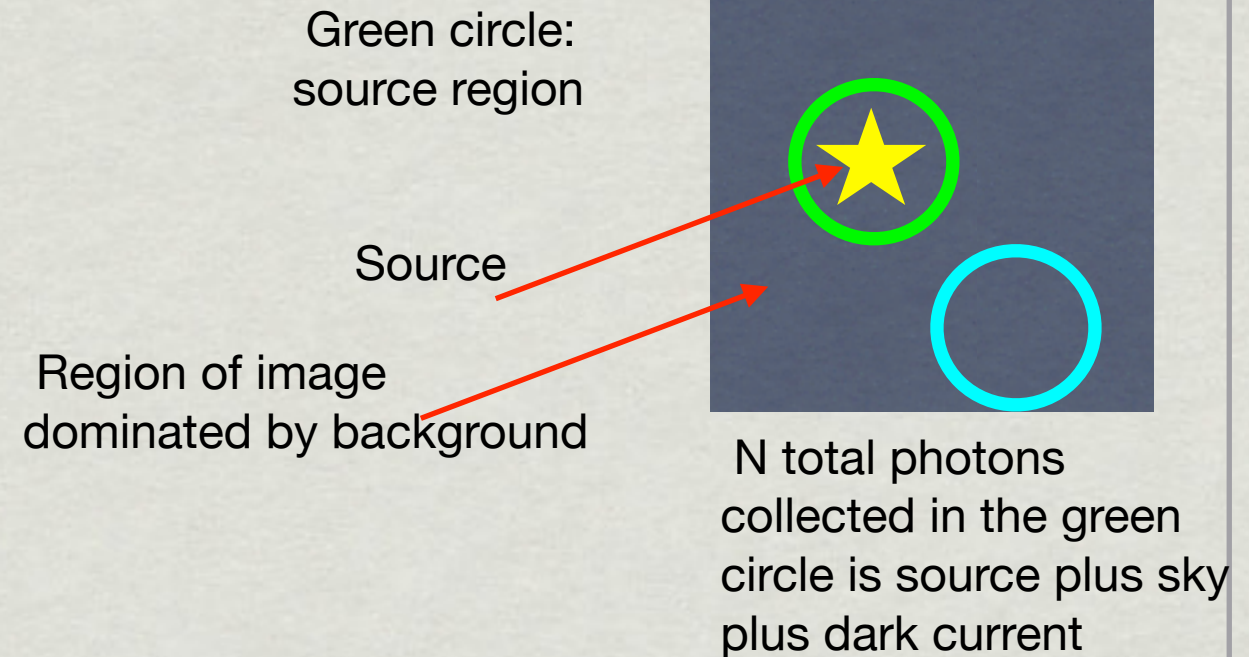
Denominator is the “noise”, the uncertainty in your measurement

- quadrature sum of the random noise processes that matters most for optical and NIR astronomy data. Adjust for radio, submm, other wavelengths
- does not include systematics, which can easily dominate any measurement error or uncertainty
  - examples: unknown gain, unstable bias or background,...

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$S_{\gamma}t$  = your signal, N, the number of photons from your source

$B_{\gamma}tA_{\Omega}$  = number of photons from the background **in the source region**

$B_{\gamma}t$  = number of photons from the sky detected per square arcsecond in that region. Rate of photon arrival from the sky per square arcsecond  $B_{\gamma}$  times exposure time  $t$ .

This expression for signal to noise assumes you know  $B_{\gamma}$  perfectly

$A_{\Omega}$  = area of source region, in square arcseconds (same sky units as  $B_{\gamma}$ ).



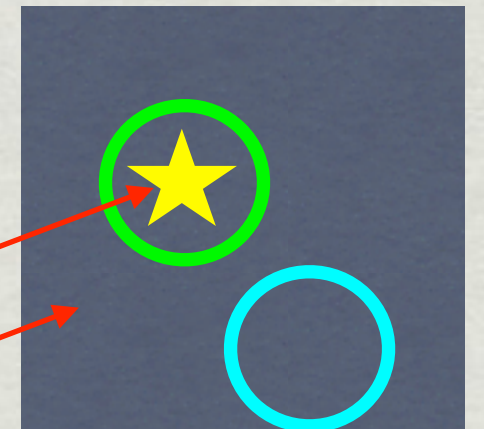
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Region of image  
dominated by background

Source



N total photons  
collected in the green  
circle is source plus sky  
plus dark current

$S_{\gamma}t$  = your signal, the number of photons from your source,  $N_s$

Written as the product of the rate of photon arrival from the source  $S_{\gamma}$  and the exposure time  $t$  you spent collecting the photons

You measure some number of counts in your source region  $N_{\text{total}}$  by adding up the counts in the green circle in the cartoon

Subtract an estimate of the background contribution by adding up the counts in a background region, the cyan circle in this cartoon.

Background estimate is  $N_B = B_{\gamma}tA_{\Omega}$

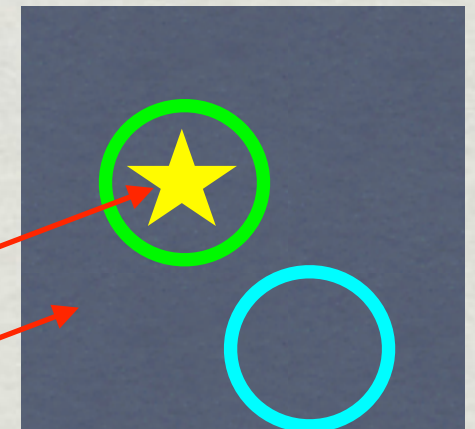
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Propagation of the uncertainties: assume you only need to include the Poisson uncertainty in the number of photons you collect in your source region

$S_{\gamma}t + B_{\gamma}tA_{\Omega} = N_{\text{total}}$  the total number of photons collected in the source region

$\sqrt{S_{\gamma}t + B_{\gamma}tA_{\Omega}}$  is the Poisson uncertainty on  $N_{\text{total}}$

That assumes you measure  $B_{\gamma}$  over a large enough area there is no additional uncertainty in your background estimate to include in the uncertainty on  $S_{\gamma}t$ .

To make this assumption valid, you should use as large a region as you can to estimate the background. We will explore that in a homework and discuss.



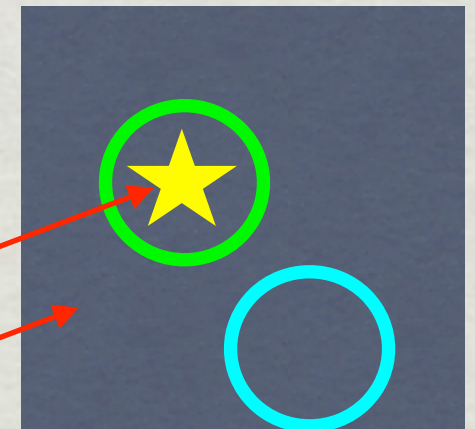
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Some detectors generate their own background.

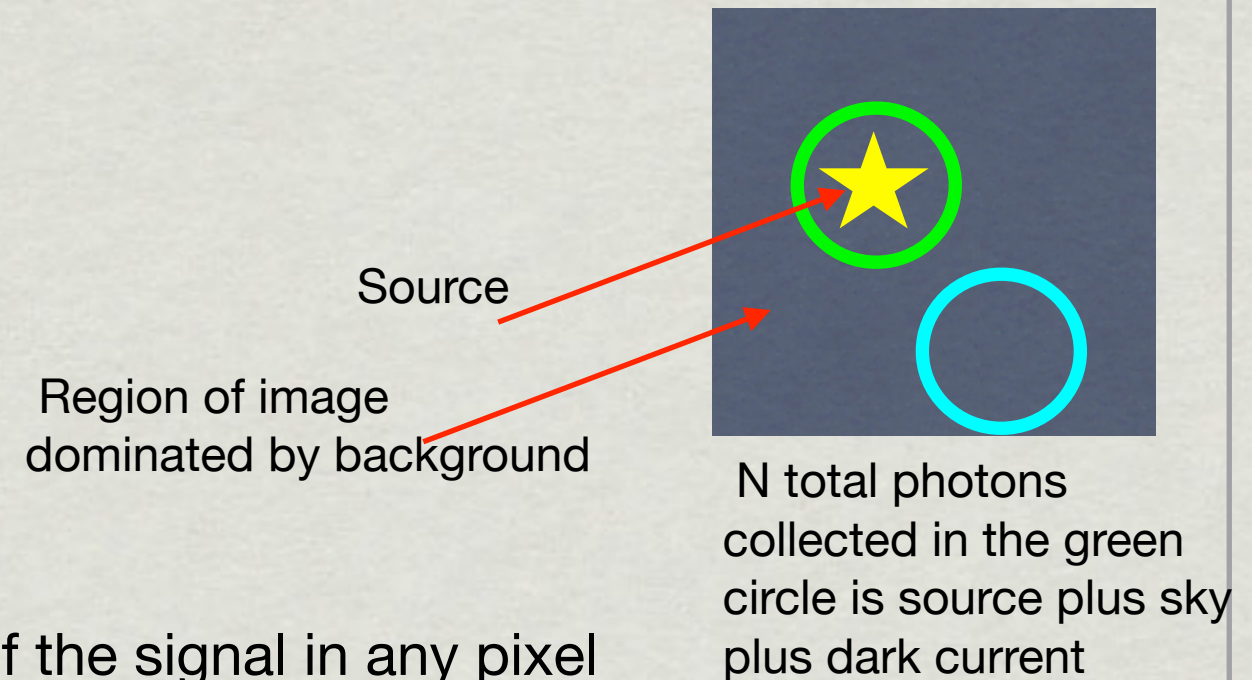
“dark current” is electrons that get a boost from thermal energy and are collected just like electrons that interact with astronomical photons

Same kind of measurement considerations apply: try to measure detector background rate well enough that you can neglect that source of uncertainty  
- impossible if that detector background is not stable — beware!

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Readnoise RN: uncertainty in the measurement of the signal in any pixel

- in photon counting detectors, units are electrons
- quantifies detector's (like a CCD or an IR detector) imperfect ability to measure the signal in any pixel
- NOT the same thing as uncertainty due to Poisson arrival statistics
- RN is per pixel. It contributes the quadrature sum of the contribution per pixel in your source region, assuming RN is the same for every pixel:

$$\sqrt{RN_1^2 + RN_2^2 + RN_3^2 + \dots + RN_N^2} = \sqrt{N_{pix} \times RN^2}$$



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Ways to improve SN:

Lower readnoise

- want source to occupy minimum number of pixels: fewer reads, more flux/pixel
- “minimum number of pixels” can only go so far
- this is what “binning” is for on detectors that can do it. Fewer reads per binned pixel
- read slower

Lower the dark current

- why you always hear about “filling dewars” and cryo-coolers

Read the detector slower (see next slide)

Reduce the sky background

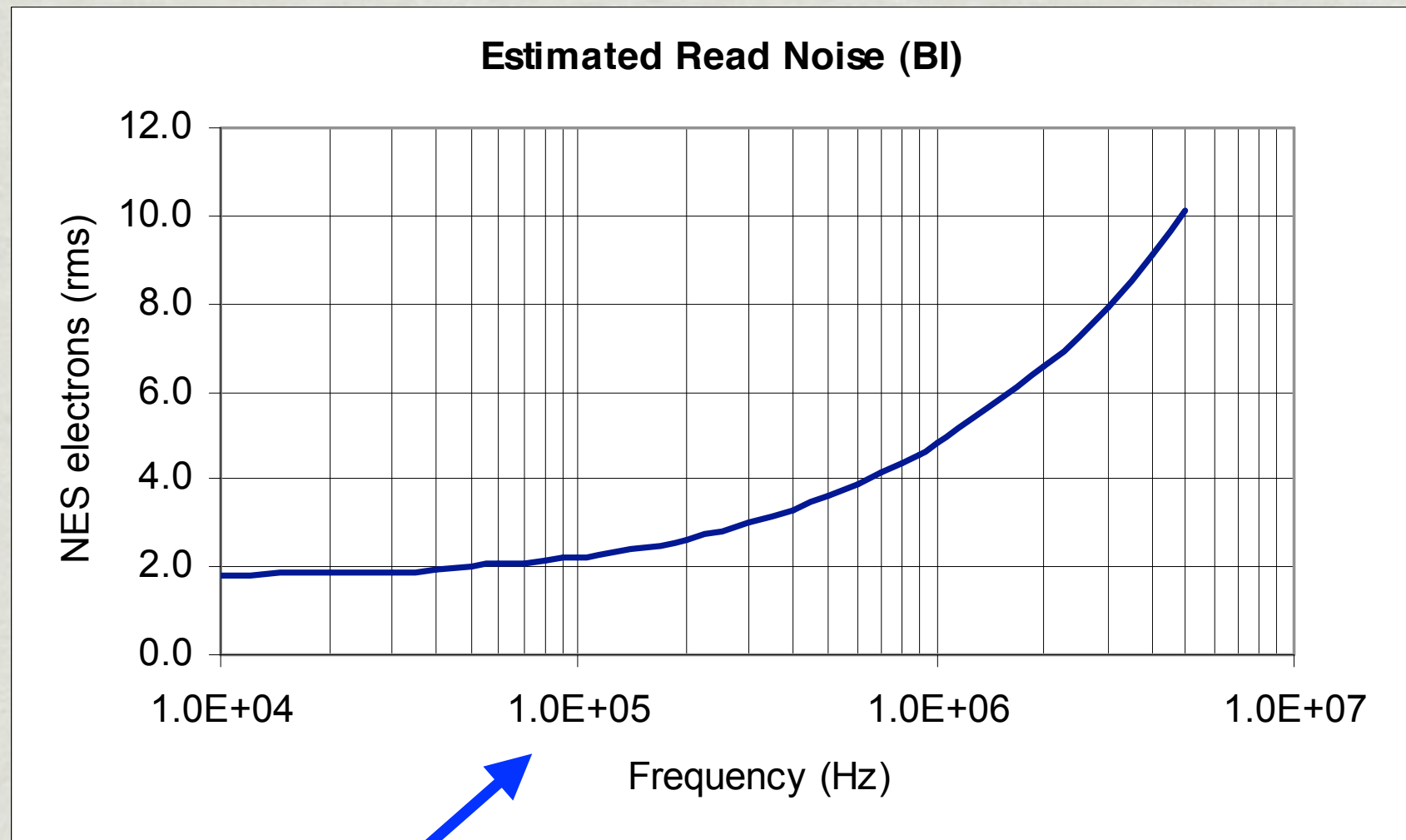
- smaller  $A_{\Omega}$  using adaptive optics
  - cool and baffle your instrument, go to space

Improve efficiency: collect more of the photons you get from the telescope

- anti-reflection coatings, better materials

Increase exposure time, take more exposures!

# CCD Readout, Double Correlated Sampling



← readout time increases  
(lower pixel rate)

100 kHz = 1e5 Hz: time to read out 2048 x 2048 pixels = 42 seconds  
(or 1 4k x 4k CCD with 4 amplifiers)

One way to lower readnoise: average over current fluctuations in the output amplifier (and any other random noise)

Penalty: increased time to read each pixel to have time to average



# Dark Current

# of electrons in conduction band set by thermal distribution,  $\sim \exp \frac{-E_{gap}}{kT}$

Density of states: at T=0 all electrons in states below the Fermi energy  
no electrons at higher energy

At larger temperatures, e- start to occupy nearby higher energy states  
At high enough temperatures, can get into conduction band

Density of e- in conduction band of Si at 300K:  $1.38 \times 10^{13} \text{ cm}^{-3}$

at 77K (conveniently, the temperature of liquid nitrogen):  $1.8 \times 10^{-18} \text{ cm}^{-3}$

This is not quite true for real Si, with impurities, but you get the idea. Cool your detector!  
IR detectors with smaller bandgaps have even bigger sensitivity to thermal e  
-> go colder!